BICEP3 focal plane design and detector performance

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ABSTRACT

BICEP3, the latest telescope in the BICEP/Keck program, started science observations in March 2016. It is a 550mm aperture refractive telescope observing the polarization of the cosmic microwave background at 95 GHz. We show the focal plane design and detector performance, including spectral response, optical efficiency and preliminary sensitivity of the upgraded BICEP3. We demonstrate 9.72\(\mu\)K\(\sqrt{s}\) noise performance of the BICEP3 receiver.

Keywords: Cosmic Microwave Background, BICEP, Keck Array, Polarization
1. INTRODUCTION

Measurements of the polarization of the Cosmic Microwave Background provide key information to further our understanding of the early universe. The ΛCDM model predicts an E-mode polarization pattern in the CMB at the level of a few $\mu$K and arc-minute B-mode polarization arises from gravitational lensing of E-mode power by the large scale structure of the universe. But inflationary gravitational waves may be a source of degree scale B-mode polarization and a detection of such signal can be used to constrain the tensor-scalar ratio $r$ and place limits on the energy scale and potential from inflation. However, several galactic mechanisms can generate B-mode foregrounds; to disentangle the cosmic signal from galactic ones, we need to probe the polarization of the CMB at multiple frequencies with high sensitivity.

The BICEP/Keck team has deployed multiple telescopes to the South Pole since 2006; we use small aperture, refracting telescopes with high sensitivity receivers to map the degree scale B-mode signal. The Keck Array is in its fifth season, currently observing at 150 GHz and 220 GHz and previously observed at 95 GHz, 150 GHz and 220 GHz optical bands with 5 BICEP2 style cameras. BICEP3 is the latest addition to this program and was first deployed to the South Pole Station in 2015. It is a 550 mm aperture, on-axis, refractive polarimeter designed to observe at 95 GHz. During the first observing season, the focal plane was only partially filled with 1152 detectors, whereas in this year’s observing season, the instrument is complete with 2560 detectors. Combining data with BICEP2/Keck, Planck and South Pole Telescope, BICEP3 is projected to set upper limits on $r < 0.03$ at 95% confidence. Ref. 4 shows an overview of the BICEP3 telescope design and observing strategy.

2. FOCAL PLANE DESIGN

BICEP3 uses the same antenna-coupled TES bolometer architecture as BICEP2/Keck, but we adapted a new modular housing design that allows us to fill larger optically illuminated areas and for easier detector repair and replacement. This section shows the designs of both the detectors and the focal plane module.

2.1 Antenna-coupled transition-edge sensor (TES)

We couple optical power to our detectors through pairs of orthogonally polarized photolithographed planar antennas in each pixel, obviating the need for horns, contacting lenses, or other bulky coupling optics. These planar antenna arrays are composed of slot sub-radiators, spaced to Nyquist sample the focal plane surface to avoid grating lobes. Waves captured by the antenna slots of a given polarization orientation in the pixel are coherently combined through a microstrip summing tree. We control the illumination pattern in each pixel through the microstrip line impedance surrounding each T-junction; in BICEP3, we use a gaussian illumination pattern that minimizes spillover onto the cold aperture stop. Optical power from the antennas passes through an on-chip band pass filter before thermally dissipating on a released bolometer island in close thermal contact to the TESes (Figure 1). Each bolometer consists of two TESes: an aluminum TES with a transition temperature $T_c \sim 1.2$ K for higher loading in lab testing and a titanium TES with $T_c \sim 0.5$ K for on-sky measurement. Ref. 5 shows design parameter and fabrication procedure of the detectors. Each detector tile has 60 pixels at 95 GHz, with 4 additional dark pixels at the corners for calibration.

2.2 Focal Plane Module

The focal plane module consists of a quartz anti-reflection coating, detector tile, niobium (Nb) quarter-wave backshort, 1st stage superconducting quantum interference device (SQUID) chips, and the readout circuit boards (Figure 2). These components are stacked together on an aluminum detector frame, aligned with a 2 mm diameter pin/slot pair at opposite side and mounted at the corners with tile clips. The detector tiles thermally sink to their aluminum frames by $\sim 500$ gold wire bonds. The front side of the aluminum frame is corrugated to suppress coupling between the edge pixels and the frame (Figure 3).

All the detector parts are enclosed in a niobium housing to control the magnetic field environment near the SQUID ammeters and multiplexers. The module is mounted to a copper heat-sinking piece at the back of the niobium housing only making thermal contact at the center of the niobium housing to avoid trapped magnetic flux during cool down. Every module unit is independently placed onto the 250 mK focal plane of the telescope.
Figure 1. This figure shows the slot antenna array. Each pixel consists of two collocated 8 by 8 orthogonally polarized antenna arrays. Top right inset is a zoom showing the slot antennas for both polarizations. Bottom inset shows the TES bolometer. The Al and Ti TESes are at the left, gold microstrip termination is at right, and a thick layer of gold in the middle ensures thermal stability.

Figure 2. Left: Solidworks model for the BICEP3 module. The AR tile, detector, backshort and mux circuit board are mounted directly onto the Aluminum frame with tile clips and aligned with pin/slot pairs. Right: mask artwork for detector lithography. Each module contains 60 optically active pixels and 4 dark channels in the corners. All detectors connect to the circuit board via wirebond on top and bottom of the tile.

We mount a low-pass metal mesh filter (developed by Cardiff University) to remove unwanted above-band radiative loading. Two 60 channel flex cables connect to the back of the module.

This new modular packaging allows easy repair and upgrades individual detector tiles, and the compact design gives more efficient use of optical area in the focal plane. Together with the faster optics design, BICEP3 has a throughput 10 times higher than that of a single Keck 95 GHz receiver (Table 1). Figure 4 shows the focal plane currently installed in BICEP3. Future improvement on this design will allow us to minimize its weight and RF/magnetic interference.

2.3 Time-domain Multiplexing and Readout

BICEP3 uses a time-domain multiplexed (TDM) system developed at NIST for the bolometer readout. The readout electronics consists of the Nyquist chips (NYQ), SQUID multiplexing chips (MUX) and the SQUID series array (SSA). The NYQ chips are used to voltage bias the detectors with a 4 mΩ shunt resistor. The chips also
Figure 3. Left: Inside of the module. A set of alumina and G-10 circuit boards interconnect the SQUID (MUX) and Nyquist (NYQ) chips with the flex-cable connectors at center. These readout chips connect to the detector tile around a \( \lambda/4 \) backshort between the detectors and printed circuit boards. A few hundred gold wire bonds directly connect between the detector tile and aluminum detector frame. Right: the front side of the module and the corrugation on the frame to suppress coupling between the edge pixel and the frame. Bottom right: Zoom of the front side to show the corrugation on the frame.

<table>
<thead>
<tr>
<th></th>
<th>Single Keck/B2</th>
<th>BICEP3</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Optics</strong></td>
<td>f/2.2</td>
<td>f/1.68</td>
</tr>
<tr>
<td><strong>Aperture</strong></td>
<td>264 mm</td>
<td>520 mm</td>
</tr>
<tr>
<td><strong>FOV</strong></td>
<td>17 degree</td>
<td>27.4 degree</td>
</tr>
<tr>
<td><strong>Throughput</strong></td>
<td>37.8 cm(^2)sr</td>
<td>381 cm(^2)sr</td>
</tr>
</tbody>
</table>

Table 1. BICEP3 and Keck throughput

have a 2 \( \mu \)H inductors to limit the bandwidth. The MUX chips contain the first stage of the SQUID multiplexer, and the SSAs provide the final SQUID amplifier stage. The NYQ and MUX chips are located inside the module cooled to 280 mK, while the SSAs are attached to the 4 K temperature stage. A Multi-Channel Electronic (MCE) system developed by the University of British Columbia controls the bias and readout of all the channels.\(^8\)

The multiplexing architecture is \( 22 \times 30 \times 5 \): 22 TESes are read out in a multiplexer row and there are 30 multiplexer columns to form a MCE unit. Each set of the NYQ-MUX chip corresponds to a signal column and 11 rows, 2 chips are connected to form the 22 row multiplexing set, and 6 of these sets are mounted inside each module. 5 modules connect to a circuit board behind the focal plane (distribution board) to group all 30 columns. The row select lines are wired in series for every 5 modules. Superconducting niobium-titanium, twisted-pair cables connect the focal plane and SSAs at 4 K. They are readout by a MCE unit attached to the cryostat at 300 K. Four independent MCE units read out all 20 modules. Figure 5 shows the block diagram of the readout schematic.

### 2.3.1 SQUID Amplifier and Multiplexer

The SQUIDs play several simultaneous roles in our readout system. They amplify the small current output of the TESes while adding noise sub-dominant to the TES itself. They transform the small \( \sim 60 \text{ m}\Omega \) impedance of the TES to levels that warm amplifiers can match. Lastly, they have sufficient bandwidth to allow multiplexing of several detectors on common readout lines.
Figure 4. Left: Fully populated BICEP3 focal plane with 20 tiles. Right: Fully populated Keck focal plane with 4 detector tiles. The detector tiles are the same size on both focal plane.

Figure 5. Readout schematic of BICEP3. Every 5 modules are grouped and connected to a distribution board behind the focal plane at 280 mK, then to the SSA at 4 K, and connected to the room temperature MCE. 4 MCEs are used to readout all 20 tiles (2560 detectors).

Each independent detector is inductively coupled to a signal SQUID array (SQ1) by an input coil and the amplifier is operated in flux-lock loop to linearize the periodic output and increase the dynamic range of the SQUIDs response. As the flux from the input coil changes in response to the TES current, a compensating flux is applied by the feedback coil to cancel it. This flux feedback serves as the output of the TES channel. The SSA provides an additional stage of amplification that provides the aforementioned impedance matching between the first stage SQUIDs and room temperature MCE, providing $\sim 1 \Omega$ dynamic resistance for a $\sim 100 \Omega$ output impedance. Figure 6 shows a simplified schematic of the SQUID amplifier system. A similar design is used BICEP2/Keck and many other experiments.

Time Domain Multiplexing is possible because the SQ1 will not generate output signal when it is biased below its critical current $I_{\text{min}}$. Each SQ1 couples a TES to a shared common readout amplifier (SSA). While the TESes are continuously biased, they are only sampled when the corresponding SQ1 channel is biased. This allows our readout system to sequentially read 22 detectors in a common column, revisiting frequently enough to nyquist sample the highest relevant frequencies in the time-stream.
Each SQ1 bias in a signal column is controlled by a superconducting-to-normal flux activated switch that biases in parallel with the SQ1 and is controlled by the 22 row-select (RS) input lines. This design differs from that in BICEP2/Keck where the RS input lines separately biased each row of SQ1s, requiring an extra per-column intermediate summing coil and SQUID (SQ2) before reaching the SSA.

The flux activated switch is designed to switch at twice the critical current of the SQ1s, allowing the switches to share the same bias line with SQ1 in BICEP3. This ultimately reduces the electrical wiring going into the cold stage of the focal plane.

Control of the MUX system and feedback-based readout of the TES data are via the room temperature Multi-Channel Electronics (MCE) systems. The multiplexing speed needs to be quick enough for the Nyquist frequency to exceed the noise bandwidth to avoid aliasing penalty. BICEP2/Keck shows the optimal multiplexing speed is 25 kHz with a $2 \mu$H bandwidth limiting inductor. The data are filtered and down sampled in the MCE before being output to the computer software. The MCE uses a fourth-order digital Butterworth filter before down-sampling by a factor of 168, the control software applies a second stage of filtering using an acausal, zero-phase-delay FIR filter to down sampled by another factor of 5, giving a final sample rate of 31.1 Hz. The full multiplexing parameters used in BICEP3 are shown in Table 2.

### 3. DETECTOR PERFORMANCE

BICEP3 was first deployed to South Pole during the 2014-15 austral summer with 9 out of 20 detector modules. This season many major improvements were made including sub-Kelvin fridge hold time, RF shielding, IR thermal filtering, and fully populating the focal plane.

#### 3.1 Detector yield

While the fully populated focal plane had 1200 optically active dual-polarized detector pairs, the final working count is 951 pairs. This is largely due to electrical opens in 7 of the multiplexing rows, likely caused by wire bond damage during cooling.
Table 2. Summary of multiplexing parameters used by BICEP3.

<table>
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<tr>
<th>Parameter</th>
<th>Value</th>
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<tbody>
<tr>
<td>Raw ADC sample rate</td>
<td>50 MHz</td>
</tr>
<tr>
<td>Row dwell</td>
<td>90 samples</td>
</tr>
<tr>
<td>Row switching rate</td>
<td>556 kHz</td>
</tr>
<tr>
<td>Number of rows</td>
<td>22</td>
</tr>
<tr>
<td>Sample-row revisit rate</td>
<td>25.3 kHz</td>
</tr>
<tr>
<td>Internal downsample</td>
<td>168</td>
</tr>
<tr>
<td>Output data rate per channel</td>
<td>150 Hz</td>
</tr>
<tr>
<td>Software downsample</td>
<td>5</td>
</tr>
<tr>
<td>Archived data rate</td>
<td>31.1 Hz</td>
</tr>
</tbody>
</table>

3.2 Beams and detector spectral response

Far-field detector response was measured using a chopped thermal source about 200 m from the telescope. Differential pointing, beam width and ellipticity were measured for each detector to characterize and control for beam systematics in analysis. Ref. 10 shows the experimental setup and beam measurement result.

The detector spectral response of BICEP3 was measured with a Martin-Puplett interferometer. BICEP3 is designed to have an optical band centered at 93 GHz with a fractional bandwidth of 27% to avoid the oxygen line at 118 GHz and below 63 GHz. Figure 7 shows the co-added measured spectra for BICEP3 with a median band center at 93.3 GHz and bandwidth of 26.3 GHz (28% fractional bandwidth). There are evidences of tile to tile non-uniformity spectral respond in this season’s measurement that require further investigation.

![BICEP3 Co-added Spectra](image)

Figure 7. Co-added BICEP3 FTS measurement.

3.3 Optical Efficiency

A flat Eccosorb sheet was placed over the cryostat window to act as a beam-filling Rayleigh-Jeans source. It can be cooled with liquid nitrogen to 77 K or left at room temperature and we use the detector load curves taken at different temperatures to determine end-to-end optical efficiency (figure 8). The median optical efficiency is 0.08 pW/K$_{eq}$, which corresponds to an efficiency $\eta \sim 24\%$.

3.4 CMB Map

We show preliminary CMB temperature and polarization maps using the first 700 hours of 2016 BICEP3 CMB data (figure 10) with detector pointing and absolute gain calibrated by correlation with a reference Planck
3.5 Preliminary Map-based NET

Noise of the telescope can be estimate by evaluating the polarization map in section 3.4. The difference of two polarization maps made with data taken in the opposite azimuth scan directions (forward/backward jackknife) was used to remove CMB signal. The per-detector NET was calculated by multiplying the noise of the jackknife polarization maps and the square root of the integration time map. Figure 9 shows the histogram of the time weighted noise, which is well described by a gaussian distribution.

The median per-detector NET is $333 \mu K \sqrt{\text{s}}$ (figure 9) and the telescope NET is $9.72 \mu K \sqrt{\text{s}}$. This is a major improvement from the first season, which had a median NET of $395 \mu K \sqrt{s}$. This is still 28% higher than Keck 95 GHz receivers at about $260 \mu K \sqrt{s}$, possibly due to excess optical loading from the telescope.

This sensitivity estimation is different than the one used in Ref. 4, which calculated the NET by evaluating the time-stream data. Data were first pair-differenced, subjected to a third-order polynomial filter to reduce the $1/f$ noise induced from the atmosphere, and condensed into a noise spectra. The noise performance of the
Figure 10. Preliminary temperature and polarization maps of the CMB made with small amount of obtained data, corresponding to 700 hours of integration time.

detector is then evaluated by taking the median of the noise spectrum in the science band from 0.1 - 1 Hz. The time-stream based calculation gives per-detector NET at 347 $\mu$K$/\sqrt{s}$ and receiver NET at 9.91 $\mu$K$/\sqrt{s}$. The map-based NET calculation is lower than the time-stream method due to smaller weighing in the low frequency noise.

4. CONCLUSION

In this proceedings, we present the design of the BICEP3 focal plane module and its readout architecture. This compact design increases the packing density of the detectors and allows more efficient use of optically illuminated area on the focal plane. The modular design makes future replacements and upgrades easier. We also show great improvement in detector performance in the second season of BICEP3, increasing detector yield from 436 to 951 polarization-sensitive pixels, reducing the per-detector NET from 395 $\mu$K$/\sqrt{s}$ to 333 $\mu$K$/\sqrt{s}$, and achieving a receiver NET of 9.72 $\mu$K$/\sqrt{s}$. 
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REFERENCES


