

Search for Lepton-Flavor-Violating Decays $D^0 \rightarrow X^0 e^\pm \mu^\mp$

J. P. Lees, V. Poireau, and V. Tisserand

*Laboratoire d'Annecy-le-Vieux de Physique des Particules (LAPP),
Université de Savoie, CNRS/IN2P3, F-74941 Annecy-Le-Vieux, France*

E. Grauges

Universitat de Barcelona, Facultat de Física, Departament ECM, E-08028 Barcelona, Spain

A. Palano

INFN Sezione di Bari and Dipartimento di Fisica, Università di Bari, I-70126 Bari, Italy

G. Eigen

University of Bergen, Institute of Physics, N-5007 Bergen, Norway

D. N. Brown and Yu. G. Kolomensky

Lawrence Berkeley National Laboratory and University of California, Berkeley, California 94720, USA

M. Fritsch, H. Koch, and T. Schroeder

Ruhr Universität Bochum, Institut für Experimentalphysik 1, D-44780 Bochum, Germany

R. Cheaib^b, C. Hearty^{ab}, T. S. Mattison^b, J. A. McKenna^b, and R. Y. So^b

*Institute of Particle Physics^a; University of British Columbia^b,
Vancouver, British Columbia, Canada V6T 1Z1*

V. E. Blinov^{abc}, A. R. Buzykaev^a, V. P. Druzhinin^{ab}, V. B. Golubev^{ab}, E. A. Kozyrev^{ab}, E. A. Kravchenko^{ab},

A. P. Onuchin^{abc}, S. I. Serebnyakov^{ab}, Yu. I. Skovpen^{ab}, E. P. Solodov^{ab}, and K. Yu. Todyshev^{ab}

*Budker Institute of Nuclear Physics SB RAS, Novosibirsk 630090^a,
Novosibirsk State University, Novosibirsk 630090^b,
Novosibirsk State Technical University, Novosibirsk 630092^c, Russia*

A. J. Lankford

University of California at Irvine, Irvine, California 92697, USA

B. Dey, J. W. Gary, and O. Long

University of California at Riverside, Riverside, California 92521, USA

A. M. Eisner, W. S. Lockman, and W. Panduro Vazquez

University of California at Santa Cruz, Institute for Particle Physics, Santa Cruz, California 95064, USA

D. S. Chao, C. H. Cheng, B. Echenard, K. T. Flood, D. G. Hitlin, J. Kim, Y. Li,

D. X. Lin, T. S. Miyashita, P. Ongmongkolkul, J. Oyang, F. C. Porter, and M. Röhrken

California Institute of Technology, Pasadena, California 91125, USA

Z. Huard, B. T. Meadows, B. G. Pushpawela, M. D. Sokoloff, and L. Sun^{*}

University of Cincinnati, Cincinnati, Ohio 45221, USA

J. G. Smith and S. R. Wagner

University of Colorado, Boulder, Colorado 80309, USA

D. Bernard and M. Verderi

Laboratoire Leprince-Ringuet, Ecole Polytechnique, CNRS/IN2P3, F-91128 Palaiseau, France

D. Bettoni^a, C. Bozzi^a, R. Calabrese^{ab}, G. Cibinetto^{ab}, E. Fioravanti^{ab}, I. Garzia^{ab}, E. Luppi^{ab}, and V. Santoro^a

INFN Sezione di Ferrara^a; Dipartimento di Fisica e Scienze della Terra, Università di Ferrara^b, I-44122 Ferrara, Italy

A. Calcaterra, R. de Sangro, G. Finocchiaro, S. Martellotti,

P. Patteri, I. M. Peruzzi, M. Piccolo, M. Rotondo, and A. Zallo

INFN Laboratori Nazionali di Frascati, I-00044 Frascati, Italy

S. Passaggio and C. Patrignani[†]
INFN Sezione di Genova, I-16146 Genova, Italy

B. J. Shuve
Harvey Mudd College, Claremont, California 91711, USA

H. M. Lacker
Humboldt-Universität zu Berlin, Institut für Physik, D-12489 Berlin, Germany

B. Bhuyan
Indian Institute of Technology Guwahati, Guwahati, Assam, 781 039, India

U. Mallik
University of Iowa, Iowa City, Iowa 52242, USA

C. Chen, J. Cochran, and S. Prell
Iowa State University, Ames, Iowa 50011, USA

A. V. Gritsan
Johns Hopkins University, Baltimore, Maryland 21218, USA

N. Arnaud, M. Davier, F. Le Diberder, A. M. Lutz, and G. Wormser
Université Paris-Saclay, CNRS/IN2P3, IJCLab, F-91405 Orsay, France

D. J. Lange and D. M. Wright
Lawrence Livermore National Laboratory, Livermore, California 94550, USA

J. P. Coleman, E. Gabathuler,[‡] D. E. Hutchcroft, D. J. Payne, and C. Touramanis
University of Liverpool, Liverpool L69 7ZE, United Kingdom

A. J. Bevan, F. Di Lodovico,[§] and R. Sacco
Queen Mary, University of London, London, E1 4NS, United Kingdom

G. Cowan
University of London, Royal Holloway and Bedford New College, Egham, Surrey TW20 0EX, United Kingdom

Sw. Banerjee, D. N. Brown, and C. L. Davis
University of Louisville, Louisville, Kentucky 40292, USA

A. G. Denig, W. Gradl, K. Griessinger, A. Hafner, and K. R. Schubert
Johannes Gutenberg-Universität Mainz, Institut für Kernphysik, D-55099 Mainz, Germany

R. J. Barlow[¶] and G. D. Lafferty
University of Manchester, Manchester M13 9PL, United Kingdom

R. Cenci, A. Jawahery, and D. A. Roberts
University of Maryland, College Park, Maryland 20742, USA

R. Cowan
Massachusetts Institute of Technology, Laboratory for Nuclear Science, Cambridge, Massachusetts 02139, USA

S. H. Robertson^{ab} and R. M. Seddon^b
Institute of Particle Physics^a; McGill University^b, Montréal, Québec, Canada H3A 2T8

N. Neri^a and F. Palombo^{ab}
INFN Sezione di Milano^a; Dipartimento di Fisica, Università di Milano^b, I-20133 Milano, Italy

L. Cremaldi, R. Godang,** and D. J. Summers
University of Mississippi, University, Mississippi 38677, USA

P. Taras
Université de Montréal, Physique des Particules, Montréal, Québec, Canada H3C 3J7

G. De Nardo and C. Sciacca
*INFN Sezione di Napoli and Dipartimento di Scienze Fisiche,
 Università di Napoli Federico II, I-80126 Napoli, Italy*

G. Raven
NIKHEF, National Institute for Nuclear Physics and High Energy Physics, NL-1009 DB Amsterdam, The Netherlands

C. P. Jessop and J. M. LoSecco
University of Notre Dame, Notre Dame, Indiana 46556, USA

K. Honscheid and R. Kass
Ohio State University, Columbus, Ohio 43210, USA

A. Gaz^a, M. Margoni^{ab}, M. Posocco^a, G. Simi^{ab}, F. Simonetto^{ab}, and R. Stroili^{ab}
INFN Sezione di Padova^a; Dipartimento di Fisica, Università di Padova^b, I-35131 Padova, Italy

S. Akar, E. Ben-Haim, M. Bomben, G. R. Bonneaud, G. Calderini, J. Chauveau, G. Marchiori, and J. Ocariz
*Laboratoire de Physique Nucléaire et de Hautes Energies, Sorbonne Université,
 Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, F-75252 Paris, France*

M. Biasini^{ab}, E. Manoni^a, and A. Rossi^a
INFN Sezione di Perugia^a; Dipartimento di Fisica, Università di Perugia^b, I-06123 Perugia, Italy

G. Batignani^{ab}, S. Bettarini^{ab}, M. Carpinelli^{ab,††}, G. Casarosa^{ab}, M. Chrzaszcz^a, F. Forti^{ab}, M. A. Giorgi^{ab},
 A. Lusiani^{ac}, B. Oberhof^{ab}, E. Paoloni^{ab}, M. Rama^a, G. Rizzo^{ab}, J. J. Walsh^a, and L. Zani^{ab}
INFN Sezione di Pisa^a; Dipartimento di Fisica, Università di Pisa^b; Scuola Normale Superiore di Pisa^c, I-56127 Pisa, Italy

A. J. S. Smith
Princeton University, Princeton, New Jersey 08544, USA

F. Anulli^a, R. Faccini^{ab}, F. Ferrarotto^a, F. Ferroni^{a,‡‡}, A. Pilloni^{ab}, and G. Piredda^{a‡}
*INFN Sezione di Roma^a; Dipartimento di Fisica,
 Università di Roma La Sapienza^b, I-00185 Roma, Italy*

C. Bünger, S. Dittrich, O. Grünberg, M. Heß, T. Leddig, C. Voß, and R. Waldi
Universität Rostock, D-18051 Rostock, Germany

T. Auyeung and F. F. Wilson
Rutherford Appleton Laboratory, Chilton, Didcot, Oxon, OX11 0QX, United Kingdom

S. Emery and G. Vasseur
IRFU, CEA, Université Paris-Saclay, F-91191 Gif-sur-Yvette, France

D. Aston, C. Cartaro, M. R. Convery, J. Dorfan, W. Dunwoodie, M. Ebert, R. C. Field, B. G. Fulsom,
 M. T. Graham, C. Hast, W. R. Innes,[‡] P. Kim, D. W. G. S. Leith,[‡] S. Luitz, D. B. MacFarlane,
 D. R. Muller, H. Neal, B. N. Ratcliff, A. Roodman, M. K. Sullivan, J. Va'vra, and W. J. Wisniewski
SLAC National Accelerator Laboratory, Stanford, California 94309 USA

M. V. Purohit and J. R. Wilson
University of South Carolina, Columbia, South Carolina 29208, USA

A. Randle-Conde and S. J. Sekula

Southern Methodist University, Dallas, Texas 75275, USA

H. Ahmed

St. Francis Xavier University, Antigonish, Nova Scotia, Canada B2G 2W5

M. Bellis, P. R. Burchat, and E. M. T. Puccio
Stanford University, Stanford, California 94305, USA

M. S. Alam and J. A. Ernst
State University of New York, Albany, New York 12222, USA

R. Gorodeisky, N. Guttman, D. R. Peimer, and A. Soffer
Tel Aviv University, School of Physics and Astronomy, Tel Aviv, 69978, Israel

S. M. Spanier
University of Tennessee, Knoxville, Tennessee 37996, USA

J. L. Ritchie and R. F. Schwitters
University of Texas at Austin, Austin, Texas 78712, USA

J. M. Izen and X. C. Lou
University of Texas at Dallas, Richardson, Texas 75083, USA

F. Bianchi^{ab}, F. De Mori^{ab}, A. Filippi^a, and D. Gamba^{ab}
INFN Sezione di Torino^a; Dipartimento di Fisica, Università di Torino^b, I-10125 Torino, Italy

L. Lanceri and L. Vitale
INFN Sezione di Trieste and Dipartimento di Fisica, Università di Trieste, I-34127 Trieste, Italy

F. Martinez-Vidal and A. Oyanguren
IFIC, Universitat de Valencia-CSIC, E-46071 Valencia, Spain

J. Albert^b, A. Beaulieu^b, F. U. Bernlochner^b, G. J. King^b, R. Kowalewski^b,
T. Lueck^b, I. M. Nugent^b, J. M. Roney^b, R. J. Sobie^{ab}, and N. Tasneem^b
Institute of Particle Physics^a; University of Victoria^b, Victoria, British Columbia, Canada V8W 3P6

T. J. Gershon, P. F. Harrison, and T. E. Latham
Department of Physics, University of Warwick, Coventry CV4 7AL, United Kingdom

R. Prepost and S. L. Wu
University of Wisconsin, Madison, Wisconsin 53706, USA
(The *BABAR* Collaboration)

We present a search for seven lepton-flavor-violating neutral charm meson decays of the type $D^0 \rightarrow X^0 e^\pm \mu^\mp$, where X^0 represents a π^0 , K_S^0 , \bar{K}^{*0} , ρ^0 , ϕ , ω , or η meson. The analysis is based on 468 fb^{-1} of e^+e^- annihilation data collected at or close to the $\Upsilon(4S)$ resonance with the *BABAR* detector at the SLAC National Accelerator Laboratory. No significant signals are observed, and we establish 90% confidence level upper limits on the branching fractions in the range $(5.0-22.5) \times 10^{-7}$. The limits are between 1 and 2 orders of magnitude more stringent than previous measurements.

PACS numbers: 13.25.Ft, 11.30.Fs

* Now at: Wuhan University, Wuhan 430072, China

† Now at: Università di Bologna and INFN Sezione di Bologna,
I-47921 Rimini, Italy

‡ Deceased

§ Now at: King's College, London, WC2R 2LS, UK

¶ Now at: University of Huddersfield, Huddersfield HD1 3DH, UK

** Now at: University of South Alabama, Mobile, Alabama 36688,

I. INTRODUCTION

Lepton-flavor-conserving charm decays such as $D \rightarrow X e^+ e^-$ or $D \rightarrow X \mu^+ \mu^-$, where X is a meson, can occur in the standard model (SM) through short-distance [1, 2] and long-distance [2] processes, which have branching fractions of order $\mathcal{O}(10^{-9})$ and $\mathcal{O}(10^{-6})$, respectively. In contrast, the lepton-flavor-violating (LFV) neutral charm decays $D^0 \rightarrow X^0 e^\pm \mu^\mp$, where X^0 is a neutral meson, are effectively forbidden in the SM because they can occur only through lepton-flavor mixing [3] and are therefore suppressed to the order $\mathcal{O}(10^{-50})$. As such, the decays $D^0 \rightarrow X^0 e^\pm \mu^\mp$ should not be visible with current data samples. However, new-physics models, such as those involving Majorana neutrinos, leptoquarks, and two-Higgs doublets, allow for lepton number and lepton flavor to be violated [4–8]. Some models make predictions for, or use constraints from, three-body decays of the form $D \rightarrow X \ell' \ell$ or $B \rightarrow X \ell' \ell$, where ℓ and ℓ' represent an electron or muon [1, 6, 7, 9–12]. Most recent theoretical work has targeted the charged charm decays $D^+ \rightarrow X^+ \ell'^+ \ell^-$. For example, Ref. [4] estimates that $\mathcal{B}(D^+ \rightarrow \pi^+ \mu^\pm e^\mp)$ can be as large as 2×10^{-6} for certain leptoquark couplings. Some models that consider LFV and lepton-number-violating (LNV) four-body charm decays, with two leptons and two hadrons in the final state, predict branching fractions up to $\mathcal{O}(10^{-5})$, approaching those accessible with current data [6–8].

The branching fractions $\mathcal{B}(D^0 \rightarrow h'^- h^+ \mu^+ \mu^-)$, where h' and h represent a K or π meson, and $\mathcal{B}(D^0 \rightarrow K^- \pi^+ e^+ e^-)$ have recently been measured to be $\mathcal{O}(10^{-7})$ to $\mathcal{O}(10^{-6})$ [13–15], compatible with SM predictions [16, 17]. The branching fractions for the decays $D^0 \rightarrow X^0 e^+ e^-$ and $D^0 \rightarrow X^0 \mu^+ \mu^-$ have not yet been measured. However 90% confidence level (C.L.) upper limits on the branching fractions do exist and are in the range $(0.3–10) \times 10^{-5}$ for $D^0 \rightarrow X^0 e^+ e^-$ and $(3.2–53) \times 10^{-5}$ for $D^0 \rightarrow X^0 \mu^+ \mu^-$ [18–21]. It is likely that one or more of these decays are a major contributor to the branching fractions of the decays $D^0 \rightarrow h'^- h^+ e^+ e^-$ or $D^0 \rightarrow h'^- h^+ \mu^+ \mu^-$, as long-distance processes are predicted to be dominant [2], and published distributions of the invariant masses $m(h'^- h^+)$ for $D^0 \rightarrow h'^- h^+ \mu^+ \mu^-$ and $D^0 \rightarrow K^- \pi^+ e^+ e^-$ indicate large yields near some of the X^0 invariant masses [13–15].

The most stringent existing upper limits on the branching fractions for the LFV four-body decays of the type $D^0 \rightarrow h'^- h^+ e^\pm \mu^\mp$ are in the range $(11.0–19.0) \times 10^{-7}$ at the 90% confidence level [22]. For the LFV decays $D^0 \rightarrow X^0 e^\pm \mu^\mp$, where X^0 is an intermediate resonance meson decaying to $h'^- h^+$, $\pi^+ \pi^- \pi^0$ or $\gamma\gamma$, the 90% C.L. limits are in the range $(3.4–118) \times 10^{-5}$ [19, 20, 23]. For the

$D^0 \rightarrow X^0 e^\pm \mu^\mp$ decays with the same final state as the $D^0 \rightarrow h'^- h^+ e^\pm \mu^\mp$ decays ($D^0 \rightarrow K_s^0 (\rightarrow \pi^+ \pi^-) e^\pm \mu^\mp$, $D^0 \rightarrow \rho^0 (\rightarrow \pi^+ \pi^-) e^\pm \mu^\mp$, $D^0 \rightarrow \bar{K}^{*0} (\rightarrow K^- \pi^+) e^\pm \mu^\mp$, and $D^0 \rightarrow \phi (\rightarrow K^+ K^-) e^\pm \mu^\mp$), the current $D^0 \rightarrow X^0 e^\pm \mu^\mp$ branching fraction upper limits, which are in the range $(3.4–8.3) \times 10^{-5}$ [19, 20, 23], are approximately 20 times less stringent than the $D^0 \rightarrow h'^- h^+ e^\pm \mu^\mp$ limits reported in Ref. [22].

In this report we present a search for seven $D^0 \rightarrow X^0 e^\pm \mu^\mp$ LFV decays, where X^0 represents a π^0 , K_s^0 , \bar{K}^{*0} , ρ^0 , ϕ , ω , or η meson, with data recorded with the *BABAR* detector at the PEP-II asymmetric-energy $e^+ e^-$ collider operated at the SLAC National Accelerator Laboratory. The intermediate mesons X^0 are reconstructed through the decays $\pi^0 \rightarrow \gamma\gamma$, $K_s^0 \rightarrow \pi^+ \pi^-$, $\bar{K}^{*0} \rightarrow K^- \pi^+$, $\rho^0 \rightarrow \pi^+ \pi^-$, $\phi \rightarrow K^+ K^-$, $\omega \rightarrow \pi^+ \pi^- \pi^0$, $\eta \rightarrow \pi^+ \pi^- \pi^0$, and $\eta \rightarrow \gamma\gamma$. The branching fractions for the signal modes are measured relative to the normalization decays $D^0 \rightarrow \pi^- \pi^+ \pi^+ \pi^-$ (for $X^0 = K_s^0, \rho^0, \omega$), $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$ ($X^0 = \bar{K}^{*0}$), and $D^0 \rightarrow K^- K^+ \pi^+ \pi^-$ ($X^0 = \phi$). For $X^0 = \pi^0$ or η , the normalization mode $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$ is used as it has the smallest branching fraction uncertainty [23] and the largest number of reconstructed candidates of the three normalization modes. Although decays of the type $D^0 \rightarrow X^0 h'^- h^+$ have momentum distributions that more closely follow those of the signal decays under study, they suffer from smaller branching fractions, greater uncertainties on their branching fractions, and reduced reconstruction efficiencies relative to the three chosen normalization modes.

The D^0 mesons are identified using the decay $D^{*+} \rightarrow D^0 \pi^+$ produced in $e^+ e^- \rightarrow c\bar{c}$ events. Although D^0 mesons are also produced via other processes, the use of this decay chain increases the purity of the D^0 samples at the cost of a smaller number of reconstructed D^0 mesons.

II. THE *BABAR* DETECTOR AND DATA SET

The *BABAR* detector is described in detail in Refs. [24, 25]. Charged particles are reconstructed as tracks with a five-layer silicon vertex detector and a 40-layer drift chamber inside a 1.5 T solenoidal magnet. An electromagnetic calorimeter comprised of 6580 CsI(Tl) crystals is used to identify and measure the energies of electrons, positrons, muons, and photons. A ring-imaging Cherenkov detector is used to identify charged hadrons and to provide additional lepton identification information. Muons are primarily identified with an instrumented magnetic-flux return.

The data sample corresponds to 424 fb^{-1} of $e^+ e^-$ collisions collected at the center-of-mass (c.m.) energy of the $\Upsilon(4S)$ resonance (10.58 GeV, on peak) and an additional 44 fb^{-1} of data collected 0.04 GeV below the $\Upsilon(4S)$ resonance (off peak) [26].

Monte Carlo (MC) simulation is used to investigate

USA

†† Also at: Università di Sassari, I-07100 Sassari, Italy

‡‡ Also at: Gran Sasso Science Institute, I-67100 L'Aquila, Italy

sources of background contamination and evaluate selection efficiencies. Simulated events are also used to validate the selection procedure and for studies of systematic effects. The signal and normalization channels are simulated with the `EvtGen` package [27]. We generate the signal channel decays uniformly throughout the three-body phase space, while the normalization modes include two-body and three-body intermediate resonances, as well as nonresonant decays. We also generate $e^+e^- \rightarrow q\bar{q}$ ($q = u, d, s, c$), Bhabha and $\mu^+\mu^-$ pairs (collectively referred to as QED events), and $B\bar{B}$ background, using a combination of the `EvtGen`, `Jetset` [28], `KK2F` [29], `AfKQed` [30], and `TAUOLA` [31] generators, where appropriate. The background samples are produced with an integrated luminosity approximately 6 times that of the data. Final-state radiation is generated using `PHOTOS` [32]. The detector response is simulated with `GEANT 4` [33, 34]. All simulated events are reconstructed in the same manner as the data.

III. EVENT SELECTION

In the following, unless otherwise noted, all observables are evaluated in the laboratory frame. In order to optimize the event reconstruction, candidate selection criteria, multivariate analysis training, and fit procedure, a rectangular area in the $m(D^0)$ versus $\Delta m = m(D^{*+}) - m(D^0)$ plane is defined, where $m(D^{*+})$ and $m(D^0)$ are the reconstructed masses of the D^{*+} and D^0 candidates, respectively. This region is kept hidden (blinded) in data until the analysis steps are finalized. The hidden region is approximately 3 times the root mean square (RMS) width of the Δm and $m(D^0)$ resolutions. Its Δm region is $0.1447 < \Delta m < 0.1462 \text{ GeV}/c^2$ for all modes. The $m(D^0)$ signal peak distribution is asymmetric due to bremsstrahlung emission, with the left-side RMS width typically 1 to $2 \text{ MeV}/c^2$ wider than the right-side. The $m(D^0)$ RMS widths vary between 5 and $21 \text{ MeV}/c^2$, depending on the signal mode.

Particle identification (PID) criteria are applied to all charged daughter tracks of the intermediate meson X^0 decays. The charged pions and kaons are identified by measurements of their energy loss in the tracking detectors, and the number of photons and the Cherenkov angle recorded in the ring-imaging Cherenkov detector. These measurements are combined with information from the electromagnetic calorimeter and the muon detector to identify electrons and muons [24, 25]. Photons are detected and their energies are measured in the electromagnetic calorimeter. For $D^0 \rightarrow \phi e^\pm \mu^\mp$, the PID requirement on the kaons from the ϕ meson decay is relaxed compared to the single-kaon modes. This increases the reconstruction efficiency for this signal mode, with little increase in backgrounds or misidentified candidates. The muon PID requirement depends on the signal mode, with tighter requirements imposed for modes with more charged pions in the final state. The PID efficiency

depends on the track momentum, and is in the range $0.87-0.92$ for electrons, $0.60-0.95$ for muons, $0.86-0.98$ for pions, and $0.84-0.92$ for kaons. The misidentification probability [35], defined as the probability that particles are identified as one flavor (*e.g.* muon) that are in reality of a different flavor (*i.e.* not a muon), is typically less than 0.03 for all selection criteria, except for the pion selection criteria, where the muon misidentification rate can be as high as 0.35 at low momentum.

We select events that have at least five charged tracks, except for $D^0 \rightarrow \pi^0 e^\pm \mu^\mp$ and $D^0 \rightarrow \eta(\rightarrow \gamma\gamma) e^\pm \mu^\mp$, which must have at least three. Two or more of the tracks must be identified as leptons. The separation along the beam axis between the two leptons at their distance of closest approach to the beam line is required to be less than 0.2 cm. The leptons must have opposite charges, and their momenta must be greater than $0.3 \text{ GeV}/c$. Electrons and positrons from photon conversions are rejected by removing electron-positron pairs with an invariant mass less than $0.03 \text{ GeV}/c^2$ and a production vertex more than 2 cm from the beam axis.

The minimum photon energy in a signal decay is required to be greater than 0.025 GeV . For the decays $D^0 \rightarrow \pi^0 e^\pm \mu^\mp$ and $D^0 \rightarrow \eta(\rightarrow \gamma\gamma) e^\pm \mu^\mp$, the momentum of the π^0 or η must be greater than $0.4 \text{ GeV}/c$ and the energy of each photon from the π^0 must be greater than 0.045 GeV . The reconstructed π^0 invariant mass for all signal decays is required to be between 120 and $160 \text{ MeV}/c^2$.

The reconstructed invariant masses of the π^0 , K_s^0 , \bar{K}^{*0} , ρ^0 , ϕ , and ω candidates are required to be within 19, 9, 76, 240, 20, and $34 \text{ MeV}/c^2$, of their nominal mass [23], respectively. For the decays $\eta \rightarrow \gamma\gamma$ and $\eta \rightarrow \pi^+\pi^-\pi^0$, the invariant mass of the η candidates must be within 47 and $35 \text{ MeV}/c^2$ of the η nominal mass, respectively. These ranges are equivalent to 3 times the reconstructed RMS widths.

Candidate D^0 mesons for the signal modes are formed from the electron or positron, muon or antimuon, and intermediate resonance candidates. For the normalization modes, the D^0 candidate is formed from four charged tracks. Particle identification is applied to all charged tracks and the D^0 candidates are reconstructed with the appropriate charged-track mass hypotheses for both the signal and normalization decays. The tracks are required to form a good-quality vertex with a χ^2 probability for the vertex fit greater than 0.005. For the decay $D^0 \rightarrow K_s^0 e^\pm \mu^\mp$, the K_s^0 must have a transverse flight distance from the D^0 decay vertex greater than 0.2 cm. A bremsstrahlung energy recovery algorithm is applied to electrons and positrons, in which the energy of photon showers that are within a small angle (35 mrad in polar angle and 50 mrad in azimuth [24]) with respect to the tangent of the initial electron or positron direction is added to the energy of the electron or positron candidate. For the normalization modes, the reconstructed D^0 meson mass is required to be in the range $1.81 < m(D^0) < 1.91 \text{ GeV}/c^2$, while for the signal modes,

$m(D^0)$ must be in the hidden $m(D^0)$ range defined above.

The candidate D^{*+} is formed by combining the D^0 candidate with a charged pion having a momentum greater than $0.1 \text{ GeV}/c$. For the normalization mode $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$, this pion is required to have a charge opposite that of the kaon. The pion and D^0 candidate are subject to a vertex fit, with the D^0 mass constrained to its known value [23] and the requirement that the D^0 meson and the pion originate from the beam spot [36]. The χ^2 probability of the fit is required to be greater than 0.005. After the application of the D^{*+} vertex fit, the D^0 candidate momentum in the c.m. system $p^*(D^0)$ must be greater than $2.4 \text{ GeV}/c$. For the normalization modes, the mass difference Δm is required to be $0.143 < \Delta m < 0.148 \text{ GeV}/c^2$, while for the signal modes the range is $0.1395 < \Delta m < 0.1610 \text{ GeV}/c^2$. The extended Δm range for the signal modes provides greater stability when fitting the background distributions.

The requirement on the number of charged tracks strongly suppresses backgrounds from QED processes. The $p^*(D^0)$ criterion removes most sources of combinatorial background, as well as charm hadrons produced in B decays, which are kinematically limited to $p^*(D^0) \lesssim 2.2 \text{ GeV}/c$ [37].

Simulated samples indicate that the remaining background arises from $e^+e^- \rightarrow c\bar{c}$ events in which charged tracks and neutral particles can either be lost or selected from elsewhere in the event to form a D^0 candidate. To reject this background, a multivariate selection based on a Boosted Decision Tree (BDT) discriminant is applied to the signal modes [38]. A common set of eight input observables is used for all modes: the momenta of the electron or positron, muon or antimuon, and reconstructed intermediate meson; the momentum of the lowest-momentum charged track or photon from the X^0 candidate; the maximum angle between the direction of D^0 daughters and the D^0 direction; the total energy of all charged tracks and photons in the event, normalized to the beam energy; the ratio $x_p = p^*(D^{*+})/\sqrt{E_{e^+e^-}^{*2} - m^2(D^{*+})}$, where $p^*(D^{*+})$ is the c.m. momentum of the D^{*+} candidate and $E_{e^+e^-}^*$ is the c.m. beam energy; and the reconstructed mass of the intermediate meson. Three additional input observables are used for the D^0 decays with ω or η decaying to $\pi^+\pi^-\pi^0$: the momentum and reconstructed mass of the π^0 candidate, and the energy of the lowest-energy photon from the π^0 . The discriminant is trained and tested independently for each signal mode, using simulated samples for the signal modes, and ensembles of data outside the hidden region and $e^+e^- \rightarrow c\bar{c}$ simulated samples for the background. Depending on the signal mode, the requirement on the discriminant output accepts between 70% to 90% of the simulated signal sample while rejecting between 50% to 90% of the background.

The cross feed to one signal mode from any other signal modes is estimated from simulated samples to be less than 4% in all cases, and typically less than 1%, assum-

ing equal branching fractions for all signal modes. The cross feed to a specific normalization mode from the other two normalization modes is predicted from simulation to be less than 0.7%, where the branching fractions are taken from Ref. [23]. In the data, no events with reconstructed normalization decays contain reconstructed signal decays.

From the data, we find that the fraction of normalization mode events with more than one candidate is 2.4%, 3.6%, and 4.4% for $D^0 \rightarrow K^- K^+ \pi^+ \pi^-$, $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$, and $D^0 \rightarrow \pi^- \pi^+ \pi^+ \pi^-$, respectively. For the signal mode with $\eta \rightarrow \pi^+ \pi^- \pi^0$, 40% of events have multiple candidates. For $\eta \rightarrow \gamma\gamma$ and ω decays, the number of events with multiple candidates is $\sim 10\%$, and for the remaining modes it is 1% to 5%. If two or more candidates are found in an event, the one with the highest D^{*+} vertex χ^2 probability is selected. After applying the best-candidate selection, the correct D^{*+} candidate in the simulated samples is selected with a probability of 95% or more for the normalization modes. For the signal modes, 70% of D^{*+} candidates are correctly selected for $\eta \rightarrow \pi^+ \pi^- \pi^0$, and between 86% and 94% for the remaining modes. After the application of all selection criteria and corrections for small differences between data and MC simulation in tracking and PID performance, the reconstruction efficiency ϵ_{sig} for the simulated signal decays is between 1.6% and 3.6%, depending on the mode. For the normalization decays, the reconstruction efficiency ϵ_{norm} is between 19.2% and 24.7%. The difference between ϵ_{sig} and ϵ_{norm} is mainly due to the minimum momentum criterion on the leptons required by the PID algorithms [25].

IV. SIGNAL YIELD EXTRACTION

The $D^0 \rightarrow X^0 e^\pm \mu^\mp$ signal mode branching fraction \mathcal{B}_{sig} is determined relative to that of the normalization decay using

$$\mathcal{B}_{\text{sig}} = \frac{N_{\text{sig}}}{N_{\text{norm}}} \frac{\epsilon_{\text{norm}}}{\epsilon_{\text{sig}}} \frac{\mathcal{L}_{\text{norm}}}{\mathcal{L}_{\text{sig}}} \frac{\mathcal{B}_{\text{norm}}}{\mathcal{B}(X^0)}, \quad (1)$$

where $\mathcal{B}_{\text{norm}}$ is the branching fraction of the normalization mode [23], and N_{sig} and N_{norm} are the fitted yields of the signal and normalization mode decays, respectively. $\mathcal{B}(X^0)$ is the branching fraction of the intermediate meson decay channel. The symbols \mathcal{L}_{sig} and $\mathcal{L}_{\text{norm}}$ represent the integrated luminosities of the data samples used for the signal ($468.2 \pm 2.0 \text{ fb}^{-1}$) and the normalization decays ($39.3 \pm 0.2 \text{ fb}^{-1}$), respectively [26]. For the signal modes, we use both the on-peak and off-peak data samples. For the normalization modes, a subset of the off-peak data is sufficient for achieving statistical uncertainties that are much smaller than the systematic uncertainties.

We perform an extended unbinned maximum likelihood fit to extract the signal and background yields for

both the normalization and signal modes [39]. The likelihood function is

$$\mathcal{L} = \frac{1}{N!} \exp\left(-\sum_{j=1}^2 n_j\right) \prod_{i=1}^N \left[\sum_{j=1}^2 n_j \mathcal{P}_j(\vec{x}_i; \vec{\alpha}_j) \right]. \quad (2)$$

We define the likelihood for each event candidate i to be the sum of $n_j \mathcal{P}_j(\vec{x}_i; \vec{\alpha}_j)$ over two hypotheses j (signal or normalization and background). The symbol $\mathcal{P}_j(\vec{x}_i; \vec{\alpha}_j)$ is the product of the probability density functions (PDFs) for hypothesis j evaluated for the measured variables \vec{x}_i of the i -th event. The total number of events in the sample is N , and n_j is the yield for hypothesis j . The quantities $\vec{\alpha}_j$ represent parameters of \mathcal{P}_j . The distributions of each discriminating variable x_i in the likelihood function is modeled with one or more PDFs, where the parameters $\vec{\alpha}_j$ are determined from fits to signal simulation or data samples.

Each normalization mode yield N_{norm} is extracted by performing a two-dimensional unbinned maximum likelihood fit to the Δm versus $m(D^0)$ distributions in the range $0.143 < \Delta m < 0.148 \text{ GeV}/c^2$ and $1.81 < m(D^0) < 1.91 \text{ GeV}/c^2$. Considering normalization and background events separately, the measured Δm and $m(D^0)$ values are essentially uncorrelated and are therefore treated as independent observables in the fits. The PDFs in the fits depend on the normalization mode and use sums of multiple Cruijff [15] and Crystal Ball [40] functions in both Δm and $m(D^0)$. The functions for each observable use a common mean. The background is modeled with an ARGUS threshold function [41] for Δm and a Chebyshev polynomial for $m(D^0)$. The ARGUS end point parameter is fixed at $0.1395 \text{ GeV}/c^2$, the Δm kinematic threshold for $D^{*+} \rightarrow D^0 \pi^+$ decays. All other PDF parameters, together with the normalization mode and background yields, are allowed to vary in the fit.

The fitted yields and reconstruction efficiencies for the normalization modes are given in Table I. Figure 1 shows projections of the unbinned maximum-likelihood fits onto the final candidate distributions as a function of Δm for the normalization modes in the range $0.143 < \Delta m < 0.148 \text{ GeV}/c^2$.

TABLE I. Summary of fitted candidate yields, with statistical uncertainties, and reconstruction efficiencies for the three normalization modes.

Decay mode	N_{norm} (candidates)	ϵ_{norm} (%)
$D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$	260870 ± 520	20.1 ± 0.2
$D^0 \rightarrow K^- K^+ \pi^+ \pi^-$	8480 ± 110	19.2 ± 0.2
$D^0 \rightarrow \pi^- \pi^+ \pi^+ \pi^-$	28470 ± 220	24.7 ± 0.2

After the application of the selection criteria, there are on the order of 100 events or fewer available for fitting in each signal mode. Each signal mode yield N_{sig} is therefore extracted by performing a one-dimensional

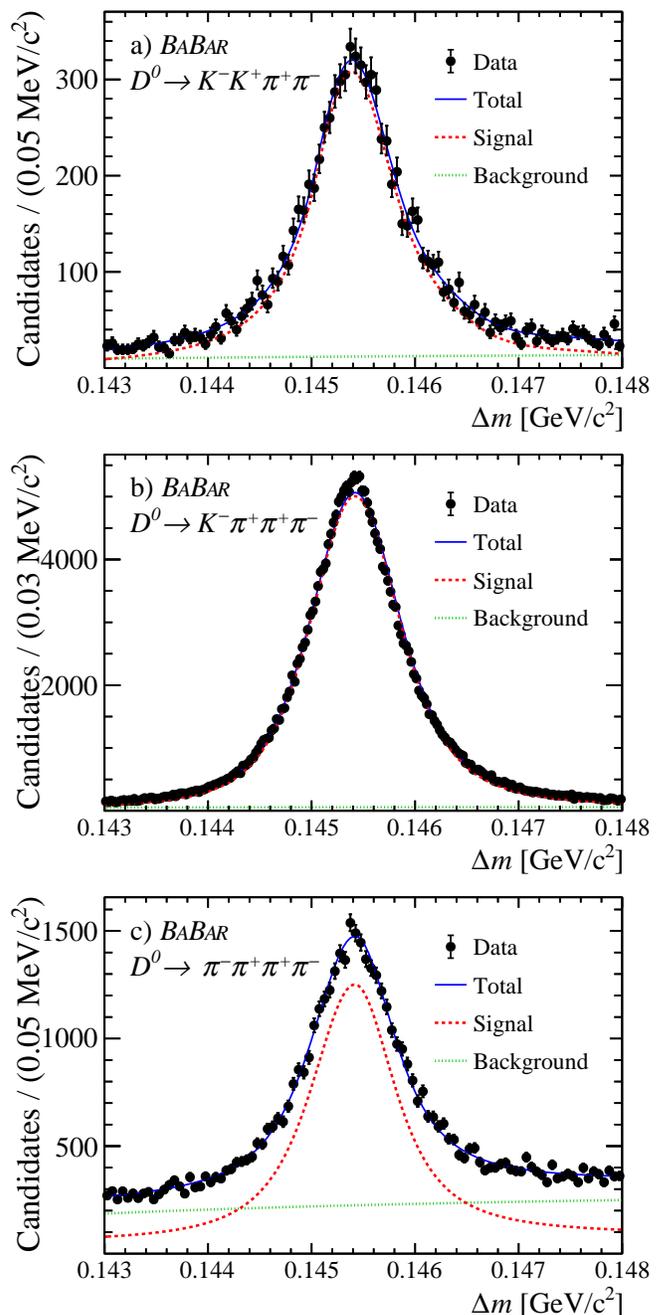


FIG. 1. Projections of the unbinned maximum-likelihood fits to the final candidate distributions as a function of Δm for the normalization modes in the range $0.143 < \Delta m < 0.148 \text{ GeV}/c^2$. The solid blue line is the total fit, the dashed red line is the signal and the dotted green line is the background.

unbinned maximum likelihood fit to Δm in the range $0.1395 < \Delta m < 0.1610 \text{ GeV}/c^2$. A Cruijff function is implemented for the signal mode PDF, except for $D^0 \rightarrow \phi e^\pm \mu^\mp$, for which two two-piece Gaussians functions are used, and $D^0 \rightarrow \rho^0 e^\pm \mu^\mp$, for which two Cruijff

functions are used. The background is modeled with an ARGUS function with the same end point used for the normalization modes. The signal PDF parameters and the end point parameter are fixed in the fit. All other background parameters and the signal and background yields are allowed to vary. Figure 2 shows the results of the fits to the Δm distributions for the signal modes.

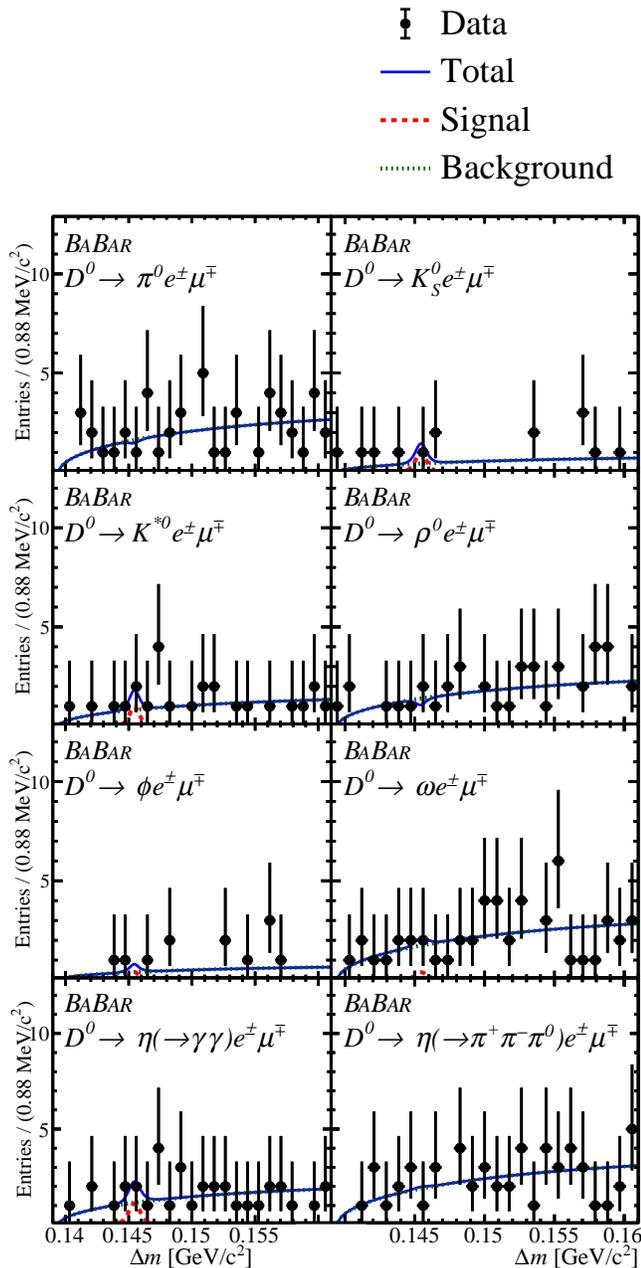


FIG. 2. Unbinned maximum-likelihood fits to the final candidate distributions as a function of Δm for the signal modes in the range $0.1395 < \Delta m < 0.1610 \text{ GeV}/c^2$. The solid blue line is the total fit, the dashed red line is the signal and the dotted green line is the background.

We test the performance of the maximum likelihood fit for the normalization modes by generating ensembles of MC samples from the normalization and background PDF distributions. The mean numbers of normalization and background candidates used in the ensembles are taken from the fits to the data. The numbers of generated background and normalization mode candidates are sampled from a Poisson distribution. All background and normalization mode PDF parameters are allowed to vary, except for the ARGUS function end point. No significant biases are observed in the fitted yields of the normalization modes. The same procedure is repeated for the maximum likelihood fits to the signal modes, with ensembles of MC samples generated from the background PDF distributions only, assuming a signal yield of zero. The signal PDF parameters are fixed to the values used for the fits to the data, and the signal yield is allowed to vary. The biases in the fitted signal yields are less than ± 0.3 candidates for all modes, and these are subtracted from the fitted yields before calculating the signal branching fractions.

To confirm the normalization procedure, the signal modes in Eq. (1) are replaced with the decay $D^0 \rightarrow K^- \pi^+$, which has a well-measured branching fraction [23]. The $D^0 \rightarrow K^- \pi^+$ decays are reconstructed using the on-peak data sample only ($424.3 \pm 1.8 \text{ fb}^{-1}$). The $D^0 \rightarrow K^- \pi^+$ decay is selected using the same criteria as used for the $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$ mode, which is used as the normalization mode for this test. The $D^0 \rightarrow K^- \pi^+$ signal yield is 1881950 ± 1380 with $\epsilon_{\text{sig}} = (27.4 \pm 0.2)\%$. Thus, we determine $\mathcal{B}(D^0 \rightarrow K^- \pi^+) = (3.98 \pm 0.08 \pm 0.10)\%$, where the uncertainties are statistical and systematic, respectively. This is consistent with the current world average of $(3.95 \pm 0.03)\%$ [23]. When the test is repeated using either $D^0 \rightarrow K^- K^+ \pi^+ \pi^-$ or $D^0 \rightarrow \pi^- \pi^+ \pi^+ \pi^-$ as the normalization mode, $\mathcal{B}(D^0 \rightarrow K^- \pi^+)$ is determined to be $(3.51 \pm 0.18 \pm 0.18)\%$ and $(4.12 \pm 0.13 \pm 0.16)\%$, respectively.

V. SYSTEMATIC UNCERTAINTIES

The systematic uncertainties in the branching fraction determinations of the signal modes arise from so-called additive systematic uncertainties that affect the significance of the signal mode yields in the fits to the data samples and from multiplicative systematic uncertainties on the luminosity and signal reconstruction efficiencies.

The main sources of the additive systematic uncertainties in the signal yields are associated with the model parametrizations used in the fits to the signal modes, the fit biases, the allowed invariant-mass ranges for the D^0 and X^0 candidates, the amount of cross feed, and the limited MC and data sample sizes available for the optimization of the BDT discriminants.

The uncertainties associated with the fit model parametrizations of the signal modes are estimated by repeating the fits with alternative PDFs. This involves

replacing the Cruijff functions with Crystal Ball functions, using a two-piece Gaussian function, and changing the number of functions used in the PDFs. For the background, the ARGUS function is replaced by a first- or second-order polynomial. The largest deviation occurs when using the Crystal Ball functions for the signal and the first-order polynomial for the background. The systematic uncertainty is taken as half this maximum deviation. The largest contribution comes from the normalization mode $D^0 \rightarrow \pi^- \pi^+ \pi^+ \pi^-$ due to the presence of increased background and greater uncertainty in the background shape. To account for potential inaccuracies in the simulation of the D^0 and X^0 invariant mass distributions, we change the mass selection ranges by $\pm 0.5\sigma$, where σ is the RMS width of the D^0 or X^0 meson.

The systematic uncertainties in the correction on the fit biases for the signal yields are taken from the ensembles of fits to the MC samples. Given the central value of the signal yield obtained from the fit in each mode, the cross feed yields from all other modes are calculated and are taken as a systematic uncertainty. To evaluate the systematic uncertainty in the application of the BDT discriminant, we vary the value of the selection criterion for the BDT discriminant output, change the size of the hidden region in data, and also retrain the BDT discriminant using a training sample with a different ensemble of MC samples. Summing the uncertainties in quadrature, the total additive systematic uncertainties in the signal yields are between 0.4 and 0.9 events.

Multiplicative systematic uncertainties are due to assumptions made about the distributions of the final-state particles in the signal simulation modeling, the model parametrizations used in the fits to the normalization modes, the normalization mode branching fractions, tracking and PID efficiencies, limited simulation sample sizes, and luminosity.

Since the decay mechanism of the signal modes is unknown, we vary the angular distributions of the simulated final-state particles from the D^0 signal decay in three angular variables, defined following the prescription of Ref. [42]. We weight the events, which are simulated uniformly in phase space, using combinations of \sin , \cos , \sin^2 , and \cos^2 functions of the angular variables. The reconstruction efficiencies calculated from simulation samples as functions of the three angles are constant, within the statistics available. The deviations of the reweighted efficiencies from the default average reconstruction efficiencies are therefore small. Half the maximum change in the average reconstruction efficiency is assigned as a systematic uncertainty.

Uncertainties associated with the fit model parametrizations of the normalization modes are estimated by repeating the fits with alternative PDFs. This involves swapping the Cruijff and Crystal Ball functions used in both Δm and $m(D^0)$. For the background, the order of the polynomials is changed and the ARGUS function is replaced by a second-order polynomial. Half the maximum change in the fitted yield is assigned as

a systematic uncertainty. The normalization modes branching fraction uncertainties are taken from Ref. [23].

For both signal and normalization modes, we include uncertainties to account for discrepancies between reconstruction efficiencies calculated from simulation and data samples of 1.0% per K_S^0 , 0.8% per lepton, and 0.7% per hadron track [43]. We include a momentum-dependent π^0 reconstruction efficiency uncertainty of 2.1% for $D^0 \rightarrow \pi^0 e^\pm \mu^\mp$ and 2.3% for $D^0 \rightarrow \omega e^\pm \mu^\mp$ and $D^0 \rightarrow \eta(\rightarrow \pi^+ \pi^- \pi^0) e^\pm \mu^\mp$. For the PID efficiencies, we assign an uncertainty of 0.7% per track for electrons, 1.0% for muons, 0.2% for charged pions, and 1.1% for kaons [25]. A systematic uncertainty of 0.4% is associated with our knowledge of the luminosities $\mathcal{L}_{\text{norm}}$ and \mathcal{L}_{sig} [26]. We assign systematic uncertainties in the range 0.8% to 1.8% to account for the limited size of the simulation samples available for calculating reconstruction efficiencies for the signal and normalization modes.

The simulation samples for the normalization modes contain a resonant structure of intermediate resonances that decay to two- or three-body final states, as well as four-body nonresonant decays. To investigate how changes in the resonant structure affect the reconstruction efficiencies, the simulation samples were generated using a four-body phase-space distribution only and the reconstruction efficiencies recalculated. The resulting changes in reconstruction efficiencies are less than the statistical uncertainties on ϵ_{norm} due to the limited size of the simulation samples, and no systematic uncertainties are assigned. The total multiplicative systematic uncertainties are between 4.7% and 6.8% for the normalization modes and between 4.2% and 7.8% for the signal modes.

Table II summarizes the contributions of the systematic uncertainties of the normalization modes to the systematic uncertainties in the signal mode branching fractions, as defined in Eq. (1). Table III summarizes the systematic uncertainties in the signal mode yields, excluding those due to the normalization modes.

TABLE II. Summary of the contributions to the systematic uncertainties on the signal mode branching fractions, as defined in Eq. (1), that arise from uncertainties in the measurement of the normalization modes.

	$\pi^- \pi^+ \pi^+ \pi^-$	$K^- \pi^+ \pi^+ \pi^-$	$K^- K^+ \pi^+ \pi^-$
PDF variation	4.6%	1.0%	1.0%
K_S^0 correction	1.0%	1.0%	1.0%
Tracking correction	3.5%	3.5%	3.5%
PID correction	0.8%	1.7%	2.6%
Luminosity	0.4%	0.4%	0.4%
Normalization \mathcal{B}	3.0%	1.8%	4.5%
Simulation size	1.0%	1.0%	0.8%
Total	6.8%	4.7%	6.6%

TABLE III. Summary of $D^0 \rightarrow X^0 e^\pm \mu^\mp$ additive and multiplicative systematic uncertainties, excluding those due to the normalization modes given in Table II.

	$X^0 = \pi^0$	K_S^0	\bar{K}^{*0}	ρ^0	ϕ	ω	η	η
$X^0 \rightarrow$	$\gamma\gamma$	$\pi^+\pi^-$	$K^-\pi^+$	$\pi^+\pi^-$	K^+K^-	$\pi^+\pi^-\pi^0$	$\gamma\gamma$	$\pi^+\pi^-\pi^0$
Additive (events):								
PDF variation	0.23	0.05	0.20	0.16	0.17	0.26	0.43	0.16
Fit bias	0.09	0.28	0.21	0.15	0.24	0.09	0.08	0.07
D^0/X^0 mass	0.30	0.04	0.05	0.07	0.07	0.07	0.04	0.23
BDT discriminant	0.83	0.68	0.71	0.30	0.06	0.35	0.27	0.58
Cross feed				0.01			0.06	
Subtotal (candidates)	0.92	0.74	0.76	0.38	0.31	0.45	0.52	0.65
Multiplicative (%):								
Angular variation	1.4	2.8	2.0	3.4	5.3	1.9	1.6	1.6
$\mathcal{B}(X^0)$ subdecay		0.1			1.0	0.8	0.5	1.2
K_S^0 correction		1.0						
Tracking correction	2.3	3.7	3.7	3.7	3.7	3.7	2.3	3.7
PID correction	2.7	2.1	3.0	2.1	3.9	3.1	2.7	3.1
π^0 correction	2.1					2.3		2.3
Luminosity	0.4	0.4	0.4	0.4	0.4	0.4	0.4	0.4
Simulation sample size	1.4	1.3	1.5	1.3	1.4	1.8	1.3	1.5
Subtotal (%)	4.2	5.4	5.4	5.6	7.8	5.7	4.2	5.6

VI. RESULTS

Table IV gives the fitted signal yields, reconstruction efficiencies, branching fractions with statistical and systematic uncertainties, 90% C.L. upper limits on the branching fractions, and previous upper limits [19, 20, 23] for the signal modes. The yields for all the signal modes are compatible with zero. We assume that there are no cancellations due to correlations in the systematic uncertainties in the numerator and denominator of Eq. (1). We use the frequentist approach of Feldman and Cousins [44] to determine 90% C.L. bands. When computing the limits, the systematic uncertainties are combined in quadrature with the statistical uncertainties in the fitted signal yields.

In summary, we report 90% C.L. upper limits on the branching fractions for seven lepton-flavor-violating $D^0 \rightarrow X^0 e^\pm \mu^\mp$ decays. The analysis is based on a sample of e^+e^- annihilation data collected with the *BABAR* detector, corresponding to an integrated luminosity of $468.2 \pm 2.0 \text{ fb}^{-1}$. The limits are in the range $(5.0 - 22.5) \times 10^{-7}$ and are between 1 and 2 orders of magnitude more stringent than previous $D^0 \rightarrow X^0 e^\pm \mu^\mp$ decay results. For the four $D^0 \rightarrow X^0 e^\pm \mu^\mp$ decays with the same final state as the $D^0 \rightarrow h'^- h^+ e^\pm \mu^\mp$ decays reported in Ref. [22], the limits are 1.5 to 3 times more stringent.

VII. ACKNOWLEDGMENTS

We are grateful for the extraordinary contributions of our PEP-II colleagues in achieving the excellent luminosity and machine conditions that have made this work possible. The success of this project also relies critically on the expertise and dedication of the computing organizations that support *BABAR*. The collaborating institutions wish to thank SLAC for its support and the kind hospitality extended to them. This work is supported by the US Department of Energy and National Science Foundation, the Natural Sciences and Engineering Research Council (Canada), the Commissariat à l’Energie Atomique and Institut National de Physique Nucléaire et de Physique des Particules (France), the Bundesministerium für Bildung und Forschung and Deutsche Forschungsgemeinschaft (Germany), the Istituto Nazionale di Fisica Nucleare (Italy), the Foundation for Fundamental Research on Matter (The Netherlands), the Research Council of Norway, the Ministry of Education and Science of the Russian Federation, Ministerio de Economía y Competitividad (Spain), the Science and Technology Facilities Council (United Kingdom), and the Binational Science Foundation (U.S.-Israel). Individuals have received support from the Marie-Curie IEF program (European Union) and the A. P. Sloan Foundation (USA).

[1] A. Paul, I. I. Bigi, and S. Recksiegel, Phys. Rev. D **83**, 114006 (2011).

[2] A. J. Schwartz, Mod. Phys. Lett. A **08**, 967 (1993).

[3] D. Guadagnoli and K. Lane, Phys. Lett. B **751**, 54 (2015).

[4] S. de Boer and G. Hiller, Phys. Rev. D **93**, 074001 (2016).

[5] S. Fajfer and N. Košnik, Eur. Phys. J. C **75**, 567 (2015).

TABLE IV. Summary of fitted signal yields N_{sig} with statistical and systematic uncertainties, reconstruction efficiencies ϵ_{sig} , branching fractions with statistical and systematic uncertainties, 90% C.L. upper limits (U.L.) on the branching fractions, and previous limits [19, 20, 23]. The additive and multiplicative uncertainties are combined to obtain the overall systematic uncertainties. The branching fraction systematic uncertainties include the uncertainties in the normalization mode branching fractions.

Decay mode	N_{sig} (candidates)	ϵ_{sig} (%)	\mathcal{B} ($\times 10^{-7}$)	\mathcal{B} 90% U.L. ($\times 10^{-7}$)	
				BABAR	Previous
$D^0 \rightarrow \pi^0 e^\pm \mu^\mp$	$-0.3 \pm 2.0 \pm 0.9$	2.15 ± 0.03	$-0.6 \pm 4.8 \pm 2.3$	8.0	860
$D^0 \rightarrow K_S^0 e^\pm \mu^\mp$	$0.7 \pm 1.7 \pm 0.7$	3.01 ± 0.04	$1.9 \pm 4.6 \pm 1.9$	8.6	500
$D^0 \rightarrow \bar{K}^{*0} e^\pm \mu^\mp$	$0.8 \pm 1.8 \pm 0.8$	2.31 ± 0.03	$2.8 \pm 6.1 \pm 2.6$	12.4	830
$D^0 \rightarrow \rho^0 e^\pm \mu^\mp$	$-0.7 \pm 1.7 \pm 0.4$	2.10 ± 0.03	$-1.8 \pm 4.4 \pm 1.0$	5.0	490
$D^0 \rightarrow \phi e^\pm \mu^\mp$	$0.0 \pm 1.4 \pm 0.3$	3.43 ± 0.04	$0.1 \pm 3.8 \pm 0.9$	5.1	340
$D^0 \rightarrow \omega e^\pm \mu^\mp$	$0.4 \pm 2.3 \pm 0.5$	1.46 ± 0.03	$1.8 \pm 9.5 \pm 1.9$	17.1	1200
$D^0 \rightarrow \eta e^\pm \mu^\mp$			$6.1 \pm 9.7 \pm 2.3$	22.5	1000
with $\eta \rightarrow \gamma\gamma$	$1.6 \pm 2.3 \pm 0.5$	2.96 ± 0.04	$7.0 \pm 10.5 \pm 2.4$	24.0	
with $\eta \rightarrow \pi^+\pi^-\pi^0$	$0.0 \pm 2.8 \pm 0.7$	2.46 ± 0.04	$0.4 \pm 25.8 \pm 6.0$	42.8	

- [6] A. Atre, T. Han, S. Pascoli, and B. Zhang, *J. High Energy Phys.* **05**, 030 (2009).
- [7] H. Yuan, T. Wang, G.-L. Wang, W.-L. Ju, and J.-M. Zhang, *J. High Energy Phys.* **08**, 066 (2013).
- [8] D. Hai-Rong, F. Feng, and L. Hai-Bo, *Chin. Phys. C* **39**, 013101 (2015).
- [9] A. Paul, A. de la Puente, and I. I. Bigi, *Phys. Rev. D* **90**, 014035 (2014).
- [10] G. Burdman, E. Golowich, J. A. Hewett, and S. Pakvasa, *Phys. Rev. D* **66**, 014009 (2002).
- [11] S. Fajfer and S. Prelovšek, *Phys. Rev. D* **73**, 054026 (2006).
- [12] S. Fajfer, N. Košnik, and S. Prelovšek, *Phys. Rev. D* **76**, 074010 (2007).
- [13] R. Aaij *et al.* (LHCb Collaboration), *Phys. Lett. B* **757**, 558 (2016).
- [14] R. Aaij *et al.* (LHCb Collaboration), *Phys. Rev. Lett.* **119**, 181805 (2017).
- [15] J. P. Lees *et al.* (BABAR Collaboration), *Phys. Rev. Lett.* **122**, 081802 (2019).
- [16] L. Cappiello, O. Cata, and G. D'Ambrosio, *J. High Energy Phys.* **04**, 135 (2013).
- [17] S. de Boer and G. Hiller, *Phys. Rev. D* **98**, 035041 (2018).
- [18] K. Kodama *et al.* (E653 Collaboration), *Phys. Lett. B* **345**, 85 (1995).
- [19] A. Freyberger *et al.* (CLEO Collaboration), *Phys. Rev. Lett.* **76**, 3065 (1996).
- [20] E. M. Aitala *et al.* (E791 Collaboration), *Phys. Rev. Lett.* **86**, 3969 (2001).
- [21] M. Ablikim *et al.* (BESIII Collaboration), *Phys. Rev. D* **97**, 072015 (2018).
- [22] J. P. Lees *et al.* (BABAR Collaboration), *Phys. Rev. Lett.* **124**, 071802 (2020).
- [23] M. Tanabashi *et al.* (Particle Data Group), *Phys. Rev. D* **98**, 030001 (2018), and 2019 update.
- [24] B. Aubert *et al.* (BABAR Collaboration), *Nucl. Instrum. Methods Phys. Res., Sect. A* **479**, 1 (2002).
- [25] B. Aubert *et al.* (BABAR Collaboration), *Nucl. Instrum. Methods Phys. Res., Sect. A* **729**, 615 (2013).
- [26] J. P. Lees *et al.* (BABAR Collaboration), *Nucl. Instrum. Methods Phys. Res., Sect. A* **726**, 203 (2013).
- [27] D. J. Lange, *Nucl. Instrum. Methods Phys. Res., Sect. A* **462**, 152 (2001).
- [28] T. Sjöstrand, *Comput. Phys. Commun.* **82**, 74 (1994).
- [29] B. F. L. Ward, S. Jadach, and Z. Was, *Nucl. Phys. Proc. Suppl.* **116**, 73 (2003).
- [30] H. Czyz and J. H. Kühn, *Eur. Phys. J. C* **18**, 497 (2001).
- [31] N. Davidson, G. Nanava, T. Przedzinski, E. Richter-Was, and Z. Was, *Comput. Phys. Commun.* **183**, 821 (2012).
- [32] P. Golonka and Z. Was, *Eur. Phys. J. C* **45**, 97 (2006).
- [33] S. Agostinelli *et al.* (GEANT 4 Collaboration), *Nucl. Instrum. Methods Phys. Res., Sect. A* **506**, 250 (2003).
- [34] J. Allison, K. Amako, J. Apostolakis, H. Araujo, P. Dubois, *et al.* (GEANT 4 Collaboration), *IEEE Trans. Nucl. Sci.* **53**, 270 (2006).
- [35] A. J. Bevan *et al.* (BABAR and Belle Collaborations), *Eur. Phys. J. C* **74**, 3026 (2014).
- [36] W. D. Hulsbergen, *Nucl. Instrum. Methods Phys. Res., Sect. A* **552**, 566 (2005).
- [37] B. Aubert *et al.* (BABAR Collaboration), *Phys. Rev. D* **69**, 111104 (2004).
- [38] Y. Freund and R. E. Schapire, *J. Comput. Syst. Sci.* **55**, 119 (1997).
- [39] J. P. Lees *et al.* (BABAR Collaboration), *Phys. Rev. D* **89**, 011102 (2014).
- [40] T. Skwarnicki, *A study of the radiative cascade transitions between the Υ' and Υ resonances*, Ph.D. thesis, Institute of Nuclear Physics, Krakow, DESY-F31-86-02, 1986.
- [41] H. Albrecht *et al.* (ARGUS Collaboration), *Phys. Lett. B* **241**, 278 (1990).
- [42] B. Aubert *et al.* (BABAR Collaboration), *Phys. Rev. D* **71**, 032005 (2005).
- [43] T. Allmendinger *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. A* **704**, 44 (2013).
- [44] G. J. Feldman and R. D. Cousins, *Phys. Rev. D* **57**, 3873 (1998).