Improving the identification of high-z Herschel sources with position priors and optical/NIR and FIR/mm photometric redshifts*


(Affiliations are available in the online edition)

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1. Introduction

Based on the IRAS, ISO, and Spitzer missions, we know that luminous infrared galaxies (LIRGs) experienced significant evolution in the past 8 Gyr, roughly decreasing their typical infrared luminosity $L(\text{IR})$ by an order of magnitude from $z \sim 1$ to $z = 0$ (Flores et al. 1999; Chary & Elbaz 2001; Pérez-González et al. 2005; Le Floc’h et al. 2005; Rodighiero et al. 2010; Magnelli et al. 2009). By $z \sim 1$, they dominated the star formation activity of the Universe, being responsible for more than $50\%$ of the cosmic SFR density up to at least $z \sim 3$, and playing an important role in the formation of massive galaxies at $z \sim 1.5$, which exhibit high star formation efficiencies (Pérez-González et al. 2008; Reddy et al. 2008; Damen et al. 2009), strong dust obscured nuclear activity (Daddi et al. 2007; Donley et al. 2008), or both. Beyond $z \sim 1$, sub-mm surveys have also detected a population of LIRGs and (mostly) ULIRGs, identified as very active stages in the formation of massive galaxies (Smail et al. 1997; Hughes et al. 1998; Chapman et al. 2003, 2005; Pope et al. 2006).

Although very useful for measuring the obscuration of optical light by interstellar dust (very abundant in the most extreme starbursts), the analysis of IR data to understand the formation and evolution of galaxies in the Universe has been classically hampered mainly by two issues. First, the sensitivity of IR telescopes has increased significantly from one mission to the next (and also their angular resolution, although at a lower rate), but there are still problems correctly identifying all sources detected in IR surveys, most noticeably at longer wavelengths where confusion remains severe (e.g., at 70/160 $\mu$m in the Spitzer/MIPS surveys and in the SCUBA or LABOCA sub-mm data). Moreover, finding counterparts and getting reliable redshifts from spectroscopy for high redshift (U)LIRGs is challenging since a good fraction of them are extremely faint at optical/NIR wavelengths (Ivison et al. 2002; Wilson et al. 2008; Frayer et al. 2004). A second problem for the study of $0 < z \lesssim 4$ (U)LIRGs is the availability of only one photometric data point (MIPS 24 $\mu$m) for most of them. The fits to dust emission models to calculate $L(\text{IR})$, and from this derive a SFR (Kennicutt 1998), need a considerable extrapolation, and thus the SFR is subject to significant systematic uncertainties (Papovich et al. 2006; Daddi et al. 2007; Barro et al. 2010). In addition, the limited FIR data already available for galaxies at $z \gtrsim 2$ indicate that they have different dust properties from those observed for nearby (U)LIRGs (Chapman et al. 2005; Papovich et al. 2007; Rigby et al. 2008; Rieke et al. 2009).

The ESA Herschel Space Observatory (Pilbratt et al. 2010) alleviates these two issues. Thanks to its sensitivity and angular resolution at $100 < \lambda < 500 \mu$m, it enables us to robustly characterize high-$z$ galaxies. In addition to the populations of IR galaxies newly discovered by Herschel, it will supply up to six photometric data points in the IR range for $0 < z < 4$ galaxies detected by MIPS at 24 $\mu$m. These data can be used to constrain the fits to the dust emission templates. Even more novel and relevant is the combination of Herschel observations with the MIPS

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* Herschel is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.
data (and also IRAC) to diminish the effects of source confusion in IR surveys. The new Herschel data indeed allow an easier and more reliable identification of counterparts from one band to the adjacent one as we move to shorter wavelengths. In addition, the Herschel bands effectively fill the gap between MIPS and (sub-)mm surveys (with SCUBA2, LABOCA, or AzTEC). By combining all these datasets, we can now use the several IR photometric data points to estimate a photometric redshift based on dust emission alone, and compare it with the estimations based on UV-to-NIR data. This can be regarded as an extension of the radio-to-IR photo-z method (Carilli & Yun 1999; Dunne et al. 2000; Rengarajan & Takeuchi 2001; Blain et al. 2003; Aretxaga et al. 2003; Daddi et al. 2009), benefiting from a finer sampling of the emission from dust at different temperatures. The procedure can be used to improve the cross-correlation with optical/NIR galaxy samples, or even to obtain the only redshift measurement possible for very faint (high-z) IR-bright optically-faint sources. These sources may be numerous in the Herschel surveys, because of its impressive sensitivity at wavelengths probing the peak of the dust emission, which can be brighter than the optical emission by orders of magnitude.

In this letter, we perform a preliminary analysis of the Herschel Lensing Survey (HLS, PI Egami) data taken for the Bullet cluster to show the utility of position priors and photometric redshifts obtained from UV-to-NIR and IR-to-mm data to improve the identification of Herschel and sub-mm sources with those detected at shorter wavelengths. A correct identification of IR sources will be extremely relevant to all Herschel cosmological surveys, but is of critical importance for clusters, because of their crowded nature and the possible presence of high-z lensed and distorted sources with very faint optical counterparts that may even only be detectable in the FIR/mm range.

2. Data and catalogs

The reduction of the Herschel SPIRE (Griffin et al. 2010) and PACS (Poglitsch et al. 2010) data of the Bullet cluster (a ∼ 06:58, δ ∼ −55:57) is described in Egami et al. (2010) and Rex et al. (2010), jointly with optical-to-mm ancillary data used to estimate photometric redshifts, stellar masses, and SFRs. Here we describe our cataloging procedure for the five Herschel bands: PACS 100 and 160 μm, and SPIRE 250, 350, and 500 μm.

Given the remarkable depth of the Spitzer MIPS 24 μm images, PACS catalogs have been compiled with a position prior technique, involving a list of MIPS sources and a PSF fitting analysis (see Pérez-González et al. 2008). To improve the reliability of the method, we first aligned all Herschel images to the WCS of the MIPS 24 μm map using the wstools software (Mink 1999), obtaining matching uncertainties smaller than one pixel for each Herschel band. The MIPS 24 μm catalog in the Bullet cluster field has a 5σ detection level of $S_{24} = 85 \pm 30 \mu Jy^1$, and a density of 5.3 sources/arcmin$^2$ above that threshold (7.2 sources/arcmin$^2$ with $S_{24} > 50 \mu Jy$).

The similar PSF sizes of the MIPS 24 μm and the PACS 100 μm bands (∼6′′ and ∼8′′, respectively) make the identification and extraction of sources straightforward for this band. In addition, we also considered that some sources may be too faint for MIPS yet still be detectable by PACS. To account for this population, we detected sources directly in the PACS data using SExtractor (Bertin & Arnouts 1996), obtaining a sample of 43 sources above the 5σ level ($S_{100} = 5.5 \pm 0.7 \mu Jy$).

1 Calculated from several artificial apertures formed with random sky pixels; see Appendix A in Pérez-González et al. (2008).

The cross-correlation of this sample with the MIPS catalog (using a search radius of 3′′) revealed that all but one of the robust PACS detections are detected at 24 μm, and 5 of them had 2 MIPS counterparts within the search radius. The only non-detection at 24 μm is, in fact, a group of three very faint 100 μm sources within 10′′, which are detected as one single source by SExtractor in the PACS image, but are individual sources at 24 μm. The 5 multiple detections were sources for which the slightly higher resolution and greater depth of MIPS allowed the deblending of close IR emitters. This cross-correlation also found that the WCS alignment between both images has an rms smaller than 1′′.

For the 160 μm image, it is expected that more 24 μm sources are merged to form single sources. For this reason, before measuring photometry, we removed sources from the main catalog that could not be separated at the resolution of the PACS 160 μm band, i.e., sources closer than 6′′, which is half of the PSF FWHM, keeping only the brightest source at 24 μm for each group of merged sources. The master catalog combined the MIPS source list with the direct detections in the 160 μm image above the 5σ level ($S_{160} [100] = 10 \pm 1 \mu Jy$), all the latter being actually detected in the MIPS data. Merged sources account for almost 20% of the master catalog, with 85% (15%) of this sub-sample having 2 (3) objects within 6′′. Just considering sources directly detected at 160 μm, the fraction of them presenting multiple identifications in the MIPS catalog is higher than 40%. The WCS accuracy between the MIPS and PACS160 images is inferior to 1.5′′.

The final purged list of sources was used to fit PSFs at the given positions using the daophot package in IRAF, allowing for one pixel centering offsets. The results of this PSF fitting algorithm are flux densities in a given aperture (we found the optimum values to be 5′′ and 7.5′′ for the PACS green and red channels, respectively), to which an aperture correction must be applied. The values we applied are 1.84 ± 0.08 and 1.93 ± 0.08 for 100 and 160 μm, respectively. We note that these aperture corrections are larger than those published by the PACS Team, but we found that their PSF matched ours (constructed with the same Bullet cluster data) for large radii, but was brighter for similar and smaller radii than that selected for the PSF fitting method.

The SPIRE catalogs were also compiled with a method based on priors and PSF fitting. The photometry was carried out in circular apertures (of radii 12′′, 18′′, and 22′′ for 250, 350, and 500 μm, respectively), applying a calibration based on the beam sizes (assuming they are Gaussians). For the 250 μm channel, we combined the list of sources detected by MIPS at 24 μm and the direct detection in the SPIRE data using SExtractor and a 5σ threshold ($S_{[250]} = 12 \pm 2 \mu Jy$). Other objects within 9″ (half of the PSF FWHM) of a given candidate source were purged, only the position of the brightest galaxy at 24 μm being kept. For the region covered simultaneously by MIPS and the SPIRE 250 μm channel, only 2 SPIRE sources out of 261 were not detected by MIPS, but both of them turned to be $F(24) \sim 60 \mu Jy$ detections at the 3σ level. Around half of the SPIRE sources have 2 or more MIPS counterparts within 9″. The WCS alignment between the MIPS and SPIRE250 images has an rms of ∼4″.

The catalog for the SPIRE 350 μm channel was constructed from the list of detections at 250 μm. Sources within 14″ were merged into the same object (about 40% of the total catalog, virtually all of them being a doublet), only the position of the brightest one at 250 μm being kept. A direct detection in the 350 μm image ($S_{[350]} = 17 \pm 3 \mu Jy$) found that 5% of sources are not directly linked to a 250 μm emitter, but virtually all of...
Table 1. Properties of the HLS Bullet cluster catalogs.

<table>
<thead>
<tr>
<th></th>
<th>PACS100</th>
<th>PACS160</th>
<th>SPIRE250</th>
<th>SPIRE350</th>
<th>SPIRE500</th>
</tr>
</thead>
<tbody>
<tr>
<td>A''</td>
<td>28.5</td>
<td>30.4</td>
<td>434</td>
<td>438</td>
<td>416</td>
</tr>
<tr>
<td>S (μm)a</td>
<td>5.5 (3.0)</td>
<td>10 (6)</td>
<td>12 (7)</td>
<td>17 (10)</td>
<td>18 (11)</td>
</tr>
<tr>
<td>D''</td>
<td>2.0 (5.0)</td>
<td>2.0 (4.2)</td>
<td>1.1 (1.8)</td>
<td>0.5 (1.2)</td>
<td>0.2 (0.6)</td>
</tr>
<tr>
<td>Md</td>
<td>12% (24)</td>
<td>19% (24)</td>
<td>48% (24)</td>
<td>42% (250)</td>
<td>48% (350)</td>
</tr>
</tbody>
</table>

Notes. (a) Surveyed area in arcmin²; (b) 5σ (3σ) flux density detection levels in mJy; (c) source densities for each significance level in sources/arcmin²; (d) fraction of merged sources and wavelength (in μm) used to obtain position priors.

Table 1 indicates the surveyed areas, detection thresholds, the corresponding source densities, and fraction of merged sources for our catalogs of the PACS and SPIRE HLS data. An important (and potentially dangerous) step in our photometric procedure based on position priors is the blending of nearby non-resolved sources, which provides only the position of the brightest galaxy in the parent catalog to the PSF fitting algorithm. We note that the MIPS 24 μm flux is dominated by the emission of warm dust and/or PAHs while the SPIRE bands are dominated by the emission of cold dust. Thus, it is probable that bright IR emitters are not the dominant sources at FIR wavelengths (see Fig. 1). In the following section, we illustrate a method to help in the robust identification and/or deblending of the Herschel sources.

3. Identification of high-z Herschel sources

Figure 1 shows thumbnails in several bands of two interesting sources detected by Herschel in the Bullet cluster field: the southern tip of the arc at z = 3.54 described in Mehlert et al. (2001), and a sub-mm galaxy detected by LABOCA (source #10 in Johansson et al. 2010). With these two examples, we demonstrate that the use of IRAC and MIPS data helps to correctly identify the Herschel sources with optical/NIR counterparts. In the case of the arc, the photometric method based on the use of the MIPS 24 μm catalog as a prior helped in the detection and measurement processes for the PACS 160 μm and SPIRE bands, where the source is very faint and almost indistinguishable from the background. In the PACS red band, our method allows a deblending of the source in the two knots detected by MIPS and linked to the arc. In the case of the LABOCA source, although it is barely detected in the optical, the position of the source can be followed as we move to redder wavelengths and larger PSFs, allowing a more reliable identification in spite of the presence of a close neighbor to the west. The brightest source at 24 μm is indeed already brighter than its companion at IRAC wavelengths (most noticeably, at 8.0 μm), and this nearby source disappears in PACS. This illustrates the benefits of the use of Spitzer data as priors to obtain the most probable identifications of Herschel and (sub-)mm emitters at shorter wavelengths.

Figure 2 shows the optical-to-mm SEDs of the two examples of high-z galaxies. Photometric redshifts for all 24 μm sources in the Bullet cluster field were estimated from the UV-to-MIR data (including IRAC fluxes) using the technique described in Pérez-González et al. (2005, 2008). Briefly, the SEDs compiled from aperture-matched photometry are fitted with stellar population models and AGN templates from Polletta et al. (2007). A minimization algorithm is used to search for the optimal redshifted template fitting the data, and the most probable
emission templates from Chary & Elbaz (2001), Dale & Helou (2002), and Rieke et al. (2009). Using these two fits, we estimated relevant parameters such as the stellar mass and the SFR. The inset plots depict the photometric redshift estimation and a SFR value (given in the plots with no correction for magnification factor). The general validity of our method for the whole sample of Herschel sources will be tested thoroughly in forthcoming papers by evaluating the quality of our photo-z’s. From a preliminary comparison of the optical and FIR photo-z’s for sources detected by MIPS, PACS, SPIRE, and a (sub)-mm instrument, we are encouraging that we measure an average $\Delta z/(1+z) = \eta = 0.09$ with 10% of catastrophic outliers ($\eta > 0.2$), comparable to the typical goodness of IR-based photo-z’s (Aretxaga et al. 2003).

Fig. 2. SEDs of the $z = 3.24$ LIRG MIPS06:58:31.1−55:56:39.2 (magnification factor $\sim 21$) and the $z = 2.5$ HyLIRG MIPS06:58:45.3−55:58:46.5, detected in the MIR/FIR by Spitzer and Herschel, and the (sub)-mm by LABOCA and AzTEC. The blue line shows the best-fit stellar population model for the photometry up to 8 $\mu$m, which provides an estimate of the photometric redshift and the stellar mass. The cyan, green, and red lines show the best-fit dust emission model for the MIR-to-mm photometry (from Chary & Elbaz 2001; Dale & Helou 2002; and Rieke et al. 2009), which gives an independent photometric redshift estimation and a SFR value (given in the plots with no correction for magnification factor). The inset plots depict the photometric redshift normalized probability distributions for the optical/NIR and the MIR/mm estimations.

The photometric redshift is calculated by integrating the redshift probability function (see also Barro et al. 2010). Photometric redshifts were also estimated from the IR data alone by comparing the Spitzer, Herschel, and (sub-mm) fluxes with dust emission templates from Chary & Elbaz (2001), Dale & Helou (2002), and Rieke et al. (2009). Using these two fits, we estimated relevant parameters such as the stellar mass and the SFR. The Rieke et al. (2009) template most closely matched the data for the arc corresponds to a local $L(\text{IR}) = 10^{11.5} L_\odot$ galaxy. This value is consistent with the average for the arc based on the three template libraries, $L(\text{IR}) = 10^{11.40\pm0.04} L_\odot$, once a magnification factor of $\sim 21$ is applied (Bradač et al. 2006). The LABOCA source qualifies as a hyper-LIRG based on the best fits to all template libraries (if it is not lensed), and the best-fit template from Rieke et al. is that of a local galaxy with $10^{12.75} L_\odot$, although our source seems to have more prominent PAH emission.

Figure 2 illustrates the power of using independent estimates of the photometric redshift based on UV-to-NIR data and IR data alone to validate the individual values (see the photo-z probability distributions). The procedure is also useful for achieving more reliable identifications of the IR emitters. For example, for the LABOCA source at $z \sim 2.5$, we calculated photometric redshifts based on UV-to-NIR data for the three close neighbors detected by MIPS within $\sim 10''$ of the central source. The galaxy to the north lies at $z = 0.3$ (a cluster member), that to the south at $z = 1.9$, and the closest companion has a photo-z compatible with the central source ($z = 2.3 \pm 0.3$), thus implying that the two sources are a possible interacting pair.