

**Supplementary information**

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**Analysis of the limitations in the oxygen reduction activity of transition metal oxide surfaces**

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## Supplementary Information

### Analysis of the Limitations in the Oxygen Reduction Activity of Transition Metal Oxide Surfaces

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## Supplementary Discussion

### High-throughput experiments

While the automation of the serial electrochemistry enables such unprecedented catalyst screening campaigns, several artifacts can occur, prompting our development of automated quality control measures to mitigate their influence of the characterization of catalyst activity. Intermittent electrochemical contact, typically from a bubble blocking contact to any 1 of the 3 electrodes, as well as oscillations or undesirably large working electrode contact area from electrolyte flow deregulation, cause erroneous measurements. Quality control requirements include (i) measured potential must be within 5 mV of intended potential for all but 25 data points acquired during a given technique, (ii) given the consistency of the shapes of CVs, PCA reconstruction of each CV using 9 components, which explained 99% of the variance in the entire dataset, was compared to the measured CV with a requirement of root mean squared difference of no more than 0.01 mA/cm<sup>2</sup>, which removes CVs with anomalous shapes due to intermittent cell issues, and (iii) exploiting prior knowledge that catalytic activity is smooth on the scale of 10 at.% within a given composition space, each measured current density was required to be within 2  $\sigma$  of the composition-smoothed value, where the predicted value and uncertainty were determined by training a Matern Gaussian process regressor ( $\nu = 1.5$ ,  $\alpha = 0.01$ , and  $\text{length\_scale} = 10$  at.%) on the other compositions in the composition space. Measurements and associated current densities meeting these requirements are tabulated in the tables in <https://data.caltech.edu/records/1632>, which include both the full sample list including duplicate compositions as well as the set of catalyst activities after de-duplication using the median

current density for duplicate compositions. This latter table is the source for all data in **Figure 1** and **Supplementary Figures 1** and **3**. Note that these high-throughput TMO samples are expected to have a mixture of different Miller facets since many stable facets (with both low- and higher-index) of complex TMOs appear to have similar calculated surface energies. Therefore, the overall ORR activity of each sample should consist of the activity contribution from various surfaces and sites.

Importantly, the Pt baseline comparisons in **Figure 1** and **Supplementary Figure 1** (from Ref.<sup>1</sup>) used geometric surface area with smooth films whose microscopic surface area was confirmed to be similar *via* CO-stripping experiments. For the non-Pt catalysts in high-throughput experiments, interpretation of surface area measurements such as CO stripping are not standard. For TMO catalysts deposited with similar elements, loading, and thermal processing, we previously characterized electrochemical surface area *via* the electrochemical capacitance, finding that the specific surface area ratio can be up to a factor of 10.<sup>2</sup> As a result, while **Figure 1** and **Supplementary Figure 1** make comparisons on a geometric surface area basis, the analogous comparison on a microscopic surface area basis would result in the Pt baseline remaining approximately the same value while the TMO current densities being lowered by up to a factor of 10, further highlighting the inability of the TMOs to approach the activity of Pt.

## **Supplementary Methods**

### **Statistical analysis on previous literature**

Data for the statistical analysis of the TMOs for acidic ORR were acquired from the representative literatures published over the past twelve years (**Figure 1** and **Supplementary**

**Figures 1-2).**<sup>3-32</sup> These literatures include pure TMOs, oxynitride, oxycarbide, and oxidized carbonitride (CNO) TMOs. For comparison, the corresponding values of pure Pt were acquired from Ref.<sup>1</sup>. Note that the literatures with potential Pt-contamination and the partially oxidized Pt-alloy systems (which possess metallic Pt as their active sites) are excluded from the statistics.

## Computational methods

The density functional theory (DFT) calculated binding energy data consist of our calculations and the previously reported data from our group.<sup>33-36</sup> For the transition metal (111) surfaces (Pt, Pd, Au, Ag, Cu, Rh, Ir, and Ni), 4-layer, 4×4 slabs were considered for calculations. Spin-polarization was employed in the calculations for Ni(111). A (3×3×1)  $k$ -point mesh was used to sample the Brillouin zone of these transition metal surfaces using the Monkhorst-Pack method.<sup>37</sup> The atop-site adsorption of oxygen on transition metal surfaces were relaxed by constraining the x-y direction of O\*, where the O\* was originally placed at the center of the metal-atop position before structural relaxation. The  $E_{HO^*}$  vs  $E_{O^*}$  scaling relation for transition metals used for microkinetic modeling was calculated by picking the most favorable site for each adsorbate (the grey line in **Figure 3b**). The cations of the TMO structures considered in this study include Zr, Hf, Ti, Sn, Ir, Pt, Rh, Nb, Re, Mn, Fe, V, Ni, Co, Pb, Pd, Ru, Sb, Mo, Sr, and La. For the TMOs containing Mn, Fe, Co, Ni, or Mo, spin-polarized DFT+U method<sup>38</sup> was adopted with Hubbard U values in line with the *MaterialsProject*.<sup>39</sup> The electronic energy and structural relaxation were converged to 10<sup>-5</sup> eV and 0.05 eV/Å, respectively. For the rest of the structures, spin-polarizations were tested and used when necessary. All of the bulk structures were obtained from the *MaterialsProject* database.<sup>39</sup> Bulk optimizations were performed for each system

before their surfaces were cut, with the lattice parameters fully relaxed with a  $k$ -point density of at least  $100 \text{ \AA}^{-3}$  (reciprocal lattice volume). For the ORR binding energy calculations, the bottom half layers were fixed to the bulk positions, and the rest layers were relaxed. A vacuum spacing of at least  $12 \text{ \AA}$  was placed in the direction perpendicular to the surface. We used the same  $k$ -point mesh grid as the bulk but setting the  $z$ -direction to be one because of the large vacuum and slab thickness. For all of the computation parameters above, stricter criteria (*i.e.*, larger kinetic cutoff, larger  $k$ -point mesh grid, lower convergent force, thicker layer, and larger unit cell) were tested for each system; no significant difference was found in the binding energies and adsorption geometries. The crystal structure manipulation and input generation were performed using Python Materials Genomics (pymatgen)<sup>40</sup> and Atomic Simulation Environment (ASE)<sup>41</sup> packages. The electronic binding energy data classified by the type of host cation are shown in **Supplementary Figure 6**. The data with the HO binding energy higher than  $\sim 2 \text{ eV}$  are excluded from our analysis. For the free energy diagram analysis shown in **Figure 2**, (111) was used as an example because it is one of the lowest energy facets of  $\text{ZrO}_2$  and  $\text{HfO}_2$ .<sup>42,43</sup>

The electronic binding energies were calculated using the total energies of  $\text{H}_2$  and  $\text{H}_2\text{O}$  as the energy references:

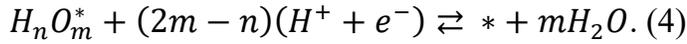
$$E_{O^*} = E_{tot} - E_{bare} - E_{\text{H}_2\text{O}} + E_{\text{H}_2}, \quad (1)$$

$$E_{\text{HO}^*} = E_{tot} - E_{bare} - E_{\text{H}_2\text{O}} + \frac{1}{2}E_{\text{H}_2}, \quad (2)$$

$$E_{\text{HOO}^*} = E_{tot} - E_{bare} - 2E_{\text{H}_2\text{O}} + \frac{3}{2}E_{\text{H}_2}, \quad (3)$$

where  $E_{tot}$  is the total energy of the surface with adsorbate,  $E_{bare}$  is the total energy of a bare surface,  $E_{\text{H}_2\text{O}}$  is the total energy of a  $\text{H}_2\text{O}$  molecule in vacuum, and  $E_{\text{H}_2}$  is the total energy of a

H<sub>2</sub> molecule in vacuum. For most of the TMO surfaces considered in our calculations, the metal-atop was found to be the adsorption site of ORR adsorbate. The reaction free energy diagrams on TMOs were calculated with the computational hydrogen electrode (CHE) method,<sup>44</sup> with the entropic, zero-point energy, and solvation corrections shown in Refs.<sup>44,45</sup>. The free energy data for Pt, Pd, and Au were acquired from Ref.<sup>46</sup>. Surface Pourbaix diagrams were developed to show the most favorable surface state of a TMO as a function of pH and U<sub>SHE</sub>, based on the representation of the adsorption of possible hydrogenated and oxygenated species in ORR:



Then the free energies shown in the surface Pourbaix diagrams ( $G_{SP}$ ) were calculated based on the method described by Refs.<sup>47-49</sup> and the entropic and zero-point energy corrections from Refs.<sup>44,49</sup>:

$$G_{SP} = G_{bare} + mG_{H_2O} - G_{tot} - (2m - n)\left(\frac{1}{2}G_{H_2} - U_{SHE} - 2.303k_B T * pH\right), \quad (5)$$

where  $m$  and  $n$  respectively represent the numbers of oxygen and hydrogen atoms adsorbed on a stoichiometric TMO surface,  $k_B$  is the Boltzmann constant, and  $T$  is the temperature.

The O\*-induced charge density differences ( $\Delta\rho$ ) were plotted using the following equation:

$$\Delta\rho = \rho_{tot} - \rho_{bare} - \rho_{O^*}, \quad (6)$$

where  $\rho_{tot}$  is the charge density of the surface with the adsorbed O,  $\rho_{bare}$  is the charge density of the bare surface, and  $\rho_{O^*}$  is the charge density of the atomic oxygen.

## Kinetic modeling methods

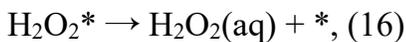
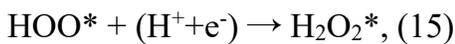
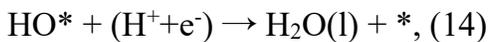
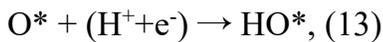
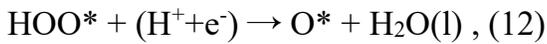
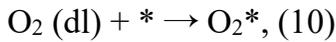
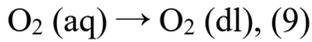
The microkinetic modeling of ORR volcano was based on the method by Hansen *et al.*<sup>50</sup> and Kelly *et al.*<sup>51</sup> using the CatMAP package.<sup>52</sup> The adsorbate coverages were determined by the steady-state approximation method. The rates for intermediate steps were calculated using Equation 7:

$$\text{rate} = k_f \prod \theta_{\text{reac}} - k_r \prod \theta_{\text{prod}}, \quad (7)$$

where  $\theta_{\text{reac}}$  and  $\theta_{\text{prod}}$  are the coverages of reactants and products, respectively. The rate constant  $k$  was calculated as the function of reaction prefactor  $A$  ( $\text{s}^{-1}$ ), activation free energy  $G_a$ , Boltzmann constant  $k_B$ , and reaction temperature  $T$ :

$$k = A \text{Exp}\left[-\frac{G_a}{k_B T}\right], \quad (8)$$

The intermediate reactions considered in the modeling are shown in Reactions 9-16:



Reaction 9 stands for the diffusion of aqueous  $O_2$  via a Nernstian diffusion layer, with the rate of  $8 \times 10^5 \text{ s}^{-1}$  according to the calculations by Hansen *et al.*<sup>50</sup> Reaction 10 represents the adsorption of  $O_2$  on the surface. Reaction 16 represents the desorption of formed  $H_2O_2$  on the surface. Prefactors of both Reactions 10 and 16 were set as  $1 \times 10^8 \text{ s}^{-1}$  due to the solvent reorganization to accommodate  $O_2$  and  $H_2O_2$ .<sup>51</sup> Reactions 11-15 are the proton-electron transfer steps, with the energy of proton-electron pair represented by the energy of half of a  $H_2$  molecule according to the CHE method.<sup>44</sup> Reactions 11-14 represent the standard associative pathway for  $4e^-$  ORR. Reactions 11 and 15 represent the  $2e^-$  ORR process. For Reaction 12, where the O-O is broken along with the protonation, data from Ref.<sup>34</sup>, where  $G_a$  were calculated for the backward reaction, were reformulated to find  $G_a = 0.8(G_{O^*} - G_{HOO^*}) + 1.35 + 0.42U$ . For all other proton-transfers, which did not include any other bond-breaks, we used an intrinsic barrier of 0.26 eV and assumed 0.5 electrons had transferred at the transition state.<sup>53</sup> Prefactors for all of the proton-electron transfer steps were set as  $1 \times 10^9 \text{ s}^{-1}$  to account for solvent reorganization.<sup>51</sup> Previous studies with explicit solvation models have shown that many TMO surfaces are hydrophobic, with the water layer floating above the surface and further stabilizing the adsorption of  $HOO^*$  and  $HO^*$ .<sup>45,53,54</sup> Zenkin *et al.* proposed an empirical theory that TMOs with lower metal electronegativity (*e.g.*, Zr and Hf oxides) tend to be more hydrophobic,<sup>55</sup> in good agreement with many previous experimental measurements.<sup>56-58</sup> Since the most widely reported stable TMOs for acidic ORR have relatively low electronegativity (*e.g.*, Zr, Nb, and Hf oxides), we consider these as hydrophobic TMOs. With the adsorption of  $HO^*$  or  $HOO^*$  under reaction conditions, due to the formed hydrogen bond with the water layer, the hydrophobicity of a TMO surface will be reduced. But since this will further strengthen the binding of  $HO^*$  and  $HOO^*$ ,<sup>45</sup>

water is still less likely to block the surface site. Therefore, H<sub>2</sub>O molecule is not expected to block the surface sites in our modeling. We note that whether TMOs can be hydrophobic or hydrophilic also depends on surface chemistry, which is still an interesting open question in the field. Dissociation reactions have been shown to have little effect on kinetics of ORR, and were therefore ignored.<sup>51</sup> Scaling relations for O\* and HOO\* were input using the scaling relations as discussed in the maintext, with the entropic, zero-point energy, and solvation corrections shown in previous Refs.<sup>44,45</sup>. Scaling relations for all other adsorbates were obtained from Kelly *et al.*<sup>51</sup> Then they were corrected for the RHE scale and electric field using Equation 17:<sup>51</sup>

$$G_{ads} = G_{ads}^{\vec{E}=0} + \mu\vec{E} - \frac{\alpha}{2}\vec{E}^2 - neU_{RHE}, \quad (17)$$

where  $G_{ads}^{\vec{E}=0}$  is the binding energy of an adsorbate calculated without external electric field.

Our recent study showed how electric field effects can predict pH effects in ORR catalysts.<sup>51</sup>

This is because the electric field is dependent on the absolute potential of the electrode, which can be approximated with a parallel-plate capacitor model, as shown in Equation 18:

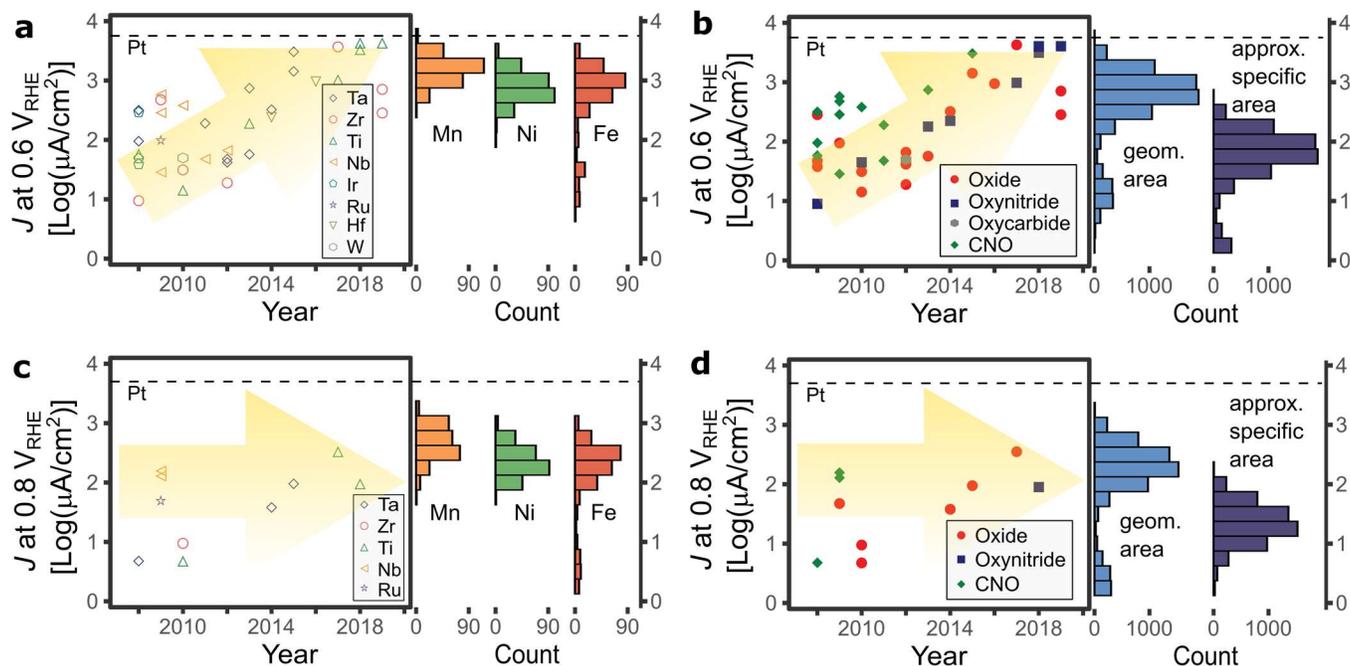
$$\vec{E} = \frac{C_H}{\epsilon\epsilon_0}(U - U_{PZC}), \quad (18)$$

where  $U_{PZC}$  refers to the potential of zero charge of an electrode. There is little to no data on the potential of zero charge of TMOs, meaning our model cannot predict rates at specific pH for a given  $U_{RHE}$ . However, because  $U_{RHE}$  differs from potential vs the standard hydrogen electrode ( $U_{SHE}$ , and absolute potential scale) by only pH, as shown in Equation 19, we can generally predict that higher pH leads to more negative fields, and vice versa.

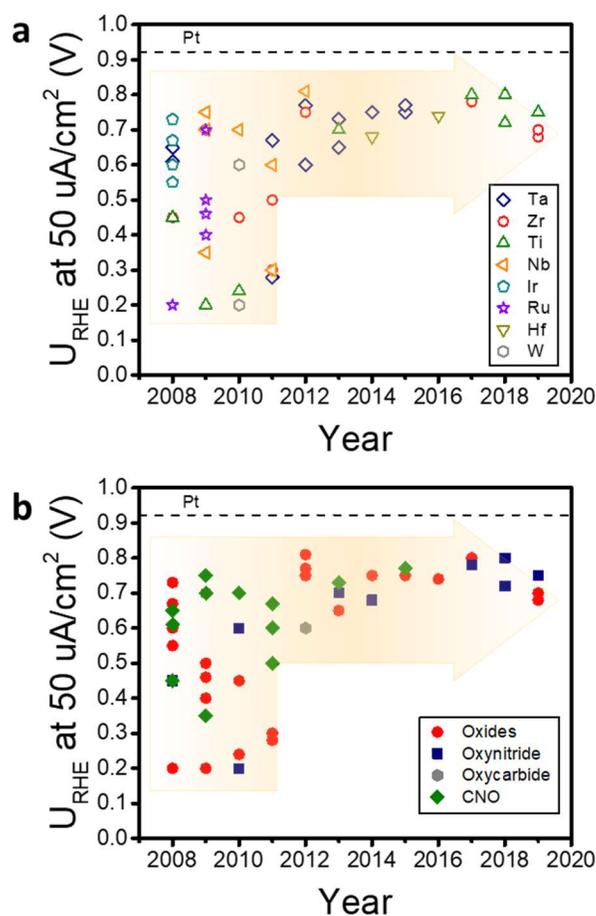
$$U_{SHE} = U_{RHE} - k_B T \ln(10) \times pH, \quad (19)$$

The values of  $\mu$  and  $\alpha$  were acquired from the second order polynomial fitting for the energy of each adsorbate across the range of electric fields, using Equation 17 where  $\vec{E}$  is the applied electric field in the calculations using the Quantum Espresso package.<sup>59</sup> The volcano modeling for metal used the field effects calculations on Au(100) as described previously.<sup>51</sup> For the modeling of TMO volcanos, electric field effect calculations on ZrO<sub>2</sub>(111) and ZrO<sub>2</sub>(100) were considered for higher and low-index TMOs, respectively, with the reason that ZrO<sub>2</sub> is one of the most typical TMOs for acidic ORR and it behaves similar to other widely studied early TMOs.<sup>5</sup> Our electric field effect calculations on other TMOs indicate that the tuning trends of ORR adsorbate bindings are general (**Supplementary Figure 7**), which do not qualitatively alter with different types of TMO surfaces and in turn will not change the qualitative trends of the kinetic volcano. Since a more realistic solvation configuration on TMOs is with the water layer floating above the TMO surfaces due to the hydrophobic nature of many stable ORR TMOs (*e.g.*, Zr, Nb, and Hf oxides, as discussed earlier), the possibility of water blocking the reaction sites is not considered. Calculations with more explicit water molecules were performed with applied fields; no significant change was found in the overall configurations. Our recent study showed that this modeling method leads to excellent agreement with the pH-dependent trends of polarization curves and rotating ring disk electrode current experiments for Pt(111), Au(111), and Au(100).<sup>51</sup> Therefore, given the neglectable change in the water layer configuration with applied field, and the high consistency between our modeling results and experiments, field effect is not expected to qualitatively influence the reaction mechanism. Besides, though water molecules may bind differently to TMOs and to metals, we did not consider an aqueous electrolyte in the present study due to the current good agreement between theory and experiments, and the tremendous

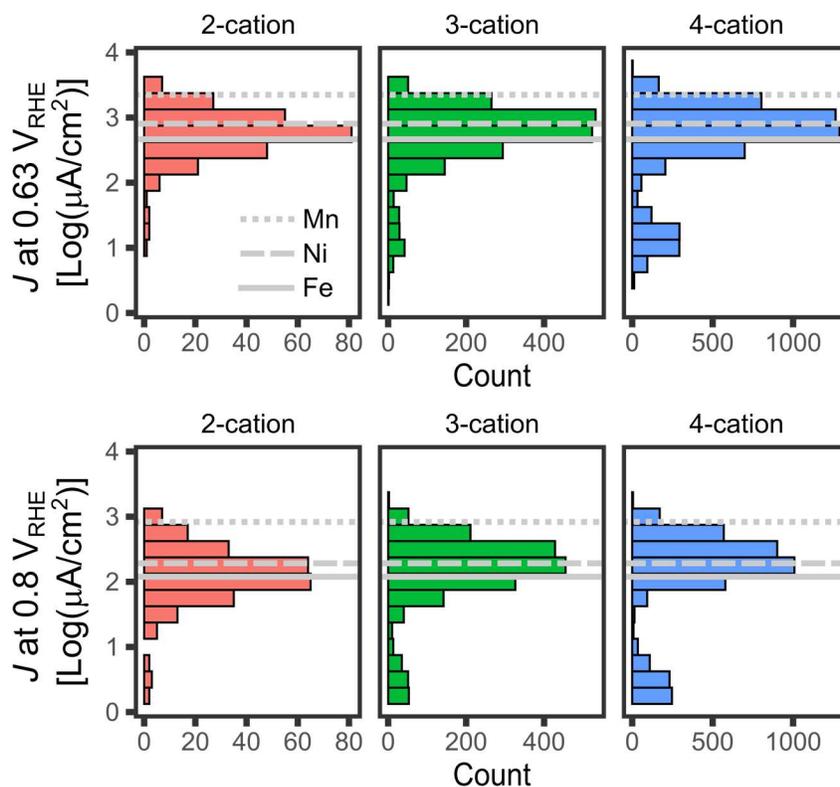
computational cost for liquid-surface interface simulations. In addition, aqueous and double-layer oxygen were set to the energy of 5.19 eV at 0 V without external electric field. The mole fraction of oxygen was set as  $2.34 \times 10^{-5}$ , which corresponds to 1 atm O<sub>2</sub> in equilibrium with water.<sup>51</sup> To more clearly see the effects of each possible rate-determining step, we also included rate-determining step analysis. We used rates solved previously to describe the left and right hand side of the metal volcano and also calculated the rate of Reaction 12 using steady-state assumption (**Supplementary Figure 8**).



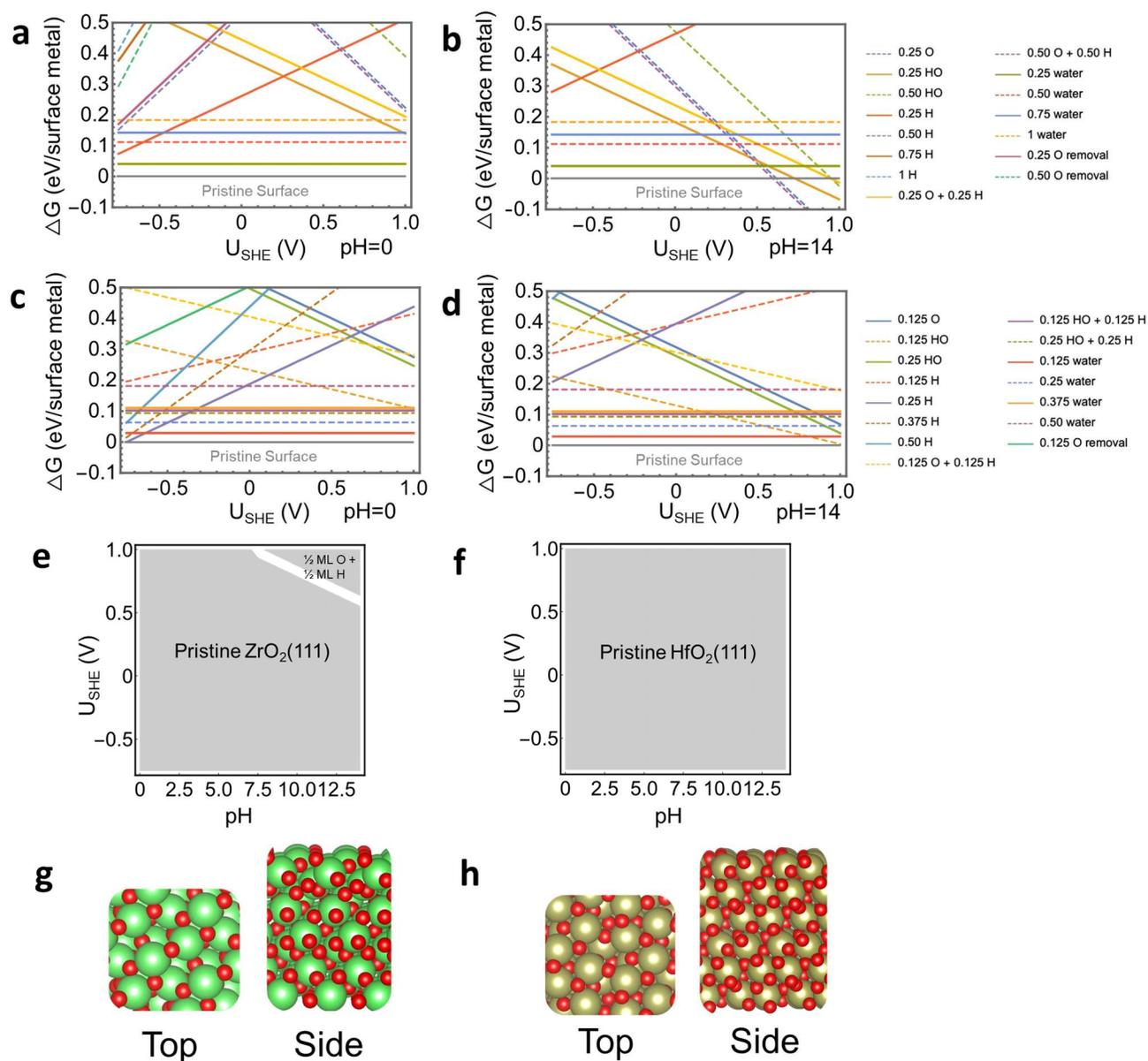
**Supplementary Figure 1. Summary of the performance of TMO catalysts for ORR.** In each subfigure, literature values published over the last twelve years for acidic ORR are shown in a scatter plot and the high-throughput alkaline ORR screening from the present work is shown as histograms. Current densities measured at the potential of  $0.6 V_{\text{RHE}}$  (or  $0.63 V_{\text{RHE}}$ ) are shown labelled by cation system in a) and by host anion in b). Current densities measured at the potential of  $0.8 V_{\text{RHE}}$  are shown labelled by cation system in c) and by host anion in d). For the catalyst compositions synthesized with more than 50% Mn, Ni, or Fe, the geometric current density histograms for each element are shown in a) and c). For the high-throughput screening results, the full dataset of TMOs is shown in b) and d), by geometric current density on the left (blue) and then by specific current density assuming an approximate specific surface area ratio of 10 (purple). The dashed lines are the values of pure Pt catalysts for acidic ORR reported by Ref.<sup>1</sup>. Catalysts with current density below  $1 \mu\text{A}/\text{cm}^2$  are not shown.



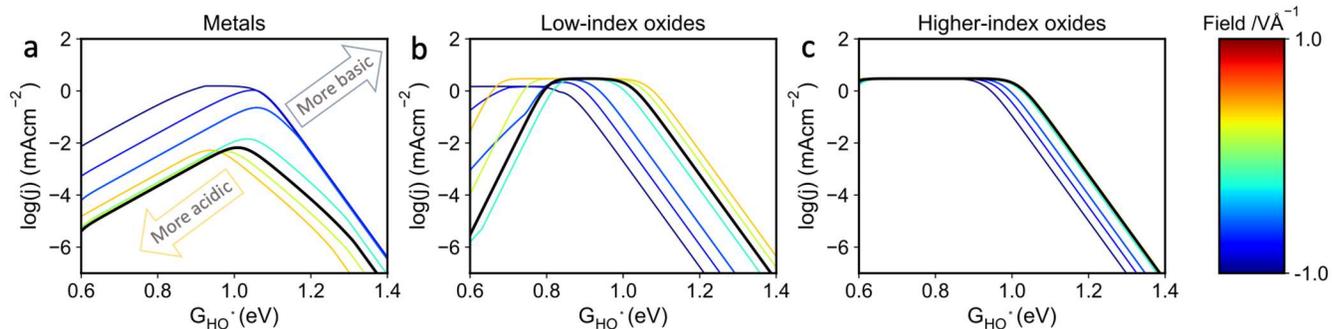
**Supplementary Figure 2. Summary of the performance of TMO catalysts for acidic ORR published over the last twelve years.**  $U_{\text{RHE}}$  measured at 50  $\mu\text{A}/\text{cm}^2$  with the data sorted by a) cation system and b) host anion. The dashed line is the value of pure Pt catalysts for acidic ORR reported by Ref.<sup>1</sup>. Note that the literatures with potential Pt-contamination and the partially oxidized Pt-alloy systems (which possess metallic Pt as their active sites) are excluded from the statistics.



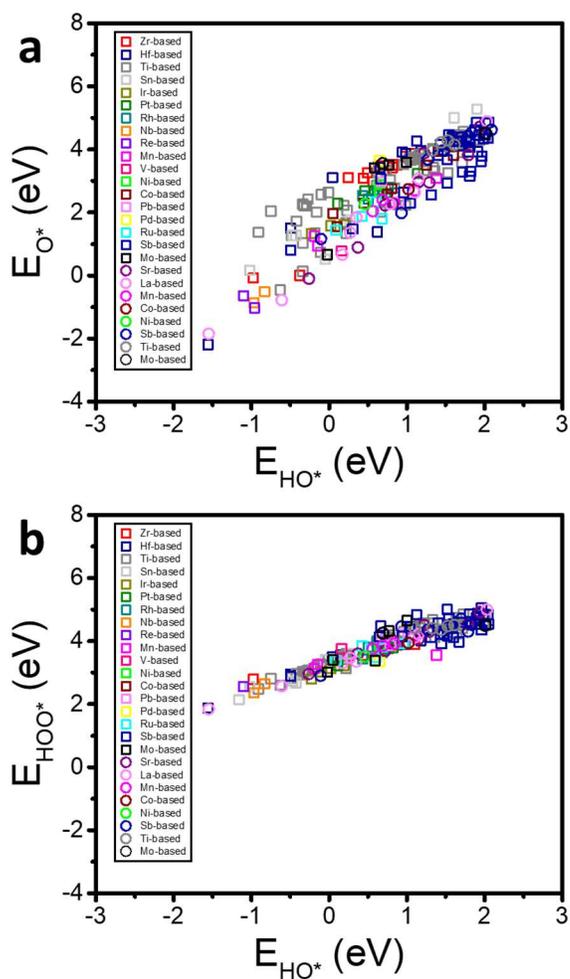
**Supplementary Figure 3. Summary of high-throughput screening of 7798 unique TMO compositions.** The values for the Mn, Ni, and Fe oxide samples are indicated by horizontal lines. The activity of the catalysts containing 2, 3, and 4 cations are shown in the respective histograms for 0.63 (top) and 0.8 (bottom) V<sub>RHE</sub>. Catalysts with current density below 1 μA/cm<sup>2</sup> are not shown.



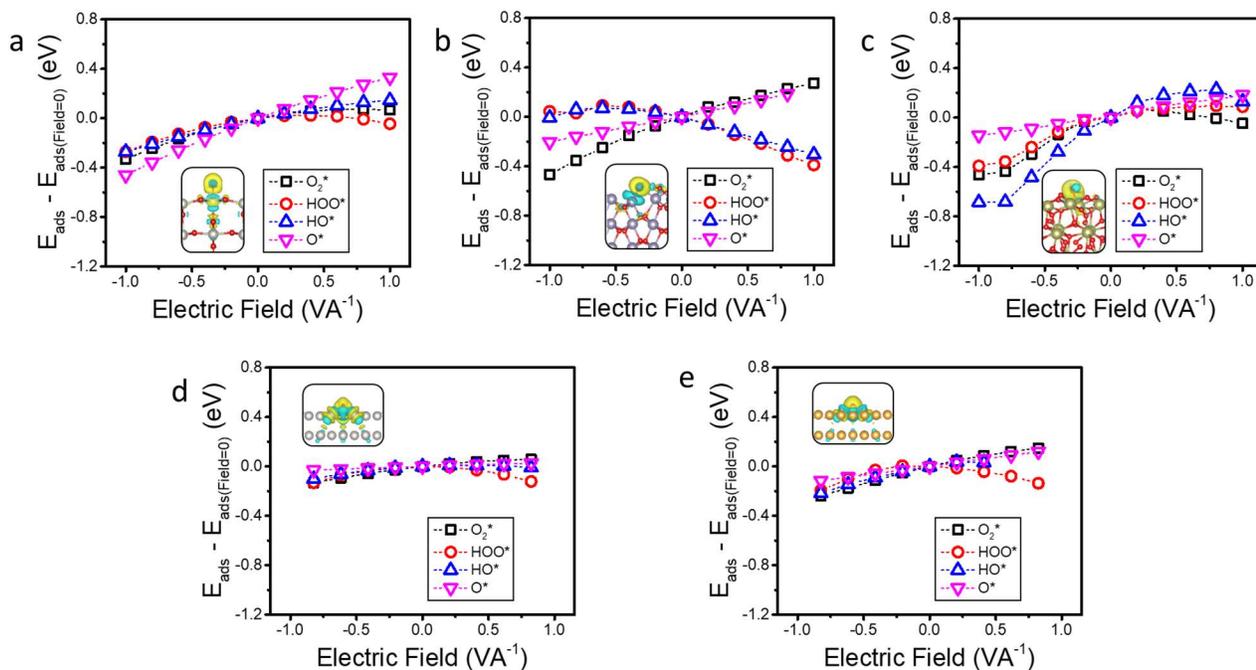
**Supplementary Figure 4. Calculated surface Pourbaix diagrams for  $ZrO_2$  and  $HfO_2(111)$ .** a-b) Surface Pourbaix diagrams of  $ZrO_2(111)$  at the pH of 0 and 14. c-d) Surface Pourbaix diagrams of  $HfO_2(111)$  at the pH of 0 and 14. e-f) 2D surface Pourbaix diagram of  $ZrO_2$  and  $HfO_2(111)$ . At  $U_{RHE}=0.80$  V, the corresponding  $U_{SHE}$  values are 0.8 and -0.03 V respectively for the pH of 0 and 14. g-h) Top and side view of the pristine surfaces of  $ZrO_2$  and  $HfO_2(111)$  considered for calculations. Red, green, and brown spheres represent O, Zr, and Hf atoms, respectively.



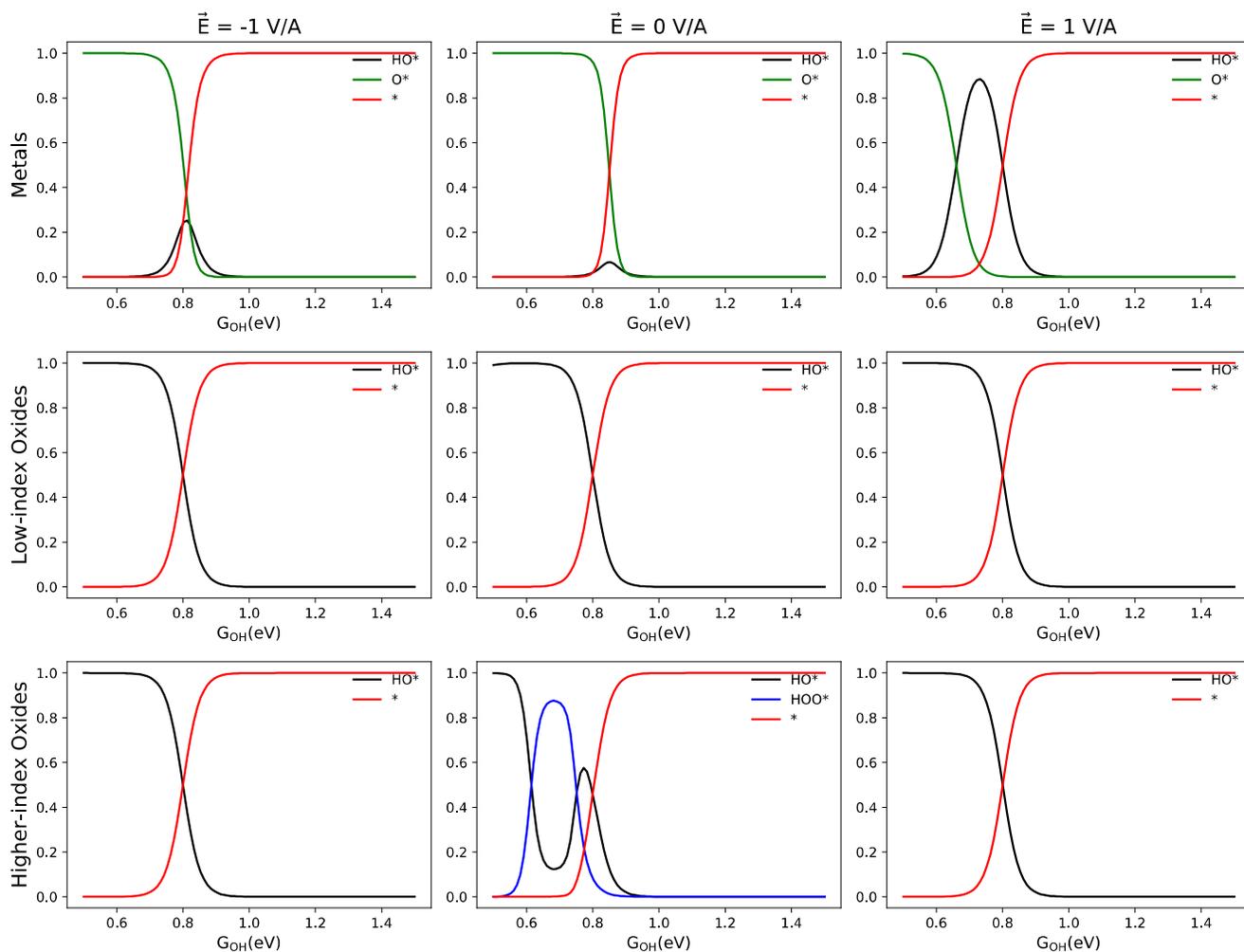
**Supplementary Figure 5. pH-dependent kinetic volcano models for the  $2e^-$  ORR process at  $0.6 \text{ V}_{\text{RHE}}$ .** a) Volcano activity models for transition metal surfaces; b,c) volcano activity models for b) low- and c) higher-index TMO surfaces. The higher and lower electric fields respectively represent more acidic and basic conditions.



**Supplementary Figure 6. ORR adsorbate binding energy data used in this study.** a)  $E_{HO^*}$  vs  $E_{O^*}$  and b)  $E_{HO^*}$  vs  $E_{HOO^*}$  scaling relations. Square and circle points represent higher- ( $h^2+k^2+l^2 > 1$ ) and low-index ( $h^2+k^2+l^2 = 1$ ) surfaces, respectively. The plots consist of the data from Refs.<sup>33–36</sup>, and the data calculated by our current research team.



**Supplementary Figure 7. Electric field effects on the ORR adsorbate bindings.** a) PdO<sub>2</sub>(110), b) SnO<sub>2</sub>(110), c) HfO<sub>2</sub>(111), d) Pt(111),<sup>51</sup> and e) Au(111).<sup>51</sup> Insets show the charge density differences induced by the adsorption of atomic oxygen. The yellow and teal colors in the isosurfaces represent electron charge gain and loss, respectively. Red, silver, purple, brown, grey, and yellow spheres represent O, Pd, Sn, Hf, Pt, and Au, respectively.



**Supplementary Figure 8. Simulated steady-state coverages of ORR adsorbates.** Fractional coverages of each adsorbate with at least 1% coverage of surface sites across the binding energy range. Coverages shown for metals, low-index TMOs, and higher-index TMOs at -1, 0, and 1 V/Å fields.

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