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34 **ABSTRACT**

35 Chemical alloying is a powerful approach to tune the electronic structure of semiconductors,
36 and has led to the synthesis of ternary and quaternary two-dimensional (2D) dichalcogenide
37 semiconductor alloys (*e.g.* MoSSe₂, WSSe₂ etc.). To date, most studies have been focused on
38 determining the chemical composition by evaluating the optical properties, primarily via
39 photoluminescence and reflection spectroscopy of these materials in the 2D monolayer limit.
40 However, a comprehensive study of alloying in multilayer films with direct measurement of
41 electronic structure, combined with first principles theory, is required for a complete understanding
42 of this promising class of semiconductors. We have combined first-principles density functional
43 theory calculations with experimental characterization of MoS_{2(1-x)}Se_{2x} (where *x* ranges from 0 to
44 1) alloys using X-ray photoelectron spectroscopy to evaluate the valence and conduction band
45 edge positions in each alloy. Moreover, our observations reveal that the valence band edge energies
46 for molybdenum sulfide/selenide alloys increase as a function of increasing selenium
47 concentration. These experimental results agree well with the results of density functional theory
48 calculations showing a similar trend in calculated valence band edges. Our studies suggest that
49 alloying is an effective technique for tuning the band edges of transition-metal dichalcogenides,
50 with implications for applications such as solar cells and photoelectrochemical devices.

51 ■ INTRODUCTION

52 Two-dimensional (2D) transition-metal dichalcogenides (TMDCs) with the general formula MX_2
53 ($M = \text{Mo}, \text{W}; X = \text{S}, \text{Se}, \text{Te}$) have been widely explored as a material system for nanoelectronics,
54 optoelectronics, and catalysis. Their layered structure and van der Waals bonding determine their
55 electron transport properties and gives rise to strong exciton confinement.¹⁻⁴ In particular, TMDCs
56 exhibit layer-dependent optoelectronic properties from a direct energy gap (monolayer) to an
57 indirect energy gap (multilayer and bulk), which makes this class of semiconductors versatile for
58 optoelectronic devices. TMDCs are also catalytically active for the hydrogen evolution reaction
59 (HER), due to the presence of chalcogen vacancies and active edge sites, as well as suitable band
60 edge positions relative to the reduction potential for the HER.⁵⁻¹² For instance, MoS_2 is well-
61 known for its superior HER performance in activity and stability, and after selenium incorporation,
62 $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ exhibits further enhanced catalytic HER performance.^{13,14} To date many studies have
63 focused on the optoelectronic properties of monolayer TMDCs, while a comprehensive study of
64 the ground state electronic structure of few-layer TMDC alloys remains lacking. Earlier theoretical
65 studies reveal that the conduction band edge of monolayer MoS_2 is well aligned to the energy level
66 for CO_2 reduction (CO_2R) to methane, while the valence band edge of MoSe_2 is well-aligned for
67 water oxidation.¹⁵⁻¹⁹ Thus, the band edge position of MoS_2 is used as an indicator to predicting its
68 catalytic activity in (photo)electrochemical CO_2R . Synthesis of TMDC substitutional alloys
69 containing Mo, W, S, and Se has been achieved via chemical vapor deposition (CVD) or chemical
70 vapor transport (CVT).²⁰⁻³¹ Several studies have focused on the synthesis, electronic structure, and
71 optical properties of monolayer 2D MoS_2 , MoSe_2 , and MoSSe .³²⁻³⁵ Despite much progress, studies
72 of the influence of TMDC alloy composition on electronic properties are very limited. Although
73 several studies have used first-principles calculations to understand the electronic structure and
74 optical properties of a variety of 2D van der Waals (vdW) heterostructures,^{20,36-41} they have
75 focused on the electronic band structure of monolayer TMDCs. Few-layer TMDCs and their alloys
76 have rarely been examined. Furthermore, most of these prior theoretical studies are independent
77 and have failed to link electronic structure calculations with a direct experimental measurement of
78 the band positions of TMDC alloys. Therefore, a combined experimental and theoretical
79 investigation of the effect of elemental composition on electronic band structure is needed to
80 expand our knowledge of few-layer TMDC materials and enable future optoelectronic and
81 photoelectrochemical applications.

We employ a combined experimental and theoretical approach to the synthesis and characterization of 4-layer (4L) $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ ($0 < x < 1$) alloys and relate chemical composition to the ground state electronic structure of the alloy. Our study provides direct experimental measurements of the valence band (VB) and conduction band (CB) positions of few-layer MoSSe alloys in the multilayer thickness regime supported by first-principles calculations. Significantly, our theoretical calculations consider three distinct atomic configurations of the synthesized MoSSe heterostructures that have not been considered previously. Our correlated experimental and theoretical results demonstrate that all three configurations exhibit an overall upward trend in VB and CB position as the Se fraction increases in the alloys. Overall, our results suggest that chemical alloying enables the systematic tuning of the band edges of molybdenum sulfide/selenide alloys, with implications for a variety of catalytic applications.^{32,42–47}

■ METHODS

2.1 Synthesis of MOCVD-grown MoS_2 . The synthesis of monolayer to few-layer MoS_2 was carried out in a 5 inch (outer diameter) hot-wall quartz tube furnace. We synthesized monolayer MoS_2 on a 4-inch Si wafer with 298 nm thick thermal SiO_2 using metal-organic chemical vapor deposition (MOCVD). Molybdenum hexacarbonyl (MHC) and diethyl sulfide (DES) were selected as chemical precursors for Mo and S, respectively, and then introduced to the furnace by flowing an Ar carrier gas through the bubblers. H_2 and Ar background gases were introduced into the chamber as the carrier gases using separate lines. In the growth process, a 4-inch Si/ SiO_2 wafer was placed on the quartz substrate holder at the center zone of the furnace. We controlled the chamber pressure at 5 torr, growth temperature of 600 °C and growth time of 3 hrs. The flow rates of precursors were 10 sccm for the Ar flowing through the MHC bubbler and 30 sccm for the Ar flowing through the DES bubbler, 85 sccm for H_2 , and 500 sccm for Ar, which were regulated by individual mass flow controllers (MFCs). The bubbler pressures for the precursors were 250 Torr for MHC and 200 Torr for DES. In the 3hr growth time, we were able to achieve the continuous multilayer MoS_2 growth. MHC: STREM 42-1350, DES: Sigma Aldrich 107247.

2.2. Thermal Selenization of As-grown MoS_2 to Form $\text{MoS}_{2(1-x)}\text{Se}_{2x}$. The selenization process was performed under a three-zone quartz tube furnace with 4.5-inch diameter. Selenium powder (Sigma-Aldrich, -100 mesh, 99.99% trace metal basis) was used as the precursor for CVD reaction. In each selenization experiment, the selenium powders with the fixed amount (1.18 g) were placed

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3 113 in a quartz boat at the upstream side of the furnace and then heated up to 300 °C with flowing a
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5 114 H₂/N₂ (1:9) gas mixture at the pressure of 400 torr. The selenium powders were volatile under
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7 115 300°C and became selenium vapors, which were carried with N₂/H₂ mixture gases to the
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9 116 downstream. Each as-grown quadrilayer MoS₂ on SiO₂/Si sample was cut into 1 x 1 cm², and
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11 117 placed on the quartz substrate holder at the center of a hot wall furnace for selenization at 600 to
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13 118 800 °C for 15~60 minutes. When temperature > 600°C, the selenium vapors reacted with H₂ to
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15 119 form H₂Se(g), which further reacted with MoS₂ samples and form MoS_{2(1-x)}Se_{2x}. The control of
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17 120 selenium concentration in MoS_{2(1-x)}Se_{2x} ($x = 0.12, 0.18, 0.23, 0.26, \text{ and } 1$) depended on various
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19 121 selenization temperature and time, which were 600°C, 650°C, 700°C, 750°C for 15 minutes and
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21 122 800°C for 60 minutes, respectively. After synthesis, the furnace was naturally cooled down to
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23 123 room temperature.

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25 124 **2.3 X-ray Photoelectron Spectroscopy Analysis.** XPS measurements were performed using
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27 125 a Kratos Axis Ultra (and Surface Science M Probe) system with a base pressure of < 1x10⁻⁹ mtorr.
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29 126 A monochromatic Al K_α ($\hbar\omega = 1486.69$ eV) source with a power of 150 W was used for all
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31 127 measurements.

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33 128 **2.3.1. Chemical Compositions of MoS_{2(1-x)}Se_{2x}.** We performed the X-ray photoelectron
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35 129 spectroscopy (XPS) to characterize the chemical compositions and valence-band positions of
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37 130 MoS_{2(1-x)}Se_{2x}. The information we acquired here are macroscopic, which averaged the local
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39 131 information since the spot size was about several hundred microns. Core levels were calibrated
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41 132 to Au 4f (84 eV) which was deposited on the SiO₂ substrate and the corresponding binding energy
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43 133 of the Si⁴⁺ 2p core level (102.7 eV) was used as a reference for all subsequent measurements.
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45 134 Overview scans were collected with a pass energy of 160eV whereas detailed scans were
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47 135 performed with a pass energy of 20 eV. CasaXPS was used for background removal (Shirley type)
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49 136 and peak fitting. For the calculation of atomic ratios and the peak intensities, both obtained from
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51 137 peak fitting, were corrected by the analyzer transmission function (TF) and the respective atomic
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53 138 sensitivity factors (ASF) of the core levels.

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55 139 **2.3.2 Valence Band Maximum (VBM) Measurement.** The valence band edge is the region
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57 140 close to zero binding energy (maximum kinetic energy). The valence band maximum position
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59 141 (VBM) was determined by fitting a fermi distribution to the edge. The function of the Fermi-Dirac
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142 distribution is described as follows:
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$$f(E) = \frac{A}{1 + e^{\frac{E - VBM}{k_B T \cdot B}}}$$

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6
7

8 145 where k_B is Boltzmann's constant, T is the absolute temperature, A as a step height
9 146 parameter, B as peak broadening, and VBM as the valence band maximum.⁴⁸
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11 148

12 149 **2.3.3 Work Function Measurement** The work function of the sample was determined by
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14 150 measuring the secondary electron cutoff of the sample while biasing the sample with a stabilized
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16 151 voltage source (+9V) to overcome the work function of the analyzer. The spectra were then
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18 152 corrected by the applied bias, and the position of the secondary electron cutoff edge was
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20 153 determined by the interception of a linear fit to the background and a linear fit to the edge. In an
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22 154 intensity vs. kinetic energy plot, this value at the intersection will directly yield the value of work
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24 155 function; for binding energy scale, the value has to be converted by the photon energy.

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26 156 **2.4. Optical Band Gap Measurement.** In order to reduce possible damage during the film
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28 157 transfer process and obtain the information directly from the sample, we perform ellipsometry (J.A.
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30 158 Woollam Co. model VASE) measurements with thin film on glassy carbon support or SiO₂/Si
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32 159 substrate. The angle of incident was chosen to be 65°, 70°, and 75°. We then extract the optical
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34 160 constant n (refractive index) and k (extinction coefficient) of the synthesized MoS_{2(1-x)}Se_{2x} samples.
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36 161 The optical band gaps are defined at the rise of the extinction coefficient.

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38 162 **2.5. Materials and Morphological Characterizations.** Raman and room temperature
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40 163 photoluminescence (PL) measurements in ambient condition using Renishaw M1000 Micro
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42 164 Raman Spectrometer System with 514 nm laser excitation with 100X objective and 1% power
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44 165 density (spot size is about 1 μ m). For wavelength calibration, we measured the Si peak at 520.7
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46 166 cm⁻¹ on an undoped silicon wafer as the reference. The emitted Stokes Raman signal was collected
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48 167 by a 0.85 N.A. of 100x objective from Leica in 2400 lines/mm grating for Raman measurement
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50 168 and 1800 lines/mm grating for PL measurement. Scanning electronic microscopy (SEM) images
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52 169 were obtained with a FEI Nova NanoSEM 450 microscope.

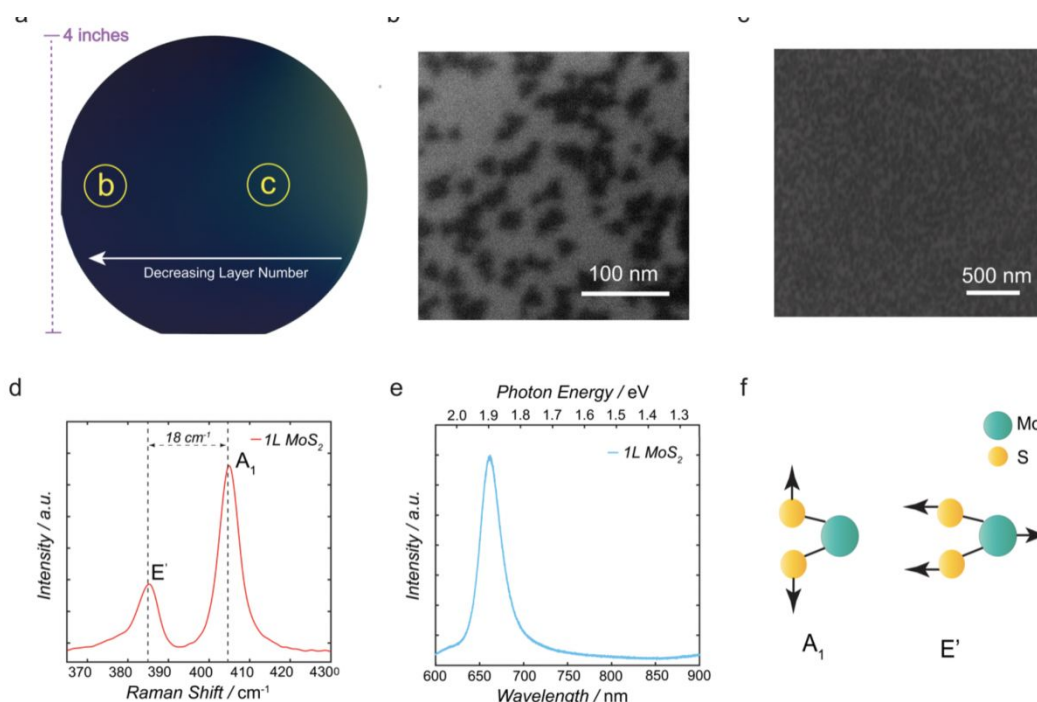
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54 170 **2.6 DFT Calculations.** First-principles ionic and electronic structure calculations were
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56 171 performed using density functional theory (DFT) with a plane-wave basis and projector augmented
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58 172 wave (PAW) potentials⁴⁹ as implemented by the Vienna *ab-initio* simulation package (VASP).^{50,51}
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60 173 All calculations are performed in the generalized gradient approximation as implemented by
174 Perdew, Burke and Ernzerhof (PBE),⁵² with additional van der Waals corrections using rev-vdW-

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3 175 DF2 to incorporate dispersion forces.⁵³ For homogeneous alloys and vertical heterostructures a
4 176 plane-wave energy cut-off of 750 eV was used with a 10x10x1 Gamma-centered k-point mesh for
5 177 structural relaxations, band structures and slab calculations. The most stable homogeneous alloys
6 178 were determined by considering all symmetrically distinct arrangements of S and Se in a 2x2x1
7 179 bulk supercell. A quadrilayer slab with 20 Å of vacuum was used for structural relaxation of the
8 180 vertical heterostructures. All forces were converged to < 0.001 eV/Å. Calculations of the
9 181 electrostatic potential along the c-axis of quadrilayer slabs with 20 Å of vacuum for each system
10 182 were performed to determine absolute band edges. The absolute position of the Fermi energy
11 183 relative to the electrostatic energy of the vacuum region was taken as the valence band maximum
12 184 (VBM). The conduction band minimum (CBM) was determined by adding the bulk band structure
13 185 to the slab VBM. For small lateral domain calculations bulk structural relaxations were performed
14 186 with a plane-wave energy cut-off of 750 eV used with a 2x2x1 Gamma-centered k-point mesh on
15 187 a 10x10x1 supercell of MoS₂ with regions of MoSe₂ of varying size. All forces were converged
16 188 to < 0.01 eV/Å. A quadrilayer slab calculation with 15 Å of vacuum was performed using a 600
17 189 eV energy cut-off and a 2x2x1 k-point mesh to determine the absolute positions of the VBM and
18 190 CBM. An empirical correction scheme was further applied to the calculated band edges.
19 191 The differences between the experimental band edges and DFT band edges for pristine
20 192 MoS₂ and MoSe₂ were calculated, then linearly interpolated across the alloys as a
21 193 function of Se content x in order to generate empirical corrections for the DFT band edges.
22 194

195 ■ RESULTS AND DISCUSSION

196 The synthesis of two-dimensional molybdenum disulfide (MoS₂) thin-films was performed via
197 metal-organic chemical vapor deposition (MOCVD) on four-inch silicon (Si) wafers with 298-nm-
198 thick thermal oxide (SiO₂) (Figure 1a). The MoS₂ film thickness varied as a function of distance
199 across the wafer substrate from multilayered films (> 5 layers) to sub-monolayer flakes. The
200 microstructure of as-grown MoS₂ films was examined under scanning electron microscopy (SEM).
201 SEM images from different portions of the substrate are shown in Figure 1b and Figure 1c, where
202 sub-monolayer MoS₂ islands can be observed with lateral sizes of approximately 20 nm.
203 Previously, many studies have reported the optical properties, such as band gap of monolayer

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3 204 MoS₂ is about 1.8 eV and few-layer samples range from 1.2-1.6 eV. We further characterized the
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5 205 MoS₂ film thickness and its crystallinity using Raman spectroscopy and photoluminescence (PL)
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7 206 as shown in Figure 1d and Figure 1e, respectively. Figure 1d shows the Raman spectrum obtained
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9 207 from monolayer MoS₂ films excited by 514 nm continuous wave laser irradiation at room
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11 208 temperature. The separation between two signature vibrational modes is used to identify the MoS₂
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13 209 film thickness: the A₁ (out-of-plane) and E' (in-plane) vibrational modes for sub-monolayer (<1L)
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15 210 and monolayer (1L), while the A_{1g} and E_{2g} modes are used for multilayer (ML) MoS₂, respectively.
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17 211 These two A₁(A_{1g}) and E'(E_{2g}) vibrational modes are illustrated in Figure 1f. The A₁ and E'
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19 212 vibrational modes for 1L MoS₂ are located at 404.8 and 386.8 cm⁻¹ with peak frequency (cm⁻¹)
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21 213 separations of around 18 cm⁻¹, while those separations for ML MoS₂ are around 19-24 cm⁻¹,
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23 214 depending on the film thickness.⁵⁴ Two Raman modes, E_{2g}¹ and A_{1g}, exhibit a dependence on layer
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25 215 number, with the frequency of the former decreasing and that of the latter increasing with
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27 216 thickness.⁵⁵ From the observed separation of 18 cm⁻¹ in Figure 1d, we estimate that our MoS₂ films
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29 217 in the thinnest regions (b) are (sub)monolayer in thickness. In region (c), the films likely exist as
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31 218 quadrilayers (4L) based on the wavenumber difference between the A₁(A_{1g}) and E'(E_{2g}) Raman
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33 219 peaks, as shown in Figure S1. Figure 1e illustrates a room temperature photoluminescence (PL)
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35 220 spectrum of 1L MoS₂ on the SiO₂/Si substrate, where we observe a strong PL signal at 661.8 nm
36
37 221 corresponding to the A₁ direct exciton transition at 1.85 eV (670 nm).¹⁰ This PL data is consistent
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39 222 with our Raman spectroscopy observations (Figure 1d) for 1L MoS₂.⁵⁷ The strong PL intensity
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41 223 observed is consistent with the high radiative efficiency expected for 1L MoS₂ due to the direct
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43 224 energy gap.⁵⁴⁻⁵⁶
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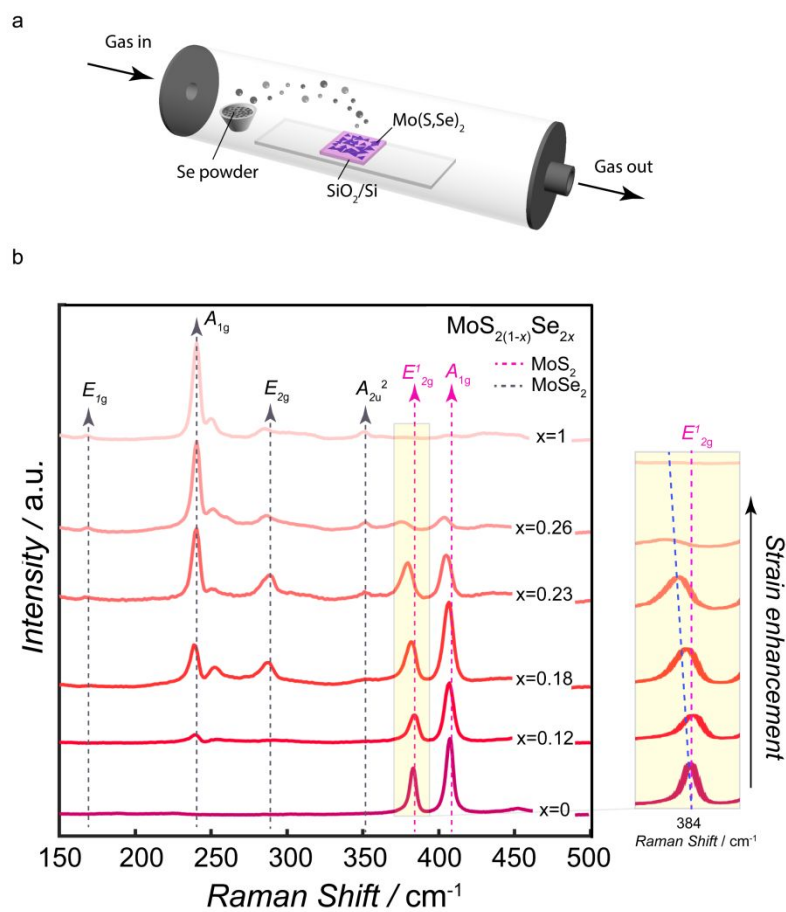
229 **Figure 1.** Metal organic chemical vapor deposition (MOCVD)-grown monolayer and multilayer
 230 MoS₂ compounds and their characterizations. (a) Monolayer and few-layer MoS₂ flakes were
 231 grown on a 4-inch SiO₂/Si substrate. (b) Scanning electron microscopy (SEM) images show the
 232 monolayer MoS₂ and (c) multilayer MoS₂ flakes homogeneously deposited on the silicon wafer.
 233 (d) Raman and (e) photoluminescence spectra were acquired under 514-nm laser excitation. (f)
 234 The in-plane (E') and out-of-plane (A₁) vibration modes that exist in MoS₂ compounds.

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237 To prepare MoS_{2(1-x)}Se_{2x} compounds, we use an established high-temperature selenization
 238 method by flowing selenium vapors reacting with hydrogen in a mixed stream to form H₂Se(g)
 239 under a N₂ environment. Figure 2a shows the experimental setup used for the selenization of
 240 quadrilayer MoS₂ to form various MoS_{2(1-x)}Se_{2x} alloys (x ranges from 0, 0.12, 0.18, 0.23, 0.26 to
 241 1), where x denotes the atomic fraction of selenium in the alloy. By control of the synthesis
 242 temperature, pressure, time, as well as the amount of selenium powder, MoS_{2(1-x)}Se₂ with higher
 243 selenium fractions (0.26 < x < 1) can be further obtained. More synthetic details of the selenization
 244 processes are described in the Methods section. Figure 2b shows the trends in the Raman spectra
 245 for MoS_{2(1-x)}Se_{2x} films as a function of selenium composition. The Raman data indicates a clear
 246 trend associated with the transition with increasing Se composition from pure MoS₂ to alloyed

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3 247 MoSSe and subsequently to pure MoSe₂. The formation of MoSSe alloys is inferred from the
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5 248 presence of Raman peaks located at 287 cm⁻¹ and 239 cm⁻¹ (denoted as E' and A₁ modes of MoSe₂,
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7 249 respectively), which suggest the formation of Mo-Se bonds, along with the weaker signals for
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9 250 MoS₂ at 404.8 cm⁻¹ and 386.8 cm⁻¹ (assigned to A₁ and E' modes of MoS₂) as the Se composition
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11 251 increases.⁵⁸⁻⁶¹ In MoS_{2(1-x)Se_{2x}} alloys for $x = 0.12$ to 0.26 , we observed both MoS₂ and MoSe₂
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13 252 peaks coexisting in the Raman spectra, which clearly shows the evidence of Mo-S and Mo-Se bond
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15 253 formation. Most importantly, no MoS₂ characteristic Raman peaks are shown in MoSe₂ (when $x=1$
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17 254 in MoS_{2(1-x)Se_{2x}}), which indicates selenization process is complete and homogeneous. The
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19 255 signatures associated with MoSSe alloys are accompanied by MoS₂ and MoSe₂ Raman signals,
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21 256 which suggests that some portions of the MoS₂ film are selenized to MoSSe. By increasing the
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23 257 selenium composition, full conversion to MoSe₂ films is achieved. Notably, we observe an obvious
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25 258 red-shift of the E'_{2g} mode of MoS₂ with increasing selenium incorporation, indicating the presence
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27 259 of strain in the resultant alloy.^{62,63} As the Se content increases in strained regions of the film, we
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29 260 expect that the Mo-S mode would be softened by the interactions between Se and S. On the
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31 261 contrary, in such regions the Se-Mo modes would be strengthened, leading to a red shift in their
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33 262 modal frequencies.⁶⁴
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Figure 2. Synthesis of $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ from pristine MoS_2 . (a) Experimental setup of the vapor-phase thermal selenization process. (b) Raman spectroscopy was utilized to examine the formation of $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ compounds. It shows an evolution of structure and chemical composition from MoS_2 to MoSSe and eventually MoSe_2 .

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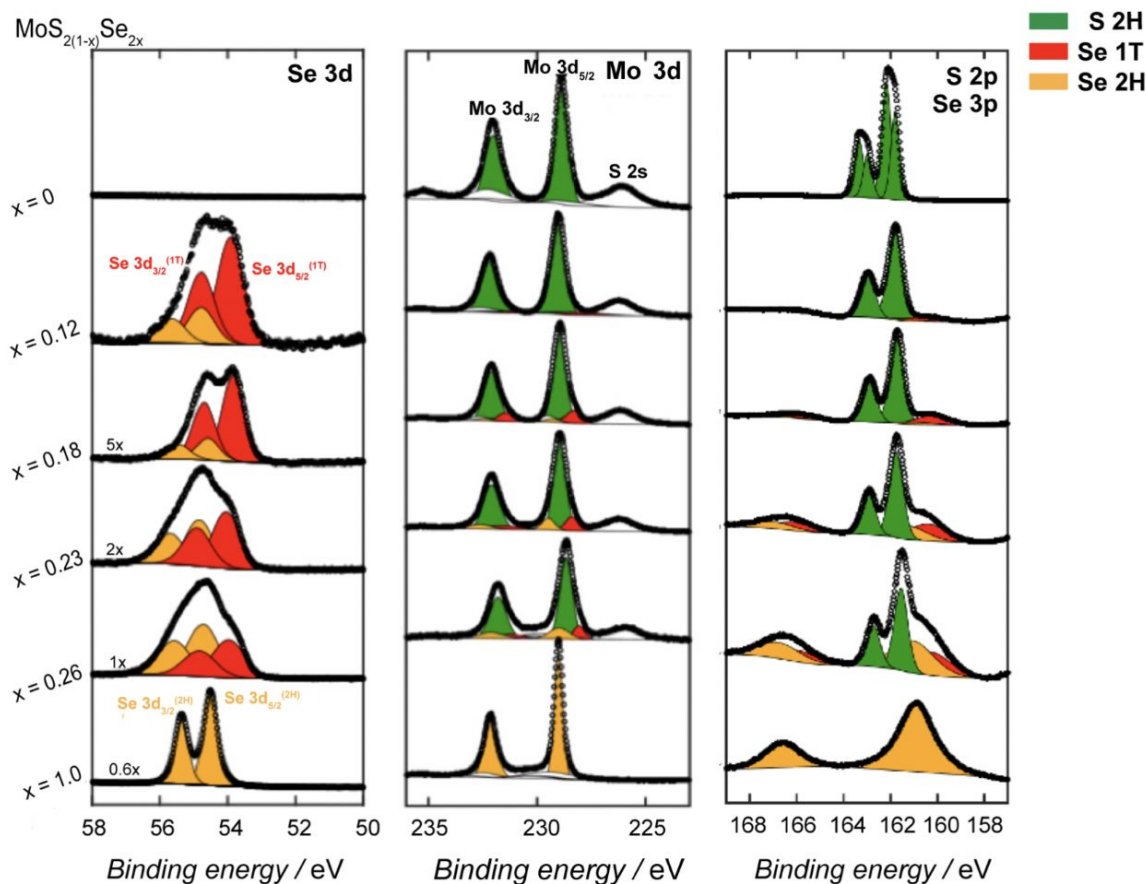
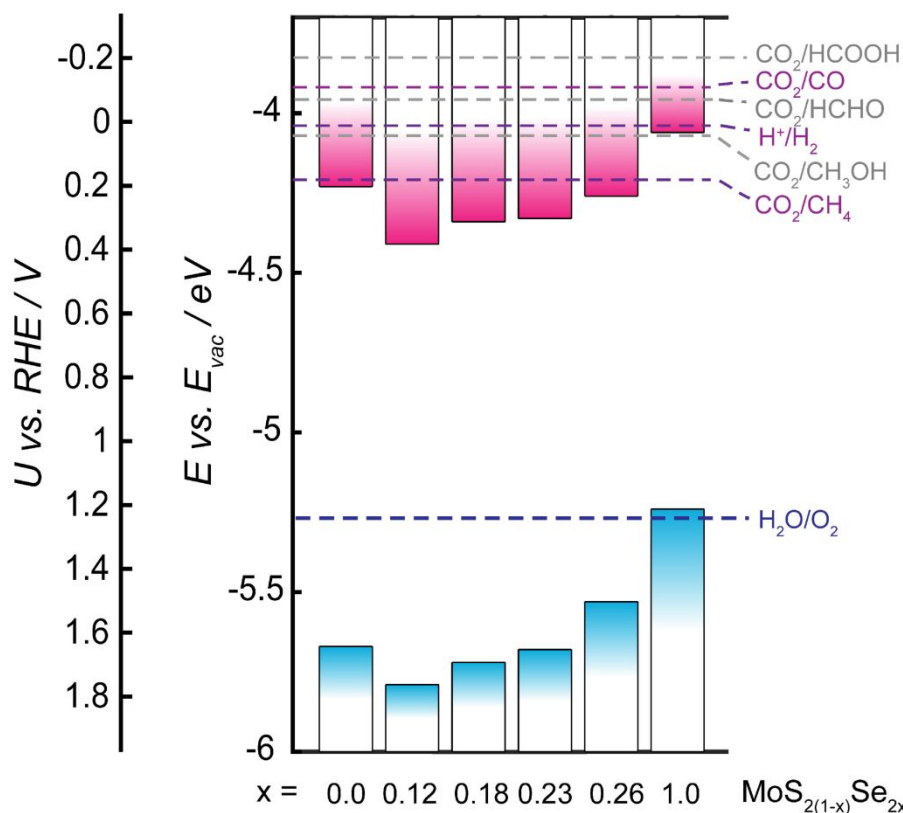


Figure 3. Chemical composition evolution of $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ as a function of Se concentrations monitored through XPS.

To probe the chemical composition as well as the phases of $\text{MoS}_{2(1-x)}\text{Se}_2$ heterostructures, we performed X-ray photoelectron spectroscopy (XPS) for $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ alloys, measuring the Mo 3d, Se 3p, Se 3d and S 2p core levels, as shown in Figure 3. The Se composition in few-layer MoSSe was determined by XPS measurements in which the X-ray penetration depth can generally reach several nanometers to acquire signals from the entire sample. We measured the Se concentrations at various spots in each alloy and averaged the obtained XPS signals over the entire sample. More details of MoSSe alloy compositions are listed in Table S1 in the Supporting Information. No significant change is observed in the Mo 3d core level for MoSSe alloys, yet the Se 3p core level spectra change in both peak shape and energy. As expected, Se 3p peak intensities increase as more selenium is incorporated into the $\text{MoS}_{2(1-x)}\text{Se}_{2x}$ alloys. This XPS result is consistent with the

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3 286 Raman measurements in Figure 2. The 2H/1T phase transition has been reported to occur in MoS₂
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5 287 upon sodium or lithium ion-intercalation, thermal annealing, laser irradiation, or mechanical
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7 288 strain.^{65–67} In the present work, we attribute the partial phase transition from 2H to 1T phase in
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9 289 MoS_{2(1-x)Se_{2x}} to the high temperature selenization process.^{65–67} To distinguish between the 2H and
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11 290 1T phases, we examine the X-ray photoemission spectra in the region of coexistence of the 2H and
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13 291 1T phases, focusing particularly on the Se 3p core level spectra in selenized MoS_{2(1-x)Se_{2x}} alloys
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15 292 at $x = 0.12$ to 0.26 . The Se 3p-2H peaks are located at 55.51 eV and 54.65 eV, while Se 3p-1T
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17 293 peaks are located at 54.69 eV and 53.83 eV binding energy, respectively. Moreover, since the S 2p
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19 294 and Se 3p core levels share a similar binding energy range, the contributions for the Se 3p-1T and
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21 295 Se 3p-2H peaks have been included along with their normalized intensities. The crystal structures
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23 296 of the 2H and 1T phase of MoS₂ exhibit different symmetry, with trigonal prismatic coordinated
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25 297 structure for 2H-MoS₂, and an octahedral coordinated structure for 1T-MoS₂ and the less common
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27 298 distorted 1T'-MoS₂.^{68,69} At present, this evidence of a phase transition from MoS₂ to MoS_{2(1-x)Se_{2x}}
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29 299 indicates the breaking of the original symmetry of MoS₂ after selenization, which could potentially
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31 300 change the electronic properties of the alloy materials.^{68,69} Notably, although a clear Raman
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33 301 signature of the 1T phase of MoS_{2(1-x)Se_{2x}} is lacking in our samples, our XPS spectra (Figure 3)
34
35 302 indicate that the 1T phase is generated upon selenization. The absence of an observed 1T phase in
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37 303 Raman spectra versus observation of the 1T phase in XPS can be attributed to the higher sensitivity
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39 304 of XPS to the 1T phase, which has been shown in earlier studies.⁵⁷ The 1T phase of MoX₂ (X = S,
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41 305 Se) is known to be metastable in an air ambient atmosphere, and transformation back into the
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43 306 more stable 2H phase is thermodynamically preferred.^{70,71} As to the electronic properties, earlier
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45 307 studies show that the metallic 1T phase of the MoSSe alloy exhibits a higher electrical conductivity
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47 308 as well as enhanced catalytic performance for energy conversion applications.^{69,72–74}
48
49 309 Corresponding impacts of the 2H → 1T phase transition on the band edge of MoX₂ electronic
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51 310 structures measurements are also expected.^{67,75–79} Due to the coexistence of two phases, we expect
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53 311 the interfaces between two the phases to comprise both in-plane and out of plane heterostructures,
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55 312 which further alters and complicates electron transport and the corresponding electronic structures
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57 313 in our alloy samples.
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320 **Figure 4.** Experimental measurements of electronic band structures of few-layered MoS_{2(1-x)}Se_{2x}
 321 films. All the energetic levels are relative to the reduction potentials (vs. RHE) for aqueous CO₂
 322 reduction reaction (pH = 6.8). The valence band edges were derived by XPS and optical band gaps
 323 are measured by ellipsometry

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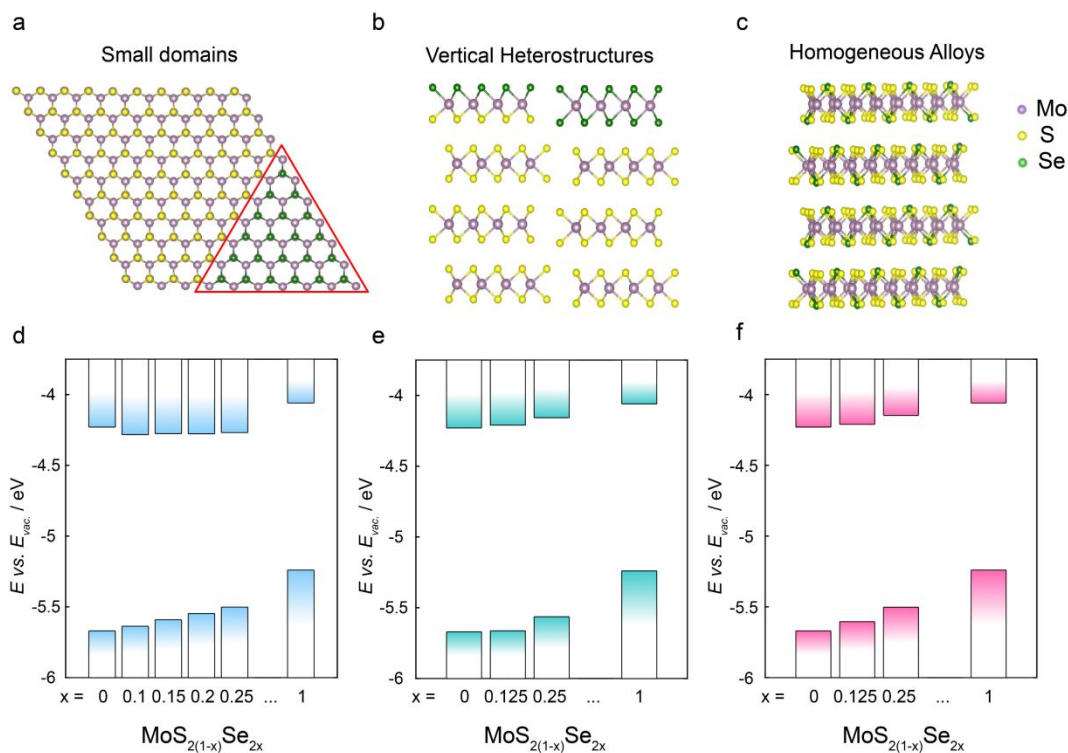
326 To gain insight into the electronic structure, the valence band maxima of MoS_{2(1-x)}Se_{2x} were
 327 measured as shown in Figure S2 and Figure S3. Figure 4 illustrates the experimentally-derived
 328 energy level diagram depicting the valence band maximum (VBM) and the conduction band
 329 minimum (CBM) for our synthesized quadrilayer MoS_{2(1-x)}Se_{2x} alloys (where x ranges from 0 to
 330 1) with respect to the vacuum level, E_{vac} (defined as 0 eV). The MoS_{2(1-x)}Se_{2x} films are synthesized
 331 via selenization of the MoS₂ films (see Methods). In this study, our focus was to evaluate the band
 332 structure of thicker films for the eventual development of devices, since multiple layers (*e.g.* 4L)
 333 are expected to offer more robust stability needed for device applications. Therefore, here we use
 334 4L MoS₂ as the starting material to ensure a continuous film of sufficient thickness on the substrate.
 335 The thickness of the MoS₂ film is indicated by the energy difference of the A₁(A_{1g}) and E'(E_{12g})

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3 336 Raman peaks as shown in Figure S1. In addition, we have depicted the band edge positions relative
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5 337 to the standard reduction potentials, and relative to the reversible hydrogen electrode (RHE), for
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7 338 several CO₂ reduction reactions at a solution of pH = 7 to help visualize the potential utility of
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9 339 MoS_{2(1-x)}Se_{2x} alloys for CO₂R applications. All the band edge positions in Figure 4 were
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11 340 experimentally derived from X-ray photoelectron spectroscopy (XPS) of valence band position
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13 341 and ellipsometry measurements to estimate the bandgap in our synthesized quadrilayer MoS₂₍₁₋
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15 342 _{x)}Se_{2x} films (see Methods and Figure S2 and Figure S3). From these data, we observe an increase
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17 343 in the valence band and conduction band edges in MoS_{2(1-x)}Se_{2x} alloys with increasing Se
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19 344 composition. This increase in band edge energy qualitatively parallels the trend predicted by
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21 345 previous calculations for monolayer transition metal dichalcogenides,^{80,81} and our measurements
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23 346 show lower valence band and conduction band edges for the alloys relative to MoS₂ and MoSe₂.
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25 347 We note that all measured experimental values are relative to the vacuum level.

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27 348 In our synthesized MoSSe alloys, we observed the coexistence of both 1T and 2H phases.
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29 349 Despite the distinct electronic properties of these two phases,^{75,82} the VB and CB positions of the
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31 350 alloy films are expected to be dictated by their overall chemical composition, regardless of the size
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33 351 of individual domains within the MoSSe alloys. We emphasize that the acquired VB and CB data
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35 352 is the result of accumulated signals averaged from several alloyed domains in the XPS and
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37 353 ellipsometry measurements due to the relatively big spot sizes sampled by these two measurements
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39 354 (about 100 μm and 1-5 mm, respectively). The materials characterization was conducted on the
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41 355 entire film and averaged together to elucidate band edge positions. Considering the relatively close
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43 356 correspondence between experiment (Fig. 4) and theory (see Fig. 5), we expect that the 1T impurity
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45 357 does not significantly contribute to the VB and CB edge position.

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47 358 In general, the electronic structure of MoS_{2(1-x)}Se_{2x} alloys could vary significantly based on the
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49 359 chemical composition, phase distribution, and atomic arrangement in the alloy. Based on our
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51 360 experimental observations, we hypothesized that MoS_{2(1-x)}Se_{2x} alloys could have three possible
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53 361 atomic structures and we performed first-principles density functional theory (DFT) calculations
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55 362 to predict band edge positions (Figure 5) corresponding to each of these structures. The first
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57 363 structure shown in Figure 5a assumes a structure comprising MoSe₂ domains within an MoS₂
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59 364 matrix. The second structure in Figure 5b shows that of a ‘vertical heterostructure’ with selenium
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365 atoms dominating the top layer while the rest of layers underneath has sulfur atoms as the dominant
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chalcogen species. The third possible structure is a homogeneous alloy, with sulfur and selenium

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3 367 distributed uniformly throughout the layers. For vertical heterostructures and homogeneous alloys,
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5 368 the conduction band edge becomes progressively more positive as a function of increasing
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7 369 selenium doping. However, in the case of MoSe₂ domains embedded in MoS₂, the conduction band
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9 370 edge varies less significantly with increasing Se content and is in all cases more negative than the
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11 371 conduction band edge of the binary alloy structures. This result of the upward trend of the
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13 372 conduction band edges is consistent with prior studies of lateral heterostructures of monolayer
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15 373 MoS₂ and MoSe₂.^{83,84} These theoretically-predicted trends suggest that our samples may consist of
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17 374 a combination of different types of MoS_{2(1-x)}Se_{2x} structures, which is consistent with our
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19 375 experimental observations of the conduction band edges shown in Figure 4. Based on the
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21 376 calculation results, it is expected to see the upward trend of the VB edges with increasing Se
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23 377 concentration; however, we note that the VB edge positions determined from experiment do not
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25 378 match perfectly with those predicted from DFT calculations. Although the overall trend is in
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27 379 accordance with theoretical predictions, the $x = 0.12$ sample is clearly an outlier. The discrepancy
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29 380 between experiment and theory for this sample, compared to pure MoS₂, is likely attributable to
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31 381 the coverage of adsorbates (oxygen and/or hydroxyl groups) on the surface of the alloy. The MoS₂
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33 382 sample is prepared in a single step via MOCVD as opposed to the alloyed samples, which require
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35 383 two steps of thermal cycling at elevated temperature to prepare. Since the alloying selenization
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37 384 procedure is performed via chalcogen ion exchange at an elevated temperature, this process may
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39 385 lead to the introduction of additional defects that serve as sites for adsorption of water or oxygen
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41 386 from the ambient.^{85,86} It is known that the adsorbates on MoS₂ can lead to higher work function
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43 387 values as compared to pristine MoS₂.⁸⁷ This aspect appears to be the most plausible explanation
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45 388 for why the $x = 0.12$ alloyed sample has a higher work function than the MoS₂ sample. Despite
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47 389 this small discrepancy between experiment and theory, the overall trend of the shifting VB position
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49 390 with increasing Se content is consistent with theory, confirming that selenization is a viable
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51 391 strategy for tuning the band edges of TMDC materials. We emphasize that this discrepancy
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53 392 between experiment and theory is minor and doesn't affect our general conclusions.
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396 **Figure 5.** The calculated band edges of three different hypothesized MoS₂(1-x)Se_{2x} structures by
 397 using first-principles density functional theory (DFT) with van der Waals (vdW) corrections (rev-
 398 vdW-DF2) using the Vienna *ab-initio* simulation package (VASP). Three structures include (a)
 399 small domains with both MoS₂ and MoSe₂, (b) Vertical heterostructures (MoS₂ as the majority and
 400 MoSe₂ as the top layer), and (c) homogeneous alloy. (d), (e), (f) are the corresponding calculated
 401 band structures of (a), (b) and (c), respectively.

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404 Further insights into the crystalline structure of MoS₂(1-x)Se_{2x}, were gained from the
 405 comparison of Raman spectra with DFT results as shown in Figure 2 and Figure 5, respectively.
 406 In Figure 5, the calculated band edge positions of three different proposed structures of MoS₂(1-
 407 x)Se_{2x} all increase as a function of Se concentration for homogeneous alloys and vertical
 408 heterostructures while lateral domains demonstrate a downward shift of the conduction band edges
 409 relative to the pristine structures. Additionally, two sets of composition-dependent vibrational
 410 modes (Figure 2b) were observed in the Raman spectra from pure MoS₂ (at higher frequency, 350-
 411 410 cm⁻¹) and MoSe₂ (at lower frequency, 200-300 cm⁻¹), rather than a single set of peaks for
 412 ternary MoSSe phases. Such two-mode behavior shown in Raman spectra indicate that our MoS₂(1-
 413 x)Se_{2x} compounds are likely composed of coexisting MoS₂ and MoSe₂ domains (both horizontally

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3 414 and vertically) instead of a single-phase, atomically-homogeneous, substitutional MoSSe alloy.
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5 415 The above results are consistent with the experimental band edge data in Figure 4. Incorporation
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7 416 of Se into pristine MoS₂ not only changes the macrostructure of MoS_{2(1-x)}Se_{2x}, but also generates
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9 417 local strain in the alloy itself, further tailoring the local electronic structure.^{88,89} Figure S4 exhibits
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11 418 the calculated the bond angles of pristine MoS₂, MoSe₂ and MoS_{1.5}Se_{0.5} alloys ($x = 0.25$, denoted
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13 419 as 25% small domains to be representative of the ternary alloy) to show the local elemental bonding
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15 420 of Mo to chalcogens (S and Se). Compared with pristine MoS₂ and MoSe₂ alloys, the MoS_{1.5}Se_{0.5}
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17 421 alloy consisting of small lateral domains of MoSe₂ exhibits observable bond angle changes at the
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19 422 domain interface in the bond angles denoted as A(A'), B, and C, which implies geometric changes
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21 423 in local atomic structure. Furthermore, the heterojunctions in MoS_{2(1-x)}Se_{2x} alloys also introduce
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23 424 strain at the MoS₂/MoSe₂ interface. A local phase transition is often accompanied by the existence
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25 425 of local strains, particularly in monolayer or few-layer transition metal dichalcogenide materials.⁶⁷
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27 426 This type of strain has been demonstrated to alter local electronic structure and adsorption energy
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29 427 of the material, ultimately influencing the catalytic activity of TMDC alloys.^{90,91} Our correlation
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31 428 of the electronic structure of a series MoS_{2(1-x)}Se_{2x} compounds suggests potential applications for
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33 429 catalysis such as in electrochemical CO₂R. Nevertheless, additional calculations of the adsorption
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35 430 energy, solvation effects, structural phase inhomogeneity, kinetic barriers and activation energy
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37 431 for CO₂R are needed to evaluate the potential catalytic performance of these materials, and will be
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39 432 the subject of ongoing studies.
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435 ■ CONCLUSIONS

436 In summary, we have synthesized a series of few-layer MoS_{2(1-x)}Se_{2x} alloys through chemical
437 alloying of MOCVD-grown MoS₂ with high-temperature selenization process. Experimental
438 characterization of the MoS_{2(1-x)}Se_{2x} band edges via XPS and ellipsometry indicates that their
439 valence and conduction band edges shift positively with increasing selenium content. The resulting
440 upward shift of their band edges relative to the electrochemical reduction potentials for CO₂
441 reduction suggests their promise as (photo)electrocatalysts for the CO₂RR. In addition to tuning
442 of band edge energies, local strains within the alloy film may lead to altered catalytic activity.
443 Finally, our experimental observation of band edge tuning through selenium incorporation in MoS₂,

444 corroborated by calculations, suggest a new approach for tailoring the optoelectronic properties of
445 TMDC alloys.

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447 ■ ASSOCIATED CONTENT

448 Supporting Information

449 Raman characterizations of MOCVD-grown MoS₂ on 4-inch SiO₂/Si substrate. Valence band and
450 work function spectra measured by XPS. Extinction coefficient of different MoS_{2(1-x)}Se_{2x} samples.
451 Simulated bond angles of pristine MoS₂, MoSSe and MoSe₂ between various bonds, such as S-S,
452 Se-Se, and S-Se in small domain case as well as comparison of calculated bond angles in various
453 compounds. All the chemical compositions of MoS_{2(1-x)}Se_{2x} in this work.

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