ZTFJ0038+2030: A Long-period Eclipsing White Dwarf and a Substellar Companion

Jan van Roestel1, Thomas Kupfer2, Keaton J. Bell3, Kevin Burdge1, Przemek Mróz1, Thomas A. Prince1, Eric C. Bellm4, Andrew Drake1, Richard Dekany3, Ashish A. Mahabal1,6, Michael Porter5, Reed Riddle5, Kyung Min Shin7, David L. Shupe8, and S. R. Kulkarni1

1 Division of Physics, Mathematics, and Astronomy, California Institute of Technology, Pasadena, CA 91125, USA; jvanroes@caltech.edu
2 University of Kansas, Lawrence, KS 66045, USA
3 DIRAC Institute, Department of Astronomy, University of Washington, Seattle, WA-98195, USA
4 DIRAC Institute, Department of Astronomy, University of Washington, Seattle, WA-98195, USA
5 California Institute of Technology, Pasadena, CA 91125, USA
6 Center for Data Driven Discovery, California Institute of Technology, Pasadena, CA 91125, USA
7 EPIC212235321
8 IPAC, California Institute of Technology, 1200 E. California Blvd., Pasadena, CA 91125, USA

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Abstract

In a search for eclipsing white dwarfs using the Zwicky Transient Facility lightcurves, we identified a deep eclipsing white dwarf with an orbital period of 10.4 hr and an undetected substellar companion. We obtained high-speed photometry and radial velocity measurements to characterize the system. The white dwarf has a mass of $0.50 \pm 0.02 \, M_\odot$ and a temperature of $10,900 \pm 200 \, K$. The companion has a mass of $0.059 \pm 0.004 \, M_\odot$ and is a brown dwarf. It has a radius of $0.0783 \pm 0.0013 \, R_\odot$, and is one of the physically smallest transiting brown dwarfs known and likely old, $\gtrsim 8 \, Gyr$. The ZTF discovery efficiency of substellar objects transiting white dwarfs is limited by the number of epochs and as ZTF continues to collect data we expect to find more of these systems.

Unified Astronomy Thesaurus concepts: White dwarf stars (1799); Eclipsing binary stars (444); Brown dwarfs (185); Extrasolar gaseous giant planets (509)

1. Introduction

Low-mass, hydrogen-rich objects (brown dwarfs and giant planets) that orbit a white dwarf in short orbital periods are rare. There are currently 10 confirmed white dwarfs with brown dwarf companions with measured orbital periods. These are: GD1400 ($P = 9.98 \, hr$; Farihi & Christopher 2004; Dobie et al. 2005; Burleigh et al. 2011), WD1037-349 ($P = 116 \, minutes$; Burleigh et al. 2006; Maxted et al. 2006; Casewell et al. 2015), NLTT 5306 ($P = 101.9 \, minutes$; Steele et al. 2013), WD0837+185 ($P = 4.2 \, hr$; Casewell et al. 2012), SDSS J141126.20+200911.1 ($P = 121.73 \, minutes$; Beuermann et al. 2013; Littlefair et al. 2014), SDSS J155720.77+091624.6 ($P = 2.27 \, hr$; Farihi et al. 2017), SDSS J120515.80-024222.6 ($P = 71.2 \, minutes$; Parsons et al. 2017a), SDSSJ123127.14+004312.9 ($P = 72.5 \, minutes$; Parsons et al. 2017a), EPIC212235321 ($P = 68.2 \, minutes$; Casewell et al. 2018) and WD1032+011 ($P = 2.21 \, hr$; Casewell et al. 2020a).

Most of these systems have short orbital periods. This is a detection bias: they have a higher eclipse probability, stronger reflection effect, and the white dwarf has a higher radial-velocity amplitude. For example, confirmation of GD1400 as a binary requires the white dwarf radial velocity to be measured, which requires a large telescope and high-resolution spectrograph, even for a system as bright ($G = 15.2$) as GD1400 (Burleigh et al. 2011).

There have also been many searches to find exoplanets ($M \lesssim 13 \, M_{\text{jup}}$) orbiting white dwarfs, e.g., Faedi et al. (2011), Fulton et al. (2014), van Sluijs & Van Eylen (2018), Dame et al. (2019), Rowan et al. (2019), but none of them found any candidates. The first white dwarf with a Jupiter-mass companion, WD 1145+017, which has an orbital period of $\approx 10 \, days$, was only recently discovered by Ginskiic et al. (2019). The planet is slowly evaporating and this material is accreted by the white dwarf. Soon after, Vanderburg et al. (2020) discovered the first transiting giant planet orbiting a white dwarf, WD J0914+1914, which has an orbital period of 1.4 days. After carefully analyzing optical and infrared lightcurves of the grazing eclipse, they conclude that the companion is a giant planet with a mass of $\lesssim 14 \, M_{\text{jup}}$.

Because the white dwarf went through a giant phase in the past, the companion must have migrated or have been formed after the giant phase. For brown dwarfs, common-envelope evolution is the commonly accepted scenario (Ivanova et al. 2013). If a brown dwarf is close enough to a star that is ascending the red giant branch (RGB) or asymptotic giant branch (AGB) (for $0.200–1000 \, R_\odot$, 1–5 au), the giant star engulfs the brown dwarf and forms a common envelope. Depending on the masses and initial orbital separation, the system can survive this process and ends up as a short-period binary (Casewell et al. 2018, 2020a, 2020b). However, the binary can also merge during a common-envelope event or the companion gets evaporated as soon as the hot white dwarf emerges from the common envelope (Nelemans & Tauris 1998; Soker 1998; Bear & Soker 2011). Particularly lower-mass objects (giant planets) are expected to merge or evaporate during this process.

To form white dwarf–planet systems, alternative pathways have been proposed. These are the formation of second-generation planets from gas around the white dwarf (e.g., Perets 2010) or capture and/or inward migration of distant planets (e.g., Stephan et al. 2020). In addition, recent work suggests that giant planets can survive common-envelope evolution (Lagos et al. 2021).

With the goal of studying the population of white dwarfs with brown dwarf or giant planet companions in close orbits, we searched for deep eclipsing white dwarfs using Zwicky Transient Facility lightcurves (Bellm et al. 2019; Graham et al. 2019; Masci et al. 2019; Dekany et al. 2020). We used the
combined PSF photometry and alert photometry lightcurves of white dwarfs (Gentile Fusillo et al. 2019) using the box least squares algorithm (Kovács et al. 2002). For more details, see van Roestel et al. (2021).

In this paper, we present the first result, the discovery of ZTFJ003855.0+203025.5 (ZTFJ0038+2030 see Table 1), an eclipsing white dwarf with a brown dwarf companion with an orbital period of 10 hr. ZTFJ0038+2030 shows a complete eclipse in the ZTF g and r band with a short eclipse duty cycle. It also showed no excess luminosity in the Gaia observational H–R diagram and no infrared excess in Pan-STARRS-y or WISE W1 and W2, which indicates the companion is cold and either a brown dwarf or giant planet and not an M dwarf. Because of these properties, we prioritized it for follow-up observations to determine the nature of the companion.

We obtained follow-up photometry and spectroscopy (Section 2), which we used the characterize the system (Section 3). We present the mass, radius, and temperature measurements in Section 4. We compare this binary system with other white dwarfs with substellar companions, and discuss the implications of this discovery for future searches for giant exoplanets around white dwarfs with ZTF (Section 5). The final section summarizes the paper.

2. Follow-up Data

2.1. CHIMERA Fast Cadence Photometry

We obtained high-speed photometry in the g and z filters using CHIMERA (see Table 2). CHIMERA (Harding et al. 2016) is a dual-channel photometer that uses frame-transfer, electron-multiplying CCDs mounted on the Hale 200 inch (5.1 m) Telescope at Palomar Observatory (CA, USA). The pixel scale is 0.28 pixel^{-1} (unbinned). We used the conventional amplifier and used 2×2 binning on most nights to reduce the readout noise. Each of the images was bias-subtracted and divided by twilight flat fields. We used the ULTRACAM pipeline to do aperture photometry (Dhillon et al. 2007). We used an optimal extraction method with a variable aperture of 1.5 the FWHM of the seeing (as measured from the reference star). A differential lightcurve was created by simply dividing the counts of the target by the counts from the reference star. Timestamps of the images were determined using a GPS receiver.

2.2. ESI

We used the Echellette Spectrograph and Imager (ESI, Sheinis et al. 2002) mounted at Keck II to obtain medium-resolution spectra (R ≈ 6000). The wavelength range is from 4000 to 10000 Å. CuAr exposures were taken at the beginning of the night. The spectra were reduced using the MAKEE pipeline following the standard procedure: bias subtraction, flat-fielding, sky subtraction, order extraction, and wavelength calibration. We did not attempt to flux-calibrate the spectra.

2.3. Archival Photometry

To be able to study the spectral energy distribution, we obtained photometry data from multiple other survey telescopes (see Table 1): UV data from GALEX (Bianchi et al. 2017), optical data from Gaia eDR3 (Brown et al. 2020a) and Pan-STARRS (Chambers et al. 2016), UKIRT Hemisphere Survey J-band data (Dye et al. 2018), and far-infrared data from WISE (Marocco et al. 2021). No near-infrared photometry H and K data are available. We used zero-points for each of the filters to convert the magnitudes to a flux.

3. Analysis

3.1. Ephemeris

We determine the ephemeris by measuring the mid-eclipse time from the CHIMERA g lightcurve. We then use the best model from the Chimera g data and use it to fit all ZTF data. In addition, we noticed that there is one non-detection on 2012 November 01 in Palomar Transient Factory data (out of 94 observations). We add this epoch with half the eclipse duration as uncertainty as a prior (BJD_TDB = 2, 456, 232.8854 ± 0.0018). This results in an ephemeris of:

\[
\text{BJD(TDB)} = 2, 459, 045.985194(2) + 0.431 920 8(14)
\]

3.2. Spectral Energy Distribution

To determine the white dwarf and companion temperature, we fit the observed spectral energy distribution with a model that combines white dwarf spectral models with spectral models of substellar companions (see Figure 1). We use a grid of DA white dwarf models by Koester (2009) and use bilinear interpolation to be able to generate a model for any temperature and surface gravity value. For the companion, we use models from Phillips et al. (2020) with log g = 5.5 and a fixed radius based on the lightcurve result. We use the extinction law by Fitzpatrick (1999) to account for any dust extinction. To compare the model spectra with the data, we convolve the

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Table 1

<table>
<thead>
<tr>
<th>Property</th>
<th>Value</th>
</tr>
</thead>
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<tr>
<td>R.A.</td>
<td>00°38’55.0’’</td>
</tr>
<tr>
<td>Decl.</td>
<td>00°30’26.1’’</td>
</tr>
<tr>
<td>distance</td>
<td>138.3±1.4 pc</td>
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<tr>
<td>E_B-V</td>
<td>0.03±0.02</td>
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<tr>
<td>G_V</td>
<td>17.70</td>
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<tr>
<td>G_B</td>
<td>17.76±0.01</td>
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<tr>
<td>G_R</td>
<td>17.63±0.01</td>
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<tr>
<td>G_A</td>
<td>20.37±0.24</td>
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<tr>
<td>G_N</td>
<td>18.58±0.06</td>
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<tr>
<td>ZTF_g</td>
<td>17.70±0.02</td>
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<tr>
<td>ZTF_r</td>
<td>17.78±0.03</td>
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<tr>
<td>ZTF_i</td>
<td>17.95±0.03</td>
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<tr>
<td>PS_g</td>
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<tr>
<td>PS_r</td>
<td>17.786±0.002</td>
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<tr>
<td>PS_i</td>
<td>17.931±0.006</td>
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<tr>
<td>PS_z</td>
<td>18.093±0.005</td>
</tr>
<tr>
<td>PS_y</td>
<td>18.18±0.02</td>
</tr>
<tr>
<td>UKIRT-J</td>
<td>17.71±0.06</td>
</tr>
<tr>
<td>WISE-W1</td>
<td>17.87±0.11</td>
</tr>
<tr>
<td>WISE-W2</td>
<td>17.63±0.29</td>
</tr>
</tbody>
</table>

Note. Gaia eDR3 data were used (Brown et al. 2020b) with the geometric distance from Baier-Jones et al. (2021). "\( \pm \)" indicates that the magnitudes are in the Vega system; other magnitudes are in the AB system.

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van Roestel et al.

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http://www.astro.caltech.edu/~ipac_staff/tab/makee/
model spectra with the filter response curves\(^\text{11}\) (Rodrigo et al. 2012; Rodrigo & Solano 2020). We use Gaussian priors on the parallax using the Gaia eDR3 data, the radius estimate from the Gaia eDR3 data, and a systematic uncertainty on the parallax \((\Delta \pi)\) estimated from the Gaia eDR3 data. The two remaining free parameters are the amplitude \((K_1)\) and a systematic velocity \((\gamma)\). We use the emcee (Foreman-Mackey et al. 2013) to determine the best value and uncertainty: \(K_1 = 24.2 \pm 1.4 \text{ km s}^{-1}\) (Figure 2).

3.4. Lightcurve Modeling

We modeled the high-cadence lightcurves using the package ellec (Maxted 2016). We use a spherical star to model the white dwarf and use Roche-lobe geometry for the companion. The free parameters for this model are the mid-eclipse time \((t_0)\), inclination \((i)\), mass-ratio \((q)\), the radii divided by the semimajor axis of both objects \((r_{1,2} \equiv R_{1,2}/a)\), the semimajor axis \((a)\), and the surface brightness ratio \((J_{\text{e.c.}})\).

We used a number of fixed parameters in the binary model. First, we use the orbital period obtained from the ZTF data (Section 3.1). For limb-darkening of the white dwarf, we use tabulated values by Claret et al. (2020) for \(T=10,000 \text{ K}\) and \(\log(g) = 8.0\).

In addition, we imposed two restrictions on the white dwarf based on the zero-temperature white dwarf mass–radius relation by Eggleton as reported in Verbunt & Rappaport (1988). The first is that it cannot be smaller than a zero-temperature white dwarf. The second constraint is a Gaussian prior with a relative size of 5% compared to a zero-temperature white dwarf to limit the maximum size of the white dwarf. As a

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\(^{11}\) \text{http://svo2.cab.inta-csic.es/theory/fps/}

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**Figure 1.** The spectral energy distribution of the ZTFJ0038+2030. Markers show GALEX, Gaia DR3, Pan-STARRS, median ZTF griz, and WISE data. The best-fit model is shown in gray, with the largest possible contribution by the companion as a dotted line. The substellar companion does not contribute significantly, even in the WISE W1 and W2 bands in the infrared. We can therefore only derive an upper limit to the companion temperature. The inset shows the averaged and velocity-corrected ESI spectra, showing from top to bottom: Hδ, Hγ, Hβ, and Hα.

**Table 2**

<table>
<thead>
<tr>
<th>Date</th>
<th>UT</th>
<th>Tele./Inst.</th>
<th>(N_{\text{exp}})</th>
<th>Exp. time (s)</th>
<th>Wavelength</th>
</tr>
</thead>
<tbody>
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<td>2020-07-15</td>
<td>11:22-11:57</td>
<td>P200/CHIMERA</td>
<td>400</td>
<td>5.0</td>
<td>(g)</td>
</tr>
<tr>
<td>2020-07-15</td>
<td>11:22-11:57</td>
<td>P200/CHIMERA</td>
<td>400</td>
<td>5.0</td>
<td>(z)</td>
</tr>
<tr>
<td>2020-09-12</td>
<td>12:17-12:38</td>
<td>Keck/ESI/Echellete</td>
<td>2</td>
<td>600</td>
<td>4000–10000 Å</td>
</tr>
<tr>
<td>2020-09-12</td>
<td>14:16-14:37</td>
<td>Keck/ESI/Echellete</td>
<td>3</td>
<td>600</td>
<td>4000–10000 Å</td>
</tr>
</tbody>
</table>
final constraint, we use a Gaussian prior on the radial velocity amplitude \( K_1 \) of the white dwarf (see Section 3.3).

To find the most probable parameter values and uncertainties, we again use emcee.

4. Results

We measured the binary properties by analyzing ZTF lightcurves (Section 3.1, used to determine the orbital period), the spectral energy distribution (Section 3.2, used to derive the temperatures), phase-resolved spectroscopy (Section 3.3, used to measure the white dwarf radial velocity semi-amplitude) and high-cadence \( g \)- and \( z \)-band lightcurves (used to measure the mass and radii of both components). The results are summarized in Table 3 and the posterior of the lightcurve modeling is shown in the Appendix.

The mass of the companion, which is mostly set by the radial velocity semi-amplitude measurement, is \( M_2 = 0.0593 \pm 0.004 \, M_\odot \) and a radius of \( R_2 = 0.0783 \pm 0.0012 \, R_\odot \). The \( z \)-band surface brightness ratio limits the temperature of the companion to \( \lesssim 1550 \, K \) and the spectral energy distribution constrains the temperature further to \( \lesssim 750 \, K \).

The mass of the white dwarf is \( 0.50 \pm 0.02 \, M_\odot \) and the radius is \( R_1 = 0.01429 \pm 0.00020 \, R_\odot \). This is consistent with the white dwarf M–R relation, which is what we enforced using a prior. The temperature of the white dwarf is \( T_1 = 10,900 \pm 200 \, K \) and a surface gravity is \( \log g = 7.83 \pm 0.01 \). This is slightly different than reported by Gentile Fusillo et al. (2019) \( (10,290 \pm 210 \, K, \log g = 7.94 \pm 0.07) \) but is within two standard deviations.

The orbital separation of the binary system is \( a = 1.987 \pm 0.027 \, R_\odot \) and the inclination of this system is \( i = 89^\circ 71 \pm 0^\circ 13 \).

5. Discussion

5.1. The Nature of the Substellar Companion

In Figure 3, we plot the mass and radius of the companion and compare it to models by Marley et al. (2018) and other binaries with substellar objects. The measured mass and radius agree with models of 10 Gyr old brown dwarfs with \( Z \gtrsim 0 \) abundances. The models predict a temperature of \( \sim 800 \, K \), which is consistent with the limit we derived from the SED, \( \lesssim 750 \, K \) (see Figure 1). If we assume a solar abundance or lower, the age of the brown dwarf (and therefore the system) is \( \gtrsim 8 \, Gyr \).

The initial to final mass relation for the white dwarfs suggests that the white dwarf progenitor was approximately a 1–2 \( M_\odot \) main-sequence star. This corresponds to a main-sequence lifetime of 2–10 Gyr (Catalán et al. 2008; Marigo 2013; Cummings et al. 2018). The cooling age of the white dwarf is approximately \( \sim 400 \, Myr \) (Koester 2009). Given the age of the brown dwarf, the white dwarf progenitor mass was likely closer to the lower bound of the mass range, \( \approx 1 \, M_\odot \).
Table 3

| $p'$ (d) | $0.4319208 \pm 0.014$ |
| $i_0$ (BJD$_{TDB}$) | 2.459, 045.985194(2) |
| $q$ | $0.11677 \pm 0.0068$ |
| $i$ (°) | $89.71 \pm 0.11$ |
| $r_1$ | $0.007195 \pm 0.000075$ |
| $r_2$ | $0.03934 \pm 0.00013$ |
| $a$ ($R_{\odot}$) | $1.987^{+0.022}_{-0.022}$ |
| $J_x$ | $\lesssim 0.000035$ |
| $J_z$ | $\lesssim 0.00014$ |

| $M_1$ ($M_{\odot}$) | $0.505^{+0.024}_{-0.018}$ |
| $M_2$ ($M_{\odot}$) | $0.0593^{+0.0056}_{-0.0039}$ |
| $R_1^2$ ($R_{\odot}$) | $0.01429^{+0.00017}_{-0.00011}$ |
| $R_2$ ($R_{\odot}$) | $0.0783^{+0.0011}_{-0.00011}$ |
| log($g_1$) (cgs) | $7.832^{+0.013}_{-0.013}$ |
| log($g_2$) (cgs) | $5.425^{+0.002}_{-0.003}$ |
| $K_1$ (km s$^{-1}$) | $24.4^{+1.4}_{-1.4}$ |
| $K_2$ (km s$^{-1}$) | $208.4^{+4.4}_{-2.9}$ |
| $p_2$ (g cm$^{-2}$) | $174^{+9}_{-11}$ |
| $K_1^2$ (km s$^{-1}$) | $24.4^{+1.4}_{-1.4}$ |
| $K_2$ (km s$^{-1}$) | $208.4^{+4.4}_{-2.9}$ |
| $T_1$ (K) | $10900 \pm 200$ |
| $T_2$ (K) | $\lesssim 750$ |
| distance (pc) | $139 \pm 2$ |

Note. The top section lists $ellc$ model parameters; the middle section shows the binary parameters derived from the $ellc$ fit. We fixed the orbital period ($i'$) and for the radius of the white dwarf ($R_1$) and radial velocity amplitude ($K_1$) we used a prior ($i'$). The temperatures, in the bottom section of the table, have been determined by modeling the SED. We use the 95% percentile to determine upper limits.

Compared to other substellar objects that are eclipsing white dwarfs, the mass and radius do not stand out and are similar to other brown dwarfs. This object does stand out because of its orbital period, which at 10 hr is an order of magnitude larger than the three other known brown dwarfs orbiting white dwarfs. This means that the amount of irradiation by the white dwarf is relatively low. Using a simple blackbody approximation (Littlefair et al. 2014), we estimate that the temperature of the brown dwarf is only increased by $\sim 50$ K due to irradiation by the white dwarf. This fact, and the system’s relative brightness, make it a good prototype system for long-period white-dwarf–brown-dwarf systems.

5.2. Formation History

Since the companion is a brown dwarf and not a giant planet, standard common-envelope evolution can explain the formation of this system. Given that the mass of the white dwarf is $\geq 0.47 M_{\odot}$, the white dwarf very likely has a CO core (see, e.g., Marigo 2013 and also Parsons et al. 2017b for observational evidence) which allows for two formation scenarios. In the first scenario, the white dwarf could have formed during a common-envelope phase on the AGB after helium-core exhaustion. The second scenario is that the common envelope happened at the tip of the RGB, just after the helium flash (Han et al. 2003) which would result in a white dwarf with a mass close to $0.47 M_{\odot}$. In that scenario the white dwarf would have emerged from the common envelope as a hot subdwarf (sdB) and appeared as an HW Vir system before it evolved into a white dwarf with a brown dwarf companion after helium exhaustion in the sdB. Several sdB + brown dwarf systems are known, although typically seen with shorter orbital periods (e.g., Geier et al. 2011; Schaffenroth et al. 2015, 2018, 2019).

The white dwarf will slowly cool down and the period will slowly decrease due to gravitational wave radiation. It will take $\sim 135$ Gyr to reach an orbital period of $\sim 40$ minutes (Rappaport et al. 2021), at which point the white dwarf will be $\sim 1000$ K. Roche-lobe overflow will commence and the system becomes a cataclysmic variable (Littlefair et al. 2003). The accretion flow will heat up the white dwarf again while the period increases. This will slowly drain the brown dwarf and the system ends up as a “period-bouncer”; a very low accretion-rate CV with an orbital period of $\approx 90$ minutes (e.g Pala et al. 2018).

5.3. Implications for Searches for Brown Dwarfs and Giant Planets Orbiting White Dwarfs

Here, we briefly discuss the detection efficiency of our search and the occurrence rate of white dwarfs with transiting substellar objects. A detailed simulation is beyond the scope of this paper and we limit ourselves to an order of magnitude estimate only.

To find ZTFJ0038+2030 we searched the ZTF lightcurves of the Gaia white dwarf catalog by Gentile Fusillo et al. (2019) which contains 486,641 candidate white dwarfs over the entire sky. There are 129,148 white dwarfs brighter than 20 mag with more than 80 epochs in their ZTF lightcurve. Based on the number of epochs in these lightcurves, we estimate an average recovery efficiency of 15%–25% (the probability of getting 7–5 in-eclipse points). We note that we did recover the other three known eclipsing WD–BD systems (Figure 3) that show longer eclipse duty cycles and are therefore easier to find than ZTFJ0038+2030.

With the discovery of ZTFJ0038+2030 and the discovery of its planet-candidate (Vanderburg et al. 2020), there are now two known long-period ($\gtrsim 10$ hr) transiting substellar objects around white dwarfs; the first is most likely a giant planet and the second a brown dwarf. Based on just these two detections of long-period transiting objects ($\gtrsim 10$ hr), the occurrence rate of planets and brown dwarfs is the same order of magnitude. If we also consider the non-transiting objects, the white-dwarf–brown-dwarf GD1400 and the white-dwarf–planet system WD 1145+017, we reach a similar conclusion. However, we note that these latter two systems were discovered due to an infrared excess and peculiar white dwarf emission lines, which are both methods that are heavily biased. In order to systematically measure the occurrence rate of objects spanning masses of $\approx 0.01–0.07 M_{\odot}$, a white dwarf transit search is needed. Such a survey allows us to measure the mass distribution of white dwarf companions and determine which of the formation channels are important in the formation of these objects.

Currently, the ZTF detection efficiency is limited by the number of epochs available per white dwarf. As more epochs are obtained, ZTF will be able to identify narrower eclipses, which means that longer period systems can be identified. Based on the recovery efficiency of ZTFJ0038+2030, we estimate that ZTF will find another 3–6 similar-sized objects at

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12 A large-scale, precise ($\lesssim 1 \text{ km s}^{-1}$) spectroscopic survey to find radial velocity changes of white dwarfs would also work.
long orbital periods ($\gtrsim 10$ hr) as it keeps accumulating more data. Other surveys like Gaia (Gaia Collaboration et al. 2016), ATLAS (Tonry et al. 2018), and BlackGEM (Groot 2019) can be used to find similar systems over the whole sky. In the near future, the Vera C. Rubin observatory (Ivezić et al. 2019) is expected to find many white dwarfs with exoplanets, possibly down to Earth-sized objects (Agol 2011).

6. Summary and Outlook

Using ZTF photometry and Gaia and Pan-STARRS data, we discovered an eclipsing binary composed of a white dwarf and a substellar companion with an orbital period of 10 hr. We obtained follow-up photometry and spectroscopy and measured the binary parameters. We showed that the substellar
companion is a $\gtrsim 8$ Gyr old, small brown dwarf with a mass of $0.06 M_\odot$, and the white dwarf a $0.50 M_\odot$, CO white dwarf. ZTFJ0038+2030 has a much longer orbital period than most known white-dwarf–brown-dwarf systems, and is very similar to non-transiting white-dwarf–brown-dwarf GD1400. However, GD1400 shows an infrared excess (which is how it was discovered) while ZTFJ0038+2030 does not because the brown dwarf is older and has a lower temperature.

ZTFJ0038+2030 is relatively bright and because the brown dwarf suffers minimal irradiation, it can be used to study the brown dwarf atmosphere, while the mass and radius can be measured even more precisely by obtaining additional high-speed photometry and phase-resolved spectra. It is also a useful target for eclipse timing to find circumbinary objects (e.g., NN Ser, Marsh et al. 2014) as brown dwarfs are not expected to show eclipse time variations due to Applegate’s mechanism (Applegate 1992; Bours et al. 2016). The discovery of this system demonstrates that ZTF data can be used to find substellar objects in long periods orbiting white dwarfs and we expect ZTF to find more of these systems as more data are obtained.

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Based on observations obtained with the 200 inch Hale Telescope at the Palomar Observatory as part of the Zwicky Transient Facility project. The Hale telescope is operated by the Caltech Optical Observatories.

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Facilities: PO:1.2 m (ZTF), Hale (CHIMERA), Keck:II (ESI).


Appendix

Appendix Information

Figure 4 shows the posterior distribution of the Chimera lightcurve modelling.
References


ORCID iDs

Jan van Roestel @ https://orcid.org/0000-0002-2626-2872
Thomas Kupfer @ https://orcid.org/0000-0002-6540-1484
Keaton J. Bell @ https://orcid.org/0000-0002-0656-032X
Kevin Burdge @ https://orcid.org/0000-0002-7226-836X
Przemek Mróz @ https://orcid.org/0000-0001-7016-1692
Thomas A. Prince @ https://orcid.org/0000-0002-8850-3627
Eric C. Bellm @ https://orcid.org/0000-0001-8018-5348
Richard Dekany @ https://orcid.org/0000-0002-5884-7867
Ashish A. Mahabal @ https://orcid.org/0000-0002-2242-0244
Reed Riddle @ https://orcid.org/0000-0002-0387-370X
Kyung Min Shin @ https://orcid.org/0000-0002-1486-3582
David L. Shupe @ https://orcid.org/0000-0003-4401-0430
S. R. Kulkarni @ https://orcid.org/0000-0001-5390-8563

Figure 4. The posterior distribution of the fit to the lightcurve data using ellc and emcee.