Combination of large nanostructures and complex band structure for high performance thermoelectric lead telluride

Yanzhong Pei, Nicholas A. Heinz, Aaron LaLonde and G. Jeffrey Snyder*

Received 9th June 2011, Accepted 11th July 2011
DOI: 10.1039/c1ee01928g

The complexity of the valence band structure in p-type PbTe has been shown to enable a significant enhancement of the average thermoelectric figure of merit \(zT\) when heavily doped with Na. It has also been shown that when PbTe is nanostructured with large nanometer sized Ag\(_2\)Te precipitates there is an enhancement of \(zT\) due to phonon scattering at the interfaces. The enhancement in \(zT\) resulting from these two mechanisms is of similar magnitude but, in principle, decoupled from one another. This work experimentally demonstrates a successful combination of the complexity in the valence band structure with the addition of nanostructuring to create a high performance thermoelectric material. These effects lead to a high \(zT\) over a wide temperature range with peak \(zT > 1.5\) at \(T > 650\) K in Na-doped PbTe/Ag\(_2\)Te. This high average \(zT\) produces 30% higher efficiency (300–750 K) than pure Na-doped PbTe because of the nanostructures, while the complex valence band structure leads to twice the efficiency as the related \(n\)-type La-doped PbTe/Ag\(_2\)Te without such band structure complexity.

Thermoelectric (TE) applications have attracted increasing interest in the last decade as a means to combat the ever growing rate of energy consumption throughout the world. The two main applications for thermoelectric materials are power generation and devices.

Thermoelectric generators, which directly convert heat into electricity are being considered for industrial or automotive waste heat recovery. However, the relatively low efficiency of thermoelectric materials has limited their use to niche applications. Effective strategies for high performance thermoelectric materials during the past half-century have depended largely on the reduction of heat conduction by the crystal lattice and recently on the proposed band structure engineering for improvements in the electronic properties. The present study shows the effort of combining these two independent research directions, namely reducing the lattice thermal conductivity by scattering phonons through nanostructuring and improving the electronic properties by exploiting complex band structure, as a route to improve the thermoelectric performance. The addition of Ag\(_2\)Te in PbTe creates nanosized scattering centers for heat carrying phonons while further Na-doping enables multiple bands for enhanced electronic properties. In this way, an enhancement of 30%–100% on thermoelectric efficiency is achieved as compared with the same material having either multiple bands or nanosized scattering centers alone, enabling a synergistic route in the search of high performance thermoelectric materials and devices.
can be rationalized to be from the multi valley valence band structure". The second valence band exists slightly below the principal valence band. Because of the complicated band structure, heavy and precise doping of the valence band is required to reach the high \( zT \) doping level.\(^4\)

Because this influence on the valence band structure is purely an electronic effect, there is good potential to further enhance \( zT \) for PbTe:Na by reducing the lattice thermal conductivity.

It has been demonstrated in numerous reports that nanostructures in PbTe effectively scatter the long wavelength heat carrying phonons to lower the lattice thermal conductivity.\(^8,11-12\)

This approach led to a peak \( zT \) of 1.4–1.7 in both p- and n-type PbTe materials.

Attempts to nanostructure p-type PbTe:Tl with TISbTe\(_2\) have not led to higher \( zT \) because of reduced mobility despite lower thermal conductivity.\(^17\)

Among these reports, PbTe/Ag\(_2\)Te nanocomposites prepared from an equilibrium precipitation differ because of the relatively large size of the nanoscale Ag\(_2\)Te precipitates and in addition, the resultant composite is an intrinsic semiconductor due to the self-compensated substitutional and interstitial Ag\(^+\) pairs.\(^16,18,19\)

The lack of extrinsic dopants enables a direct observation of the \( \kappa_L \)-reduction due to the embedded nanostructures because there is very little contribution to \( \kappa \) due to the electrons. Moreover, the large Ag\(_2\)Te nanostructures were found to reduce the \( \kappa_L \) of PbTe at high temperatures with an equal or greater effect than those of smaller nanosized particles,\(^16\) while the small nanoparticles reduce the thermal conductivity better than the large particles at low temperatures.

The PbTe matrix in the PbTe/Ag\(_2\)Te nanocomposites can be doped n-type with La\(^+\) or excess Ag\(^+\) which allows as an independent control of carrier density while maintaining the composite microstructure.

Additionally, introducing embedded Ag\(_2\)Te nanoparticles into the matrix may mechanically strengthen these brittle materials for practical applications.

The present work shows our effort toward Na-doping PbTe/Ag\(_2\)Te nanocomposites and the resulting thermoelectric properties. The low lattice thermal conductivity due to the Ag\(_2\)Te nanoinclusions and the excellent electronic properties due to the complex valence band structure is combined to result in a \( zT \) of 1.5 at high temperatures. Moreover, there are significant improvements in the average \( zT \) and thermoelectric efficiency throughout the entire temperature range studied, as compared with analogous materials without nanostructures or band structure complexity, or with small nanostructures.

P-type nanocomposites were obtained by Na-doping (PbTe:Na/Ag\(_2\)Te). PbTe/Ag\(_2\)Te nanocomposites with a composition of (PbTe)\(_{0.945}\)(Ag\(_2\)Te)\(_{0.055}\) (Fig.1) were pre-synthesized according to the reported method\(^16\) and subsequently used as the starting material for making PbTe:Na/Ag\(_2\)Te together with appropriate amounts of Na and Te metals. The nominal concentration of Na \([\text{Na}_{x}\text{PbTe\(_{1-x}\)}\text{Ag\(_2\)Te}]_{0.055}\) is \( x = 0-3 \) at.% and the samples for this study were synthesized with the same method used previously;\(^16\) including sealing, melting (Step 1 in Fig. 1), quenching, annealing (Step 2 followed by Step 3) and hot pressing. Phase components were checked using X-ray diffraction and scanning electron microscopy (SEM) equipped with an energy dispersive spectrometer (EDS).

Fig. 1 Phase diagram of PbTe-Ag\(_2\)Te system and the procedure used for material preparation. A low-temperature annealing leads to the precipitation of Ag\(_2\)Te from the high temperature supersaturated alloy phase.

Hot pressed disk-shape samples with relative densities of 98% or higher were used for the measurements. Details on measuring the transport properties were given elsewhere,\(^4,16\) where the heat capacity \( C_p \) in \( \text{K}_B \text{ per atom} = 3.07 + 4.7 \times 10^{-4} \times (T/\text{K}-300)\)\(^{17,19,21}\) was determined by drop calorimetry and believed to be accurate for several lead chalcogenides.\(^4,22,23\) The thermal conductivity for most of the recently reported high \( zT \) PbTe materials at 600–800 K, was determined using a heat capacity of \( C_p \) or close to, the Dulong-Petit approximation by \( \pm 5\% \).\(^4,10,14,24-26\)

It should be noted that this equation determines a value ~10% higher than the Dulong-Petit law (3\( k_B \text{ per atom} \) at \( T > 700 \text{ K} \)), and this method is used to recalculate all the literature data used for comparison in this study. The uncertainty for each measurement of \( S \), \( \sigma \) and \( \kappa \) is ~5%, resulting in a combined error of ~20% in the determined \( zT \) values.

Annealing the high temperature supersaturated solid solution phase (Step 2) of PbTe/Ag\(_2\)Te\(^27,28\) in the low temperature, two-phase region (Step 3) creates homogenously distributed Ag\(_2\)Te nanoinclusions in a PbTe matrix.\(^16\) Fig. 2a shows the typical nanostructure for the obtained as-cast ingots. The plate-like Ag\(_2\)Te precipitates have preferred orientation parallel with \(<001>\) directions in the PbTe matrix, consistent with previous reports.\(^16,19,29\) The Ag\(_2\)Te precipitates are nanometer sized in one direction but micron-size in the other directions, which is larger than those generally observed in PbTe/AgSbTe\(_2\) and its analogs.\(^5,30\) No Ag\(_2\)Te precipitates smaller than ~50 nm could be found in PbTe/Ag\(_2\)Te system with similar synthesis conditions.\(^16\)

Similar to PbTe:Na,\(^4\) Na is found to be an effective p-type dopant in PbTe/Ag\(_2\)Te as indicated by the Hall coefficient (\( R_H \)) and Seebeck coefficient measurements. However, the doping solubility of Na in PbTe/Ag\(_2\)Te is found to be much smaller than in pure PbTe, since the maximum measured room temperature Hall density (\( \rho_H = 1/eR_H \), \( e \) is the electron charge) in PbTe:Na/Ag\(_2\)Te is much smaller than that in PbTe:Na, this is presumably due to the presence of a small amount of soluble Ag\(_2\)Te (~1)%\(^4\) in PbTe matrix reducing the solubility of Na. Even though the nominal concentrations of Na are comparable in PbTe:Na/Ag\(_2\)Te and PbTe:Na,\(^4\) the measured \( \rho_H \) in PbTe:Na/Ag\(_2\)Te does
not exceed $4 \times 10^{19}$ cm$^{-3}$, while $p_H$ in PbTe:Na can be as high as $14 \times 10^{19}$ cm$^{-3}$. Since it was found that the optimal $p_H$ was close to $10^{20}$ cm$^{-3}$, due to the dominance of the heavy band behavior in p-PbTe, the most heavily doped PbTe:Na/Ag$_2$Te samples with nominal Na content of 1%, 2% and 3% having $p_H$ of 2.5, 3.1 and $3.7 \times 10^{19}$ cm$^{-3}$ are used for the following discussions and marked as 2.5e19, 3.1e19 and 3.7e19, respectively.

As has been well established in p-PbTe, the room temperature Seebeck coefficient versus Hall density (Pisarenko plot, solid line in Fig. 3a) provides a powerful characterization of the band structure for PbTe. The majority of published measurements on bulk p-type PbTe fall on this Pisarenko line, as do the present results, within the measurement uncertainty (Fig. 3a). Similar to the findings in La-doped PbTe/Ag$_2$Te (PbTe:La/Ag$_2$Te), neither doping with Na nor introducing Ag$_2$Te alters the band structure of the PbTe matrix.

The flattening of $S$ at high doping levels ($p_H > 3 \times 10^{19}$ cm$^{-3}$) can be well understood by the complex valence band structure shown in the inset of Fig. 3a. Furthermore, the energy of the light valence band reduces with increasing temperature and moves below the heavy valence band at $\sim 450$ K, making the heavy valence band the dominant band for hole transport at high temperatures where the $zT$ peaks. The electronic transport, optical spectroscopy and other properties of p-PbTe can be well described by this two-valence-band mode.

Rather than the Seebeck coefficient being proportional to absolute temperature (e.g. n-type PbTe:La/Ag$_2$Te in Fig. 3b), a significant increase in $S$ (Fig. 3b) is observed, particularly at high temperatures, due to the increasing contribution of the heavy band carriers resulting from the temperature dependent band structure and Fermi broadening. The analogous n-type material (PbTe:La/Ag$_2$Te) has a much lower Seebeck coefficient than p-type PbTe:Na/Ag$_2$Te, due to the lack of conduction band complexity. These PbTe:Na/Ag$_2$Te materials show a roughly unchanged Seebeck coefficient with respect to PbTe: Na (without Ag$_2$Te) at similar doping levels, which further supports the above discussion ($S$ vs. $p_H$) that Ag$_2$Te inclusions do not affect the band structure of PbTe.

While the Seebeck coefficient is similar, the PbTe:Na/Ag$_2$Te composites have lower charge carrier mobility compared with PbTe:Na, due to both the increased scattering of carriers at the phase boundaries and the increased point defect scattering in the
The reduced electrical resistivity of PbTe:Na/Ag$_2$Te nanocomposites is higher than that of PbTe:Na (Fig. 4) at low temperatures, when the doping levels are nearly the same ($\rho_H = 3.6 \times 10^{19}$ cm$^{-3}$ for PbTe:Na vs. $\rho_H = 3.7 \times 10^{19}$ cm$^{-3}$ for PbTe:Na/Ag$_2$Te). Similar to PbTe:Na$^4$ (and n-PbTe$^{31,32}$), the electrical resistivity ($\rho$) rises faster than normally expected for a system dominated by acoustic scattering ($\rho \propto T^{-1.5}$). This can be explained by the increasing effective mass of carriers because of the temperature dependent mass of the light bands as well as the transition of holes from the light to the heavy band that has lower mobility.$^{13,31–33}$ Increasing Hall density results in lower electrical resistivity in PbTe:Na/Ag$_2$Te nanocomposites, as shown in Fig. 4.

The thermal conductivity ($\kappa$) is reduced by $\approx 50\%$ in roughly the entire measured temperature range due to the Ag$_2$Te nanoinclusions (Fig. 5). It is important to note that the Hall density of all the samples was similar and found to be $\sim 3.5 \times 10^{19}$ cm$^{-3}$. To avoid the ambiguity of bipolar thermal conductivity and for the sake of clarity, only heavily doped samples are shown in Fig. 5.

The observed $\kappa$-reduction is only partially attributed to the presence of the nanoinclusions. The reduced electrical conductivity and thus a reduced electronic component to the thermal conductivity ($\kappa_E$) also contributes to lower $\kappa$. It is difficult to accurately estimate the $\kappa_E$ via the Wiedemann-Franz law ($\kappa_E = L_oT$) because of the difficulty in determining the Lorenz number ($L$) in p-PbTe due to the complex valence band structure and the band non-parabolicity.$^{31,32,39,40}$ For simplicity, an estimation of $L$ is made using a single parabolic band (SPB) model assuming an acoustic phonon scattering mechanism,$^{41}$ which results in an $L$ with a deviation of less than $10\%$ when compared to a more rigorous single non-parabolic band and multiple band model calculation.$^{7,16,39,42}$ And the same method of $L$-estimation is used to recalculate the literature data for the following discussion.

By subtracting the electronic component from the total thermal conductivity, the lattice component ($\kappa_L$) is calculated and shown in Fig. 5. It can be expected that the introduction of nanoinclusions effectively reduces the lattice thermal conductivity due to the enhanced scattering of phonons at boundaries.$^{8,9,16,19}$ The extremely low $\kappa_L$ of $\approx 0.5$ W/m-K at $T > 600$ K is approaching the theoretical minimum value of 0.36 W/m-K.$^{16,43}$ It should be noted that the total thermal conductivity of PbTe:Na/Ag$_2$Te composites is significantly lower than the lattice thermal conductivity of PbTe:Na$^4$. As compared with other PbTe nanocomposites having smaller structure features, such as PbTe:Na/SrTe$_2$ and PbTe/NaSbTe$_2$ (SALT) and their analogs, where the $\kappa_L$ is now recalculated with the same estimation of both $C_p$ and $L$, PbTe with relatively large Ag$_2$Te precipitates has even lower lattice thermal conductivity at high temperatures, similar to that observed in n-type materials.$^{16}$ This is presumably due to at least partial precipitate dissolution of small nanoparticles observed at high temperatures,$^{44}$ while the large Ag$_2$Te precipitates in our current and previous study$^{16,19}$ are obtained by annealing at high temperature (Fig. 1) and thus are likely to be more stable.

When compared with n-type L$_3$-doping$^{16}$ (using the same estimation of $C_p$), p-type Na-doping in PbTe/Ag$_2$Te nanocomposites has significantly higher $\varepsilon T$ (Fig. 6a) at low temperatures because of the favorable electronic effect of the complex band structure.$^{45}$ Alternatively, compared with PbTe:Na having similar $\rho_H$ of $3.6 \times 10^{19}$ cm$^{-3}$, introducing Ag$_2$Te nanoinclusions significantly reduces the thermal conductivity (Fig. 5) and thus increases $\varepsilon T$ (Fig. 6a). This is further evidenced by the excellent agreement between the measured $\varepsilon T$ of the PbTe:Na/Ag$_2$Te sample with $\rho_H = 3.7 \times 10^{19}$ cm$^{-3}$ and the projected $\varepsilon T$ curve for PbTe:Na/Ag$_2$Te.
PbTe:Na with \( p_H = 3.6 \times 10^{19} \text{ cm}^{-3} \) (dashed orange), which is generated by simply assuming the \( \kappa_L \) value to be that of PbTe:Na/Ag\(_2\)Te with \( p_H = 3.7 \times 10^{19} \text{ cm}^{-3} \). However, it is unlikely that these materials are fully optimized as the optimal carrier density, as observed in PbTe:Na with \( p_H = 9 \times 10^{19} \text{ cm}^{-3} \) (cyan curve),\(^4\) was not obtained in this study due to the small dopant solubility.

The current effort of combining both complex band structure and nanostructures enables a \( zT \) higher than 1.5 at \( T > \sim 650 \text{ K} \). More importantly, as shown in Fig. 6b, both the average \( zT \) (\( zT_{\text{ave}} \)) and the device \( zT \) that is calculated from the theoretically available power generation efficiency (\( \eta \))\(^{45-48} \) (between 300 K and 750 K), are found to be enhanced by 40–100\% in PbTe:Na/Ag\(_2\)Te, when compared to PbTe:La/Ag\(_2\)Te\(^6\) and PbTe:Na\(^4\) having similar \( p_H \). Here, the estimations of device \( zT \) also take into account the thermoelectric compatibility effect.\(^{45-48} \) It should also be note that the \( zT_{\text{ave}} \) enables a fairly accurate evaluation of the relative change on the efficiency\(^{48} \) even though real device \( zT \) is lower than \( zT_{\text{ave}} \) due to the variable compatibility factor.\(^{46-48} \) Even though PbTe:Na/Ag\(_2\)Te composite is not fully optimized for carrier concentration, 15–20\% higher \( zT \) and \( zT_{\text{ave}} \) have already been achieved as compared with the best PbTe:Na (\( p_H = 9 \times 10^{19} \text{ cm}^{-3} \))\(^6\) and PbTe:Na/SrTe\(^9\) having smaller and endotaxial nanostructures, indicating the effectiveness of nanostructures for high performance thermoelectrics.

In summary, band structure complexity and nanostructure effects are simultaneously considered as an effective approach for improving thermoelectric performance. These combined approaches have been demonstrated in PbTe:Na/Ag\(_2\)Te resulting in a peak \( zT \) higher than 1.5, and more importantly, this results in significant enhancements of the average \( zT \) and thermoelectric efficiency. Further study of the combination of carrier density and nanostructure control should result in even higher thermoelectric performance in similar PbTe materials.

The authors gratefully acknowledge NASA-JPL and DARPA for funding support of this work.

**References**