Measurement of Intrinsic Dirac Fermion Cooling on the Surface of the Topological Insulator Bi$_2$Se$_3$ Using Time-Resolved and Angle-Resolved Photoemission Spectroscopy


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We perform time- and angle-resolved photoemission spectroscopy of a prototypical topological insulator (TI) Bi$_2$Se$_3$ to study the ultrafast dynamics of surface and bulk electrons after photoexcitation. By analyzing the evolution of surface states and bulk band spectra, we obtain their electronic temperature and chemical potential relaxation dynamics separately. These dynamics reveal strong phonon-assisted surface-bulk coupling at high lattice temperature and total suppression of inelastic scattering between the surface and the bulk at low lattice temperature. In this low temperature regime, the unique cooling of Dirac fermions in TI by acoustic phonons is manifested through a power law dependence on carrier density.

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Electrons on the surface of a three-dimensional topological insulator (TI) are massless Dirac fermions with linear energy-momentum dispersion [1–3]. These electronic systems are promising for novel applications ranging from spin-based field effect transistors [4], ultrafast opto-spintronic devices [5] and high-speed quantum information processors [6], whose performance depends critically on the dynamics of hot carriers. Unlike the case in graphene [7], helical Dirac fermions in a TI interact not only with phonons but also with an underlying bulk reservoir of electrons [3]. Therefore, it is important to understand separately their coupling mechanisms to each of these degrees of freedom. However, high frequency optical conductivity [8] and ultrafast optical experiments [9,10] do not directly separate bulk and surface signals nor do they distinguish different relaxation channels.

Time- and angle-resolved photoemission spectroscopy (TrARPES) is a powerful technique to study carrier dynamics with energy and momentum resolution. A recent TrARPES study has shown that persistent surface carrier dynamics with energy and momentum resolution. A recent TrARPES is a powerful technique to study carrier dynamics following photoexcitation over a range of lattice temperature and total suppression of inelastic scattering between the surface and the bulk at low lattice temperature. In this low temperature regime, the unique cooling of Dirac fermions in TI by acoustic phonons is manifested through a power law dependence on the surface temperature decay rate on carrier density.

Figure 1: (a) Raw ARPES spectra taken immediately after the pump excitation. The ARPES spectra taken immediately after the pump excitation is consistent with prior studies [16]. However, it can be seen from the difference of momentum-integrated spectra that the equilibrium has not been fully reached within the measurement time window [Fig. 1(e)]. The ARPS spectra taken immediately after the pump excitation [Fig. 1(b)] resemble a thermalized electron distribution at elevated temperature for both SS and CB. This hot distribution cools down progressively at longer time delays [Figs. 1(c) and 1(d)]. However, it can be seen from the difference of momentum-integrated spectra $\Delta I(E,t)=I(E,t)-I(E,t=5\text{ps})$ that the equilibrium has not been fully reached within the measurement time window [Fig. 1(e)].

In order to quantitatively describe the relaxation dynamics of the hot bulk and surface populations in Fig. 1, we separately analyze the spectral intensities for the surface and the bulk. Figure 2(a) shows $I(E,t)$ for SS and CB at $t=0.25\text{ ps}$ obtained in their respective energy-momentum region [Fig. 1(a)]. We fit $I(E,t)$ for both SS and CB with the following equation [17] to extract their respective electronic temperatures $T_e(t)$ and chemical potentials $\mu(t)$ as a function of time:

$$I(E,t) = A(t) \int_{-\infty}^{\infty} dE [f_{FD}(E,T_e(t),\mu(t))D(E)]G(E-\epsilon,w(t)).$$

(1)
Here, $A(t)$ is a scaling factor [17], $f_{FD}$ is the Fermi-Dirac distribution, $G(E, w)$ is a Gaussian function with energy independent width $w$ that includes the instrumental resolution and time-dependent broadening due to increased scattering rate after photoexcitation [18]. The time independent $D_t$, which is a product of density of states (DOS) and photoemission matrix element [19], is obtained through a global best fit for all $t$ [12]. Both the SS and CB fittings agree very well with data for all $t$ [Fig. 2(b)] (also see [12]). The fact that both SS and CB populations can be described by Fermi-Dirac distributions at $t = 0.25$ ps [Fig. 2(a)] suggests that intraband thermalization of SS and CB electrons is established within the instrumental resolution time ($\sim 200$ fs).

Figure 3 shows the obtained electronic temperatures and chemical potentials for both the SS and CB. We focus on the temperature first. As shown in Figs. 3(a), SS (CB) temperature $T_{CR}^{SS}$ ($T_{CR}^{CB}$) starts at the lattice temperature $T_L = 15$ K at $t = -5$ ps. The maximum temperature increase of $\sim 1100$ K at $t = 0$ ps allows us to estimate an electron specific-heat coefficient $\gamma$ of $2$ mJ/mol·K$^2$ (assuming 50 nm penetration depth of the pump pulse [11]). The relaxation time agrees with previous optical studies [10].

The analysis of the time dependence of the chemical potential reveals carrier population dynamics both within and between the bands. In general, chemical potential is a function of both temperature and particle number $N$ [20]. In the case where $N$ is fixed, we derive [12] a relationship between chemical potential and our measured temperature assuming linear and parabolic dispersion for the SS and CB bands [20] respectively. The SS and CB chemical potentials derived this way ($\mu_{N}^{SS}$ and $\mu_{N}^{CB}$) are plotted in Fig. 3(b) along with the chemical potentials directly obtained by fitting our data to Eq. (1) ($\mu^{SS}$ and $\mu^{CB}$). Remarkably, a good overlap between $\mu^{SS}$ and $\mu_{N}^{SS}$ as a function of $t$ is observed at $T_L = 15$ K [Figs. 3(b) green] indicating constant SS particle number $N^{SS}$. $\mu^{CB}$ is significantly higher than $\mu_{N}^{CB}$ [Figs. 3(b) red], suggesting instantaneous $N^{CB}$ increase upon photoexcitation. The increasing $N^{CB}$ shows that photoexcitation from deeper valence bands primarily populates CB, consistent with a related study that shows bulk dominant direct optical transition [11]. The fact that $N^{SS}$ stays constant indicates that there is no net particle transfer from CB to SS.

In order to understand whether the transfer of carriers between surface and bulk can happen at higher lattice temperature, we now investigate the chemical potentials and temperature dynamics at 300 K [Figs. 3(c) and 3(d)]. The most striking feature is that $\mu^{SS}$ is much bigger than $\mu_{N}^{SS}$ [Fig. 3(d)], indicating a strong increase of surface carrier density in contrast to 15 K [Fig. 3(b)]. To find where the extra surface carriers come from, we notice $\mu^{CB}$ is lower at 300 K than at 15 K [Fig. 3(b)–3(d) solid red traces]. By calculating the change of particle numbers after photoexcitation of SS and CB ($\Delta N^{SS}$ and $\Delta N^{CB}$) [12], we see that the increase of $\Delta N^{SS}$ at 300 K has a slower rise time than $\Delta N^{CB}$.

FIG. 1 (color). TrARPES spectra of Bi$_2$Se$_3$. (a)-(d) Energy-momentum spectra sliced through 3D data volume along $\Gamma - M$ direction at -5 ps (a), 0.25 ps (b), 2.5 ps (c) and 15 ps (d) taken at 15 K on an $E_D = 0.28$ eV sample. (e) Momentum-integrated difference spectra subtracting the spectrum at $t = -5$ ps. The blue dashed lines in (a) define the boundary between conduction bulk band (CB) and surface state (SS) for the analysis performed in the text.

FIG. 2 (color). Fitting of momentum-integrated spectra. (a) Momentum-integrated spectra for SS (green) and CB (red) and their respective fits (purple and orange lines) at $t = 0.25$ ps for the data shown in Fig. 1(b). (b) Constant energy slices of the difference spectra as a function of $t$ at 0.1 eV for the same data and fit. The spectra are normalized over the prepump intensity $I_0 = I(t = -5$ ps).
at 15 K [Fig. 3(d) inset]. This shows that SS particle number increase is not due to direct optical transition. Rather it is a result of transfer from CB to SS through phonon-scattering at elevated lattice temperature due to higher phonon-scattering rate above Debye temperature ($\Theta_D = 182$ K [21]) [22]. The evidence for such phonon-assisted CB-SS scattering can also be seen in the temperature dynamics [Figs. 3(a)–3(c)]. $T^SS_e$ and $T^CB_e$ only equilibrate at 300 K within the probed time window, indicating that phonons scatter carriers, transferring energy and mediating the thermalization between SS and CB at 300 K. Since inelastic $e$-$e$ scattering is mainly responsible for ultrafast thermalization [22] within a single band, the lack of thermalization between SS and CB at 15 K [Figs. 3(a)] suggests that inelastic $e$-$e$ scattering time between SS and CB exceeds the measured time window. (See [12] for more $T^SS_e$ vs $T^CB_e$ at 15 K.) Such suppression of interband inelastic $e$-$e$ scattering is likely due to the kinematics constraints [23].

Having shown that interactions between SS and CB are suppressed at 15 K [Figs. 3(a) and 3(b)], we proceed to investigate their respective intraband cooling mechanisms. Because different mechanisms can be distinguished based on carrier density dependence, as demonstrated in other Dirac materials [24], we measure the intraband cooling dynamics of $\text{Bi}_2\text{Se}_3$ over a wide range of dopings [12]. To vary the doping, we used samples from two different batches with high and low carrier concentrations along with surface doping effect [25,26]. The surface doping effect has been commonly observed in topological insulators such as $\text{Bi}_2\text{Se}_3$ [25] and is attributed to either adsorption or moving of Se vacancies towards the surface [26]. Figure 4(a) shows that the surface temperature has two decay components $\tau^SS = 0.7 \pm 0.1$ ps [Fig. 4(b)] matches with a previous ultrafast study, which has been attributed to an optical phonon-mediated intra-band cooling of CB [10]. The slow component $\tau^CB$ exhibits no discernible dependence on $E_D$ [Fig. 4(b)].

We use the density dependence of the slow component to reveal its microscopic origin. We extract $\tau^SS$ through a biexponential fit of $T^SS_e(t)$ [12] and plot it as a function of doping $E_D$ in Fig. 4(b) inset. Such exponential decay of the electronic temperature in a Dirac system has been previously predicted for graphene [24] and also observed in graphite [28] and recently $\text{Bi}_2\text{Se}_3$ [11]. The SS cooling rate scales as $1/\tau^SS \propto E_D^{3+0.5}$, which is in good agreement

![Figure 3](image-url)
FIG. 4 (color). Electron density dependent intraband cooling of surface Dirac fermions. (a) Electronic temperature as a function of $t$ of surface electrons with different $E_D$ at 15 K. Different colors correspond to $E_D$ in matching color in the inset. (b) $T^{CB}$ at 15 K. Inset: Inverse cooling time of the slow component of $T_e$ as a function of $E_D$ at 15 K for SS (green triangles) and CB (red squares). $1/\tau^S_2 \approx E_D^2$ at 15 K.

with the theoretical calculations of cooling of 2D Dirac fermions via acoustic phonons [24]. In stark contrast, $1/\tau^C_2$ exhibits little dependence on doping, which agrees with theoretical predictions that acoustic phonon-mediated carrier cooling in 3D metals is density independent [29]. Consistent with the phonon-assisted surface-bulk coupling presented earlier, the cooling rate at 300 K lies in between that of the SS and CB at 15 K [12]. It is worth mentioning that in the low doping regime where $\mu < k_B T_e$, acoustic phonon cooling of Dirac fermions will be further suppressed and is predicted to have a different decay dynamics [24]. As a result, the doping dependence of the slow decay component may deviate from the power at low doping levels.

In conclusion, we have directly visualized different scattering channels of surface and bulk electrons on a TI with TrARPES. At high temperature, phonons are largely responsible for scattering electrons between the surface states and the conduction bulk band. At low temperature, the coupling between surface and bulk is suppressed. We further reveal the surface cooling rate follows a power law dependence on the carrier density, which is a signature of acoustic phonon-mediated cooling of 2D Dirac fermions.

This suggests that by tuning to the charge neutrality point one can enter a regime where acoustic phonon-scattering of Dirac fermions is completely eliminated and transport properties are determined solely by structural and chemical defects [30,31]. An exciting possibility is the creation of long-lived hot photocarriers that can be used for high-efficiency photothermoelectric applications [5,32].

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