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## HYPERCONTRACTIVITY OF THE SEMIGROUP OF THE FRACTIONAL LAPLACIAN ON THE $n$ -SPHERE

RUPERT L. FRANK AND PAATA IVANISVILI

ABSTRACT. For  $1 < p \leq q$  we show that the Poisson semigroup  $e^{-t\sqrt{-\Delta}}$  on the  $n$ -sphere is hypercontractive from  $L^p$  to  $L^q$  in dimensions  $n \leq 3$  if and only if  $e^{-t\sqrt{n}} \leq \sqrt{\frac{p-1}{q-1}}$ . We also show that the equivalence fails in large dimensions.

### 1. INTRODUCTION

1.1. **Poisson semigroup on the sphere.** Let

$$\mathbb{S}^n = \{x \in \mathbb{R}^{n+1} : \|x\| = 1\}$$

be the unit sphere in  $\mathbb{R}^{n+1}$ , where  $\|x\| = \sqrt{x_1^2 + \dots + x_{n+1}^2}$  for  $x = (x_1, \dots, x_{n+1}) \in \mathbb{R}^{n+1}$ . Let  $\Delta$  be the Laplace–Beltrami operator on  $\mathbb{S}^n$ . We will be working with spherical polynomials  $f : \mathbb{S}^n \rightarrow \mathbb{C}$ , i.e., finite sums

$$f(\xi) = \sum_{d \geq 0} H_d(\xi),$$

where  $H_d$  satisfies

$$\Delta H_d = -d(d+n-1)H_d.$$

The heat semigroup  $e^{t\Delta}$  is defined by  $e^{t\Delta}f = \sum_{d \geq 0} e^{-d(d+n-1)t} H_d$ . The hypercontractivity result for the heat semigroup on  $\mathbb{S}^n$  states that for any  $1 \leq p \leq q < \infty$ , any integer  $n \geq 1$ , and any  $t \geq 0$  we have

$$(1) \quad \|e^{t\Delta}f\|_q \leq \|f\|_p \quad \text{for all } f \quad \text{if and only if} \quad e^{-tn} \leq \sqrt{\frac{p-1}{q-1}},$$

where  $\|f\|_p^p = \|f\|_{L^p(\mathbb{S}^n, d\sigma_n)}^p = \int_{\mathbb{S}^n} |f|^p d\sigma_n$ , and  $d\sigma_n$  is the normalized surface area measure of  $\mathbb{S}^n$ . The case  $n = 1$  was solved independently in [9] and [10], and the general case  $n \geq 2$  was settled in [7]. We remark that the condition  $e^{-tn} \leq \sqrt{\frac{p-1}{q-1}}$  in (1) is different from the classical hypercontractivity condition  $e^{-t} \leq \sqrt{\frac{p-1}{q-1}}$  in Gauss space due to Nelson [8], and on the hypercube due to Bonami [2]. The appearance of the extra factor  $n$  in (1) can be explained from the fact that the spectral gap (the smallest nonzero eigenvalue) of  $-\Delta$  equals  $n$ .

In [7] the authors ask what the corresponding hypercontractivity estimates are for the Poisson semigroup on  $\mathbb{S}^n$ . As pointed out in [7], there are two natural

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Poisson semigroups on  $\mathbb{S}^n$  one can consider: 1)  $e^{-t\sqrt{-\Delta}}f$ , and 2)  $P_r f = \sum t^d H_d$ ,  $r \in [0, 1]$ . Notice that when  $n = 1$  both of these semigroups coincide (with  $r = e^{-t}$ ). It was conjectured by E. Stein that

$$\|P_r f\|_q \leq \|f\|_p \quad \text{if and only if} \quad r \leq \sqrt{\frac{p-1}{q-1}}$$

holds on  $\mathbb{S}^n$  for all  $n \geq 1$ . Besides the case  $n = 1$  mentioned above, the case  $n = 2$  was confirmed in [4], and the general case  $n \geq 2$  in [1].

The question of hypercontractivity for the semigroup  $e^{-t\sqrt{-\Delta}}$  on  $\mathbb{S}^n$  for  $n \geq 2$ , however, has remained open. Since the spectral gap of  $\sqrt{-\Delta}$  equals  $\sqrt{n}$ , it is easy to see that a necessary condition for the estimate  $\|e^{-t\sqrt{-\Delta}}f\|_q \leq \|f\|_p$  is  $e^{-t\sqrt{n}} \leq \sqrt{\frac{p-1}{q-1}}$ ; see Section 2.1. One might conjecture that this necessary condition is also sufficient. Surprisingly, it turns out the answer is positive in small dimensions and negative in large dimensions.

**Theorem 1.1.** *Let  $1 < p < q$ ,  $n \geq 1$ , and  $t \geq 0$ . Then*

$$(2) \quad (i) \|e^{-t\sqrt{-\Delta}}f\|_q \leq \|f\|_p \quad \text{for all } f \quad \text{implies} \quad (ii) \quad e^{-t\sqrt{n}} \leq \sqrt{\frac{p-1}{q-1}}.$$

Moreover, (ii) implies (i) in dimensions  $n \leq 3$ . Finally, for any  $q > \max\{2, p\}$ , there exists  $n_0 = n_0(p, q) \geq 4$  such that (ii) does not imply (i) in dimensions  $n$  with  $n \geq n_0$ .

It remains an open problem to find a necessary and sufficient condition on  $t > 0$  in dimensions  $n \geq 4$  for which the semigroup  $e^{-t\sqrt{-\Delta}}$  is hypercontractive from  $L^p(\mathbb{S}^n)$  to  $L^q(\mathbb{S}^n)$ .

## 2. PROOF OF THEOREM 1.1

**2.1. The necessity part (i)  $\Rightarrow$  (ii).** We recall this standard argument for the sake of completeness. Let  $f(\xi) = 1 + \varepsilon H_1(\xi)$  where  $H_1$  is any (real) spherical harmonic of degree 1, i.e.,  $\Delta H_1 = -nH_1$ . Then  $e^{-t\sqrt{-\Delta}}f(\xi) = 1 + \varepsilon e^{-t\sqrt{n}}H_1(\xi)$ . As  $\varepsilon \rightarrow 0$ , we obtain

$$\begin{aligned} & \int_{\mathbb{S}^n} |1 + \varepsilon e^{-t\sqrt{n}}H_1(\xi)|^q d\sigma_n \\ &= \int_{\mathbb{S}^n} \left( 1 + q\varepsilon e^{-t\sqrt{n}}H_1(\xi) + \frac{q(q-1)}{2}\varepsilon^2 e^{-2t\sqrt{n}}H_1^2(\xi) + O(\varepsilon^3) \right) d\sigma_n \\ &= 1 + \frac{q(q-1)}{2}\varepsilon^2 e^{-2t\sqrt{n}}\|H_1\|_2^2 + O(\varepsilon^3). \end{aligned}$$

Thus,

$$(3) \quad \|e^{-t\sqrt{-\Delta}}f\|_q = 1 + \frac{q-1}{2}\varepsilon^2 e^{-2t\sqrt{n}}\|H_1\|_2^2 + O(\varepsilon^3).$$

Similarly, we have

$$(4) \quad \|f\|_p = 1 + \frac{p-1}{2}\varepsilon^2\|H_1\|_2^2 + O(\varepsilon^2).$$

Substituting (3) and (4) into the inequality  $\|e^{-t\sqrt{-\Delta}}f\|_q \leq \|f\|_p$ , and taking  $\varepsilon \rightarrow 0$  we obtain the necessary condition  $e^{-2t\sqrt{n}} \leq \frac{p-1}{q-1}$  which coincides with (ii) in (2).

**2.2. The sufficiency part (ii)  $\Rightarrow$  (i) in dimensions  $n = 1, 2, 3$ .** Our goal is to show that if  $1 < p < q$  and if  $t \geq 0$  is such that  $e^{-t^2\sqrt{n}} \leq \frac{p-1}{q-1}$ , then

$$(5) \quad \|e^{-t\sqrt{-\Delta}}f\|_q \leq \|f\|_p \quad \text{in dimensions } n = 1, 2, 3.$$

The case  $n = 1$  was confirmed in [10]. In what follows we assume  $n \in \{2, 3\}$ . First we need the fact that the heat semigroup  $e^{t\Delta}$  has a nonnegative kernel. Indeed, for each  $t > 0$  there exists  $K_t : [-1, 1] \rightarrow [0, \infty)$  such that

$$e^{t\Delta}f(\xi) = \int_{\mathbb{S}^n} K_t(\xi \cdot \eta)f(\eta)d\sigma_n(\eta),$$

where  $\xi \cdot \eta = \sum_{j=1}^{n+1} \xi_j \eta_j$  for  $\xi = (\xi_1, \dots, \xi_{n+1})$  and  $\eta = (\eta_1, \dots, \eta_{n+1})$ , see, for example, Proposition 4.1 in [7]. Next, we recall the subordination formula

$$(6) \quad e^{-x} = \frac{1}{\sqrt{\pi}} \int_0^\infty e^{-y-x^2/(4y)} \frac{dy}{\sqrt{y}} \quad \text{valid for all } x \geq 0,$$

By the functional calculus, we deduce that the Poisson semigroup  $e^{-t\sqrt{-\Delta}}$  has a positive kernel with total mass 1. The latter fact together with the convexity of the map  $x \mapsto |x|^p$  for  $p \geq 1$  implies that  $\|e^{-t\sqrt{-\Delta}}\|_p \leq \|f\|_p$  for all  $t \geq 0$ . Thus, it suffices to verify (5) for those  $t \geq 0$  for which  $e^{-2t\sqrt{n}} = \frac{p-1}{q-1}$ .

Next we claim that it suffices to verify (5) only for the powers  $p, q$  such that  $2 \leq p \leq q$ . Indeed, assume (5) holds for  $2 \leq p \leq q$ . By duality and the symmetry of the semigroup  $e^{-t\sqrt{-\Delta}}$  we obtain  $\|e^{-t\sqrt{-\Delta}}f\|_{p'} \leq \|f\|_{q'}$  where  $p' = \frac{p}{p-1}$ ,  $q' = \frac{q}{q-1}$ ,  $1 < q' \leq p' \leq 2$ . Notice that  $\frac{p-1}{q-1} = \frac{q'-1}{p'-1}$ , thus we extend (5) to all  $p, q$  such that  $1 < p \leq q \leq 2$ . It remains to extend (5) for those powers  $p, q$  when  $p \leq 2 \leq q$ . To do so, let  $p \leq 2 \leq q$ , and let  $t \geq 0$  be such  $e^{-2t\sqrt{n}} = \frac{p-1}{q-1}$ . Choose  $t_1, t_2 \geq 0$  so that  $t = t_1 + t_2$  and  $e^{-2t_1\sqrt{n}} = p-1$  and  $e^{-2t_2\sqrt{n}} = \frac{1}{q-1}$ . Then we have

$$\|e^{-t\sqrt{-\Delta}}f\|_q = \|e^{-t_2\sqrt{-\Delta}}(e^{-t_1\sqrt{-\Delta}}f)\|_q \leq \|e^{-t_1\sqrt{-\Delta}}f\|_2 \leq \|f\|_p.$$

In what follows we assume  $2 \leq p \leq q$ . We will use a standard argument to deduce the validity of the hypercontractivity estimate from a log Sobolev inequality. Nonnegativity of the kernel for the Poisson semigroup combined with the triangle inequality implies  $|e^{-t\sqrt{-\Delta}}f| \leq e^{-t\sqrt{-\Delta}}|f|$  for any  $f$ . Thus by continuity and standard density arguments we can assume that  $f \geq 0$ ,  $f$  is not identically zero, and  $f$  is smooth in (5).

The equality  $e^{-2t\sqrt{n}} = \frac{p-1}{q-1}$  implies  $q = 1 + e^{2t\sqrt{n}}(p-1)$ . Fix  $p \geq 2$  and consider the map

$$\varphi(t) = \|e^{-t\sqrt{-\Delta}}f\|_{q(t)} > 0, \quad t \geq 0,$$

where  $q(t) = 1 + e^{2t\sqrt{n}}(p-1)$ . If we show  $\varphi'(t) \leq 0$ , then we obtain  $\varphi(t) \leq \varphi(0) = \|f\|_p$ , and this proves the sufficiency part. Let  $\psi(t) = \ln \varphi(t)$ . We have

$$\frac{q^2}{q'}\psi'(t) = -\ln \left( \int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}}f)^q d\sigma_n \right) + \frac{\int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}}f)^q \left( \ln(e^{-t\sqrt{-\Delta}}f) + \frac{q^2}{q'} \frac{\partial_t e^{-t\sqrt{-\Delta}}f}{e^{-t\sqrt{-\Delta}}f} \right) d\sigma_n}{\int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}}f)^q d\sigma_n}.$$

Clearly  $\psi' \leq 0$  if and only if

$$\begin{aligned} & \int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}} f)^q \ln(e^{-t\sqrt{-\Delta}} f)^q d\sigma_n - \int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}} f)^q d\sigma_n \ln \left( \int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}} f)^q d\sigma_n \right) \\ & \leq \frac{q^2}{q'} \int_{\mathbb{S}^n} (e^{-t\sqrt{-\Delta}} f)^{q-1} \sqrt{-\Delta} (e^{-t\sqrt{-\Delta}} f) d\sigma_n. \end{aligned}$$

Let  $g = e^{-t\sqrt{-\Delta}} f \geq 0$ . Then we can rewrite the previous inequality as

$$(7) \quad \int_{\mathbb{S}^n} g^q \ln g^q d\sigma_n - \int_{\mathbb{S}^n} g^q d\sigma_n \ln \left( \int_{\mathbb{S}^n} g^q d\sigma_n \right) \leq \frac{q^2}{2(q-1)\sqrt{n}} \int_{\mathbb{S}^n} g^{q-1} \sqrt{-\Delta} g d\sigma_n,$$

where we used the fact that  $q' = 2(q-1)\sqrt{n}$ . Since  $e^{-t\sqrt{-\Delta}}$  is contractive in  $L^\infty(\mathbb{S}^n)$  with a nonnegative, symmetric kernel, it follows that the validity of the estimate (7) for  $q = 2$  implies (7) for all  $q \in [2, \infty)$ ; see, e.g., Theorem 4.1 in [3].

Let  $g = \sum_{k \geq 0} H_k$  be the decomposition of  $g$  into its spherical harmonics. Then the estimate (7) for  $q = 2$  takes the form

$$\int_{\mathbb{S}^n} g^2 \ln g^2 d\sigma_n - \int_{\mathbb{S}^n} g^2 d\sigma_n \ln \left( \int_{\mathbb{S}^n} g^2 d\sigma_n \right) \leq \sum_{k \geq 0} 2 \sqrt{\frac{k(k+n-1)}{n}} \|H_k\|_2^2.$$

It follows from Beckner's conformal log Sobolev inequality [1] (which is a consequence of Lieb's sharp Hardy–Littlewood–Sobolev inequality [6]) that for any smooth nonnegative  $g = \sum_{k \geq 0} H_k$  we have

$$\int_{\mathbb{S}^n} g^2 \ln g^2 d\sigma_n - \int_{\mathbb{S}^n} g^2 d\sigma_n \ln \left( \int_{\mathbb{S}^n} g^2 d\sigma_n \right) \leq \sum_{k \geq 0} \Delta_n(k) \|H_k\|_2^2$$

with  $\Delta_n(k) = 2n \sum_{m=0}^{k-1} \frac{1}{2m+n}$ . Thus, the estimate (5) is a consequence of the following lemma.

**Lemma 2.1.** *Let  $n \in \{2, 3\}$ . Then for all integers  $k \geq 1$  one has*

$$n \sum_{m=0}^{k-1} \frac{1}{2m+n} \leq \sqrt{\frac{k(k+n-1)}{n}}.$$

*Proof.* We first check the inequality for  $k \leq 3$  by direct computation. Indeed, the case  $k = 1$  is an equality. The case  $k = 2$  can be checked as follows,

$$1 + \frac{n}{2+n} = \frac{2+2n}{2+n} \leq \sqrt{\frac{2+2n}{n}},$$

which is true because  $n(2+2n) \leq (2+n)^2$  holds for  $n = 2, 3$ . The case  $k = 3$  can be checked similarly:

$$\frac{2+2n}{2+n} + \frac{n}{4+n} \leq \sqrt{\frac{6+3n}{n}}$$

holds for  $n = 2, 3$  (notice that this inequality fails for  $n = 4$ ).

Next, we assume  $k \geq 4$ . We have

$$\sum_{m=0}^{k-1} \frac{1}{m + \frac{n}{2}} = \frac{2}{n} + \sum_{m=1}^{k-1} \frac{1}{m + \frac{n}{2}} \leq \frac{2}{n} + \int_0^{k-1} \frac{1}{x + \frac{n}{2}} dx = \frac{2}{n} + \ln \left( \frac{k + \frac{n}{2} - 1}{\frac{n}{2}} \right).$$

Thus it suffices to show

$$\frac{2}{n} + \ln\left(\frac{k + \frac{n}{2} - 1}{\frac{n}{2}}\right) - \frac{2}{n} \sqrt{\frac{k(k+n-1)}{n}} \leq 0.$$

Notice that the left hand side, call it  $h(k)$ , is decreasing in  $k$ . Indeed, we have

$$h'(k) = \frac{1}{\frac{n}{2} + k - 1} - \frac{2k + n - 1}{n\sqrt{k(k+n-1)}} \leq \frac{1}{\frac{n}{2} + k - 1} - \frac{1}{\sqrt{kn}} \leq \frac{1}{2\sqrt{\frac{n}{2}(k-1)}} - \frac{1}{\sqrt{kn}} \leq 0.$$

On the other hand, we have for  $n = 2, 3$ ,

$$h(4) = \frac{2}{n} + \ln\left(\frac{6+n}{n}\right) - \frac{2}{n} \sqrt{\frac{12+4n}{n}} \leq 0.$$

Indeed, if  $n = 2$ ,  $h(4) = 1 + 2\ln 2 - \sqrt{10} < 0$ , and if  $n = 3$ ,  $h(4) = \frac{2+3\ln 3-4\sqrt{2}}{3} < 0$ .  $\square$

**2.3. Counterexample to (ii)  $\Rightarrow$  (i) in high dimensions.** Let  $\lambda := \frac{n-1}{2}$ , and let  $C_d^{(\lambda)}(x)$  be the Gegenbauer polynomial

$$(8) \quad C_d^{(\lambda)}(x) = \sum_{j=0}^{\lfloor \frac{d}{2} \rfloor} (-1)^j \frac{\Gamma(d-j+\lambda)}{\Gamma(\lambda)j!(d-2j)!} (2x)^{d-2j},$$

where  $\lfloor \frac{d}{2} \rfloor$  denotes the largest integer  $m$  such that  $m \leq \frac{d}{2}$ , and  $\Gamma(x)$  is the Gamma function. Notice that if we let  $Y_d(\xi) = C_d^{(\lambda)}(\xi \cdot e_1)$ , where  $e_1 = (1, 0, \dots, 0) \in \mathbb{R}^{n+1}$ , then  $Y_d(\xi)$  is a spherical harmonic of degree  $d$  on  $\mathbb{S}^n$ . In particular, for  $t \geq 0$  such that  $e^{-2t\sqrt{n}} = \frac{p-1}{q-1}$ , the estimate  $\|e^{-t\sqrt{-\Delta}} f\|_{L^q(\mathbb{S}^n)} \leq \|f\|_{L^p(\mathbb{S}^n)}$  applied to  $f = Y_d(\xi)$  is equivalent to the estimate

$$(9) \quad \frac{\|Y_d\|_q}{\|Y_d\|_p} \leq e^{t\sqrt{d(d+n-1)}} = \left(\frac{q-1}{p-1}\right)^{\frac{1}{2}\sqrt{\frac{d(d+n-1)}{n}}}.$$

Next, we need

**Lemma 2.2.** *For any  $d \geq 0$  we have*

$$(10) \quad \lim_{n \rightarrow \infty} \frac{\|Y_d\|_{L^q(\mathbb{S}^n, d\sigma_n)}}{\|Y_d\|_{L^p(\mathbb{S}^n, d\sigma_n)}} = \frac{\|h_d\|_{L^q(\mathbb{R}, d\gamma)}}{\|h_d\|_{L^p(\mathbb{R}, d\gamma)}},$$

where  $d\gamma(y) = \frac{e^{-y^2/2}}{\sqrt{2\pi}} dy$  is the standard Gaussian measure on the real line, and  $h_d(x)$  is the probabilistic Hermite polynomial

$$(11) \quad h_d(x) = \sum_{j=0}^{\lfloor \frac{d}{2} \rfloor} \frac{(-1)^j d!}{j!(d-2j)!} \frac{x^{d-2j}}{2^j}.$$

*Proof.* Indeed, notice that

$$(12) \quad \|Y_d\|_p^p = \int_{\mathbb{S}^n} |C_d^{(\lambda)}(\xi \cdot e_1)|^p d\sigma_n(\xi) = \int_{-1}^1 |C_d^{(\lambda)}(t)|^p c_\lambda(1-t^2)^{\lambda-\frac{1}{2}} dt,$$

where  $c_\lambda = \frac{\Gamma(\lambda+1)}{\Gamma(\frac{\lambda}{2})\Gamma(\lambda+\frac{1}{2})}$ . In particular, after the change of variables  $t = \frac{s}{\sqrt{2\lambda}}$  in (12), and multiplying both sides in (12) by  $(d!/(2\lambda)^{d/2})^p$  we obtain

$$\left(\frac{d!}{(2\lambda)^{d/2}}\right)^p \|Y_d\|_p^p = \int_{\mathbb{R}} \left| \frac{d!}{(2\lambda)^{d/2}} C_d^{(\lambda)} \left( \frac{s}{\sqrt{2\lambda}} \right) \right|^p \frac{c_\lambda}{\sqrt{2\lambda}} \left(1 - \frac{s^2}{2\lambda}\right)^{\lambda-\frac{1}{2}} \mathbb{1}_{[-\sqrt{2\lambda}, \sqrt{2\lambda}]}(s) ds,$$

where  $\mathbb{1}_{[-\sqrt{2\lambda}, \sqrt{2\lambda}]}(s)$  denotes the indicator function of the set  $[-\sqrt{2\lambda}, \sqrt{2\lambda}]$ . Notice that by Stirling's formula for any  $j \geq 0$ , and any  $d \geq 0$  we have

$$(13) \quad \lim_{\lambda \rightarrow \infty} \frac{1}{\lambda^{d-j}} \frac{\Gamma(d-j+\lambda)}{\Gamma(\lambda)} = 1.$$

Therefore, (11) and (8) together with (13) imply that for all  $s \in \mathbb{R}$  we have

$$\lim_{\lambda \rightarrow \infty} \frac{d!}{(2\lambda)^{d/2}} C_d^{(\lambda)} \left( \frac{s}{\sqrt{2\lambda}} \right) = h_d(s).$$

Invoking Stirling's formula again we have

$$\lim_{\lambda \rightarrow \infty} \frac{c_\lambda}{\sqrt{2\lambda}} \left(1 - \frac{s^2}{2\lambda}\right)^{\lambda-\frac{1}{2}} \mathbb{1}_{[-\sqrt{2\lambda}, \sqrt{2\lambda}]}(s) = \frac{e^{-s^2/2}}{\sqrt{2\pi}} \quad \text{for all } s \in \mathbb{R}.$$

Finally, to apply Lebesgue's dominated convergence theorem it suffices to verify that for all  $s \in \mathbb{R}$  and all  $\lambda \geq \lambda_0$  we have the following pointwise estimates

$$\begin{aligned} a) \quad & \frac{c_\lambda}{\sqrt{2\lambda}} \left(1 - \frac{s^2}{2\lambda}\right)^{\lambda-\frac{1}{2}} \mathbb{1}_{[-\sqrt{2\lambda}, \sqrt{2\lambda}]}(s) \leq C e^{-s^2/2} \\ b) \quad & \frac{d!}{(2\lambda)^{d/2}} C_d^{(\lambda)} \left( \frac{s}{\sqrt{2\lambda}} \right) \leq c_1(d)(1+|s|)^{c_2(d)}, \end{aligned}$$

where  $\lambda_0, C, c_1(d), c_2(d)$  are some positive constants independent of  $\lambda$  and  $s$ .

To verify a) it suffices to consider the case  $s \in [-\sqrt{2\lambda}, \sqrt{2\lambda}]$ . Since  $\lim_{\lambda \rightarrow \infty} \frac{c_\lambda}{\sqrt{2\lambda}} = \frac{1}{\sqrt{2\pi}}$  it follows that  $\frac{c_\lambda}{\sqrt{2\lambda}} \leq C$  for all  $\lambda \geq \lambda_0$ , where  $\lambda_0$  is a sufficiently large number.

Next, the estimate  $(1 - \frac{s^2}{2\lambda})^{\lambda-1/2} \leq C' e^{-s^2/2}$  for  $s \in [-\sqrt{2\lambda}, \sqrt{2\lambda}]$  follows if we show that  $(1 - \frac{t}{2\lambda}) \ln(1-t) \leq C''/\lambda - t$  for all  $t := \frac{s^2}{2\lambda} \in [0, 1]$  where  $C''$  is a universal positive constant. The latter inequality follows from  $\ln(1-t) \leq -t$  for  $t \in [0, 1]$ .

To verify b) it suffices to show that for all  $\lambda \geq \lambda_0 > 0$  and all integers  $j$  such that  $d \geq j \geq 0$  one has

$$\frac{1}{\lambda^{d-j}} \frac{\Gamma(d-j+\lambda)}{\Gamma(\lambda)} \leq C(d-j),$$

where  $C(d-j)$  depends only on  $d-j$ . The latter inequality follows from (13) provided that  $\lambda \geq \lambda_0$  where  $\lambda_0$  is a sufficiently large number.

Thus, it follows from the Lebesgue's dominated convergence theorem that

$$\lim_{n \rightarrow \infty} \frac{d!}{(n-1)^{d/2}} \|Y_d\|_{L^p(\mathbb{S}^n, d\sigma_n)} = \|h_d\|_{L^p(\mathbb{R}, d\gamma)}.$$

The lemma is proved.  $\square$

Now we fix  $q > \max\{p, 2\}$  and, in order to prove the failure of (ii)  $\Rightarrow$  (i) for all sufficiently large  $n$ , we argue by contradiction and assume that there is a sequence of dimensions  $\{n_j\}_{j \geq 1}$  going to infinity such that (ii)  $\Rightarrow$  (i) in Theorem 1.1 does hold. Then, by combining (9) and (10) we have

$$(14) \quad \frac{\|h_d\|_{L^q(\mathbb{R}, d\gamma)}}{\|h_d\|_{L^p(\mathbb{R}, d\gamma)}} \leq \left(\frac{q-1}{p-1}\right)^{\frac{\sqrt{d}}{2}}.$$

On the other hand, a consequence of the main result in [5] and the assumption  $q > \max\{p, 2\}$  is that

$$\lim_{d \rightarrow \infty} \left( \frac{\|h_d\|_{L^q(\mathbb{R}, d\gamma)}}{\|h_d\|_{L^p(\mathbb{R}, d\gamma)}} \right)^{1/d} = \left( \frac{q-1}{\max\{p, 2\}-1} \right)^{\frac{1}{2}},$$

which is in contradiction with (14).

**Remark 2.1.** Let  $B(x, y)$  be the Beta function. The estimate (9) for  $p = 2$  and  $q = 4$  takes the form

$$(15) \quad \int_{-1}^1 |C_d^{(\frac{n-1}{2})}(t)|^4 (1-t^2)^{\frac{n-2}{2}} dt \leq 9 \sqrt{\frac{d(d+n-1)}{n}} \frac{(n-1)^2 B(1/2, n/2)}{d^2 (2d+n-1)^2 B^2(n-1, d)},$$

where we used the fact that  $\|Y_d\|_{L^2(\mathbb{S}^n)}^2 = \frac{n-1}{d(2d+n-1)B(n-1, d)}$ . The numerical computations show that the inequality (15) already fails for  $d = 7$  and  $n = 13$ .

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