

Confirmation of the Long-Period Planet Orbiting Gliese 411 and the Detection of a New Planet Candidate

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ABSTRACT

We perform a detailed characterization of the planetary system orbiting the bright, nearby M dwarf Gliese 411 using radial velocities gathered by APF, HIRES, SOPHIE, and CARMENES. We confirm the presence of a signal with a period near 2900 days that has been disputed as either a planet or long-period stellar magnetic cycle. An analysis of activity metrics including H_α and $\log'R_{HK}$ indices supports the interpretation that the signal corresponds to a Neptune-mass planet, GJ 411 c. An additional signal near 215 days was previously dismissed as an instrumental systematic, but our analysis shows that a planetary origin cannot be ruled out. With a semi-major axis of 0.5141 ± 0.0038 AU, this candidate's orbit falls between those of its companions and skirts the outer edge of the habitable zone. It has a minimum mass of $4.1 \pm 0.6 M_\oplus$, giving a radial velocity amplitude of 0.83 ± 0.12 m s⁻¹. If confirmed, this would be one of the lowest-amplitude planet detections from any of these four instruments. Our analysis of the joint radial velocity data set also provides tighter constraints on the orbital parameters for the previously known planets. Photometric data from *TESS* does not show any signs of a transit event. However, the outermost planet and candidate are prime targets for future direct imaging missions and GJ 411 c may be detectable via astrometry.

1. INTRODUCTION

The discovery of planets orbiting nearby stars has traditionally relied on the radial-velocity (RV) detection method. Out of the 89 known planets within 10 pc of the Solar System, 86 were discovered—and each has been observed—using this method. In contrast, less than a fifth of all known planets were discovered using Doppler spectroscopy.¹

Many of these radial-velocity detections were made through systematic monitoring such as the Eta-Earth Survey (Howard et al. 2009), which used the HIRES spectrograph at the W.M. Keck Observatory (Vogt et al. 1994) to find low-mass ($\sim 3 - 30 M_\oplus$) planets around 230 of the closest GKM stars. The California Planet Search (CPS) group has continued similar observing campaigns, including the APF-50 Survey (Fulton et al. 2016; Fulton 2017), which uses the robotic Automatic Planet Finder (APF; Radovan et al. 2014; Vogt et al. 2014) to observe the 51 brightest and least chromospherically active stars in the Eta-Earth Survey at high cadence. Most of these stars have been observed over baselines spanning more than one or two decades, offering the opportu-

nity to detect long-period planets. Moreover, high-resolution spectroscopy of nearby stars can be obtained with shorter exposures, allowing many precise RV measurements that make it easier to detect low-amplitude signals. As a consequence, these surveys have played an important role in understanding the frequency of low-mass planets, along with their characteristics (Howard et al. 2010; Mayor et al. 2011; Fulton 2017).

Gliese 411, also known as Lalande 21185 or HD 95735, is one of the brightest and closest M dwarfs to our Solar System (Lépine & Gaidos 2011; Gaidos et al. 2014), making it a frequent target in the search for exoplanets. Early astrometric studies suggested the presence of an orbital companion with estimates of a minimum mass ranging from 0.1 to 0.3 M_\odot (van de Kamp & Lippincott 1951; Lippincott 1960). Also using astrometry, Gatewood (1996) later claimed that a planet with a minimum mass of 0.9 M_J was orbiting the star. None of these claims have since been verified. Over the past five years, several RV searches for planets around GJ 411 have been published, beginning with Butler et al. (2017) who used HIRES spectroscopy to identify a 9.9 day signal corresponding to a minimum mass of 3.8 M_\oplus . This signal, however, has been disputed by subsequent analyses (Díaz et al.

¹ <https://exoplanetarchive.ipac.caltech.edu>

2019). Using RVs from the SOPHIE spectrograph (Perruchot et al. 2008), Díaz et al. (2019) discovered GJ 411 b, a planet with a minimum mass of $3 M_{\oplus}$ on a 12.95 day orbit. Stock et al. (2020) subsequently combined the SOPHIE data with RVs from the CARMENES spectrograph (Quirrenbach et al. 2014, 2018) to confirm this planet. They also identified a 2900 day periodicity, which they attributed to long-period magnetic activity cycles. The California Legacy Survey also detected GJ 411 b and the 2900 day signal but instead attributed the latter to a planet with a minimum mass of $18 M_{\oplus}$, GJ 411 c (Rosenthal et al. 2021).

In this paper, we combine additional RVs gathered by the Eta-Earth and APF-50 Surveys with data from Díaz et al. (2019), Stock et al. (2020), and Rosenthal et al. (2021) to conduct a detailed search for planets around GJ 411. We find that the disputed long-period signal is inconsistent with the star’s magnetic activity cycle and favor the planetary interpretation. We investigate another signal identified by Rosenthal et al. (2021) as an instrumental systematic and show that a planetary origin cannot be ruled out. Additionally, we provide new constraints on the orbital parameters of GJ 411 b and the most detailed characterization of the system to date. In Section 2, we describe our observations and RVs. We provide an updated analysis of the star’s properties in Section 3. In Section 4, we describe our analysis of the RVs. We then discuss the transit, direct imaging, and astrometric detection prospects in Section 5, and place the system in the broader exoplanet population. Lastly, in Section 6, we conclude our results.

2. OBSERVATIONS

2.1. HIRES Spectroscopy

We obtained high-resolution spectra of GJ 411 with the HIRES spectrograph located on the Keck I telescope (Vogt et al. 1994) at Mauna Kea, HA. The star was observed with a resolving power of $R \approx 60,000$ over wavelengths ranging roughly from 300 nm to 1000 nm. Spectra were wavelength-calibrated using a warm iodine-gas cell and then modeled using a deconvolved stellar spectral template, an atlas iodine spectrum, and an approximate line spread function. We derived RVs following the techniques described by Butler et al. (1996). We additionally calculated the $\log R'_{\text{HK}}$ and H_{α} indices for each spectra, following the procedures outlined in Isaacson & Fischer (2010) and Robertson et al. (2014).

While our observations originally ranged from June 1997 to January 2021, HIRES received a new CCD and other optical improvements in 2004. Before these upgrades, charge transfer efficiency in the CCD would blur the dense spectral lines found in the spectra of M dwarfs like GJ 411, contaminating the RVs. Consequently, we only include the HIRES data gathered after these upgrades, leaving 299 spectra beginning in November 2004. Five observations had low counts,

Table 1. Radial Velocities of Gliese 411

BJD -2450000	RV m s^{-1}	Uncertainty m s^{-1}	H_{α}	$\log R'_{\text{HK}}$	Instrument
8955.7945	-0.6962	1.5564	0.0917	-5.5499	APF
8955.8042	-1.1892	1.4823	0.0909	-5.4362	APF
8955.8106	1.4112	1.5855	0.0910	-5.5268	APF
9185.0451	-2.3714	1.8003	0.0895	-5.4849	APF
9185.0502	-2.4402	1.9729	0.0891	-5.5193	APF
9185.0553	-0.0376	1.7513	0.0893	-5.5403	APF
9192.9958	-3.4514	1.7958	0.0904	-5.3846	APF

NOTE—This table is available in its entirety in machine-readable form.

resulting in decreased signal to noise and inflated uncertainties; we discarded these measurements, giving a total of 294 radial velocities. Our HIRES observations up to 2014 are shared with Butler et al. (2017) but processed through a separate pipeline. Additionally, all RVs collected prior to April 2020 are shared with Rosenthal et al. (2021).

2.2. APF Spectroscopy

We collected 435 spectra of GJ 411 using the Automated Planet Finder (APF; Radovan et al. 2014; Vogt et al. 2014) between March 2014 and February 2021. The APF is a robotic telescope located at Lick Observatory on Mt. Hamilton, CA. It is equipped with the Levy Spectrograph, a high-resolution echelle spectrometer that achieves $R \approx 100,000$ over wavelengths ranging from 374.3 nm to 980 nm. Spectra passed through a warm iodine-gas cell for wavelength calibration and were modeled the same way as the HIRES spectra in Section 2.1. Radial velocities were then derived with the pipeline described by Fulton et al. (2015), which originates from the Butler et al. (1996) pipeline used for the HIRES observations. We also calculated the $\log R'_{\text{HK}}$ and H_{α} indices for each APF spectrum. 13 of the observations had low counts and were excluded from our analysis, leaving 422 RVs. Similar to the HIRES observations, all APF data from before April 2020 were included in Rosenthal et al. (2021)’s analysis. The new APF and HIRES RVs are provided along with their uncertainties and respective activity indices in Table 1.

2.3. SOPHIE Spectroscopy

Díaz et al. (2019) collected 155 spectra of GJ 411 between October 2011 and June 2018 using the SOPHIE spectrograph (Perruchot et al. 2008, 2011) located on the 1.93 m reflector telescope at the Haute-Provence Observatory in France. The observations were made in high-resolution mode ($R \approx 75,000$) and span wavelengths ranging from 387.2 nm

Table 2. Stellar Parameters for Gliese 411

Parameter	Value	Units	Source
R.A.	11 03 19.43	hh mm ss	(1)
Dec.	+35 56 55.15	dd mm ss	(1)
μ_α	-580.057 ± 0.026	mas yr ⁻¹	(1)
μ_δ	-4776.589 ± 0.030	mas yr ⁻¹	(1)
Parallax	392.753 ± 0.032	mas	(1)
Distance	2.54613 ± 0.00021	pc	(1)
T_{eff}	3719^{+21}_{-31}	K	(2)
Rotation Period	56.16 ± 0.27	days	(3)
Mass	$0.3899^{+0.0106}_{-0.0061}$	M_\odot	(2)
Radius	$0.3685^{+0.0081}_{-0.0054}$	R_\odot	(2)
Age	$8.047^{+3.958}_{-4.523}$	Gyr	(2)
[Fe/H]	$-0.3621^{+0.0872}_{-0.0687}$...	(2)
log g	$4.895^{+0.008}_{-0.010}$	cgs	(2)
log R'_{HK}	$-5.47^{+0.11}_{-0.09}$...	(2)

NOTE—(1) Gaia Early Data Release 3 (Gaia Collaboration et al. 2016a, 2020); (2) This Work; (3) Díaz et al. (2019)

to 694.3 nm. In addition to radial velocities, they derive observables such as $\log R'_{\text{HK}}$ and H_α indices for each spectra, although they note the latter are influenced by an unknown systematic present in those data for other intensely observed stars. We obtained these measurements using the Vizier catalogue access tool. The original RVs provided had not been corrected for instrumental zero-point changes, however, the appropriate offsets were calculated from a set of constant RV stars (Courcol et al. 2015; Hobson et al. 2018) and had been provided along with the rest of the publicly available data. After applying these offsets, there still appeared to be a discrepancy between the RVs and the data presented in Díaz et al. (2019); we concluded that they were not corrected for secular acceleration. We accounted for this discrepancy using the model described by Kürster et al. (2003), which relies on the star’s measured distance, proper motion, and absolute radial velocity. Informative priors on the distance and proper motion were placed using data from Gaia Early Data Release 3 (see Table 2; Gaia Collaboration et al. 2016a, 2020). After fitting this model, the residuals are used throughout the remainder of this work as the SOPHIE RVs. 2 radial velocities had uncertainties greater than three times the median uncertainty and were removed, leaving 153 data points.

2.4. CARMENES Spectroscopy

Stock et al. (2020) collected 321 spectra of GJ 411 using the CARMENES spectrograph (Quirrenbach et al. 2014, 2018), installed at the 3.5 m telescope of the Calar Alto Observatory in Spain. Radial velocities were derived from the

VIS channel, which has a resolution of $R \approx 95,000$ and a wavelength coverage of 520 nm to 960 nm. We accessed the publicly available RVs through Vizier, along with the respective H_α indices, although we note these are calculated using a different method from the APF and HIRES indices (Schöfer et al. 2019). The differential line width (dLW) and chromatic index (CRX) for each observation were also provided. Several RVs had large error bars and we removed any where the uncertainties were greater than three times the median uncertainty, leaving 309 data points.

2.5. TESS Photometry

GJ 411 was observed by NASA’s TESS mission (Ricker et al. 2015) during sector 22, between February 19 and March 17 of 2020. We downloaded the short-cadence data file from the Mikulski Archive for Space Telescopes (MAST) and analyzed the pre-search data conditions simple aperture photometry (PDCSAP) light curve, which has been corrected for instrumental systematics (Stumpe et al. 2012; Smith et al. 2012; Stumpe et al. 2014).

3. STELLAR PROPERTIES

Gliese 411 is an M1.9 star (Mann et al. 2015) located in the constellation of Ursa Major. The star is one of our closest neighbors at a distance of 2.5451 ± 0.0002 pc (Gaia Collaboration et al. 2016a, 2020). It also has a V-band magnitude of 7.49, making it the third-brightest red dwarf (Lépine & Gaidos 2011; Gaidos et al. 2014).

Following the procedures outlined in Petigura et al. (2017) and Dedrick et al. (2020), which utilize the publicly available packages SpecMatch-Empirical (Yee et al. 2017) and isoclassify (Huber 2017; Berger et al. 2020), we derive the effective temperature, mass, radius, age, metallicity, and log-gravity for the star from our HIRES and APF spectra. Using the available values from each instrument, we also find that the star has a median $\log R'_{\text{HK}}$ of -5.47 ± 0.10 , suggesting that it experiences very low activity levels. However, previous photometric observations have shown rotational modulation with a period of 56.15 ± 0.27 days (Díaz et al. 2019). Each of these stellar properties can be found in Table 2.

4. ANALYSIS OF RADIAL VELOCITY SIGNALS

4.1. Signal Discovery

We generate ΔBIC periodograms using RVSearch (Rosenthal et al. 2021) to identify periodic signals in each data set along with the combined RVs, shown in Figure 1. This type of search is preferred to generalized Lomb-Scargle periodograms (GLS; Zechmeister & Kürster 2009) because it accounts for instrument-specific parameters such as jitter and offset, allowing us to more accurately search the data from all four instruments simultaneously. Several peaks with false alarms probabilities (FAPs) under 0.1% appear in the

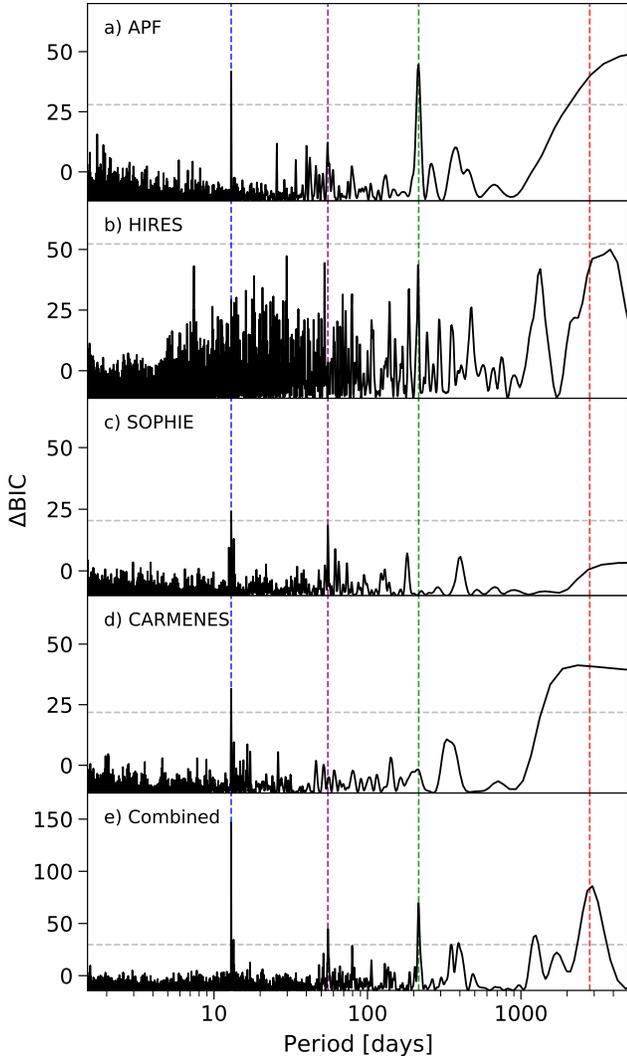


Figure 1. Δ BIC periodograms for each data set (first four panels) and the combined radial velocities (bottom panel). The horizontal dashed grey lines mark the 0.1% FAP threshold for each periodogram. The vertical dashed lines mark the periods of interest, with blue corresponding to GJ 411 b (~ 12.9 days), red denoting the long-period, disputed signal (~ 2900 days), green marking the candidate signal (~ 215.2 days), and purple indicating the rotation period (~ 55.3 days).

periodogram of the combined data, with the most prominent located at 12.9 days and corresponding to the known planet GJ 411 b. This signal is easily identified separately in the APF, SOPHIE, and CARMENES data but surprisingly is not visible in the HIRES periodogram. It is possible that HIRES, the lowest-resolution instrument out of all four, blurred the dense spectral lines of the M dwarf, contributing to increased noise and uncertainty in the RVs. This, combined with the data set’s relatively low cadence, could hide a low-amplitude ($< 2 \text{ m s}^{-1}$) signal observed by the other instruments such as GJ 411 b.

The second peak is located at 2923 days, corresponding to the disputed signal identified by [Stock et al. \(2020\)](#) and [Rosenthal et al. \(2021\)](#). In Section 4.2, we investigate whether the peak is caused by stellar magnetic cycles. The APF and CARMENES periodograms show significant power near this periodicity but plateau because their baselines are too short to constrain any long-period signals. While insufficient precision prevents HIRES from making a strong detection of the signal, a somewhat resolved peak appears in the corresponding periodogram. Two marginally significant peaks located near 340 days and 420 days appear to be one-year aliases of this long-period signal and disappear, along with a peak near 1300 days, when fitting a Keplerian to a period near 2900 days. This 1300 day peak is not well-distinguished from the long-period signal and we conclude it is likely a harmonic or alias given that our data set’s baseline is unable to well-constrain the 2900 day periodicity.

The third peak in the combined RVs is located at 215.2 days and was also identified by [Rosenthal et al. \(2021\)](#), but was dismissed as an APF systematic, primarily because it was not detected in the HIRES data. We note that while not significant, there is a resolved peak exactly at this period in the HIRES periodogram, although the signal is not independently detected in the SOPHIE or CARMENES data either. We further investigate this periodicity’s origins in Section 4.3 and will refer to it as the candidate signal.

A fourth peak is found at 55.3 days, which we note is approximately one-fourth of 215 days and could easily be mistaken for a harmonic of the candidate signal. However, considering that [Díaz et al. \(2019\)](#) found the stellar rotation period to be 56.16 days using well-sampled photometry, and that this peak is surrounded by a forest of smaller peaks, this most likely corresponds to rotational modulation and is consistent with the known distribution of rotation periods for cool stars ([McQuillan et al. 2014](#)). Out of the four data sets, this modulation is most significant in the SOPHIE RVs.

4.2. Isolating the Magnetic Activity Cycle

[Stock et al. \(2020\)](#) first identified a long-period signal of 2852 ± 568 days using the SOPHIE and CARMENES data. They primarily use activity metrics derived from the CARMENES spectra to vet the origins of this periodicity, claiming that GLS periodograms show elevated power near 2900 days. We reproduce these periodograms for the CARMENES H_{α} , CRX, and dLW indices along with observables from the other instruments in Figure 2. While periodograms of the APF, SOPHIE, and CARMENES activity indices are all elevated near 2900 days, they also plateau and do not resolve any peaks, indicating that they have baselines far too short for a periodogram analysis to constrain the origins of a long periodicity. [Stock et al. \(2020\)](#) also identify several peaks near 1400 days, approximately half of the sig-

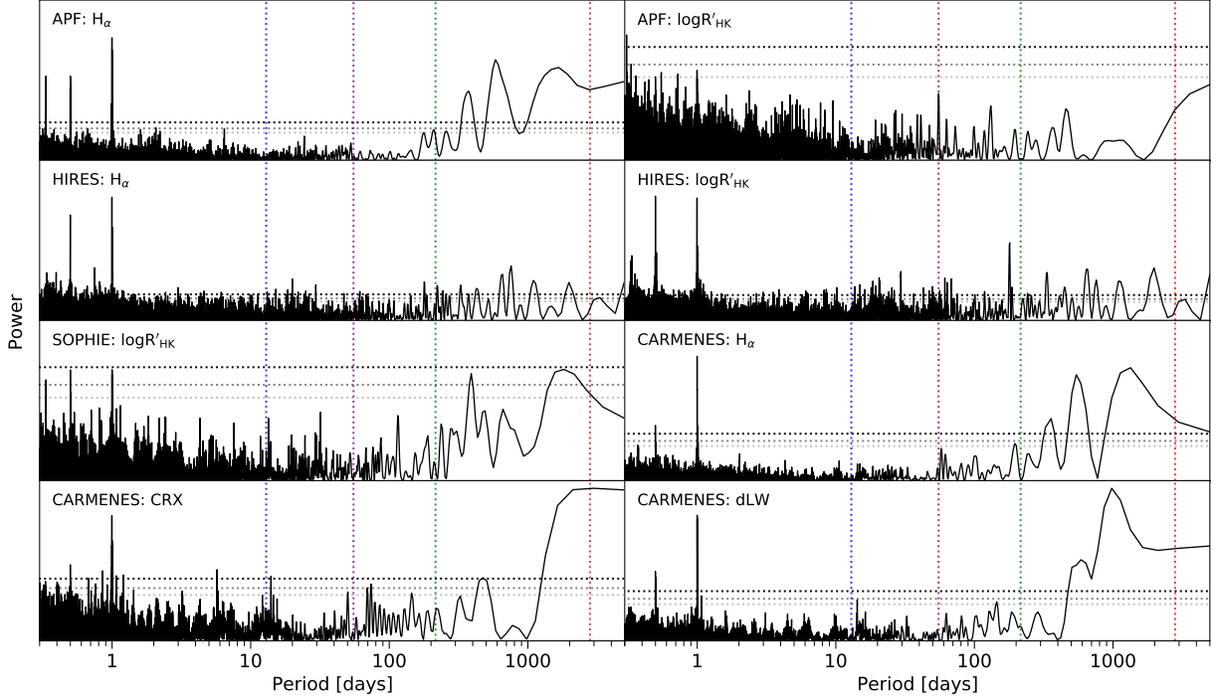


Figure 2. GLS periodograms of the activity indices for each instrument. The three horizontal dashed lines in each panel represent the 1%, 0.5%, and 0.01% FAP levels for the corresponding periodogram. As with Figure 1, the vertical dashed lines represent the signals of interest, with blue corresponding to GJ 411 b, purple representing the rotational modulation, green marking the candidate signal, and red denoting the long-period, disputed signal. No significant or resolved peaks appear near these four periodicities, however periodograms for the APF, SOPHIE, and CARMENES indices do show elevated power near the long-period signal of ~ 2900 days. This structure likely appears because the baselines for these data sets are too short to constrain any long periods.

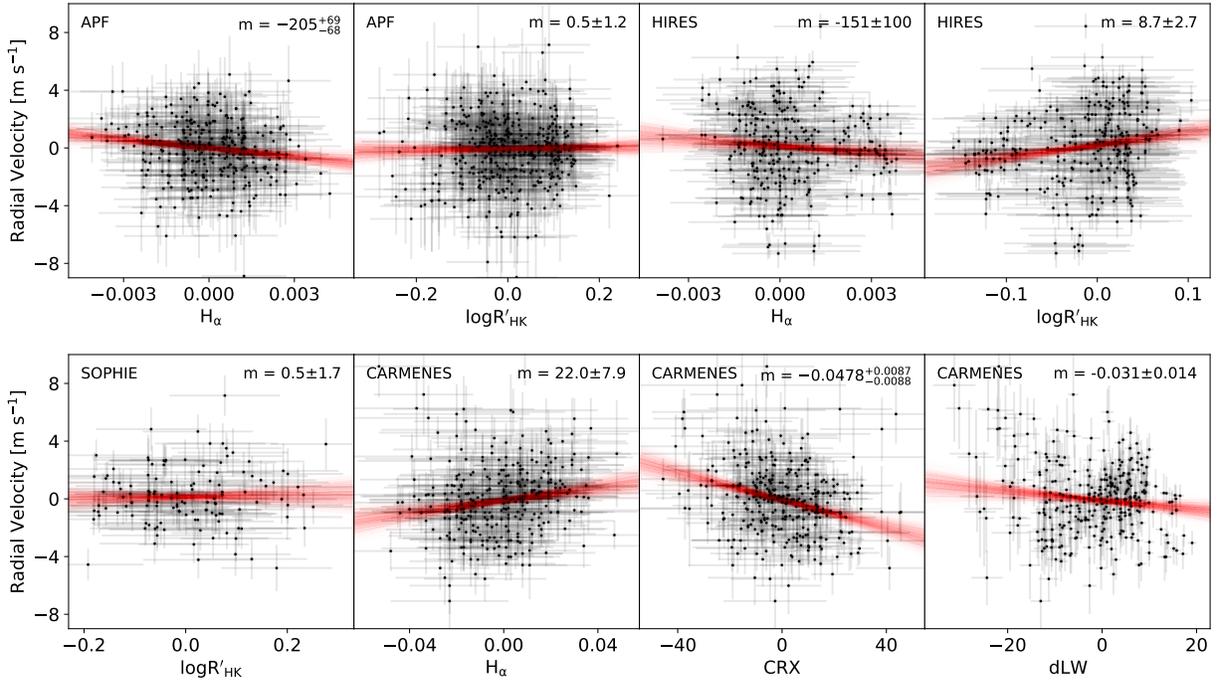


Figure 3. Plots of the radial velocities versus activity indices for each data set. Signals corresponding to GJ 411 b, the candidate, and stellar rotation have been subtracted from the RVs using the model presented in Section 4.4. Samples of linear models from the MCMC posterior for each pair of variables are shown as red lines. The 68.3% confidence interval for the slope is given in the top-right corner of each subplot. No correlations are observed for the APF and HIRES $\log R'_{\text{HK}}$ values while the rest of the correlations appear to be weak and can be explained by non-astrophysical signals or stellar magnetic cycles that span even longer than 2900 days.

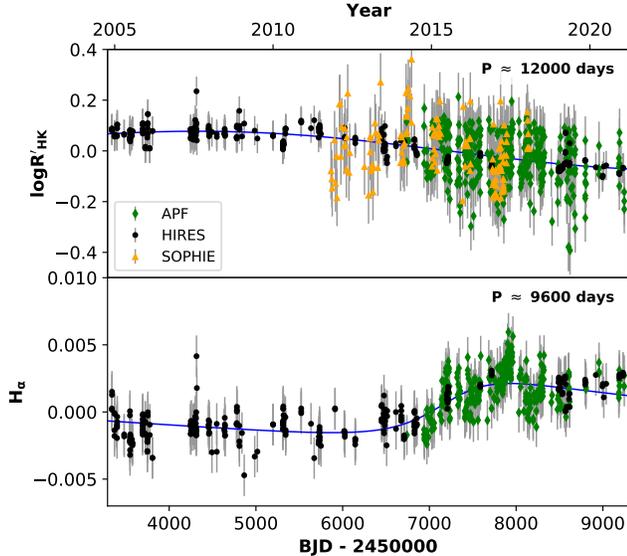


Figure 4. The time series data and best-fit Keplerian models for the $\log R'_{\text{HK}}$ (top) and H_{α} (bottom) activity indices. We observe less than one cycle of the long-period signals contained in both, whereas our radial velocity data spans nearly two cycles of the 2900 day signal (see Figure 8 for comparison). This indicates that the disputed signal does not correspond to stellar magnetic activity and instead originates from a planet designated as GJ 411 c.

nal in question. While our periodogram analysis also shows these peaks for the CARMENES H_{α} and dLW metrics, this period is suspiciously close to the length of time spanned by the CARMENES data.

Rosenthal et al. (2021) independently identified a signal of 3190 ± 185 days, consistent with the long periodicity from Stock et al. (2020), using their HIRES and APF RVs. They note that their S_{HK} indices for HIRES, which span a baseline longer than the disputed signal, show signs of a long-term trend. However, periodograms of the indices reveal no power near 2900 days. Similarly, our periodograms of the HIRES H_{α} and $\log R'_{\text{HK}}$ indices do not contain significant power in this region. This indicates that any detectable magnetic activity cycle must be much longer and lends credence towards a planetary origin.

Because our periodogram analysis is limited by the baseline of each instrument, we also search for correlations between the radial velocities and stellar activity metrics for each data set, following the procedure outlined in Robertson et al. (2014). We are most interested in whether or not the long period is caused by activity, so we remove the signals corresponding to GJ 411 b, the candidate, and stellar rotation from our RVs using the model presented in Section 4.4. We also reject any outlying activity indices with a three-sigma clip and then offset each data set such that their respective means are zero, mitigating degeneracies between slope and intercept when searching for linear trends. To determine whether

the RVs and activity metrics are correlated, we fit a linear model accounting for uncertainties in both variables along with intrinsic scatter (Hogg et al. 2010) using a Markov chain Monte Carlo (MCMC) process. If a slope of zero is within one sigma of the resulting posterior, we consider there to be no correlation. Plots showing representative samples of the posteriors for each activity index can be found in Figure 3. Only two indices, the APF and SOPHIE $\log R'_{\text{HK}}$ values, do not show any correlation. However, we note that all of the correlations appear to be weak, with the actual data showing large amounts of scatter. Further, these relationships seem to be independent of the 2900 day signal. While periodograms of the HIRES H_{α} and $\log R'_{\text{HK}}$ values show no power near 2900 days, both of the indices are correlated with the RVs, implying some other signal must be driving these relationships, such as stellar magnetic cycles even longer than 2900 days. Although the CARMENES H_{α} values are derived differently from the APF and HIRES H_{α} values, both are related to the line width and we expect them to behave similarly. We see that the APF and HIRES H_{α} indices are negatively correlated with the RVs while the CARMENES H_{α} indices are positively correlated, meaning some of the observed relationships could be non-astrophysical. While unable to provide a conclusive answer, our correlation analysis provides no strong evidence in favor of stellar activity as an origin of the disputed signal.

Visually, there is a long-term trend present in the $\log R'_{\text{HK}}$ and H_{α} indices. We use `RadVel` (Fulton et al. 2018) to jointly model the APF, HIRES, and SOPHIE $\log R'_{\text{HK}}$ values with a single Keplerian. The best-fit model returns a period of $\sim 12,000$ days, however, because this signal spans such a long time, MCMC fits are poorly constrained and fail to converge. Similarly, we fit a Keplerian model to the APF and HIRES H_{α} indices, returning a best-fit period of $\sim 9,600$ days. These models and the corresponding time series are shown in Figure 4. Interestingly the H_{α} and $\log R'_{\text{HK}}$ indices appear to be anti-correlated, consistent with observations for other low-activity M dwarfs (Gomes da Silva et al. 2011). It is apparent that any long-period magnetic activity cycles must be several times longer than the 2900 day signal and could be driving the weak correlations observed between the activity metrics and RVs. Consequently, we agree with Rosenthal et al. (2021)’s interpretation that the disputed long-period signal originates from a planet, GJ 411 c.

4.3. Investigating the Origin of the Candidate Signal

Rosenthal et al. (2021) detected a 214.59 ± 0.82 day signal in a joint analysis of their APF and HIRES data. This signal was primarily dismissed as a systematic because it was prominent within the APF RVs but went undetected in the HIRES data. As shown in Figure 1 and discussed in Section 4.1, our HIRES periodogram does show a peak at

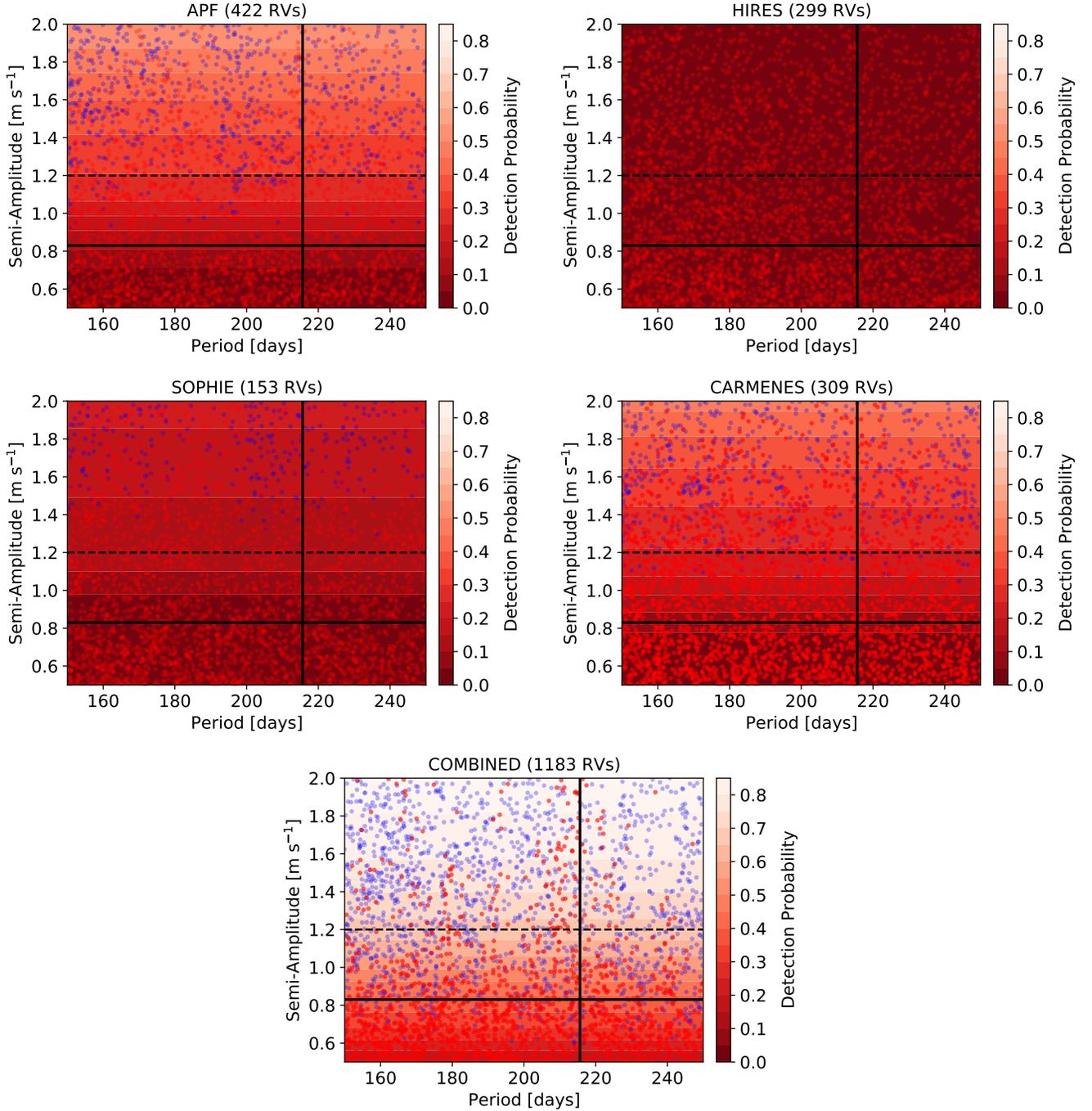


Figure 5. Completeness maps for the APF (top left), HIRES (top right), SOPHIE (middle left), CARMENES (middle right), and complete data sets (bottom). Signals that were recovered by *RVSearch* are denoted by blue points while signals that were not detected are given by red points. The contours represent the probability that a planet with the corresponding period and amplitude are detected. The vertical black line marks the period of the candidate signal while the horizontal solid black line gives the amplitude from our adopted model (see Section 4.4). The dashed horizontal black line marks the elevated amplitude from models that neglect the stellar rotation signal. For both amplitudes, HIRES and SOPHIE are unlikely to detect the candidate signal while the APF and combined RVs are more likely to make a detection than CARMENES.

215 days, however it does not fall above the 0.001 FAP threshold and is not flagged by the *RVSearch* algorithm. The CARMENES and SOPHIE periodograms do not show any power near this period, however, the peak at 215 days

appears to be more prominent in the combined data set than in the APF periodogram with a FAP of 2.56×10^{-10} versus 1.38×10^{-6} . If we only include the HIRES or CARMENES RVs with the APF RVs, the FAP of the 215 day periodicity

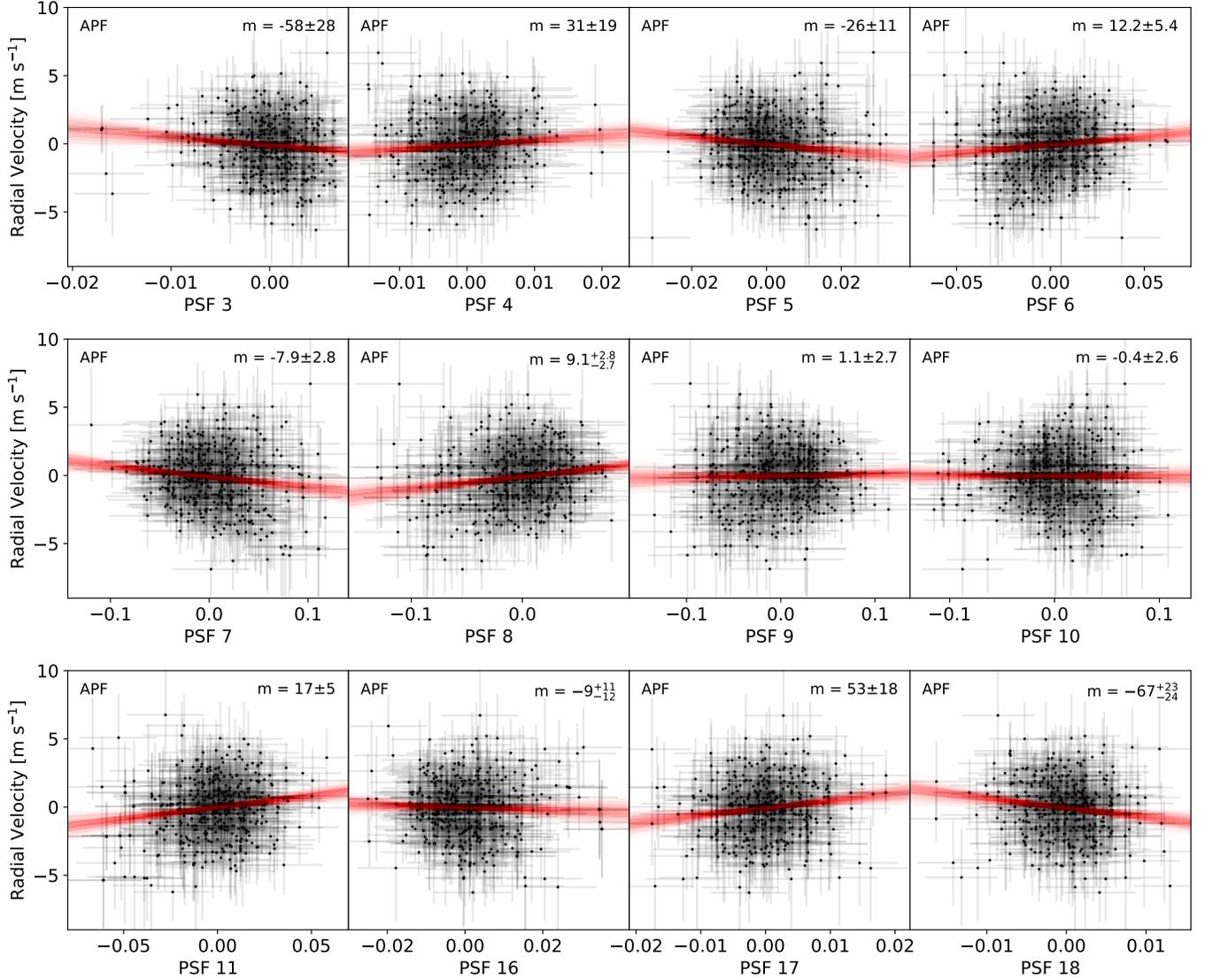


Figure 6. Plots of the radial velocities versus PSF parameters for the APF data set. Signals corresponding to GJ 411 b, GJ 411 c, and stellar rotation have been subtracted from the RVs using the model presented in Section 4.4. Similar to Figure 3, samples of linear models from the MCMC posterior for each pair of variables are shown as red lines while the 68.3% confidence interval for the slope is given in the top-right corner of each subplot. Most of the PSF parameters show some correlation with the RVs, however they appear to be weak and are not strictly caused by the candidate signal.

still improves, suggesting that the signal is present in each of these three data sets. Combining the SOPHIE RVs with the APF data marginally decreases the peak’s FAP, however this is likely because the rotation signal at 55 days becomes much stronger. Visually, the 215 day peak continues to rise further above the noise floor in the periodogram.

To determine whether the signal could be present but go undetected in three of our four RV data sets, we perform injection recovery tests. For each instrument, we randomly generate 3000 signals with periods drawn from a uniform distribution between 150 and 250 days, amplitudes drawn from a uniform distribution between 0.5 and 2 m s^{-1} , eccentricities drawn from a beta distribution described by Kipping

(2013), and arguments of periastron and times of periastron passage drawn from uniform distributions. Each Keplerian is then separately injected into our radial velocities and if `RVSearch` is able to recover a significant peak near the expected period, we consider the signal detectable. Using our results, we create completeness maps revealing the detection limits (Howard & Fulton 2016) for each instrument, shown in Figure 5. Our adopted model in Section 4.4 gives an amplitude of $0.83 \pm 0.12 \text{ m s}^{-1}$, which corresponds to detection probabilities of nearly zero in the HIRES and SOPHIE data. The CARMENES radial velocities have a $\sim 7.5\%$ chance of detecting the signal whereas the APF RVs have a $\sim 11\%$ chance. If we neglect the rotation signal in our modeling,

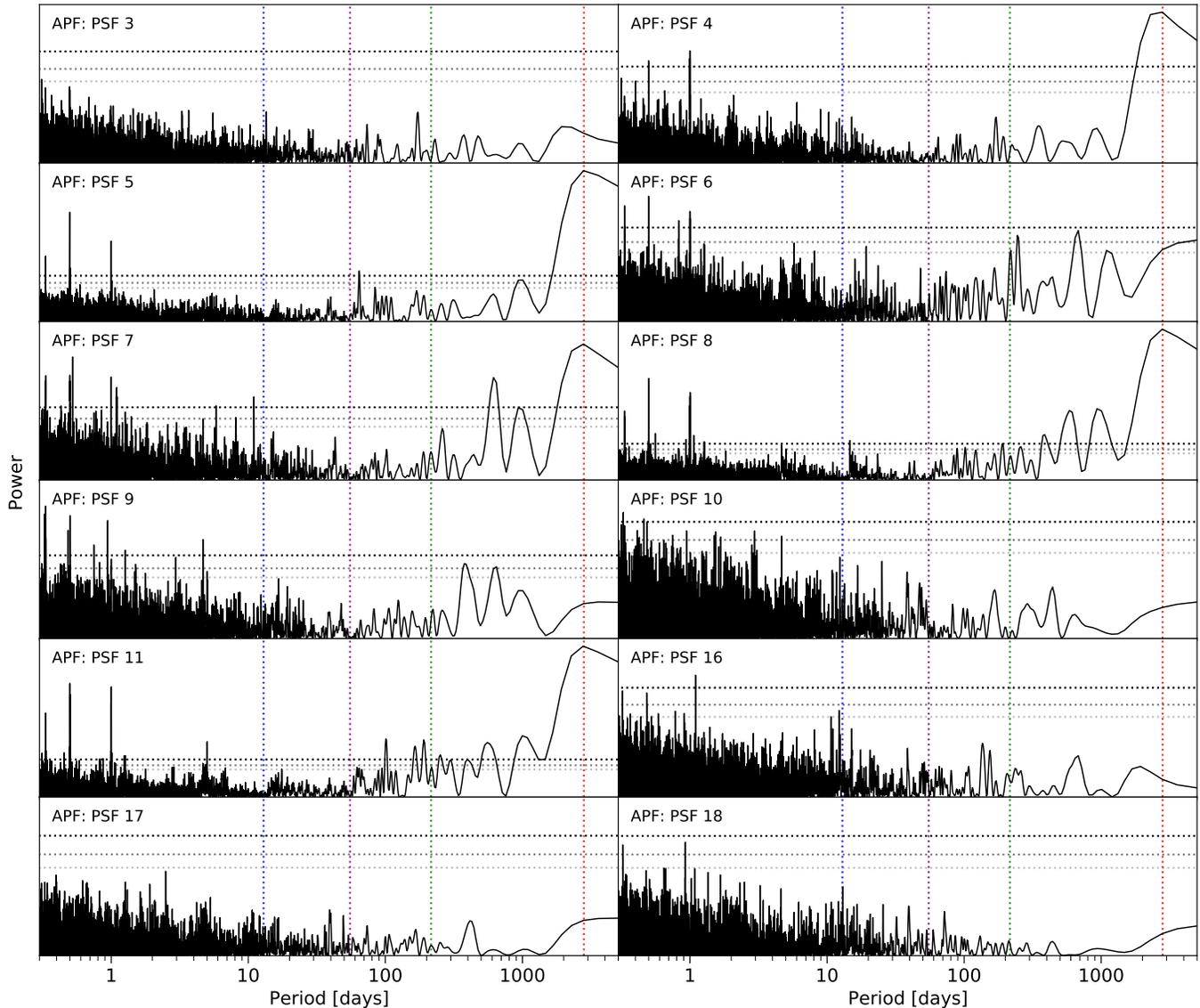


Figure 7. GLS periodograms of the APF PSF parameters. As with Figure 2, the three horizontal dashed lines in each panel represent the 1%, 0.05%, and 0.01% FAP levels for the corresponding periodogram. Additionally, the vertical dashed lines represent the periods of interest, with blue marking GJ 411 b, purple giving the rotational modulation, green denoting the candidate signal, and red representing GJ 411 c. No peak near 215 days appears for any of the PSF parameters, suggesting that the candidate signal is not driven by a systematic.

we find that the signal has a larger amplitude of 1.22 ± 0.16 m s^{-1} (see Section 4.4), however the APF data set is still the most likely to detect the 215 day periodicity. It is feasible that the candidate signal is real and present in all four data sets despite only being detected by the APF.

Rosenthal et al. (2021) also noted that the candidate signal had a high eccentricity, with their $\text{RV}_{\text{Search}}$ output giving values of 0.52 ± 0.18 . This was used to further dismiss the signal as a systematic. In contrast, our modeling returns an eccentricity of 0.15 ± 0.14 , consistent with a circular orbit. We attribute this discrepancy to the stellar rotation signal; when fitting three-Keplerian models and neglecting rotational modulation, we also find elevated eccentricities for

the candidate signal. It is likely that the stellar rotation period, which falls near a harmonic of 215 days, is interfering with the 215 day signal near its maxima and minima, driving artificially high amplitudes and eccentricities.

Common sources of systematics such as instrumental hysteresis should affect the point spread function (PSF) parameters used to model the APF spectra in addition to the radial velocities. Following a procedure similar to that in Section 4.2, we search for correlations between the APF RVs and each PSF parameter. This time, however, we remove the signals corresponding to GJ 411 b, GJ 411 c, and stellar rotation from the RVs, allowing us to primarily search for a 215 day systematic. Additionally, each of the PSF parameters

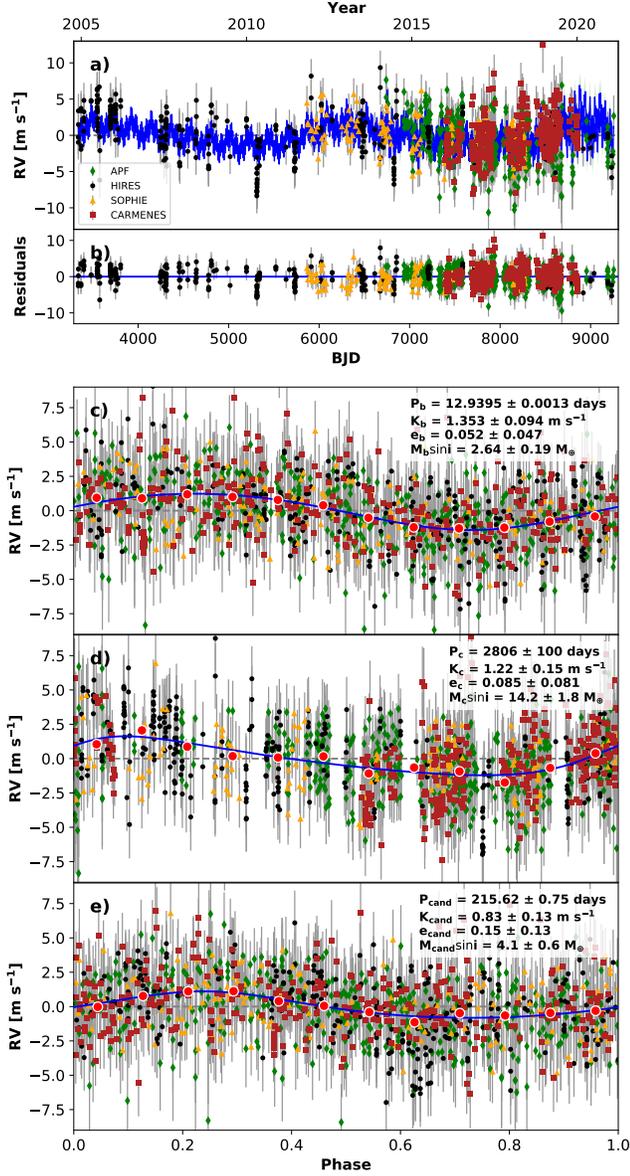


Figure 8. GJ 411 radial velocities with error bars representing observational uncertainties and instrumental jitter values added in quadrature. Panel a: Full time series with the best fit three-planet model and mean GP shown in blue. Panel b: residuals of the best fit model and mean GP. Panel c: RVs phase folded to the ephemeris of GJ 411 b. The Keplerian orbits of all other planets and the rotation signal have been subtracted and the Keplerian fit for the individual planet is shown in blue. Red dots represent the binned velocities. Panels d and e are the same, except they represent the orbits of GJ 411 c and the candidate signal, respectively.

are tightly clustered around some central value, but outliers continue to affect linear models of the data. We perform an initial three-sigma clip on each PSF parameter and then remove any points further than 5 median absolute deviations from the mean. In some cases, the PSF parameters show too little variation to fit a model and are excluded from our anal-

ysis. As seen in Figure 6, nine out of twelve PSF parameters show correlation with the RVs, however the actual data have very large scatters. We also note that the results are highly dependent on the method used to reject outliers. GLS periodograms of the PSF parameters in Figure 7 show that none have significant power near 215 days, confirming that any correlations between the RVs and PSF parameters are not associated with the candidate signal. Ultimately, we find little evidence that a systematic drives the 215 day periodicity in the APF data. However, given that we cannot independently detect the signal in any of the other three data sets, we continue to consider it to be a candidate planet.

4.4. Modeling Signals of Interest

We model the three planetary and candidate signals with Keplerian orbits using the Python package *exoplanet* (Foreman-Mackey et al. 2021a). Additionally, we choose to represent the rotational modulation with a Gaussian process (GP; Rasmussen & Williams 2006; Roberts et al. 2012), which are known to represent the effects of active regions on the surface of a star (Haywood et al. 2014; Rajpaul et al. 2015). Further, GPs give us the flexibility needed to account for changes in rotational modulation between the four instruments, particularly as they cover different wavelengths. In preliminary models, most GPs, including simple periodic kernels, over-fit our data and drive the reduced- χ^2 statistic to values well below one. This is because a combination of cadence and precision prevents us from constraining the high-frequency variation in the rotation signal. However, we can build a simpler model that continues to represent the the long-term variation of the signal by setting the hyper-parameter that dictates the short-term behavior to arbitrarily large values, effectively removing high-frequency variations from our GP. We settle on the simple harmonic oscillator kernel implemented through *celerite2* (Foreman-Mackey et al. 2017; Foreman-Mackey 2018) with τ , the dampening timescale of the process, set to values many times the baseline of our full data set.

No linear or quadratic terms are included in our model because they are highly covariant with the signal for GJ 411 c. However, we include jitter (ξ) and RV-offset (γ) terms for each instrument. We additionally place priors on several MCMC step parameters (period, amplitude, eccentricity, and jitter) to keep them within physically possible bounds. The final log-likelihood for this model is

$$\mathcal{L} = -\frac{1}{2} \mathbf{r}_\theta^T K_\alpha^{-1} \mathbf{r}_\theta - \frac{1}{2} \ln \det K_\alpha - \frac{N}{2} \ln(2\pi), \quad (1)$$

where \mathbf{r}_θ is the vector of residuals from our mean function, calculated using the three Keplerians, K_α is the kernel function, and N is the number of data points. *celerite2* approximates the kernel function using its power spectral den-

Table 3. Model Parameters and Derived Planetary Properties

Parameter	Value	Units
P_b	$12.9395^{+0.0013}_{-0.0012}$	days
T_{conj_b}	$2456305.31^{+0.28}_{-0.25}$	JD
$\sqrt{e} \cos \omega_b$	$-0.09^{+0.19}_{-0.16}$...
$\sqrt{e} \sin \omega_b$	0.03 ± 0.18	...
K_b	$1.353^{+0.093}_{-0.094}$	m s^{-1}
P_c	2806^{+110}_{-94}	days
T_{conj_c}	2457136^{+79}_{-87}	JD
$\sqrt{e} \cos \omega_c$	$0.04^{+0.2}_{-0.21}$...
$\sqrt{e} \sin \omega_c$	$-0.01^{+0.29}_{-0.27}$...
K_c	1.22 ± 0.15	m s^{-1}
P_{cand}	$215.62^{+0.76}_{-0.73}$	days
$T_{\text{conj}_{\text{cand}}}$	$2456353.3^{+9.9}_{-13.0}$	JD
$\sqrt{e} \cos \omega_{\text{cand}}$	$0.16^{+0.26}_{-0.29}$...
$\sqrt{e} \sin \omega_{\text{cand}}$	$-0.14^{+0.3}_{-0.26}$...
K_{cand}	$0.83^{+0.13}_{-0.12}$	m s^{-1}
ρ	$55.2^{+0.14}_{-0.15}$	days
τ	500,000	days
σ (APF)	103^{+22}_{-20}	m s^{-1}
σ (HIRES)	102^{+19}_{-16}	m s^{-1}
σ (SOPHIE)	$2.7^{+12.0}_{-1.7}$	m s^{-1}
σ (CARMENES)	49 ± 25	m s^{-1}
ξ (APF)	$0.28^{+0.26}_{-0.19}$	m s^{-1}
ξ (HIRES)	1.47 ± 0.11	m s^{-1}
ξ (SOPHIE)	1.63 ± 0.18	m s^{-1}
ξ (CARMENES)	$1.31^{+0.16}_{-0.15}$	m s^{-1}
γ (APF)	-0.57 ± 0.17	m s^{-1}
γ (HIRES)	$-0.92^{+0.2}_{-0.19}$	m s^{-1}
γ (SOPHIE)	0.06 ± 0.18	m s^{-1}
γ (CARMENES)	0.11 ± 0.19	m s^{-1}
e_b	$0.052^{+0.057}_{-0.037}$...
ω_b	$0.7^{+1.9}_{-3.0}$	radians
T_{peri_b}	$2456306.8^{+2.8}_{-4.9}$	JD
$M_b \sin i$	2.64 ± 0.19	M_{\oplus}
a_b	$0.0788^{+0.00056}_{-0.00057}$	AU
e_c	$0.08^{+0.1}_{-0.06}$...
ω_c	$-0.1^{+1.9}_{-1.8}$	radians
T_{peri_c}	2457025^{+1100}_{-960}	JD
$M_c \sin i$	14.2 ± 1.8	M_{\oplus}
a_c	$2.845^{+0.077}_{-0.067}$	AU
e_{cand}	$0.15^{+0.16}_{-0.11}$...
ω_{cand}	$-0.6^{+1.5}_{-1.2}$	radians
$T_{\text{peri}_{\text{cand}}}$	2456306^{+120}_{-36}	JD
$M_{\text{cand}} \sin i$	4.1 ± 0.6	M_{\oplus}
a_{cand}	$0.5141^{+0.0038}_{-0.0039}$	AU

sity (PSD), which is given by

$$S(\omega) = \sqrt{\frac{2}{\pi}} \frac{S_0 \omega_0^4}{(\omega^2 - \omega_0^2)^2 + \omega^2 \omega_0^2 / Q^2}, \quad (2)$$

and can be reparameterized using the simple harmonic oscillator's undamped period ($\rho = 2\pi/\omega_0$), the timescale of dampening ($\tau = 2Q/\omega_0$), and the standard deviation of the process ($\sigma = \sqrt{S_0 \omega_0 Q}$).

Using the No-U-Turns algorithm (NUTS; Hoffman & Gelman 2011) implemented through `PyMC3` (Salvatier et al. 2016), we perform an affine-invariant MCMC process to explore the parameter space. Our analysis used 16 chains and was well-converged after 15,000 draws with a final maximum Gelman-Rubin statistic (Gelman & Rubin 1992) of 1.0005. The resulting posterior is shown in Table 3 and the best fit model in Figure 8.

We also consider models that represent the rotational modulation as a simple Keplerian or completely neglect it. We find that the Keplerian is an inadequate model given that the rotation signal has vastly different amplitudes in the four different data sets (interestingly, the amplitude of the GP, σ , is the same for the APF and HIRES, which roughly cover the same wavelengths). Neglecting the rotation signal altogether returns parameters consistent with those presented in Table 3 with the exception of elevated eccentricities (0.41 ± 0.15 versus 0.15 ± 0.14) and amplitudes ($1.22 \pm 0.17 \text{ m s}^{-1}$ versus $0.83 \pm 0.12 \text{ m s}^{-1}$) for the candidate signal. As discussed in Section 4.3, the rotation signal could possibly be constructively and destructively interfering with the candidate signal, artificially elevating both the eccentricity and amplitude. Therefore, we caution that the uncertainties on e_{cand} and K_{cand} may be under-reported. We then fit a two-Keplerian model only representing the signals from GJ 411 b and c in addition to a GP representing the rotation. The results are consistent with the values given in Table 3, indicating that if the candidate signal is an APF systematic, it does not interfere with the orbital solutions for the confirmed planets.

5. DISCUSSION

5.1. Comparisons to Previous Work

By combining the APF, HIRES, SOPHIE, and CARMENES data, we have a more comprehensive data set than Díaz et al. (2019), Stock et al. (2020), or Rosenthal et al. (2021), allowing our modeling to place tighter constraints on each of the orbital parameters for GJ 411 b and c, as seen in Table 4. For GJ 411 b, our orbital period is formally inconsistent with the models from Díaz et al. (2019) and Stock et al. (2020), however, all other parameters for the planet are consistent. Interestingly, we find eccentricities better aligned with a nearly-circular orbit for this inner-most planet and obtain somewhat lower minimum masses compared to Díaz et al. (2019) and Rosenthal et al. (2021). Our orbital

Table 4. Comparisons to Previous Work

Parameter	Comparison Model	This Work	Units
Díaz et al. (2019)			
$M_b \sin i$	2.99 ± 0.46	2.64 ± 0.19	M_\oplus
a_b	0.0785 ± 0.0027	$0.0788^{+0.00056}_{-0.00057}$	AU
P_b	12.9532 ± 0.0079	$12.9395^{+0.0013}_{-0.0012}$	days
e_b	0.22 ± 0.13	$0.052^{+0.057}_{-0.037}$...
K_b	1.59 ± 0.23	$1.353^{+0.093}_{-0.094}$	m s^{-1}
Stock et al. (2020)			
$M_b \sin i$	2.69 ± 0.25	2.64 ± 0.19	M_\oplus
a_b	$0.07890^{+0.00068}_{-0.00077}$	$0.0788^{+0.00056}_{-0.00057}$	AU
P_b	12.946 ± 0.005	$12.9395^{+0.0013}_{-0.0012}$	days
e_b	$0.12^{+0.12}_{-0.09}$	$0.052^{+0.057}_{-0.037}$...
P_c	2852 ± 568	2806^{+110}_{-94}	days
Rosenthal et al. (2021)			
$M_b \sin i$	$2.803^{+0.292}_{-0.311}$	2.64 ± 0.19	M_\oplus
a_b	$0.07892^{+0.00055}_{-0.00098}$	$0.0788^{+0.00056}_{-0.00057}$	AU
e_b	$0.095^{+0.099}_{-0.066}$	$0.052^{+0.057}_{-0.037}$...
$M_c \sin i$	$18.0^{+2.9}_{-2.6}$	14.2 ± 1.8	M_\oplus
a_c	$3.1^{+0.13}_{-0.11}$	$2.845^{+0.077}_{-0.067}$	AU
P_c	3190^{+200}_{-170}	2806^{+110}_{-94}	days
e_c	$0.14^{+0.160}_{-0.095}$	$0.08^{+0.1}_{-0.06}$...
P_{cand}	$214.59^{+0.99}_{-0.64}$	$215.62^{+0.76}_{-0.73}$	days
K_{cand}	$1.45^{+0.42}_{-0.27}$	$0.83^{+0.13}_{-0.12}$	m s^{-1}

period for planet c is inconsistent with the model presented by Rosenthal et al. (2021), contributing to inconsistencies in the semi-major axis. We find minimum masses and eccentricities that are in agreement, although the central value for the minimum mass in our model is notably lower ($\sim 4 M_\oplus$) than that found by Rosenthal et al. (2021). Even though Rosenthal et al. (2021) dismissed the 215 day signal as a systematic, they provide the respective amplitude and period from their modeling, which are also shown in Table 4. While our adopted model returns a much lower amplitude for the candidate signal, this once again is likely because Rosenthal et al. (2021) neglect stellar rotation.

Butler et al. (2017) identified a planetary candidate with a period of 9.8693 ± 0.0016 days using HIRES radial velocities, which has not been confirmed by any other studies. Díaz et al. (2019) conducted an independent analysis of this HIRES data set but were unable to identify a similar signal. Notably, these two analyses used different models, with Butler et al. (2017) decorrelating the radial velocities using Ca II lines. This discrepancy could cause different signals to be observed. However, Díaz et al. (2019) found a peak located at

the same period in the window function, suggesting that the signal is not astrophysical. While processed through a different pipeline, many of our HIRES radial velocities come from the same spectra—although our analysis excludes any of the pre-2004 data. As seen in Figure 1, there is a stand-alone peak in the HIRES periodogram located before 10 days, although it does not appear to fall near 9.9 days. As emphasized by previous works, the Butler et al. (2017) signal is likely a false detection that can be attributed to systematics or modeling artifacts.

5.2. Prospects for Follow-up

5.2.1. Transit Photometry

Using the analytic relationship between mass and radius given by Weiss & Marcy (2014), we calculate the expected radius distribution for each planet and candidate from our RV model posterior. To account for the mass-inclination degeneracy, we randomly assign orbital inclinations drawn from a uniform cosine distribution to each sample. This yields radii of $1.16^{+0.75}_{-0.17} R_\oplus$, $7.22^{+4.50}_{-1.43} R_\oplus$, and $1.89^{+1.16}_{-0.41} R_\oplus$ for planets b, c, and the candidate, respectively. With the radius posteriors and the equation

$$P_{\text{transit}} = \left(\frac{R_* + R_p}{a} \right) \left(\frac{1 + e \sin \omega}{1 - e^2} \right), \quad (3)$$

we calculate the respective geometric transit probabilities to be $\sim 2\%$, $\sim 0.07\%$, and $\sim 0.3\%$.

We then calculate the fraction of orbits in our posterior with transit times falling between the start and end of TESS sector 22. Multiplying these by the geometric probabilities, we find that the probabilities of GJ 411 b, c, and the candidate transiting during our TESS data are $\sim 2\%$, $< 3 \times 10^{-4}\%$, and $\sim 0.05\%$, respectively.

As seen in Figure 9, there is no visible transit in the PD-CSAP TESS data. However, we search the light curve using

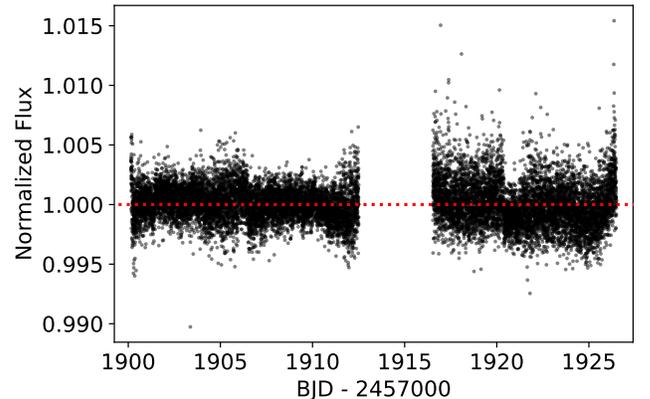


Figure 9. The normalized PDCSAP TESS light curve for GJ 411 during sector 22. The horizontal red dashed line marks a constant normalized flux of 1, from which there is little deviation and no apparent transiting event.

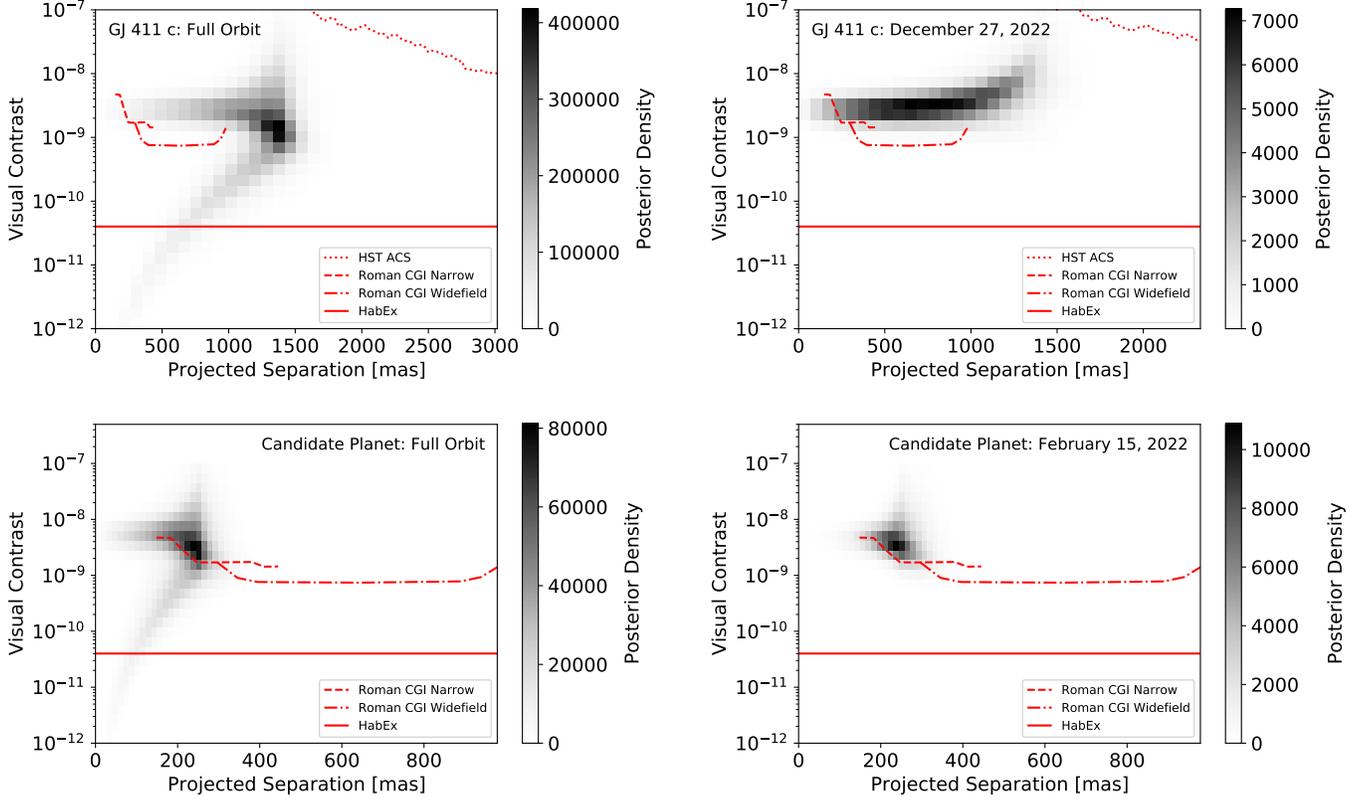


Figure 10. Two-dimensional histograms showing the joint planet contrast and projected separation posteriors for GJ 411 c (top row) and the candidate planet (bottom row). Both are approximated as a Lambertian disk with a Bond albedo of 0.5. Figures on the left show the posterior for the entire orbit while figures on the right show the posterior for a single time where the planet would be most easily observed. GJ 411 c is expected to be most easily observed on December 22, 2023 while the candidate would be best detected on February 15, 2022. The predicted post-processing detectability floor is shown for the Roman Space Telescope CGI in narrow and wide FOV modes, along with the HabEx performance requirement. We also show the detection limits for HST ACS in the upper row. Current instrumentation, such as HST ACS, do not have the sensitivity needed to detect reflected light from GJ 411 c and the candidate planet, but future missions such as Roman or HabEx may be able image the planets in the visual spectrum.

the package `Transit Least Squares` (TLS; Hippke & Heller 2019) and confirm Stock et al. (2020)’s findings that the strongest periodicity is located at 13.0 days but has a negligible signal detection efficiency of 7.1. While no apparent transits have occurred during *TESS* sector 22, during which GJ 411 b would have occulted twice, it is possible that the signal is obscured by noise. Given the expected radius, we anticipate a transit depth of ~ 1 ppt, however the RMS of the light curve is 1.7 ppt.

5.2.2. Direct Imaging

Orbiting a close, relatively bright M dwarf, GJ 411 c and the candidate planet are prime targets for imaging. To explore the prospects for a direct detection, we follow a procedure similar to that outlined by Blunt et al. (2019). We derive the projected separation posterior for the outermost planet and candidate at future epochs throughout their orbits with the expression

$$\Delta\theta = \frac{r}{d} [\cos^2(\omega + f) + \sin^2(\omega + f)\cos^2(i)]^{1/2} \quad (4)$$

(Kane & Gelino 2011). Here, r is the star-planet separation, d is the star-observer distance, ω is the argument of periastron, f is the true anomaly, and i is the orbital inclination randomly assigned in Section 5.2.1.

We then estimate the planets’ expected reflected visual light contrasts, approximating both as Lambertian disks with scattering albedos of 0.5—roughly between the albedos of Earth (~ 0.3) and Venus (~ 0.8). The phase angle, α , is calculated at each epoch for every sample in our posterior using the relation

$$\cos(\alpha) = \sin(f + \omega)\sin(i). \quad (5)$$

With the radii posteriors derived in Section 5.2.1, the expected contrast from the star is evaluated using the equation

$$\frac{F_p}{F_\star} = A_g \left(\frac{R_p}{a}\right)^2 \Phi(\alpha), \quad (6)$$

where the geometric albedo is related to the scattering albedo by

$$A_g = \frac{2}{3}A_s \quad (7)$$

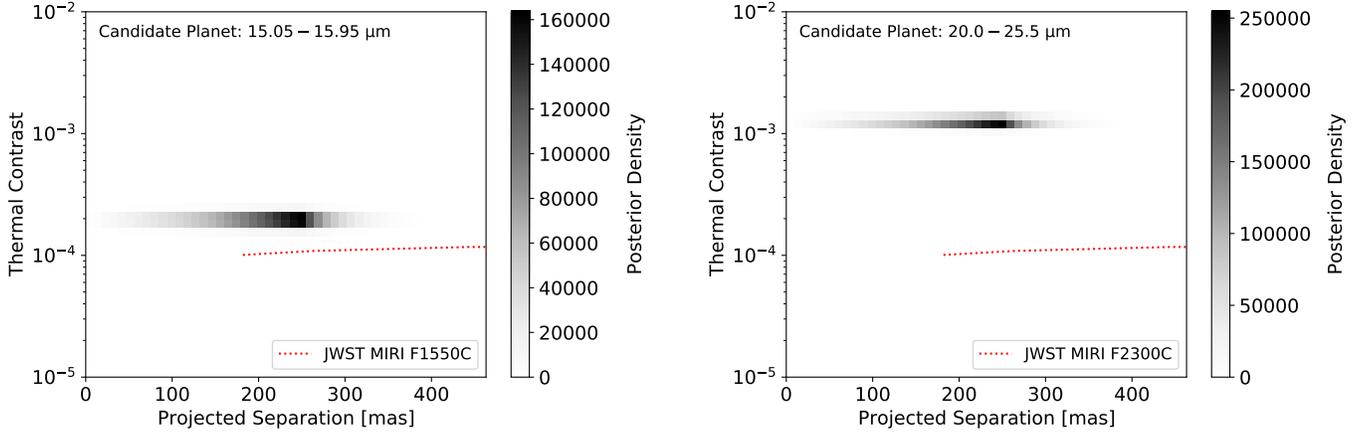


Figure 11. Two-dimensional histograms showing the detectability of the candidate planet in the thermal infrared, assuming that the planet and star are perfect black bodies and the orbit is circular. The left plot shows the contrast through the JWST MIRI F1550C filter while the right shows the contrast through the JWST MIRI F2300C filter. The red lines represent the limiting contrast for the coronagraph (Boccaletti et al. 2015). While the candidate could be too close to its star to image, it may be bright enough for JWST to make a detection.

and

$$\Phi(\alpha) = \frac{\sin(\alpha) + (\pi - \alpha)\cos(\alpha)}{\pi} \quad (8)$$

(Madhusudhan & Burrows 2012). Figure 10 shows two-dimensional histograms comparing the planet-star contrast distribution to the projected separation distribution, along with the expected detection limits for several instruments. While GJ 411 c and the candidate are not detectable by instrumentation such as the HST Advanced Camera for Surveys (ACS), which represents some of the best detection limits currently available, future space missions such as the Nancy Grace Roman Telescope will likely be able to detect both at certain points in their orbits. Rough estimates

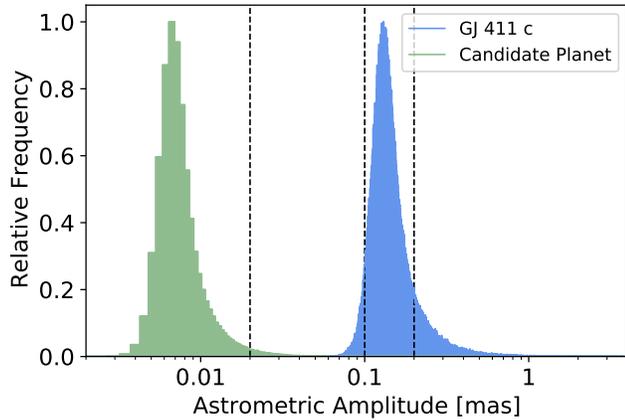


Figure 12. Histograms showing the astrometric amplitude posterior for GJ 411 c (blue) and the candidate (green). Assuming Gaia has uncertainties of 0.02 mas, the vertical dashed lines mark one-sigma, five-sigma, and ten-sigma detection limits. For an orbit to be detectable, we require the astrometric amplitude to be greater than ten-sigma. 97% and 2% of simulated orbits for GJ 411 c and the candidate planet, respectively, are detectable with Gaia.

of HabEx performance indicate that future missions may be able to detect them at most times².

The candidate planet could possibly be warm enough to also be imaged in the infrared. We derive its equilibrium temperature posterior, assuming a Bond albedo of 0.5 and a circular orbit; although the the 215 day signal is somewhat eccentric, variations in the equilibrium temperature due to orbital position are negligible. Assuming both the star and planet behave as ideal black bodies, we take the ratios of their fluxes and integrate over different band-passes. The resulting posteriors are shown for two JWST Mid-Infrared Instrument (MIRI) filters, along with the instrument’s limiting contrast curve (Boccaletti et al. 2015), in Fig. 11. While the orbit may be too small for MIRI to capture, the expected contrast is high enough to detect in the infrared.

5.2.3. Astrometry

Given the star’s proximity, GJ 411 is also a prime target for astrometry. After using the inclinations drawn in Section 5.2.1 to account for the mass-inclination degeneracy, we calculate the expected astrometric amplitude for each planet using the equation

$$A_\theta = \frac{M_p}{M_\star + M_p} \frac{a}{d}. \quad (9)$$

The resulting distributions for GJ 411 c and the candidate planet are shown in Figure 12. Assuming that each planet’s signal is distinguishable from proper motion, we consider an orbit to be detectable if its amplitude is five times greater than the expected uncertainty. Gaia Collaboration et al. (2020) gave a typical uncertainty of 0.02 mas during a sin-

² Anticipated detection limits obtained from github.com/nasavbailey/DI-flux-ratio-plot

gle epoch for stars with magnitudes below 15. By this standard, 97% of GJ 411 c’s and 2% of the candidate’s simulated orbits are detectable. While the ~ 7.5 magnitude star would likely have greater uncertainties, multiple measurements would contribute to a more significant detection.

Gaia Data Releases 1 (Gaia Collaboration et al. 2016a,b) and 2 (Gaia Collaboration et al. 2018) did not contain any observations of Gliese 411. The significance of the excess noise can indicate whether an astrometric solution behaves well, with values greater than 2 indicating a poor fit. In Early Data Release 3 (Gaia Collaboration et al. 2020), this value for GJ 411 was 36.9, with the caveat that modeling errors may inflate the significance of the excess noise for an early data release. In the case that the astrometric model poorly fits Gaia data from GJ 411, this would suggest the presence of detectable astrometric signals.

5.3. Comparative Exoplanetology

Only 2.55 pc away from the Solar System, GJ 411 is a remarkably nearby star, with only the Alpha Centauri system, Barnard’s Star, and Wolf 359 closer. While each of these stars or systems has at least one planet candidate (Anglada-Escudé et al. 2016; Ribas et al. 2018; Tuomi et al. 2019; Damasso et al. 2020), GJ 411 could potentially be the closest star with three or more planets.

Assuming circular orbits and Bond albedos of 0.5, the equilibrium temperatures for GJ 411 b, c, and the candidate—derived from the RV model posterior—are ~ 326 K, ~ 54 K, and ~ 128 K, respectively. Following the methods in Kopparapu et al. (2013), the habitable zone extends from 0.163 ± 0.003 AU to 0.308 ± 0.007 AU. Planet b and the candidate both straddle this region but never enter it throughout their full orbits.

While the majority of planets detected through imaging have masses on the scale of thousands of M_{\oplus} , recent observations of Proxima c using SPHERE at the Very Large Telescope (Gratton et al. 2020) show that the direct detection of small, nearby planets—particularly those around bright stars—may not be far out of reach. Even though the Roman telescope and potential missions like HabEx promise to image a smaller class of planets than previously possible, GJ 411 c and the candidate planet would still be unusually small and close to their star for a direct image. Consequently, the Neptune and potential super-Earth could be among the first imaged planets of their kind. Figure 13 illustrates how different both are from objects that are conventionally imaged.

It is also worth noting that GJ 411 does not show large RV variations and each signal of interest has a low amplitude (< 1.5 m s $^{-1}$). At the time of its discovery, GJ 411 b was the lowest-amplitude signal detected by SOPHIE, however, the instrument has since aided in the confirmation of HD 158259 b, which has an amplitude of ~ 1.05 m s $^{-1}$ (although the ac-

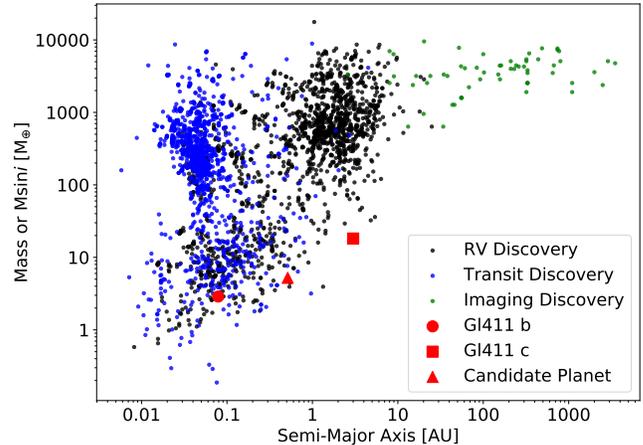


Figure 13. Mass or $M \sin i$ versus semi-major axis for planets in the NASA Exoplanet Archive (Retrieved 2021 May 2) and the GJ 411 system. GJ 411 b, c, and the candidate have relatively low minimum masses compared to other known planets at similar orbital distances. Additionally, despite GJ 411 c and the candidate possibly being imageable, they have much smaller semi-major axes and masses than planets that have been discovered using direct imaging.

tual discovery was made using *TESS* data; Hara et al. 2020). The candidate signal at 215 days has an incredibly low amplitude of 0.83 ± 0.12 m s $^{-1}$, although this value could be somewhat larger due to limitations in our modeling. While HIRES data was used in the discovery of the τ Ceti system, which does have lower amplitudes (Feng et al. 2017), this signal, if confirmed, could be one of the most sensitive detections made primarily using HIRES and APF data. Regardless, GJ 411 c is the next lowest-amplitude signal (1.22 ± 0.15) discovered with these instruments and is followed by HD 7924 d (~ 1.65 m s $^{-1}$; Fulton et al. 2015).

6. CONCLUSIONS

We perform a joint analysis of APF, HIRES, SOPHIE, and CARMENES radial velocities from the bright, nearby M dwarf Gliese 411, re-detecting 12.9, 2900, 215, and 55 day periodicities. The first of these signals corresponds to the previously detected GJ 411 b. The 2900 day signal has been contested by previous works as either a planet or long-period stellar magnetic cycles. We analyze stellar activity metrics including $\log R'_{\text{HK}}$ and H_{α} indices to conclude that this disputed signal corresponds to a planet, GJ 411 c. The 215 day signal was originally identified as an APF systematic, not showing in periodograms for other instruments. However, injection recovery tests show that the signal could have been obscured by cadence and noise in the other data sets, accounting for the non-detections. Further, the signal becomes more significant when combining RVs from all four sources. In an analysis of the PSF parameters used to fit the APF spectra, no signs of a 215 day systematic are found. While we cannot conclusively determine the cause of this signal, we consider

the planetary interpretation to be most probable. Lastly, the 55 day periodicity most likely corresponds to rotational modulation detected in previous photometric work.

We model the three planetary and candidate signals with Keplerians and represent the stellar rotation with a simplified periodic Gaussian process. Additional fits neglecting the rotational modulation do not significantly alter orbital parameters except for the candidate amplitude and eccentricity. The rotation signal, located near a harmonic of 215 days, likely interferes with the candidate signal, accounting for these discrepancies. Additionally, two-Keplerian models representing GJ 411 b and c are consistent with the three-Keplerian model, verifying that our orbital solutions are accurate even in the case that the 215 day signal does not correspond to a planet. Our adopted model generally agrees with previous models for GJ 411 b and c while providing tighter constraints on their masses and orbital radii. The candidate signal would correspond to a super-Earth falling between the other planets with a minimum mass of $4.1 \pm 0.6 M_{\oplus}$ and a semi-major axis of 0.5141 ± 0.0038 AU.

With two confirmed planets and potentially three or more orbital companions, GJ 411 is one of the closest multi-planet systems to ours. While a secondary detection through transit photometry is not promising, direct imaging with missions just as JWST, the Roman Telescope, or HabEx offer the chance to make a dual detection of GJ 411 c and to confirm the candidate planet. Additionally, astrometry through Gaia could potentially make a secondary detection of GJ 411 c. While similar observations are not possible for GJ 411 b—the inner-most planet—in the near future, learning more about the architecture of the outer system is important towards understanding its formation and evolution. Most importantly, this system could offer the opportunity to directly probe lower-mass planets.

ACKNOWLEDGMENTS

This research has made use of the NASA Exoplanet Archive, which is operated by the California Institute of Technology, under contract with the National Aeronautics and Space Administration under the Exoplanet Exploration Program. This work made use of the SIMBAD database (operated at CDS, Strasbourg, France), NASA’s Astrophysics Data System Bibliographic Services. This research has made use of the VizieR catalogue access tool, CDS, Strasbourg, France (DOI: 10.26093/cds/vizieR). The original description of the VizieR service was published in A&AS 143, 23. This work has made use of data from the European Space Agency (ESA) mission *Gaia* (<https://www.cosmos.esa.int/gaia>), processed by the *Gaia* Data Processing and Analysis Consortium (DPAC, <https://www.cosmos.esa.int/web/gaia/dpac/consortium>). Funding for the DPAC has been provided by national institutions, in particular the institutions participating in the *Gaia* Multilateral Agreement. L.M.W. is supported by the Beatrice Watson Parrent Fellowship and NASA ADAP Grant 80NSSC19K0597. We are very grateful for the donations of the Levy family that helped facilitate the construction of the Levy spectrograph on APF. Without their support the APF would not be contributing to the discovery of planets like these. Research at the Lick Observatory is partially supported by a generous gift from Google. We are grateful to the time assignment committees of the University of Hawaii, the University of California, and NASA for their generous allocations of observing time. Without their long-term commitment to RV monitoring, these planets would likely remain unknown. We acknowledge R. Paul Butler and S. S. Vogt for many years of contributing to the data presented here. Finally, the authors wish to extend special thanks to those of Hawaiian ancestry on whose sacred mountain of Maunakea we are privileged to be guests. Without their generous hospitality, the Keck observations presented herein would not have been possible.

Software: `astropy` (Astropy Collaboration et al. 2018), `emcee` (Foreman-Mackey et al. 2013), `exoplanet` (Foreman-Mackey et al. 2021b), `matplotlib` (Hunter 2007), `numpy` (Harris et al. 2020), `RadVel` (Fulton et al. 2018), `RVSearch` (Rosenthal et al. 2021)

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