Mechanically tunable optofluidic distributed feedback dye laser

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Abstract: A continuously tunable optofluidic distributed feedback (DFB) dye laser was demonstrated on a monolithic replica molded poly(dimethylsiloxane) (PDMS) chip. The optical feedback was provided by a phase-shifted higher order Bragg grating embedded in the liquid core of a single mode buried channel waveguide. Due to the soft elastomeric nature of PDMS, the laser frequency could be tuned by mechanically stretching the grating period. In principle, the mechanical tuning range is only limited by the gain bandwidth. A tuning range of nearly 60nm was demonstrated from a single dye laser chip by combining two common dye molecules Rhodamine 6G and Rhodamine 101. Single-mode operation was maintained with less than 0.1nm linewidth. Because of the higher order grating, a single laser, when operated with different dye solutions, can provide tunable light output covering the entire spectrum from near UV to near IR in which efficient laser dyes are available. An array of five DFB dye lasers with different grating periods was also demonstrated on a chip. Such tunable integrated laser arrays are expected to become key components in inexpensive advanced spectroscopy chips.

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References and links

1. Introduction

On-chip liquid dye lasers represent promising coherent light sources for ‘lab-on-a-chip’ systems in that they allow the integration of laser sources with other microfluidic and optical devices. Several groups have so far demonstrated such dye lasers by using different materials and laser cavity designs [1,2,3]. Tunable output can also be obtained using concentration or index tuning methods [3,4]. Indeed, on-chip liquid dye lasers are examples of the new class of emerging optofluidic devices, in which the integration of microfluidics with the adaptive nature of liquids enables unique performance that is not obtainable within solid state materials [5]. Recently, an optofluidic distributed feedback (DFB) dye laser was demonstrated on a monolithic poly(dimethylsiloxane) (PDMS) elastomer chip [6]. Stable single-mode operation with narrow linewidth was obtained using a phase-shifted higher order Bragg gating embedded in a single mode microfluidic channel waveguide. The use of PDMS and soft lithography results in inexpensive and easy fabrication, rapid device prototyping, biocompatibility and the opportunity to create disposable devices. In this paper, we demonstrate a tunable single-mode DFB dye laser that combines the mechanical flexibility of elastomer materials with the reconfigurability of the liquid dye gain medium.

2. Chip design

As shown in Fig. 1, an optofluidic DFB dye laser was fabricated on a monolithic PDMS chip by replica molding, an inexpensive soft lithography technique described in Ref. [6]. When the microfluidic channel is subsequently filled with liquid of higher refractive index than that of PDMS (1.406, GE RTV615), a buried channel waveguide can be formed. The channel dimensions are 2μm×3μm and the index contrast is less than 0.003 so that the waveguide supports only the fundamental TE_{00} and TM_{00} modes. Distributed optical feedback is provided by defining periodic PDMS posts inside the channel with a 3080nm period. This forms a 1cm long 15th-order Bragg grating at wavelength approximately 570nm. The PDMS posts also provide mechanical support, preventing the microfluidic channel from collapsing. An
effective $\pi/2$ phase shift is introduced at the center of the grating to ensure single frequency operation at the Bragg wavelength [7]. The gain medium, a 2mM solution of Rhodamine 6G (Rh6G) or Rhodamine 101 (Rh101) in a methanol and ethylene glycol mixture with refractive index of 1.409, is then introduced into the flow channel to form the waveguide. To pump the dye molecules, 6ns Q-switched Nd:YAG laser pulses with 532nm wavelength are focused through a cylindrical lens onto a ~100µm×1cm stripe aligned with the microfluidic channel. The fabrication, dye chemistry, and operation of the laser chip are fully compatible with silicone elastomer based microfluidics technology [8].

When a $\pi/2$ phase-shifted DFB structure is used to provide the optical feedback, the lasing wavelength is determined by the Bragg condition:

$$m\lambda_m = 2n_{\text{eff}}\Lambda$$

where $\lambda_m$ is the $m$th order resonant wavelength, $n_{\text{eff}}$ is the effective index of the guided mode and $\Lambda$ is the grating period. Given $\Lambda = 3080\text{nm}$ and $n_{\text{eff}} = 1.407$, the 15th resonant wavelength and the free spectral range (FSR) are 577.8nm and 41.3nm respectively. This large FSR ensures that at most two resonances can be simultaneously supported within the gain spectrum (typically 30-50nm wide for dye molecules). Thus single frequency operation is obtained even at high pump levels due to gain discrimination. Figure 2 shows the simulated reflectivity spectrum of the overall structure using the Rouard’s method [9]. The parameters used are: $\Lambda = 1280\text{nm} + 1800\text{nm}$, grating length $L = 1\text{cm}$, effective $\pi/2$ phase shift at the center, core index $n_{\text{core}} = 1.409$, and cladding/post index $n_{\text{clad}} = 1.406$. Also shown in Fig. 2 are the normalized measured fluorescence spectra of Rh6G and Rh101 solutions used in the lasing experiment.

Fig. 2. Simulated reflectivity spectrum of a $\pi/2$ phase shifted higher order DFB structure. The parameters used are given in the main text. Also shown are the normalized measured fluorescence spectra of Rh6G and Rh101 solutions used in the lasing experiment.

3. Wavelength tuning

The lasing wavelength can be tuned by changing either $n_{\text{eff}}$, $\Lambda$ or $m$ as has been demonstrated in conventional DFB dye lasers [10]. The effective index $n_{\text{eff}}$ can be varied by changing the core index or the cross sectional dimensions of the waveguide. However, the low Young’s modulus of PDMS (~750kPa) [11] enables a much more straight forward tuning method consisting of changing the grating period by simply stretching or compressing the elastomeric chip along the waveguide direction. Moreover, the grating order $m$ can be chosen by using different dye molecules whose emission spectra cover different spectral regions. The last two
methods were used in this work to achieve a nearly 60nm tuning range from yellow to red. As can be observed from Fig. 2, the potential tuning range for Rh6G and Rh101 is larger than 100nm covering the spectrum from 550nm to beyond 650nm. Actually, because of the multiple spectral resonances supported by the higher order grating, this laser cavity design can provide tunable output covering the entire available dye laser spectrum from 320nm to 1200nm \[12\] when suitable dye molecules and pump light are selected. With a mixture of several dye molecules, simultaneous multiple color lasing from the same cavity is also possible.

4. Results and discussion

Fig. 3. Upper: normalized laser output of the mechanically tunable optofluidic DFB dye laser. Different peaks correspond to different grating periods. The measured laser linewidth is less than 0.1nm throughout the tuning range. Lower: lasing wavelength versus the measured chip deformation. The points are the experimental data and the curve is the linear fit. The achieved single-mode tuning range for Rh6G is from 565nm to 594nm and is from 613nm to 638nm for Rh101.

To achieve mechanical tuning, the laser chip was glued to two micrometer stages with the laser region suspended in the center as shown in Fig. 1. One of the stages is a high resolution
micrometer with 1μm sensitivity which provides accurate control and quantitative measurement of the deformation of the elastomer chip. The mechanical stages allow us to both stretch and compress the chip along the channel direction, and the result of mechanical cavity tuning is summarized in Fig. 3. The points on this Figure represent experimental data whereas the curves follow the linear fit. From this plot, the single mode tuning range for Rh6G covers a range from 565nm to 594nm whereas for Rh101, the tuning range extends from 613nm to 638nm. A linear relationship between the lasing wavelength and the chip deformation was also observed in Fig. 3, indicating completely reversible elastic extension of the cavity. When the length of the central suspended region is 1cm, the total chip deformation required to traverse the above tuning ranges are about 500μm for Rh6G and 400μm for Rh101. These macroscopic deformations correspond to 28nm and 25nm distortions of the grating period, respectively. Because of the extremely large available elastic deformation of PDMS of over 120%, the ultimate tuning range is limited by the gain bandwidth of the laser dyes chosen. Only ~5% deformation was used to achieve the ~60nm tuning range demonstrated in this experiment, and we believe that an even wider tuning range from 550nm to 650nm is obtainable with improved cavity design and a more uniform mechanical loading. The tuning is continuous and completely reversible, and no noticeable degradation of the chip was observed during a 5-cycle full range tuning test. Throughout the tuning range, stable single-mode operation was maintained with measured linewidth below 0.1nm, the resolution limited by our spectrometer (Ocean Optics HR4000).

The absorbed pump thresholds of this tunable dye laser are ~150nJ and ~200nJ for Rh6G and Rh101 respectively. As expected, we observe a decrease in the laser output power as the lasing wavelength is moved away from the peak in the gain spectrum in either direction. The deformation along the channel also causes the transverse waveguide dimensions to change slightly, resulting in changes the effective index of the guided mode. However, given that the Poisson’s ratio of PDMS is approximately 0.5, the estimated effective index change is only about 1.5×10^{-5}, the effect of lateral waveguide distortion on the lasing wavelength is negligible.

![Fig. 4. Left: optical micrograph of an integrated array of five optofluidic DFB dye lasers. The grating period of each laser is given on the left. Right: normalized laser output of the array using Rh6G dye solution as the gain medium.](image)

We also fabricated an array of five DFB dye lasers on a single PDMS chip. Figure 4 shows measured lasing results from this DFB array after filling the lasers with the Rh6G dye solution. Lasers with output wavelengths spanning a ~15nm range were measured from the DFB stripes with different grating periods. The low pump threshold (<1μJ) of each optofluidic DFB dye laser makes it possible to use a single high energy pulsed laser to pump hundreds of such lasers on a chip. This opens up the possibility of building highly parallel multiplexed biosensors on a chip ranging from applications as multiple-color flow cytometers and surface plasmon resonance based sensors to Raman spectroscopy sources and compact excitation spectroscopy systems. The introduction of replica molded multi-spectral sources in PDMS fluidic systems also provides an alternative to tunable lasers for constructing compact
and inexpensive multi-wavelength scanning-less spectrometers integrated in microfluidic devices [13]. The low pump threshold enables the use of visible semiconductor laser diodes or even light emitting diodes as the pump sources to construct low-cost and compact portable spectrometers.

5. Conclusions

We have demonstrated a continuously tunable optofluidic DFB dye laser on a monolithic PDMS chip that uses simple mechanical deformation for accurate wavelength control. Single-mode operation was maintained throughout a ~60nm tuning range. Due to the higher order of the DFB grating in our devices, a single laser is capable of generating tunable output covering the spectral region from near ultra-violet to near infrared when a UV pump source is used. An integrated multi-spectral DFB laser array was also demonstrated by replica molding. Such laser arrays can be used to make highly parallel multiplexed biosensors and scanning-less spectrometers on a chip. It is important to note that these lasers are still not stand-alone devices because an external gas pressure source is still required to actuate the microfluidic valves, and an external pump laser is required to pump the dye. The gas pressure source can be eliminated by using electrokinetically driven flow, whereas an external pump light will continue to be necessary for our dye laser designs. For portable and inexpensive devices, visible semiconductor lasers or bright light emitting diodes can be used as the pump source.

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