

**Supporting Information for  
Fast and Accurate Computation of  
Nonadiabatic Coupling Matrix Elements  
(NACMEs) using the Truncated Leibniz Formula  
(TLF) and Mixed-Reference Spin-Flip  
Time-Dependent Density Functional Theory  
(MRSF-TDDFT)**

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## Description of the MRSF Method

Detailed discussion of MRSF-TDDFT can be found elsewhere.<sup>1,2</sup> The main concept is summarized in this section. Derivation of the MRSF-TDDFT methodology<sup>1</sup> is based on the density-matrix formulation of time-dependent Kohn-Sham theory.<sup>3</sup> In MRSF-TDDFT, the zeroth-order MR reduced density matrix (RDM),  $\rho_0^{\text{MR}}(x, x')$ , is defined by satisfying an idempotence relation,

$$\rho_0^{\text{MR}}(x, x') = \int \rho_0^{\text{MR}}(x, x'') \rho_0^{\text{MR}}(x'', x') dx'', \quad (\text{S1})$$

and its density,  $\rho_0^{\text{MR}}(x) = \rho_0^{\text{MR}}(x, x)$ , is same as an equiensemble density of  $M_S = +1$  and  $M_S = -1$  components of a triplet state, *e.g.*,

$$\rho_0^{\text{MR}}(x) = \frac{1}{2} \{ \rho_0^{M_S=+1}(x) + \rho_0^{M_S=-1}(x) \}. \quad (\text{S2})$$

Within the Tamm-Dancoff approximation,<sup>4,5</sup> the use of MR-RDM in the linear-response formalism yields completely decoupled two linear-response equations for singlet and triplet excited states, respectively,<sup>1</sup>

$$\sum_{rs} (A_{pq,rs}^{(k)(0)} + A'_{pq,rs}{}^{(k)}) X_{rs}^{(k)} = \Omega_{(k)} X_{pq}^{(k)}, \quad k = S, T \quad (\text{S3})$$

where  $k = S, T$  labels singlet and triplet states,  $A_{pq,rs}^{(k)(0)}$  is an orbital Hessian matrix derived by the linear response, and  $A'_{pq,rs}{}^{(k)}$  is a coupling matrix between configurations originating from different components,  $M_S = +1$  and  $M_S = -1$ , of the mixed reference.<sup>1,2</sup>  $X_{pq}^{(k)}$  and  $\Omega_{(k)}$  are the amplitude vectors and the excitation energies with respect to the reference state, respectively.

It was shown that a linear response of MRSF-TDDFT can be represented as configurations from  $M_S = +1$  and  $-1$  references depicted as those with black and red arrows, respectively, in Fig. S1.<sup>1</sup> The configurations with the red arrows are missing in the conventional SF-TDDFT; their absence leads to spin contamination of the response states. These are recovered in MRSF-TDDFT and the spin contamination of the response states is nearly eliminated.<sup>1</sup> Configurations with the blue arrows are originated from both  $M_S = +1$  and  $-1$  components of the mixed reference.

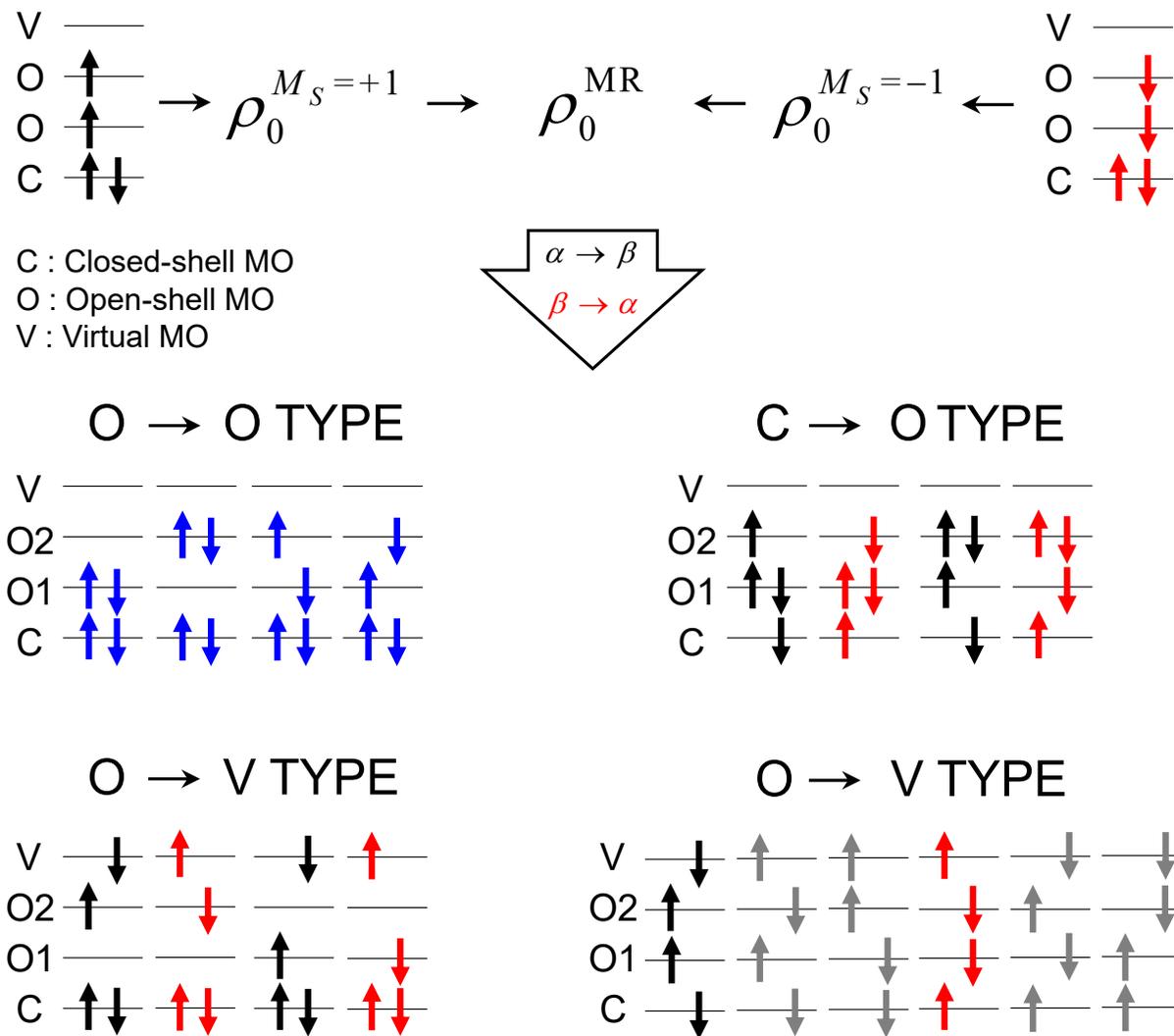


Figure S1: Electronic configurations of the  $M_S = +1$  and  $M_S = -1$  components of triplet reference in the upper panel and a complete set of configurations for MRSF-TDDFT in the lower panel. Response states are described by configurations represented with blue, black, and red arrows in MRSF-TDDFT, while with blue and black arrows in SF-TDDFT.

Although not all electronic configurations can be recovered by using the MR-RDM, the missing C  $\rightarrow$  V configurations (depicted with gray arrows in Fig. S1) represent high-lying excited states and their effect on the lower part of the excitation spectrum is insignificant.<sup>1</sup>

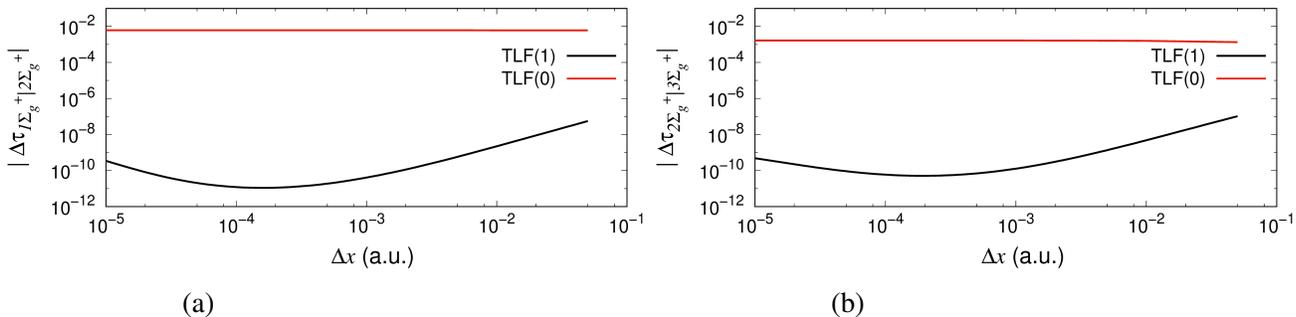


Figure S2: Comparison of the NACME,  $|\Delta\tau_{IJ}|$ , absolute errors (in a.u.) for the TLF(0) and TLF(1) algorithms for  $\text{H}_2$  (a) between  $1^1\Sigma_g^+$  and  $2^1\Sigma_g^+$  at  $R = 4.5$  a.u., (b) between  $2^1\Sigma_g^+$  and  $3^1\Sigma_g^+$  at  $R = 2.8$  a.u. The  $\Delta x$  used here is  $10^{-4}$  a.u. All the data shown were obtained with MRSF-BHHLYP.

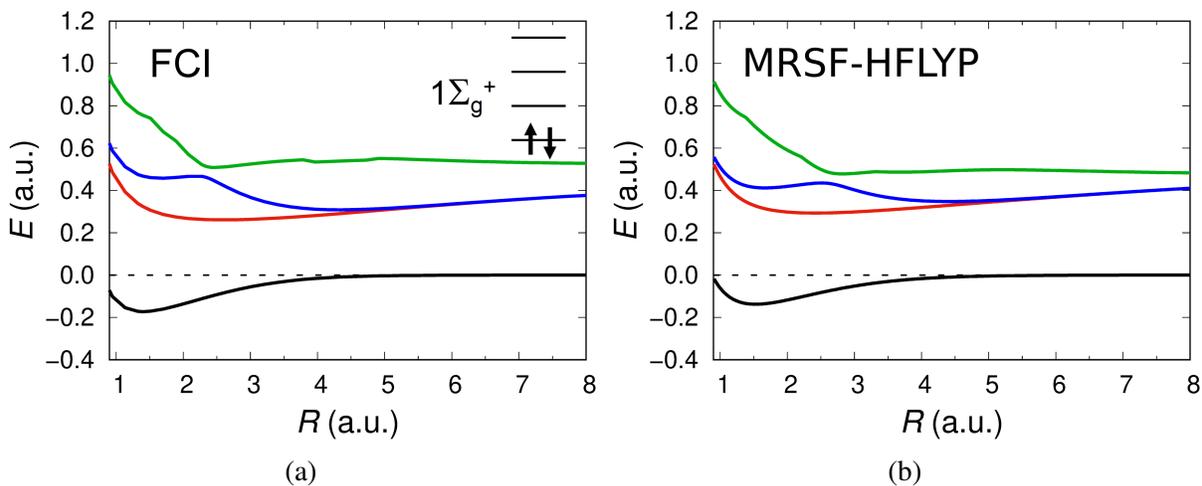


Figure S3: Comparison of the potential energy surfaces for  $^1\text{H}_2$  dissociation calculated with (a) FCI and (b) MRSF-HFLYP methods. The black, red, blue, and green curves correspond to the  $S_0$ ,  $S_1$ ,  $S_2$ , and  $S_3$  states, respectively. For the symmetry of the states as well as their configurations, see Figure 2 in the main text. The MRSF energies,  $\Delta E$ , were calculated as explained in Figure 2 in the main text.

## Computational Details

All the LR-TDDFT and the MRSF calculations reported in this work were performed by the locally modified GAMESS-US program package.<sup>6</sup> The accurate  $\text{H}_2$  dissociation curves and the NACME curves were calculated by the full configurational interaction (FCI) method implemented in the MOLPRO program package.<sup>7</sup>

The following density functionals were employed: PBE0<sup>8</sup> (0.25), BHHLYP<sup>9</sup> (0.50), STG1X<sup>10</sup> (0.85), and HFLYP<sup>10</sup> (1.00) where the amount of the exact exchange is given in parentheses. The

all-electron cc-pVTZ<sup>11</sup> basis set was employed in for the calculations of H<sub>2</sub>. For brevity, the calculations performed with the LR- or MRSF-TDDFT method and the BHHLYP functional will be abbreviated as “LR-BHHLYP” or “MRSF-BHHLYP”, respectively. The NACME calculations employed the TLF algorithm of Ref. 12 and geometry displacements (in Cartesian coordinates) of  $m \times 10^{-n}$  a.u. with  $m = 1, 5$  and  $n = 2-7$ . The SCF convergence criterion was set to  $10^{-6}$ , the convergence tolerance for Davidson algorithm was set to  $10^{-7}$ . The number of points for the radial and Lebedev grid was 99 and 590, respectively, for DFT, whereas it was 96 and 302, respectively, for TDDFT.

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