

First tidal disruption events discovered by *SRG/eROSITA*: X-ray/optical properties and X-ray luminosity function at $z < 0.6$

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ABSTRACT

We present the first sample of tidal disruption events (TDEs) discovered during the *SRG* all-sky survey. These 13 events were selected among X-ray transients detected in the $0^\circ < l < 180^\circ$ hemisphere by eROSITA during its second sky survey (2020 June 10 to December 14) and confirmed by optical follow-up observations. The most distant event occurred at $z = 0.581$. One TDE continued to brighten at least 6 months. The X-ray spectra are consistent with nearly critical accretion on to black holes of a few $\times 10^3$ to $10^8 M_\odot$, although supercritical accretion is possibly taking place. In two TDEs, a spectral hardening is observed 6 months after the discovery. Four TDEs showed an optical brightening apart from the X-ray outburst. The other nine TDEs demonstrate no optical activity. All 13 TDEs are optically faint, with $L_g/L_X < 0.3$ (L_g and L_X being the g band and 0.2–6 keV luminosity, respectively). We have constructed a TDE X-ray luminosity function, which can be fit by a power law with a slope of -0.6 ± 0.2 , similar to the trend observed for optically selected TDEs. The total rate is estimated at $(1.1 \pm 0.5) \times 10^{-5}$ TDEs per galaxy per year, an order of magnitude lower than inferred from optical studies. This suggests that X-ray bright events constitute a minority of TDEs, consistent with models predicting that X-rays can only be observed from directions close to the axis of a thick accretion disc formed from the stellar debris. Our TDE detection threshold can be lowered by a factor of ~ 2 , which should allow a detection of ~ 700 TDEs by the end of the *SRG* survey.

Key words: accretion, accretion discs – black hole physics – quasars: supermassive black holes – X-rays: galaxies – transients: tidal disruption events.

1 INTRODUCTION

Stellar disruptions by the gravitation of supermassive black holes (SMBHs) – tidal disruption events (TDEs) – provide valuable and largely unique information on relatively small ($M_{\text{BH}} \lesssim 10^8 M_\odot$) SMBHs in, usually, dormant galactic nuclei and allow us to explore various regimes of accretion on to black holes. Predicted by theorists (Hills 1975; Lidskii & Ozernoi 1979; Gurzadian & Ozernoi 1981; Rees 1988), TDEs were first discovered as soft X-ray transients (Komossa & Bade 1999) by the *ROSAT* satellite during its all-sky

survey in 1990–1991. Soft X-ray emission is one of two expected distinctive features (together with a $t^{-5/3}$ flux decline) of TDEs, since the thermal emission of an accretion disc that forms around the SMBH from the debris of the disrupted star is expected to have a characteristic temperature of $\sim 10^6$ K. Accordingly, X-ray searches remained the main channel of discovering TDEs until recently (e.g. Donley et al. 2002; Esquej et al. 2008; Maksym, Ulmer & Eracleous 2010; Khabibullin & Sazonov 2014), and the current list of TDEs discovered in X-rays comprises some 20 events (see Saxton et al. 2021 for a recent review).

In the last ~ 15 yr, a new influx of TDE discoveries has emerged from optical/ultraviolet (UV) surveys, which have provided an additional sample of TDEs that is comparable in size to the X-

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ray based one (e.g. van Velzen et al. 2011; Gezari et al. 2012; Blagorodnova et al. 2019; Holoien et al. 2019; van Velzen et al. 2020, 2021b). Interestingly, these optical/UV selected objects seem to have, on average, quite different properties compared to the X-ray selected ones. Specifically, the spectral energy distribution (SED) of the former can be described to a first approximation as thermal emission with a temperature \sim a few 10^4 K, which is at least an order of magnitude lower than expected for near-Eddington standard accretion discs (Shakura & Sunyaev 1973) around moderately massive black holes.

This apparent dichotomy in TDE SEDs resembles the distinction between type 1 and type 2 active galactic nuclei (AGNs), which has led to suggestions that we might be dealing with a similar orientation-driven effect. In particular, it was proposed (Dai et al. 2018; Curd & Narayan 2019) that X-ray rich TDEs are observed from directions close to the axis of a thick accretion disc with a powerful wind, whereas X-ray weak ones are viewed from larger inclination angles. In the latter case, the central X-ray source is obscured by the disc so that we can only see the reprocessed optical/UV emission. However, in addition to the viewing direction, other factors, in particular the black hole mass, are also likely to strongly affect TDE SEDs and light curves (Mummery 2021). In other proposed models, the conversion of the accretion disc emission to optical wavelengths occurs in the unbound part of the stellar debris (Metzger & Stone 2016; Lu & Bonnerot 2020), or the optical emission is directly powered by energy liberated during the formation of the accretion disc rather than by energy released during subsequent accretion on to the black hole (Piran et al. 2015).

It is clear that more observational data in various energy bands are needed for a better understanding of the physics of the TDE phenomenon. Therefore, the launch of the *eROSITA* telescope (Predehl et al. 2021) onboard the *SRG* observatory (Sunyaev et al. 2021) was eagerly awaited, because it was expected to find hundreds to thousands of TDEs during its revolutionary all-sky X-ray survey (Khabibullin, Sazonov & Sunyaev 2014). On 2019 July 13, *SRG* was successfully launched from the Baikonur Cosmodrome and on 2019 December 12 it started its all-sky X-ray survey from a halo orbit around the Sun–Earth L2 point. The survey is to consist of eight consecutive full scans of the sky, each lasting 6 months. Already during the first weeks of the survey a few TDE candidates were found in the *eROSITA* data through comparison with archival observations by previous X-ray missions (Khabibullin et al. 2020a, b). On 2020 June 10, the second *eROSITA* all-sky survey began, which allowed us to begin a regular search for TDEs over the sky. Specifically, TDE candidates are sought among the multitude of transient X-ray sources detected in a given *eROSITA* all-sky survey and undetected in the preceding survey. The second all-sky survey was completed on 2020 December 14, and the third survey has also been finished by now.

In this paper, we present an initial sample of 13 relatively bright TDEs discovered by *SRG/eROSITA* during its second all-sky survey at $0^\circ < l < 180^\circ$.¹ The TDE nature of these transients was suggested by their X-ray properties and then confirmed by our follow-up optical observations. Below we discuss the X-ray and optical properties of this TDE sample and use it to draw inferences on the statistical properties of TDEs in the $z < 0.6$ Universe.

In what follows we adopt a flat Λ cold dark matter cosmological model with $h = 0.70$ and $\Omega_\Lambda = 0.7$.

2 SELECTION OF TDE CANDIDATES

The initial sample of X-ray transients for this pilot study consisted of sources that were undetected in the first *eROSITA* all-sky survey (hereafter eRASS1) but were detected during the second scan (hereafter eRASS2) at a flux level exceeding at least tenfold the upper limit at a likelihood of 6 ($\approx 3\sigma$) on their flux in the 0.3–2.2 keV energy band during eRASS1. The search for transients spanned the entire period of eRASS2 (2020 June 10 to December 14). By the time of this writing, all transients found in eRASS2 have been scanned by *eROSITA* for the third time during eRASS3.

Among these X-ray transients, we filtered out objects of likely Galactic origin based on positional coincidence of the X-ray source with a star having a statistically significant ($>5\sigma$) parallax and/or proper motion in the *Gaia* astrometric catalogue (Gaia Collaboration 2016, 2021). For the majority of the remaining, potentially extragalactic, objects, a likely counterpart was readily found in archival optical and/or infrared (IR) images within the *eROSITA* localization region (of ~ 5 arcsec radius). Some of these counterparts are known AGNs or have signatures suggestive of an AGN origin, namely, archival detections by previous X-ray missions, a $W1 - W2 > 0.5$ colour (Assef et al. 2013) in the *WISE* IR all-sky survey (Wright et al. 2010), or irregular optical variability in the Zwicky Transient Facility (ZTF; Bellm et al. 2019; Graham et al. 2019; Masci et al. 2019) data many months before the *eROSITA* X-ray discovery. In addition, we analysed the *eROSITA* data and found that a number of transients had very soft X-ray spectra, suggesting a TDE origin.

For a final identification, we have been following up the presumably extragalactic *eROSITA* transients detected in eRASS2, apart from the indisputable AGN, with optical spectroscopy (Section 4). This campaign has now come to an end and allowed us to discriminate TDEs from AGNs. The new AGN discovered during this campaign will be discussed elsewhere.

As a result, we have obtained a sample of firmly established TDEs detected at $0^\circ < l < 180^\circ$ during eRASS2. The limiting X-ray flux of our sample varies across the sky depending on the sensitivity achieved at a given location in the first *SRG* sky survey. In this paper, we discuss the 13 events from this catalogue that have revealed themselves as bright transients in X-rays only (although a few of them also showed moderate activity in the optical, as will be discussed in Section 4). In a companion paper (Gilfanov et al., in preparation), we discuss the subset of *eROSITA* TDEs with prominent optical transient counterparts.

3 THE SAMPLE AND ITS X-RAY PROPERTIES

Table 1 presents our TDE sample. Specifically, the following information is provided: (1) the source name in the *eROSITA* catalogue (‘SRGE’ followed by the equatorial coordinates of the source measured in eRASS2), (2) the radius of the localization region (at the 98 per cent confidence level) in eRASS2, (3) the dates when the source was scanned during eRASS1, (4) the upper limit (3σ) on the source flux (0.3–2.2 keV) during this period, (5) the dates when the source was scanned during eRASS2, (6) the average X-ray flux during this period, (7) the dates when the source was scanned during eRASS3, and (8) the average X-ray flux or 3σ upper limit for this period. The fluxes and upper limits were determined from the measured count rates using the results of our X-ray spectral analysis, discussed in Section 3.3, namely the best-fitting parameters of eRASS2 and eRASS3 spectra by a multiblackbody accretion disc emission model if available or the corresponding eRASS2 spectral parameters in the remaining cases.

¹ Analysis of *eROSITA* data in this half of the sky is performed by the Russian *eROSITA* consortium.

Table 1. The first *SRG* TDE sample.

No.	X-ray source	R_{98}^a	eRASS1:		eRASS2:		eRASS3:	
			Dates	X-ray flux ^b	Dates	X-ray flux	Dates	X-ray flux
1	SRGE J135514.8 + 311605	4.9	2019 Dec 22–24	$<2.7 \times 10^{-14}$	2020 Jun 23–24	3.8×10^{-13}	2020 Dec 24–25	$<2.0 \times 10^{-14}$
2	SRGE J013204.6 + 122236	4.8	2020 Jan 08–09	$<2.3 \times 10^{-14}$	2020 Jul 08–09	2.8×10^{-13}	2021 Jan 07–08	$<2.5 \times 10^{-14}$
3	SRGE J153503.4 + 455056	3.8	2020 Jan 11–14	$<1.1 \times 10^{-14}$	2020 Jul 13–16	3.2×10^{-13}	2021 Jan 09–11	1.6×10^{-13}
4	SRGE J163831.7 + 534020	4.5	2020 Jan 30–04	$<6.7 \times 10^{-15}$	2020 Aug 01–08	1.2×10^{-13}	2021 Jan 22–26	$<8.4 \times 10^{-15}$
5	SRGE J163030.2 + 470125	4.0	2020 Feb 08–16	$<1.4 \times 10^{-14}$	2020 Aug 06–14	3.7×10^{-13}	2021 Jan 29–31	$<1.7 \times 10^{-14}$
6	SRGE J021939.9 + 361819	5.5	2020 Feb 05–07	$<1.8 \times 10^{-14}$	2020 Aug 07–09	2.4×10^{-13}	2021 Jan 27–28	1.5×10^{-13}
7	SRGE J161001.2 + 330121	5.1	2020 Feb 12–18	$<1.2 \times 10^{-14}$	2020 Aug 14–16	1.8×10^{-13}	2021 Feb 01–02	$<2.3 \times 10^{-14}$
8	SRGE J171423.6 + 085236	3.5	2020 Mar 18–19	$<3.0 \times 10^{-14}$	2020 Sep 18–19	1.2×10^{-12}	2021 Mar 14–16	7.7×10^{-14}
9	SRGE J071310.6 + 725627	4.0	2020 Apr 07–08	$<2.8 \times 10^{-14}$	2020 Oct 11–12	1.4×10^{-12}	2021 Apr 06–07	2.0×10^{-13}
10	SRGE J095928.6 + 643023	4.8	2020 Apr 25–26	$<2.0 \times 10^{-14}$	2020 Oct 28–29	3.9×10^{-13}	2021 Apr 28–30	3.0×10^{-13}
11	SRGE J091747.6 + 524821	5.9	2020 Apr 26–27	$<4.5 \times 10^{-14}$	2020 Oct 28–29	4.6×10^{-13}	2021 Apr 29–30	$<4.5 \times 10^{-14}$
12	SRGE J133053.3 + 734824	5.3	2020 May 01–03	$<1.9 \times 10^{-14}$	2020 Nov 02–04	2.8×10^{-13}	2021 May 05–07	$<1.9 \times 10^{-14}$
13	SRGE J144738.4 + 671821	5.4	2020 May 20–22	$<2.4 \times 10^{-14}$	2020 Nov 19–21	2.6×10^{-13}	2021 May 20–22	6.3×10^{-13}

Notes. ^aRadius of the 98 per cent localization region in eRASS2 in units of arcsec.

^bObserved fluxes or upper limits (3σ) in the 0.3–2.2 keV energy band in units of $\text{erg s}^{-1} \text{cm}^{-2}$.

Weak, marginally significant ($\approx 3.4\sigma$) X-ray emission at the position of SRGE J135514.8+311605 was registered by eROSITA during eRASS1. This emission may be associated with a quasar (which was revealed by our follow-up optical spectroscopy), J135514.83+311612.7, that is located 8 arcsec away, i.e. close to the eROSITA localization region of SRGE J135514.8+311605. In computing the upper limit on the eRASS1 flux of this TDE, we took into account all counts registered near this source.

3.1 Data reduction

eROSITA raw data were processed by the calibration pipeline developed at IKI based on the eROSITA Science Analysis Software System (ESASS,² Brunner et al. 2021) and using early in-flight calibration data. Source spectra and light curves were extracted using a circular aperture of 60 arcsec (corresponding to ≈ 90 per cent encircled energy) centred on the best-fitting X-ray source position. An annulus with the inner and outer radii of 150 and 450 arcsec around the source position was used for background extraction. Any faint sources detected in the background extraction area were masked out using a 15 arcsec circular mask.

We took advantage of the soft X-ray response of eROSITA to perform an X-ray spectral analysis of the TDEs in the 0.2–6 keV energy band. To this end, we used the five (out of seven) telescope modules that are equipped with the on-chip filter. Source detection and construction of X-ray light curves was done in the 0.3–2.2 keV energy band, using the data from all operational telescope modules.

The spectral analysis was done using XSPEC, version 12.11.1 (Arnaud 1996). The quality of spectral fits was assessed using χ^2 -statistic for the source and background spectra with Poisson statistics. To avoid the well-known bias in a profile likelihood, we bin the source spectrum so that every bin in the corresponding background spectrum contains at least five counts. The binning was done using the ftgrouppha tool of the HEASOFT package (v. 6.28). Errors for the best-fitting parameters are quoted at the 90 per cent confidence level.

²<https://erosita.mpe.mpg.de/edr/DataAnalysis/>

3.2 X-ray light curves

Fig. 1 shows the X-ray (0.3–2.2 keV) light curves of the TDEs based on eROSITA data. Specifically, these light curves consist of three flux measurements or upper limits (see Table 1) taken at 6-month intervals during eRASS1, eRASS2, and eRASS3. The fluxes and upper limits for each source were determined from the measured count rates using the spectral parameters derived from our X-ray spectral analysis (see Section 3.3). In addition, we present in Appendix A the short-term X-ray light curves of the TDEs obtained during their visits by eROSITA in eRASS2 and eRASS3, which last between ~ 1 and ~ 7 d. None of the transients have demonstrated substantial variability on these short time-scales.

12 out of the 13 TDEs have faded during the 6-month interval between their passages by eROSITA in eRASS2 and eRASS3, with 7 of them not detected anymore in the third scan. The amplitude of the flux drop varies between a factor of ~ 1.3 and $\gtrsim 20$. However, SRGE J144738.4+671821 has instead become brighter by a factor of ~ 2 in eRASS3 compared to eRASS2.

Assuming that eROSITA caught all these events at a canonical $t^{-5/3}$ TDE decay phase (Evans & Kochanek 1989), we can use the eRASS2 and eRASS3 flux measurements/upper limits to evaluate the time of onset of a given TDE relative to the date of its discovery during eRASS2, $t_{\text{det}} - t_0$. Specifically, we assume that the TDE X-ray light curve consists of an initial rise/peak (possibly supercritical) phase with duration τ and flux F_{max} and a subsequent phase when the flux (0.3–2.2 keV) decreases as

$$F(t) = F_{\text{max}} \left(\frac{t_0 - t}{\tau} \right)^{-5/3}. \quad (1)$$

The resulting estimates are presented in Table 2. The maximum flux (F_{max}) depends on the unknown τ , and we thus do not try to estimate it.

For the majority of the TDEs, $t_{\text{det}} - t_0 < 6$ months, and, moreover, $t_{\text{det}} - t_0 < 2$ months. This implies that these events started somewhere between their first and second visits by *SRG*, and in fact shortly before the eRASS2 observations. This is consistent with the non-detection of these transients in eRASS1. However, $t_{\text{det}} - t_0 > 6$ months for three events. This implies that their X-ray temporal behaviour cannot be

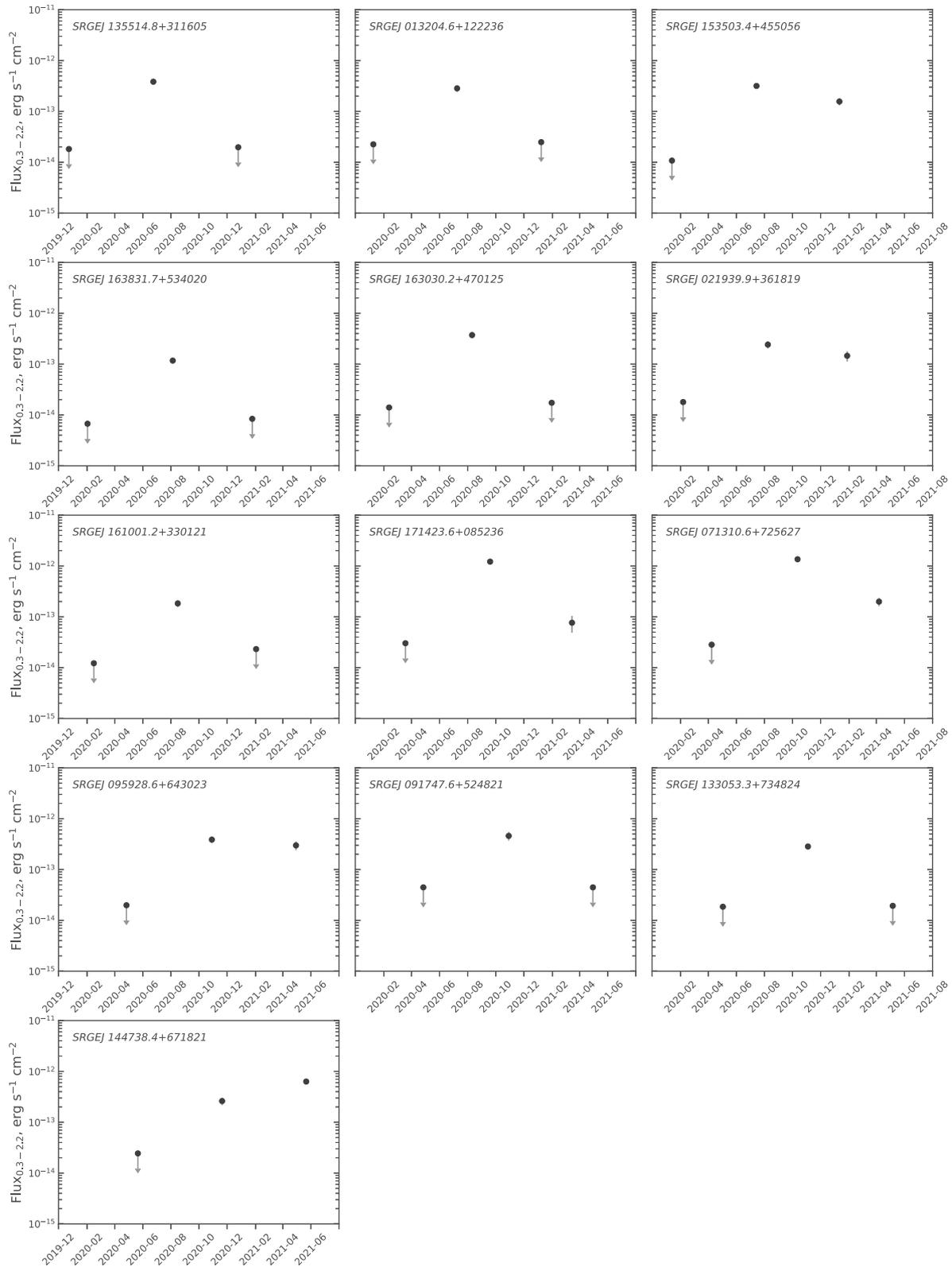


Figure 1. X-ray light curves of the TDEs obtained by eROSITA in the 0.3–2.2 keV energy range.

described by a $t^{-5/3}$ law and that eROSITA probably caught these TDEs during eRASS2 at a rise/peak phase. An even more extreme case is the already mentioned SRGEJ144738.4+671821, which has brightened between eRASS2 and eRASS3.

We stress that the above X-ray variability analysis is currently based on just two flux measurements/upper limits per TDE. As five more scans of the entire sky are planned during the SRG mission, it should be possible to add a few more data points to the X-ray light

Table 2. Time of TDE onset, assuming a $t^{-5/3}$ law.

Object	$t_{\text{det}} - t_0$, month
SRGE J135514.8+311605	<1.7
SRGE J013204.6+122236	<1.9
SRGE J153503.4+455056	11.4
SRGE J163831.7+534020	<1.1
SRGE J163030.2+470125	<1.3
SRGE J021939.9+361819	16.9
SRGE J161001.2+330121	<1.7
SRGE J171423.6+085236	5.6
SRGE J071310.6+725627	2.8
SRGE J095928.6+643023	35.3
SRGE J091747.6+524821	<2.2
SRGE J133053.3+734824	<1.5
SRGE J144738.4+671821	Undefined

curves of at least some of these TDEs and draw firmer conclusions about their long-term behaviour. Dedicated follow-up observations by other X-ray observatories can also be helpful in studying this unique TDE sample.

The above discussion was based on the assumption that the observed luminosity in the 0.3–2.2 keV band (which corresponds to a factor of $(1+z)$ harder X-ray band in the TDE rest frame) is proportional to the accretion rate, \dot{M} . In reality, when the accretion rate drops below the critical Eddington rate, the maximum temperature of the accretion disc (Shakura & Sunyaev 1973) is expected to decline as $\dot{M}^{-1/4}$. As a result, eROSITA should at some point start probing the Wien tail of the disc thermal emission, with the X-ray flux declining exponentially (Lodato & Rossi 2011). This might affect some of the eROSITA light curves presented here.

We also note that optically selected TDEs often show dramatic soft X-ray variability in contrast to the smooth $t^{-5/3}$ power-law decline of the UV/optical light curve (e.g. van Velzen et al. 2021b). In particular, TDEs ASASSN-15oi (Gezari, Cenko & Arcavi 2017) and AT 2019azh/ASASSN-19dj (Liu et al. 2019; Hinkle et al. 2021; van Velzen et al. 2021b) showed a prominent (by a factor of ~ 20 and ~ 100 , respectively) brightening in X-rays over a period of ~ 250 d since the peak of the initial optical-UV flare, which might be attributed to a delayed formation of an accretion disc. Something similar might have been occurring in SRGE J144738.4+671821, for which we see an X-ray brightening on a ~ 180 d time-scale, although there is no evidence of prominent optical activity for this TDE (see Section 4).

3.3 X-ray spectra

We analysed the X-ray spectra of the TDEs in their ‘bright’ phase, using the eROSITA data obtained during eRASS2 (Fig. 2). Four transients after their discovery in eRASS2 remained sufficiently bright (at least 30 detector counts) 6 months later (during eRASS3) to allow us to analyse their spectra obtained at this late phase.

3.3.1 Simple models

We first tried to describe the spectra by two alternative simple models, modified by Galactic and intrinsic absorption: (i) power law (TBABS*ZPHABS*ZPOWERLW) and (ii) multiblackbody accretion disc emission (TBABS*ZPHABS*ZASHIFT*DISKBB). The Galactic absorption was adopted from the HI4PI survey (HI4PI Collaboration

2016), while the redshifts of the TDEs have been measured during our optical spectroscopy program (see Section 4).

Most of the studied spectra can be described similarly well by the power-law and accretion disc models, with some intrinsic absorption ($N_{\text{H}} \sim 10^{21} \text{ cm}^{-2}$) required for the former. Intrinsic absorption does not improve the quality of approximation of the spectra by DISKBB, i.e. N_{H} is consistent with zero for this model. We thus omitted intrinsic absorption from subsequent consideration. The best-fitting parameters of the two models applied to the eRASS2 spectra are given in Table 3.

The inferred temperatures at the inner boundary of the accretion disc vary between $kT_{\text{in}} \approx 0.05$ keV for SRGE J091747.6+524821 and $T_{\text{in}} \approx 0.5$ keV for SRGE J171423.6+085236. The significantly better fit quality provided by the power-law model compared to DISKBB for SRGE J163030.2+470125 suggests the presence of a harder component in addition to the soft thermal emission in the spectrum of this object, which is possibly associated with Comptonized emission from a hot corona of the accretion disc.

3.3.2 Black hole masses and Eddington ratios

Under the assumption that the X-ray emission in TDEs is produced in a standard accretion disc (Shakura & Sunyaev 1973), the DISKBB model (Makishima et al. 1986) is more physically motivated than POWERLAW. However, the real situation is likely more complicated. First, the shape of the spectrum emergent from a standard accretion disc can significantly deviate from a sum of blackbodies (e.g. Koratkar & Blaes 1999; Davis et al. 2005), assumed in DISKBB. Second, at least some of the TDEs discussed here may have been caught by eROSITA in their early super-Eddington phase, when the accretion disc is expected to be geometrically thick, with the radial distribution of X-ray surface brightness being different than in the case of a thin disc (Shakura & Sunyaev 1973; Abramowicz et al. 1988; Watarai et al. 2000). Finally, general relativity can significantly affect the properties of emission from the inner accretion disc (e.g. Wen et al. 2021). On the other hand, at a later stage in the evolution of a TDE, the X-ray spectrum might experience a transition from a purely thermal state to a combination of thermal and Comptonized emission, as suggested by our knowledge of X-ray binary systems and AGNs and has been observed in a number of TDEs (e.g. Komossa et al. 2004; Jonker et al. 2020; Wevers 2020; Wevers et al. 2021).

We thus next tried to describe the eROSITA spectra by the OPTXAGNF model (Done et al. 2012), designed for AGN SEDs. In its simplest version, it represents the spectrum of emission from a standard (i.e. geometrically thin and optically thick) accretion disc around an SMBH, taking into account the expected deviations from a simple multitemperature blackbody shape due to the incomplete thermalization of the radiation in the disc. The main parameters of the model are black hole mass (M_{BH}), spin (a_*), and accretion rate in terms of the Eddington critical luminosity (λ_{Edd}). Essentially, this model is a modification of DISKBB, with the blackbody spectrum of each annulus of the disc modified by a temperature-dependent hardening factor ~ 2 . In addition, it takes into account the radiative efficiency of the accretion disc as a function of spin (Novikov & Thorne 1973). As a result, OPTXAGNF retains the key property of the DISKBB model that, for a given a_* , the peak energy of the spectrum is proportional to $M_{\text{BH}}^{-1/4} \lambda_{\text{Edd}}^{1/4}$ (Shakura & Sunyaev 1973), while the observed bolometric flux is of course proportional to $M_{\text{BH}} \lambda_{\text{Edd}}$ (divided by the distance squared). Therefore, by fitting a measured spectrum by OPTXAGNF it is possible to infer both M_{BH} and λ_{Edd} for an assumed a_* .

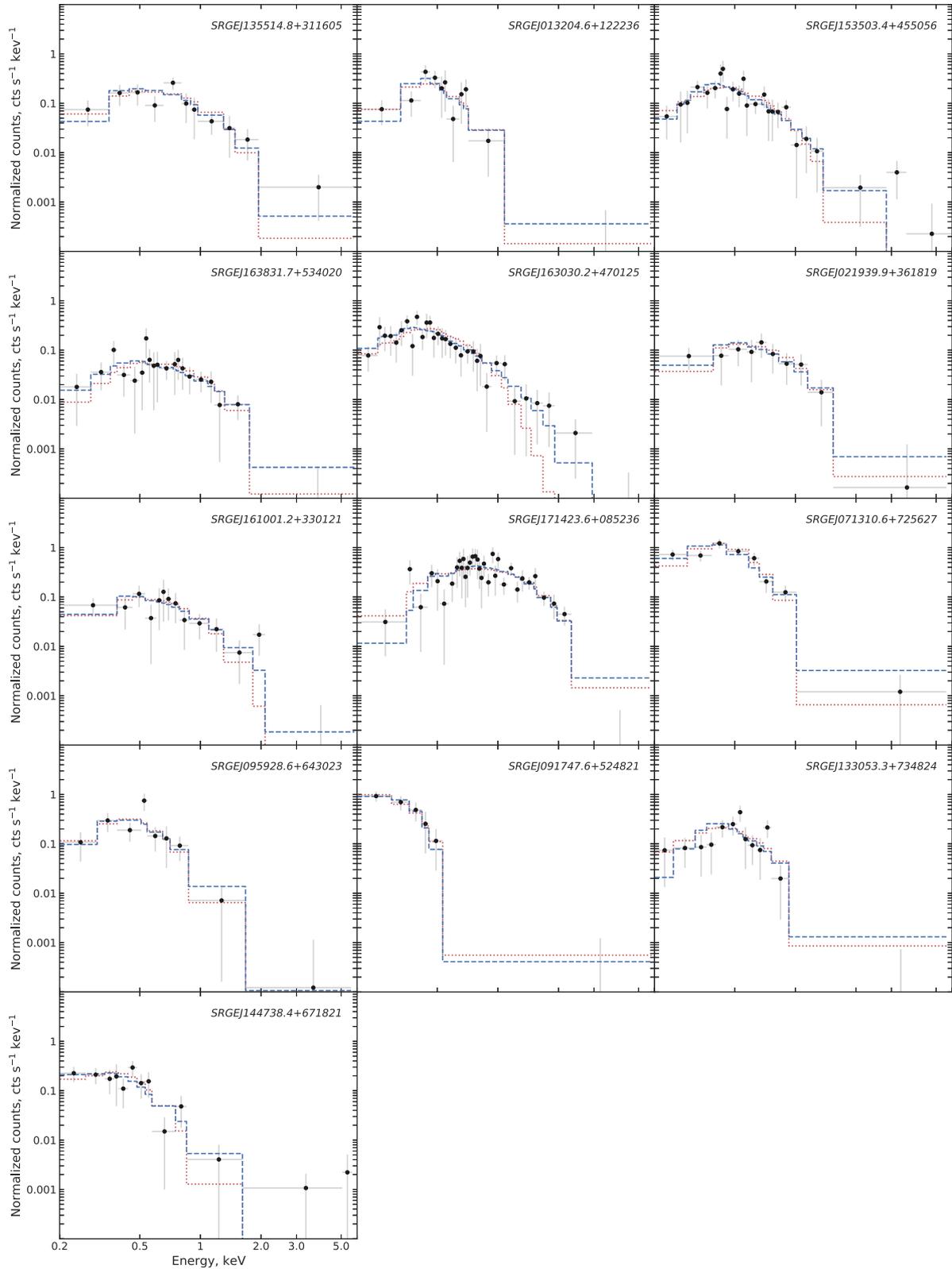


Figure 2. X-ray spectra of the TDEs obtained by eROSITA during the second SRG all-sky survey, fitted by the absorbed power-law model (blue dashed line) and the colour temperature corrected disc model (red dotted line) in the 0.2–6 keV energy range; see Tables 3 and 4 for the corresponding best-fitting parameters. The spectral channels are combined in bins with at least 1σ significance (for plotting purposes only).

Table 3. Parameters of simple models applied to eRASS2 spectra.

Absorbed power-law model: TBABS*ZPHABS*ZPOWERLW				
Object (SRGE)	N_{H} 10^{21} cm^{-2}	$K_{1 \text{ keV}}^a$	Γ	cstat/ d.o.f
J135514.8+311605	$1.9^{+2.7}_{-1.5}$	$3.1^{+2.8}_{-1.2} \times 10^{-4}$	$3.8^{+1.6}_{-1.4}$	28.5/25
J013204.6+122236	$2.3^{+3.9}_{-1.7}$	$7.2^{+7.0}_{-4.0} \times 10^{-5}$	$7.1^{+3.8}_{-2.2}$	35.7/43
J153503.4+455056	$1.2^{+1.2}_{-0.8}$	$1.5^{+0.6}_{-0.4} \times 10^{-4}$	$4.6^{+1.1}_{-0.9}$	68.4/64
J163831.7+534020	$0.7^{+2.9}_{-0.7}$	$1.6^{+1.6}_{-0.5} \times 10^{-4}$	$3.0^{+1.2}_{-0.6}$	91.1/99
J163030.2+470125	$0.6^{+0.7}_{-0.6}$	$1.9^{+0.5}_{-0.4} \times 10^{-4}$	$4.3^{+0.8}_{-0.7}$	59.8/77
J021939.9+361819	$1.8^{+5.0}_{-1.8}$	$3.4^{+7.2}_{-1.7} \times 10^{-4}$	$4.1^{+2.5}_{-1.7}$	19.5/29
J161001.2+330121	$0.7^{+1.5}_{-0.7}$	$9.5^{+4.8}_{-3.3} \times 10^{-5}$	$3.4^{+1.5}_{-1.1}$	44.9/43
J171423.6+085236	$2.9^{+2.3}_{-1.7}$	$1.1^{+0.8}_{-0.4} \times 10^{-3}$	$3.1^{+1.0}_{-0.8}$	62.8/55
J071310.6+725627	$0.4^{+0.5}_{-0.4}$	$1.9^{+0.7}_{-0.6} \times 10^{-4}$	$5.4^{+0.9}_{-0.8}$	27.5/19
J095928.6+643023	$2.4^{+2.1}_{-2.4}$	$4.6^{+148.4}_{-2.3} \times 10^{-4}$	$5.7^{+8.9}_{-2.1}$	18.6/17
J091747.6+524821	$1.6^{+1.9}_{-1.2}$	$9.5^{+91.0}_{-9.3} \times 10^{-7}$	$12.4^{+5.9}_{-4.3}$	9.9/20
J133053.3+734824	$1.7^{+4.8}_{-1.1}$	$8.7^{+7.1}_{-3.5} \times 10^{-5}$	$5.7^{+3.8}_{-1.4}$	36.3/38
J144738.4+671821	$0.3^{+0.6}_{-0.3}$	$3.0^{+2.3}_{-1.6} \times 10^{-5}$	$5.2^{+1.6}_{-1.1}$	34.9/41
Multiblackbody accretion disc model: TBABS*ZASHIFT*DISKBB				
Object (SRGE)	$R_{\text{in}} \sqrt{\cos \theta}$ km	kT_{in} eV	cstat/ d.o.f	
J135514.8+311605	$1.9^{+1.3}_{-0.8} \times 10^5$	324^{+100}_{-66}	30.3/26	
J013204.6+122236	$1.3^{+1.4}_{-0.7} \times 10^6$	121^{+31}_{-24}	36.9/44	
J153503.4+455056	$7.6^{+2.5}_{-2.5} \times 10^5$	188^{+30}_{-24}	69.2/65	
J163831.7+534020	$3.1^{+2.4}_{-1.4} \times 10^5$	413^{+128}_{-105}	93.5/100	
J163030.2+470125	$1.4^{+0.6}_{-0.4} \times 10^6$	178^{+28}_{-23}	73.6/78	
J021939.9+361819	$5.7^{+6.1}_{-3.0} \times 10^5$	288^{+105}_{-69}	20.4/30	
J161001.2+330121	$1.4^{+1.5}_{-0.8} \times 10^5$	247^{+117}_{-64}	49.4/44	
J171423.6+085236	$1.8^{+0.8}_{-0.6} \times 10^4$	527^{+126}_{-87}	60.7/56	
J071310.6+725627	$2.8^{+1.2}_{-0.9} \times 10^6$	109^{+13}_{-11}	26.3/20	
J095928.6+643023	$3.9^{+2.9}_{-2.0} \times 10^6$	164^{+49}_{-32}	19.4/18	
J091747.6+524821	$4.4^{+4.6}_{-2.2} \times 10^7$	54^{+10}_{-8}	10.1/21	
J133053.3+734824	$7.4^{+4.9}_{-3.0} \times 10^5$	150^{+31}_{-24}	30.3/39	
J144738.4+671821	$1.7^{+1.4}_{-0.8} \times 10^6$	100^{+22}_{-17}	36.3/42	

Note. ^aPower-law normalization at 1 keV in the source's rest frame, in units of photons $\text{keV}^{-1} \text{cm}^{-2} \text{s}^{-1}$.

We note, however, that OPTXAGNF is not applicable to the case of a slim accretion disc. Hence, whenever our spectral analysis implies a super-Eddington accretion rate for a given TDE, the inferred parameter values should be taken with great caution.

The best fits of the eRASS2 spectra by the OPTXAGNF model are shown in Fig. 2 in comparison with those by the power-law model. The quality of approximation by the OPTXAGNF model proves to be nearly insensitive to the black hole spin. We thus fixed this parameter at two extreme values: $a_* = 0$ (Schwarzschild black hole) and $a_* = 0.998$ (maximally rotating Kerr black hole).³ The resulting best-fitting parameters are given in Table 4. The quality of approximation by the OPTXAGNF model is nearly identical to that by the DISKBB model (Table 3), which is not surprising since both models describe multiblackbody accretion disc emission.

As seen from Table 4, the black hole mass for each of our objects may range by a factor of ~ 8 depending on the unknown spin of the black hole, with M_{BH} increasing with a_* . The corresponding accretion rates are close to the Eddington limit if the black holes are rapidly spinning ($a_* \approx 0.998$) and are supercritical if the rotation is not extreme. As already noted, the inferred parameter values, including the black hole masses, become unreliable in the latter case.

Although our X-ray spectral analysis is admittedly simplistic, we can none the less conclude that the available X-ray spectral data are consistent with the studied events being stellar disruptions by black holes with masses between $\sim 3 \times 10^3$ and $\sim 10^8 M_{\odot}$.

³ Assuming that the accretion disc rotates in the same direction as the black hole.

Table 4. Parameters of the colour temperature corrected disc model applied to eRASS2 spectra.

TBABS*OPTXAGNF				
Object (SRGE)	a_*	$\log \lambda_{\text{Edd}}$	$M (M_{\odot})$	cstat/d.o.f
J135514.8+311605	0	$1.13^{+0.25}_{-0.20}$	$3.3^{+2.7}_{-1.6} \times 10^4$	30.4/26
J013204.6+122236	0	$0.20^{+0.14}_{-0.12}$	$3.2^{+3.7}_{-1.7} \times 10^5$	36.8/44
J153503.4+455056	0	$0.73^{+0.11}_{-0.10}$	$1.5^{+0.8}_{-0.5} \times 10^5$	69.3/65
J163831.7+534020	0	$1.68^{+0.24}_{-0.21}$	$3.9^{+3.5}_{-1.9} \times 10^4$	93.5/100
J163030.2+470125	0	$0.87^{+0.11}_{-0.10}$	$2.6^{+1.3}_{-0.9} \times 10^5$	74.0/78
J021939.9+361819	0	$1.35^{+0.26}_{-0.21}$	$9.1^{+10.8}_{-5.1} \times 10^4$	20.5/30
J161001.2+330121	0	$0.54^{+0.35}_{-0.25}$	$2.9^{+3.6}_{-1.8} \times 10^4$	49.4/44
J171423.6+085236	0	$1.07^{+0.22}_{-0.18}$	$3.3^{+1.6}_{-1.2} \times 10^3$	60.7/56
J071310.6+725627	0	$0.34^{+0.07}_{-0.06}$	$7.0^{+3.3}_{-2.3} \times 10^5$	16.5/20
J095928.6+643023	0	$1.12^{+0.17}_{-0.13}$	$6.8^{+7.3}_{-3.8} \times 10^5$	19.4/18
J091747.6+524821	0	$0.21^{+0.06}_{-0.06}$	$1.2^{+1.3}_{-0.6} \times 10^7$	10.1/21
J133053.3+734824	0	$0.33^{+0.13}_{-0.12}$	$1.7^{+1.2}_{-0.7} \times 10^5$	30.2/39
J144738.4+671821	0	$-0.04^{+0.12}_{-0.11}$	$4.3^{+3.9}_{-2.0} \times 10^5$	36.4/42
J135514.8+311605	0.998	$0.22^{+0.25}_{-0.20}$	$2.5^{+2.0}_{-1.2} \times 10^5$	30.4/26
J013204.6+122236	0.998	$-0.72^{+0.14}_{-0.12}$	$2.4^{+2.8}_{-1.3} \times 10^6$	36.8/44
J153503.4+455056	0.998	$-0.19^{+0.11}_{-0.10}$	$1.2^{+0.6}_{-0.4} \times 10^6$	69.3/65
J163831.7+534020	0.998	$0.76^{+0.25}_{-0.21}$	$2.9^{+2.7}_{-1.4} \times 10^5$	93.5/100
J163030.2+470125	0.998	$-0.05^{+0.11}_{-0.10}$	$2.0^{+1.0}_{-0.7} \times 10^6$	74.0/78
J021939.9+361819	0.998	$0.43^{+0.26}_{-0.21}$	$6.8^{+8.1}_{-3.9} \times 10^5$	20.4/30
J161001.2+330121	0.998	$-0.38^{+0.36}_{-0.25}$	$2.2^{+2.7}_{-1.3} \times 10^5$	49.4/44
J171423.6+085236	0.998	$0.16^{+0.23}_{-0.18}$	$2.4^{+1.2}_{-0.9} \times 10^4$	60.8/56
J071310.6+725627	0.998	$-0.58^{+0.07}_{-0.06}$	$5.3^{+2.5}_{-1.7} \times 10^6$	26.5/20
J095928.6+643023	0.998	$0.20^{+0.17}_{-0.13}$	$5.1^{+5.5}_{-2.8} \times 10^6$	19.4/18
J091747.6+524821	0.998	$-0.71^{+0.06}_{-0.06}$	$8.9^{+9.9}_{-4.6} \times 10^7$	10.1/21
J133053.3+734824	0.998	$-0.59^{+0.13}_{-0.12}$	$1.3^{+0.9}_{-0.6} \times 10^6$	30.2/39
J144738.4+671821	0.998	$-0.96^{+0.12}_{-0.11}$	$3.2^{+2.9}_{-1.5} \times 10^6$	36.4/42

3.3.3 Late-phase X-ray spectra

As was mentioned before, we have also analysed the spectra of four TDEs obtained in eRASS3, 6 months after the discovery of these transients by eROSITA. These spectra are compared in Fig. 3 with the corresponding spectra taken during eRASS2. We similarly applied the absorbed power-law, multiblackbody accretion disc, and the colour temperature corrected disc models to these late-phase TDE spectra; see the resulting best-fitting parameters in Table 5.

Individual TDEs show significantly different evolution. In the case of SRGE J071310.6+725627 and SRGE J095928.6+643023, the black hole masses estimated from the late-phase spectra are consistent with those inferred from the early-phase ones and the observed spectral evolution can be accounted for by moderate decreases in the accretion rates.

The situation is quite different for SRGE J153503.4+455056. Here, fitting the eRASS3 spectrum by OPTXAGNF leads to a much lower black hole mass compared to the eRASS2 spectrum, which indicates that this model is inadequate for description of the observed spectral evolution. We might be witnessing a change from purely thermal accretion disc emission to a harder spectral state, where a major contribution is provided by Comptonized emission associated with a freshly formed hot corona of the accretion disc. This is also suggested by the somewhat better fit quality provided by the power-law model compared to OPTXAGNF for the eRASS3 spectrum (see Table 5).

The least clear situation is with SRGE J144738.4+671821. Here too, the eRASS3 spectrum is much harder than the eRASS2 one and the early- and late-phase M_{BH} estimates are inconsistent with each other, despite the good fit quality provided by the OPTXAGNF model

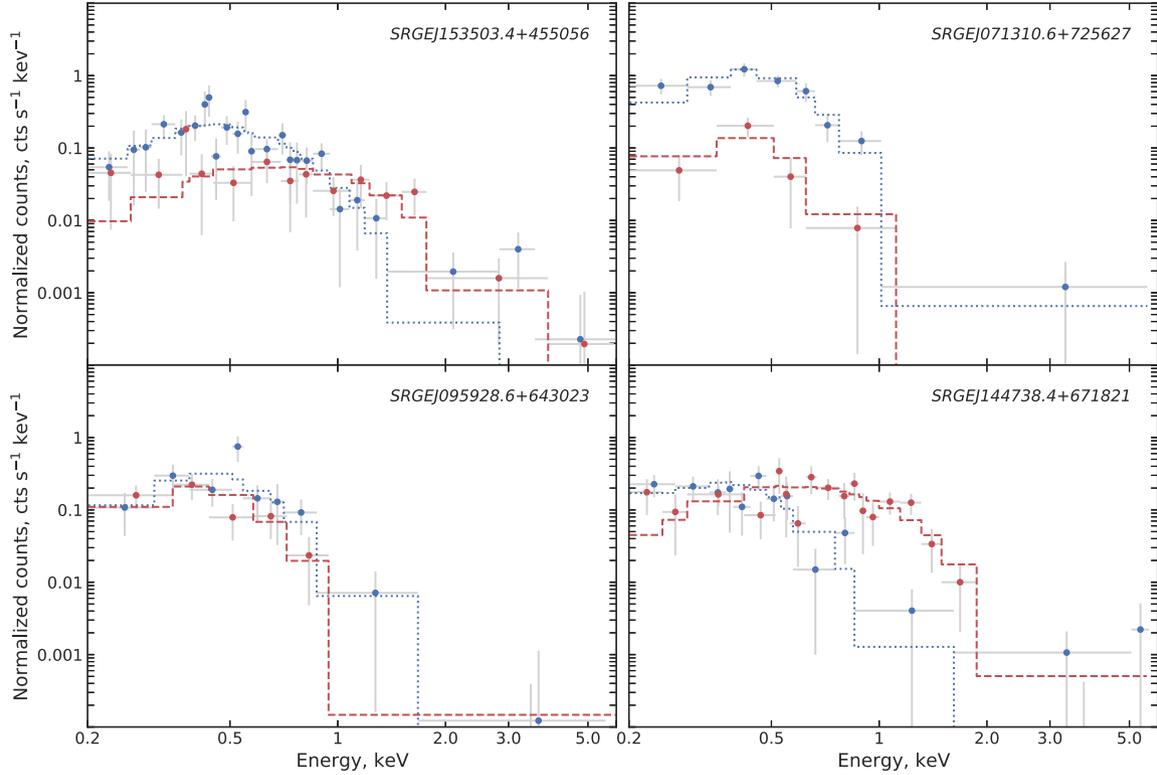


Figure 3. Comparison of the X-ray spectra of four TDEs obtained by eROSITA in the second (blue) and third (red) SRG all-sky surveys. The dotted and dashed lines show the corresponding best fits by the colour temperature corrected disc model.

Table 5. Best-fitting model parameters for eRASS3 spectra.

Absorbed power-law model: TBABS*ZPHABS*ZPOWERLW				
Object (SRGE)	N_{H} 10^{21} cm^{-2}	$K_{1\text{keV}}^a$	Γ	cstat d.o.f
J153503.4+455056	<0.5	$8.7^{+2.9}_{-2.4} \times 10^{-5}$	$2.0^{+0.6}_{-0.6}$	47.9/42
J071310.6+725627	$3.0^{+6.2}_{-2.2}$	$9.3^{+34.1}_{-9.3} \times 10^{-6}$	$10.1^{+6.0}_{-4.1}$	8.7/19
J095928.6+643023	$1.3^{+3.9}_{-1.3}$	$1.8^{+2.2}_{-0.9} \times 10^{-4}$	$6.0^{+3.3}_{-1.7}$	8.5/16
J144738.4+671821	<0.6	$2.6^{+0.7}_{-0.5} \times 10^{-4}$	$2.4^{+0.6}_{-0.3}$	34.9/27
Multiblackbody accretion disc model: TBABS*ZASHIFT*DISKBB				
Object (SRGE)	$R_{\text{in}} \sqrt{\cos \theta}$ km	kT_{in} eV	cstat/ d.o.f	
J153503.4+455056	$5.7^{+7.9}_{-4.4} \times 10^4$	499^{+640}_{-155}	52.1/43	
J071310.6+725627	$1.5^{+3.1}_{-0.9} \times 10^6$	98^{+37}_{-27}	12.6/20	
J095928.6+643023	$5.7^{+8.2}_{-3.2} \times 10^6$	138^{+42}_{-33}	10.9/17	
J144738.4+671821	$1.2^{+0.6}_{-0.4} \times 10^5$	335^{+75}_{-54}	30.5/28	
Colour temperature corrected disc model: TBABS*OPTXAGNF				
Object (SRGE)	a_*	$\log \lambda_{\text{Edd}}$	M, M_{\odot}	cstat d.o.f
J153503.4+455056	0	$1.40^{+0.85}_{-0.35}$	$9.1^{+12.9}_{-7.4} \times 10^3$	52.1/43
J071310.6+725627	0	$-0.13^{+0.18}_{-0.15}$	$3.8^{+9.0}_{-2.5} \times 10^5$	12.6/20
J095928.6+643023	0	$0.97^{+0.16}_{-0.14}$	$1.0^{+1.6}_{-0.6} \times 10^6$	10.9/17
J144738.4+671821	0	$1.04^{+0.19}_{-0.16}$	$2.3^{+1.3}_{-0.9} \times 10^4$	30.5/28
J153503.4+455056	0.998	$0.49^{+0.90}_{-0.36}$	$6.7^{+9.7}_{-3.5} \times 10^4$	52.0/43
J071310.6+725627	0.998	$-1.05^{+0.18}_{-0.15}$	$2.9^{+6.7}_{-1.9} \times 10^6$	12.6/20
J095928.6+643023	0.998	$0.05^{+0.16}_{-0.14}$	$7.7^{+12.3}_{-4.5} \times 10^6$	10.9/17
J144738.4+671821	0.998	$0.12^{+0.19}_{-0.16}$	$1.7^{+1.0}_{-0.7} \times 10^5$	30.4/28

Note. ^aPower-law normalization at 1 keV in the source's rest frame, in units of photons $\text{keV}^{-1} \text{cm}^{-2} \text{s}^{-1}$.

for both spectra. We recall that this TDE is unique in that eROSITA seems to have caught it during an usually long rising phase, so that it has become brighter in eRASS3 compared to eRASS2. Hopefully, if SRGEJ144738.4+671821 remains bright in several subsequent

SRG/eROSITA scans, we will be able to better understand its spectral evolution and nature.

4 OPTICAL/INFRARED PROPERTIES

Fig. 4 shows optical images around the studied objects from the Panoramic Survey Telescope and Rapid Response System DR1 (Pan-STARRS, PS1; Flewelling et al. 2020; Waters et al. 2020). There is a single potential optical counterpart within each eROSITA localization region. All of these candidates appear to be extended and thus can be TDE host galaxies. Their optical positions are provided in Table 11.

We carried out spectroscopy and photometry of the candidate optical counterparts of the eROSITA transients using a number of telescopes and instruments, namely the CCD-photometer (CMO RC600; Berdnikov et al. 2020) on the RC600 60-cm telescope of the Caucasus Mountain Observatory of the Sternberg Astronomical Institute (CMO SAI MSU, Russia), the ADAM low and medium resolution spectrograph (Afanasiev et al. 2016; Burenin et al. 2016) and the Andor iKon-M imaging camera on the AZT-33IK 1.6-m telescope (Kamus et al. 2002) of the Sayan Observatory (Russia), the TÜBITAK Faint Object Spectrograph and Camera⁴ (TFOSC) on the Russian-Turkish 1.5-m Telescope (RTT150) of the TÜBITAK National Observatory (Turkey), the Transient Double-Beam Spectrograph (TDS; Potanin et al. 2020) and the NBI CCD-photometer on the 2.5-m telescope of CMO SAI MSU (Russia), the SCORPIO-2 universal focal reducer (Afanasiev & Moiseev 2011) on the BTA 6-m telescope of the Special Astrophysical Observatory (Russia), and the

⁴<https://tug.tubitak.gov.tr/en/teleskoplar/rtt150-teleskope-0>

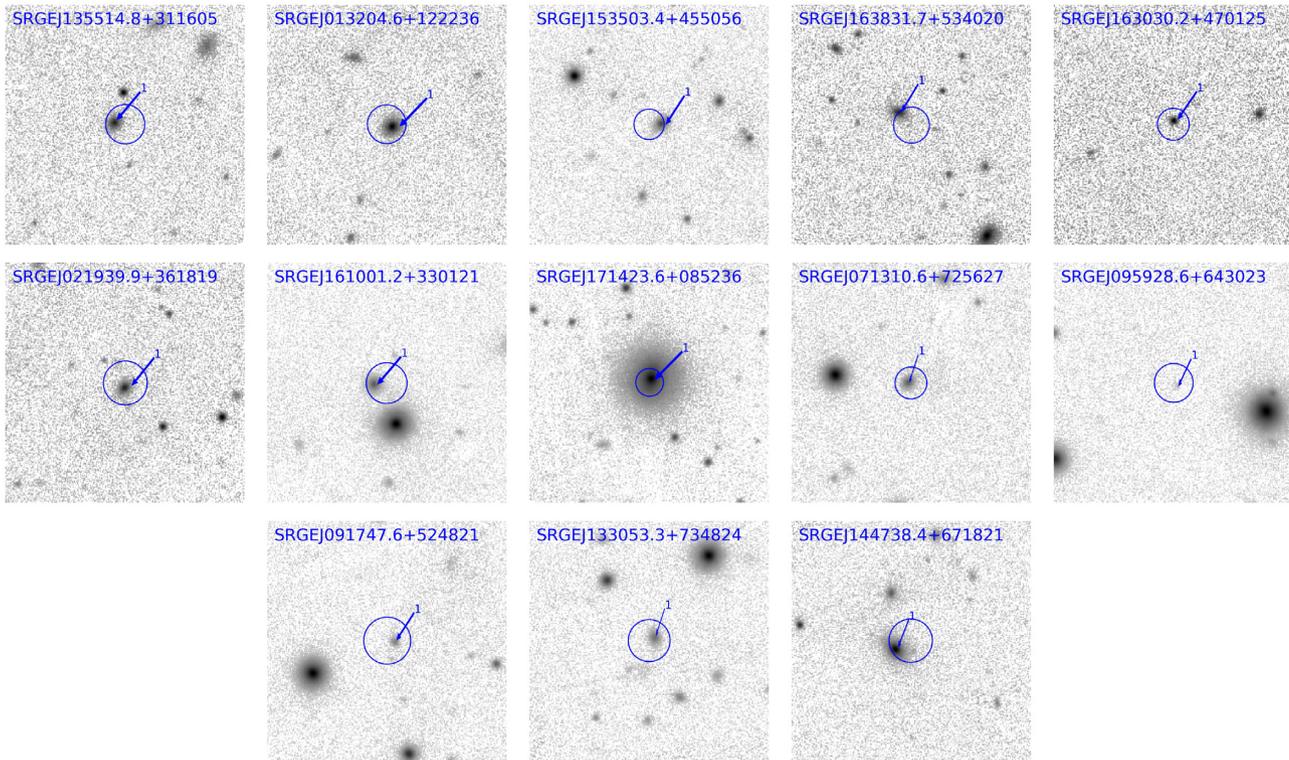


Figure 4. Pan-STARRS *i*-band 1 arcmin \times 1 arcmin images around the TDEs. In each panel, the circle shows the eROSITA localization region (with the size given in Table 1), while the arrow shows the object for which optical spectroscopy was performed.

Low Resolution Imaging Spectrograph (LRIS; Oke et al. 1995) on the Keck-I 10-m telescope (USA).

Table 6 presents a log of our photometric follow-up observations and the (*gri*) apparent magnitudes measured during these observations. Table 7 presents a log of our spectroscopic follow-up observations. Further details on the observations and data reduction are presented in Appendix B.

4.1 Optical light curves

The left-hand panels of Fig. 5 show long-term optical light curves of the TDE host galaxies constructed from our follow-up photometry (Table 6) and archival photometry provided by Pan-STARRS DR2 (PS2) and the Sloan Digital Sky Survey (SDSS; Alam et al. 2015). The right-hand panels show forced differential photometry light curves from ZTF, which cover epochs both before and after the eROSITA X-ray observations. Since the ZTF coverage at the locations of SRGEJ153503.4+455056, SRGEJ171423.6+085236, and SRGEJ144738.4+671821 is poor, we also show the Asteroid Terrestrial-impact Last Alert System (ATLAS; Tonry et al. 2018; Smith et al. 2020) forced photometry in the cyan (*c*) and orange (*o*) bands for these objects.

Three, or possibly four, TDEs exhibit optical flares. For SRGEJ153503.4+455056, we observe an optical brightening starting around 2019 July and ending before 2020 July, i.e. just before its discovery in X-rays by eROSITA on 2020 July 13–16, while during the previous visit in 2020 January eROSITA did not detect this source. For SRGEJ163030.2+470125, we observe a clear optical brightening starting in 2019 July and lasting at least until the end of 2020, whereas eROSITA discovered this transient in X-rays in 2020 August and did not detect it in 2020 February and at the

end of 2021 January. For SRGEJ091747.6+524821, we observe an optical brightening between 2020 June and 2021 March, while eROSITA discovered the X-ray transient in 2020 October and did not detect it in 2020 April and 2021 April. A fourth object, SRGEJ095928.6+643023, might also exhibit a weak optical flare around the eRASS2 visit in 2020 October–November.

Therefore, for two events (SRGEJ153503.4+455056 and SRGEJ163030.2+470125) there is evidence that the optical activity started before the X-ray activity, whereas the data for another two transients (SRGEJ095928.6+643023 and SRGEJ091747.6+524821) are consistent with the X-ray and optical activity occurring nearly concurrently.

The other nine events do not show evidence of optical flares. For SRGEJ171423.6+085236 and SRGEJ144738.4+671821, we observe moderate variability in the difference photometry light curve from 2015 July to 2021 July (ATLAS data before 2018 March are not shown in Fig. 5 for visual clarity). An amplitude of $< 20 \mu\text{Jy}$ in difference photometry corresponds to < 0.07 mag for SRGEJ171423.6+085236 and < 0.16 mag for SRGEJ144738.4+671821. Therefore, the scatter shown in Fig. 5 is probably due to imperfect image subtraction at the bright galaxy nuclei, rather than variability associated with the stellar tidal disruption. Some other events (e.g. SRGEJ135514.8+311605, SRGEJ013204.6+122236, SRGEJ071310.6+725627) show $> 3\sigma$ data points in the binned differential photometry light curves. However, those data are consistent with random fluctuations.

4.2 Constraints on optical luminosity

The observed optical flare peak magnitudes or upper limits are shown in Table 8. For the nine events without optical flares, we compute

Table 6. Log of photometric observations.

Object (SRGE)	Date	Telescope	Exposure (s), <i>gri</i>	m_g	m_r	m_i
J135514.8+311605	2020 July 21	RTT150	10 × 60, 10 × 60, 10 × 60	20.49 ± 0.07	19.60 ± 0.04	19.13 ± 0.05
J153503.4+455056	2020 July 23–27	CMO	4 × 200, 8 × 200+4 × 300, 200 + 3 × 300	20.06 ± 0.05	19.02 ± 0.03	18.55 ± 0.05
	2020 Aug 5	CMO	3 × 200, 3 × 200 + 2 × 300, 9 × 200	20.01 ± 0.05	19.02 ± 0.02	18.54 ± 0.04
	2020 Sept 17	CMO	3 × 300, 4 × 300, 3 × 300	20.09 ± 0.02	19.02 ± 0.01	18.49 ± 0.01
J163030.2+470125	2020 Oct 17	BTA	4 × 40, 4 × 40, 4 × 40	21.45 ± 0.08	20.42 ± 0.06	20.04 ± 0.06
J161001.2+330121	2021 Mar 12	RTT150	2 × 300, 3 × 300, 2 × 300	19.28 ± 0.01	18.51 ± 0.01	18.11 ± 0.01
J171423.6+085236	2020 Oct 18	CMO RC600	2 × 200, 13 × 200, 2 × 200	16.90 ± 0.03	16.02 ± 0.02	15.57 ± 0.06
	2020 Oct 21	CMO RC600	8 × 240, 3 × 240, 3 × 240	16.86 ± 0.02	16.00 ± 0.03	15.59 ± 0.02
J095928.6+643023	2021 Mar 10	AZT-33IK	18 × 90, 18 × 90, 18 × 90	23.34 ± 0.22	22.31 ± 0.11	22.23 ± 0.13

Table 7. Log of spectroscopic observations.

Object (SRGE)	Date	Telescope	Exp. (s)
J135514.8+311605	2020 Dec 17	BTA	4 × 900
	2021 Jun 7	Keck-I	435
J013204.6+122236	2021 Jul 6	Keck-I	600
J153503.4+455056	2020 Sep 16–20	CMO	21 × 1200
	2021 May 13	Keck-I	570
J163831.7+534020	2021 Apr 14	Keck-I	1250
J163030.2+470125	2020 Oct 17	BTA	4 × 900
	2021 Jun 7	Keck-I	1010
J021939.9+361819	2021 Jul 6	Keck-I	750
J161001.2+330121	2021 Apr 14	Keck-I	300
J171423.6+085236	2020 Oct 18, 21	CMO	2 × 1200
	2021 Jun 7	Keck-I	280
J071310.6+725627	2020 Nov 11–12	AZT-33IK	7 × 600
	2020 Nov 20	Keck-I	900
J095928.6+643023	2021 May 13	Keck-I	2500
J091747.6+524821	2021 Apr 14	Keck-I	600
J133053.3+734824	2021 Apr 14	Keck-I	400
J144738.4+671821	2021 Mar 11	AZT-33IK	5 × 600
	2021 Apr 14	Keck-I	300

long-term median of 3σ upper limits using ZTF and ATLAS forced differential photometry. For the four events with optical flares, we calculate maximum of the model fits shown in Fig. 5. Since the optical data have seasonal gaps, the peak of the optical emission might be missed. Therefore, we consider the actual optical peak to be 1–2 times the model maximum.

Previous studies have shown that the broad-band SED of UV or optically discovered TDEs can be described by blackbody spectra with temperatures (T_{bb}) between $\sim 1.3 \times 10^4$ and $\sim 4.0 \times 10^4$ K (van Velzen et al. 2020; Gezari 2021). Assuming three typical values of T_{bb} , we report the constraints on optical luminosities in Table 8. Although the total blackbody luminosity L_{bb} is largely model dependent, the rest-frame g -band luminosity L_g only has a weak dependence on T_{bb} .

4.3 WISE detection of a luminous infrared echo in SRGE J153503.4+455056

If the circumnuclear medium of a TDE is dusty, a significant fraction of the UV/optical radiation energy will be absorbed by dust and reprocessed to the IR (Lu, Kumar & Evans 2016). The resulting IR echoes bear important information on the dust properties at sub-pc scales in quiescent galaxies (van Velzen et al. 2016, 2021a).

By performing photometry on time-resolved AllWISE/NEOWISE coadds, Jiang et al. (2021) reported the detection of IR echoes in 8 of 23 optically selected TDEs. The ratio of their peak dust luminosity ($L_{\text{dust, peak}} \sim 10^{41}$ – 10^{42} erg s $^{-1}$) and their UV/optical luminosity ($L_{\text{bb, peak}} \sim 10^{44}$ erg s $^{-1}$) suggests a dust-covering factor $f_c = L_{\text{dust, peak}}/L_{\text{bb, peak}}$ of $\lesssim 0.01$.

In order to search for IR echoes in the SRG TDE sample, we collected AllWISE and NEOWISE-R data for the TDE hosts. Significant IR brightening was observed in SRGE J153503.4+455056 (Fig. 6). In the most recent NEOWISE epoch (from 2021 June 26 to July 1), compared with the 2010 baseline, its IR flux has increased by 150 ± 7 μ Jy in W1 and 174 ± 16 μ Jy in W2. The corresponding dust luminosity and temperature are $L_{\text{dust}} = (3.0 \pm 0.2) \times 10^{43}$ erg s $^{-1}$ and $T_{\text{dust}} = 1393^{+196}_{-145}$ K. The derived L_{dust} is a lower limit on $L_{\text{dust, peak}}$, but is already at the same level of the most luminous dust echo detected in optically selected TDEs (Jiang et al. 2021). We have constrained the peak UV/optical luminosity of SRGE J153503.4+455056 to be $L_{\text{bb, peak}} < 4 \times 10^{44}$ erg s $^{-1}$ (Table 8). Taken together, we infer a dust covering factor of $f_c > 0.1$, which is greater than typical values seen in optically selected TDEs.

Future detailed studies on the IR properties of SRG/eROSITA selected TDEs should reveal if the dust environment of X-ray selected TDEs is statistically different from that of optically selected events.

4.4 Optical spectra

Fig. 7 shows the optical spectra obtained during our follow-up program. For some objects, only a single-epoch spectrum is available, while for others we obtained a couple of spectra with an interval of several months. The achieved spectral quality is sufficient for a reliable measurement of the redshift in all cases.

Only in the spectrum of SRGE J163030.2+470125 do we see a clear indication of optical emission associated with the presumed stellar tidal disruption on top of the host galaxy light. Specifically, the first spectrum taken in 2020 October, i.e. nearly 2 months after the eROSITA discovery, exhibits a blue continuum, while this blue excess is not observed anymore in a second spectrum obtained 8 months later (in 2021 June). This spectral evolution is consistent with the brightening seen in the optical light curve between the middle of 2019 and the end of 2020.

The other three objects demonstrating optical flares in their light curves (as was discussed in Section 4.1), namely SRGE J153503.4+455056, SRGE J095928.6+643023, and SRGE J091747.6+524821, do not show any signatures in their optical spectra that could be attributed to a TDE. This is, however, fully consistent with the fact that we conducted spectroscopy for these events *after* the flares faded away.

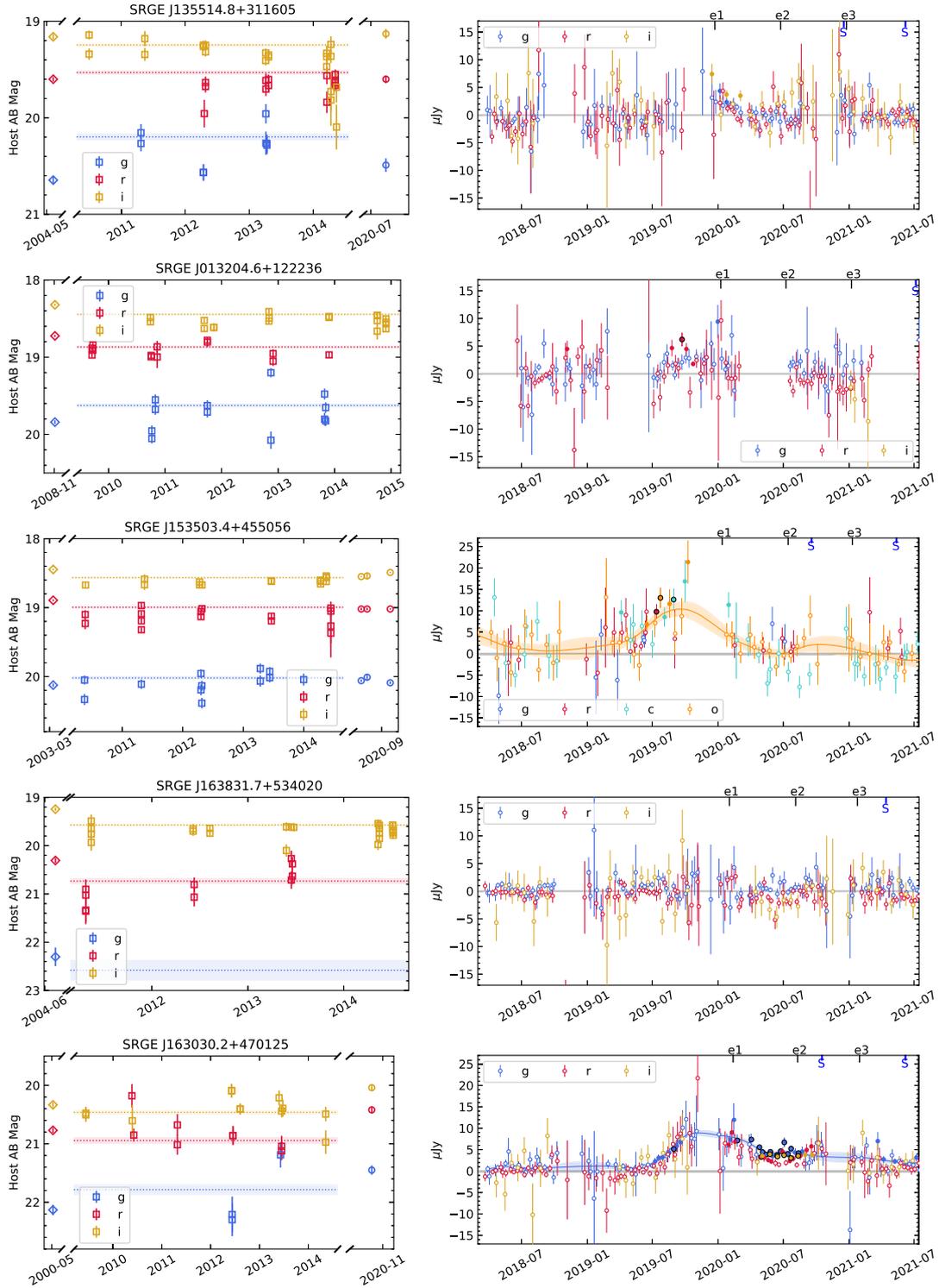


Figure 5. *Left:* Optical light curves of *SRG* TDE host galaxies, showing SDSS (diamonds), Pan-STARRS DR2 (PS2, squares), and photometry obtained by us (circles, Table 6). The horizontal dotted lines mark the PS2 (*gri*) Kron magnitudes from stacked images. *Right:* ZTF and ATLAS differential photometry performed at the centroids of the host galaxies. $<3\sigma$ data points are shown in hollow circles, $>3\sigma$ data points are shown in solid circles, and $>5\sigma$ data points are further highlighted using the black edge colour. Along the upper axis, epochs of the *SRG*/eROSITA visit (*e1*, *e2*, *e3*) and spectroscopic observations (*S*) are marked. For visual clarity, the ZTF-*g* and ZTF-*r* photometry are binned by 10 d, and the ZTF-*i* and ATLAS photometry are binned by 20 d. For the four TDEs with optical flares, we show Gaussian process model fits following procedures described in appendix B.4 of Yao et al. (2020). The models are fitted to a single band where the photometric uncertainty is the smallest or the temporal coverage is the highest.

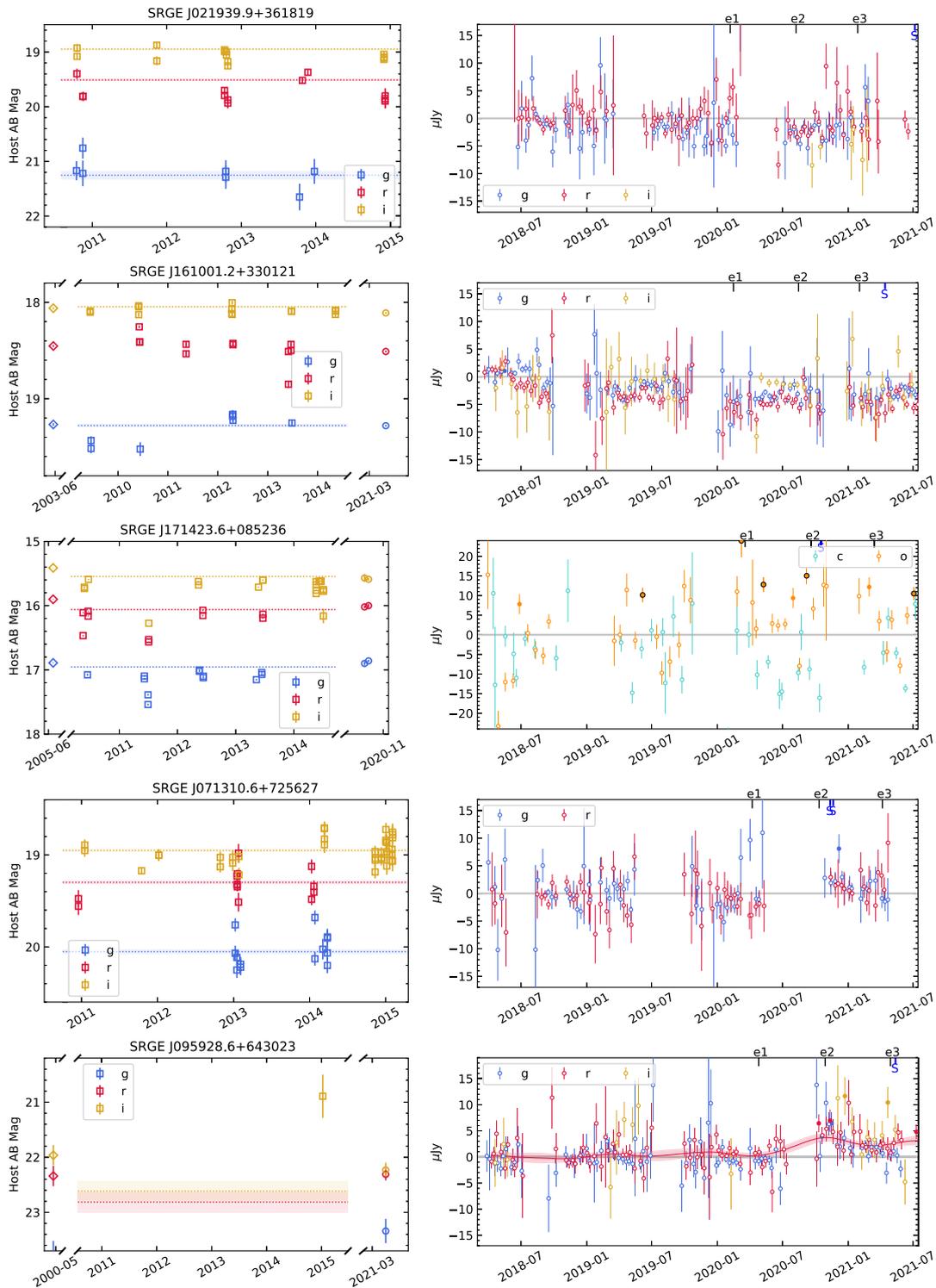


Figure 5. Continuation.

To characterize the emission content of the host galaxies, we fit the observed Keck spectra with stellar population models using PPF (Cappellari 2017). We use the MILES spectra library (Vazdekis et al. 2015), and emission lines of Balmer series, [O II], [S II], [O III], [O I], and [N II]. Among the 13 hosts, emission lines are confidently detected in 8 objects. Table 9 presents their line fluxes and classes on the Baldwin, Phillips, & Terlevich (BPT) diagnostic diagram

(Baldwin, Phillips & Terlevich 1981; Veilleux & Osterbrock 1987; Kewley et al. 2006).

In order to explore the star formation history of *SRG*-selected TDE hosts (discussed in Section 5.2), we measure the equivalent width (EW) of the Lick $H\delta_A$ index in the PPF best-fitting stellar continuum. This index was defined to capture stellar absorption from A stars (Worthey & Ottaviani 1997). We then measure the EW of the

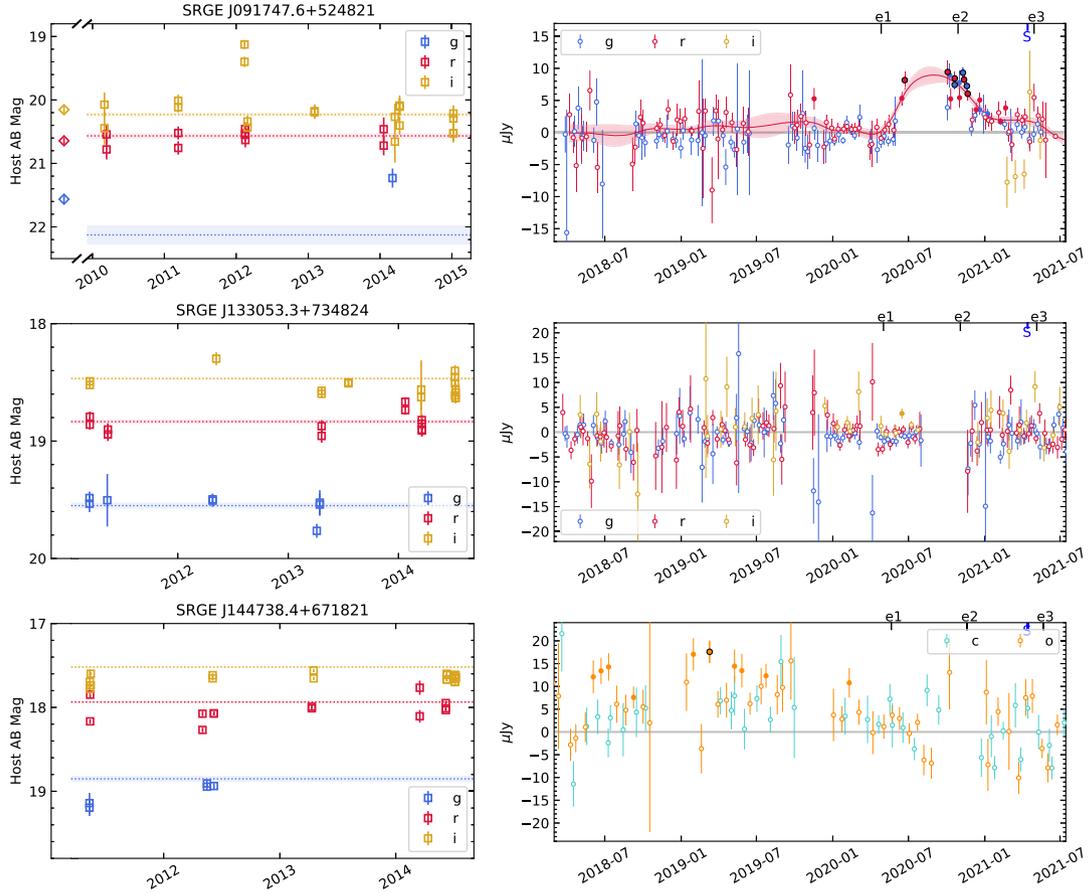


Figure 5. Continuation.

Table 8. Constraints on the optical luminosity.

Object (SRGE)	Optical flare peak AB mag	L_{bb}^a (10^{43} erg s $^{-1}$)			L_g^a (10^{43} erg s $^{-1}$)		
		1.3×10^4 K	2.5×10^4 K	4.0×10^4 K	1.3×10^4 K	2.5×10^4 K	4.0×10^4 K
J135514.8+311605	$g > 21.83$	<0.66	<1.92	<5.67	<0.39	<0.33	<0.31
J013204.6+122236	$g > 21.51$	<0.40	<1.24	<3.75	<0.22	<0.20	<0.19
J153503.4+455056	$21.33 < o < 20.58$	1.56–3.12	6.01–11.99	19.52–38.95	0.94–1.87	1.05–2.09	1.09–2.17
J163831.7+534020	$g > 21.84, r > 21.86$	<4.49	<10.24	<26.67	<3.46	<2.29	<1.91
J163030.2+470125	$21.47 < g < 20.71$	1.92–3.86	5.13–10.32	14.66–29.52	1.21–2.43	0.94–1.89	0.86–1.73
J021939.9+361819	$g > 21.43$	<3.34	<8.21	<22.78	<2.26	<1.61	<1.43
J161001.2+330121	$g > 21.68$	<0.34	<1.06	<3.18	<0.19	<0.17	<0.16
J171423.6+085236	$c > 20.04$	<0.13	<0.47	<1.51	<0.07	<0.07	<0.07
J071310.6+725627	$g > 21.41$	<0.28	<0.89	<2.69	<0.15	<0.14	<0.13
J095928.6+643023	$22.40 < r < 21.64$	1.80–3.62	5.80–11.67	17.73–35.71	1.27–2.56	1.19–2.41	1.17–2.35
J091747.6+524821	$21.49 < r < 20.74$	0.90–1.79	3.42–6.83	11.11–22.17	0.52–1.04	0.58–1.15	0.60–1.19
J133053.3+734824	$g > 21.73$	<0.42	<1.29	<3.86	<0.24	<0.21	<0.20
J144738.4+671821	$c > 20.41$	<1.04	<3.51	<10.90	<0.57	<0.56	<0.56

Note. Numbers presented in this table are corrected for Galactic extinction.

^aAssuming the optical emission can be described by a blackbody with typical temperatures of 1.3×10^4 , 2.5×10^4 , or 4.0×10^4 K, L_{bb} is the blackbody luminosity, and L_g is the rest-frame g -band luminosity.

$H\alpha$ emission using the following formula:

$$\text{EW}(H\alpha_{\text{em}}) = -[\text{EW}(H\alpha) - \text{EW}(H\alpha_{\text{abs}})], \quad (2)$$

where $\text{EW}(H\alpha)$ is the total EW of $H\alpha$ in the observed spectrum, and $\text{EW}(H\alpha_{\text{abs}})$ is the EW of the absorption line in the PPF best-

fitting stellar continuum. The $H\alpha$ line of SRGE J163831.7+534020 and SRGE J095928.6+643023 are redshifted out of the LRIS bandpass. For these two objects, we assume that the emission-line flux of $H\alpha$ is $\gtrsim 3$ \times the emission-line flux of $H\beta$. For SRGE J095928.6+643023, the $H\beta$ emission line is not confidently

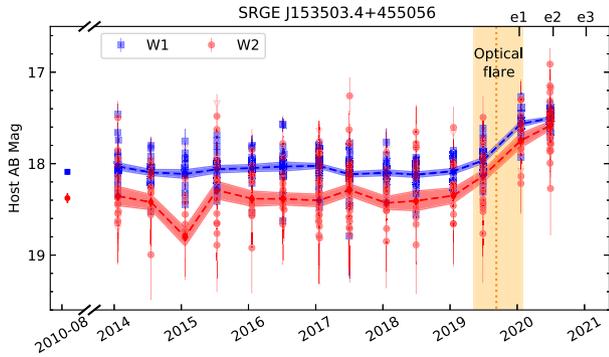


Figure 6. *WISE* W1 (3.4 μm) and W2 (4.6 μm) light curves of SRGE J153503.4+455056. The blue and red dashed lines mark the median magnitude of each *WISE* epoch. The vertical dotted line denotes the optical maximum. The vertical orange band marks the e-folding rise and decay timescales of the optical flare around the optical maximum (estimated using the Gaussian process model shown in Fig. 5). Epochs of the *SRG/eROSITA* visits are marked along the upper axis.

detected in the observed spectrum, and therefore $\text{EW}(\text{H}\alpha_{\text{em}}) \gtrsim 0 \text{ \AA}$. For SRGE J163831.7+534020, $f_{\text{H}\alpha} \gtrsim 1.9 \times 10^{-16} \text{ erg s}^{-1} \text{ cm}^{-2}$. By inspecting its observed spectrum (Fig. 7), we take the $\text{H}\alpha$ continuum flux to be $\sim 2.7 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ \AA}^{-1}$. Therefore, $\text{EW}(\text{H}\alpha_{\text{em}}) \gtrsim 7 \text{ \AA}$. Table 10 presents the measurements of $\text{EW}(\text{H}\alpha_{\text{em}})$ and $\text{EW}(\text{H}\delta_{\text{A}})$, as well as the corresponding *WISE* W1–W2 colours.

Most of the emission lines observed in the spectra are consistent with being due to star formation in the host galaxies. However, we see clear signatures of AGN activity (namely, high $[\text{N II}] \lambda 6583/\text{H}\alpha$, $[\text{S II}] \lambda \lambda 6717, 6731/\text{H}\alpha$, and $[\text{O III}] \lambda 5007/\text{H}\beta$ ratios), for a few objects, in particular SRGE J171423.6+085236, SRGE J133053.3+734824, and SRGE J144738.4+671821, which are classified according to the BPT diagram as ‘LINER’, ‘Composite’ (i.e. $[\text{H II}]/\text{LINER}$), and ‘LINER’ or ‘Seyfert’, respectively. Furthermore, although the $[\text{O III}] \lambda 5007$ and $\text{H}\beta$ emission lines are very weak or absent in some of the spectra, we can try to use the EW of the $\text{H}\alpha$ emission line as a proxy of the $[\text{O III}] \lambda 5007/\text{H}\beta$ ratio on the BPT diagram (Cid Fernandes et al. 2010) together with the $[\text{N II}] \lambda 6583/\text{H}\alpha$ ratio. Based on the values of $\text{EW}(\text{H}\alpha_{\text{em}})$ from Table 10, we infer that SRGE J153503.4+455056, SRGE J161001.2+330121, and SRGE J071310.6+725627 might also be LINERS.

Despite this evidence of AGN activity, it is unlikely that it is responsible for the X-ray transient phenomena revealed by *SRG/eROSITA* in these objects. First of all, $W1 - W2 < 0.3$ for all of the suspected AGNs (see Tables 9 and 10), which indicates that the mid-IR emission associated with an active nucleus is overwhelmed by that from the surrounding galaxy, hence the suspected AGNs cannot be luminous.

A more quantitative assessment of the AGN activity can be done based on the measured luminosity in the $[\text{O III}] \lambda 5007$ line (see Table 9). This quantity correlates with the X-ray luminosity in Seyfert galaxies (Heckman et al. 2005) and may be used as a proxy of AGN bolometric luminosity (e.g. LaMassa et al. 2010). Typically for Seyfert 1 galaxies (i.e. for AGNs whose observed spectra are not significantly affected by intrinsic absorption), the $[\text{O III}] \lambda 5007$ luminosity is $\sim 1\text{--}3$ per cent of the luminosity in the standard X-ray band (2–10 keV) (Heckman et al. 2005), although there is a substantial scatter around this mean trend. Taking into account that for typical SEDs of Seyferts/quasars (Sazonov, Ostriker & Sunyaev 2004) the luminosity in the 0.2–6 keV band (our working energy range in this study), L_X , is a factor of ≈ 2 higher than in the 2–10 keV

band, we may roughly predict the AGN contribution to the 0.2–6 keV luminosity of our objects as $L_X \sim 10^2 L_{[\text{O III}]}$. As shown in Fig. 8, the $L_X/L_{[\text{O III}]}$ ratios for all of the *SRG/eROSITA* transients with detectable $[\text{O III}] \lambda 5007$ emission are larger than 10^3 and concentrate around 10^4 .

We conclude that the few galaxies tentatively classified as LINER/Seyfert according to their optical emission-line ratios appear to be relatively low luminosity AGNs in which *SRG/eROSITA* has registered luminous soft X-ray outbursts associated with stellar tidal disruptions. The low persistent X-ray luminosity expected for these objects based on the $[\text{O III}] \lambda 5007$ flux is consistent with their non-detection by *eROSITA* during its first scan (eRASS1). However, we cannot completely rule out that AGNs do play some role in the properties of the X-ray transients discussed here (see Zabludoff et al. 2021 for a discussion of the overlap of the observed properties of TDEs and AGN), especially in the case of SRGE J144738.4+671821, which showed an atypical X-ray brightening over the 6-month period after its discovery by *eROSITA*. The presence of a few AGN in the current sample of 13 TDEs is not surprising, since roughly every second galaxy in the local Universe appears to have a weakly active nucleus (Ho 2008).

5 DISCUSSION

As has been demonstrated in the preceding sections, the totality of existing X-ray and optical data leaves little doubt that the objects under consideration are TDEs.

5.1 Properties of the SRG TDE sample

Table 11 summarizes the key properties of the discovered TDEs. Specifically, we provide for each object: (1) coordinates of the host galaxy, (2) the Galactic H I column in this direction, (3) its redshift, (4) TDE intrinsic X-ray luminosity in the 0.2–6 keV energy band, L_X , or the corresponding upper limit in eRASS2 and eRASS3, (5) TDE rest-frame g -band luminosity, L_g , or the corresponding upper limit, and (6) the volume of the Universe, V_{max} , within which the TDE could be detected during eRASS2 (see Section 5.3).

The quoted X-ray luminosities and upper limits were determined from the best-fitting DISKBB models (Tables 3 and 5) and corrected for the Galactic absorption in those cases where the available number of *eROSITA* photons had allowed us to perform a spectral analysis; otherwise the X-ray luminosities/upper limits for eRASS3 were estimated from the measured count rates adopting the best-fitting spectral model from eRASS2 (i.e. during the ‘bright phase’ of the TDE). The estimates and upper limits on L_g are adopted from the last three columns in Table 8, i.e. allowing for the characteristic blackbody temperature to range between 1.3×10^4 and 4.0×10^4 K.

Fig. 9 shows the distribution of the TDEs over redshift and X-ray luminosity. Thanks to the high sensitivity of the *SRG/eROSITA* all-sky survey, the effective ‘horizon’ of TDE observability in X-rays has moved out to $z \sim 0.6$ from $z \sim 0.15$, where it was during the *ROSAT* all-sky survey (Komossa 2015). Moreover, our current sample is based on a conservative, high detection threshold (see Section 2). The latter can be lowered in future work by a factor of ~ 2 , which should lead to the discovery of TDEs at even higher redshifts. Therefore, thanks to *SRG* we are starting to explore the TDE phenomenon beyond the low-redshift Universe.

In Section 3.3, we made an attempt to estimate the masses (M_{BH}) and Eddington ratios (λ_{Edd}) of the black holes associated with the TDEs (see Table 4) from their X-ray spectra measured by *eROSITA*,

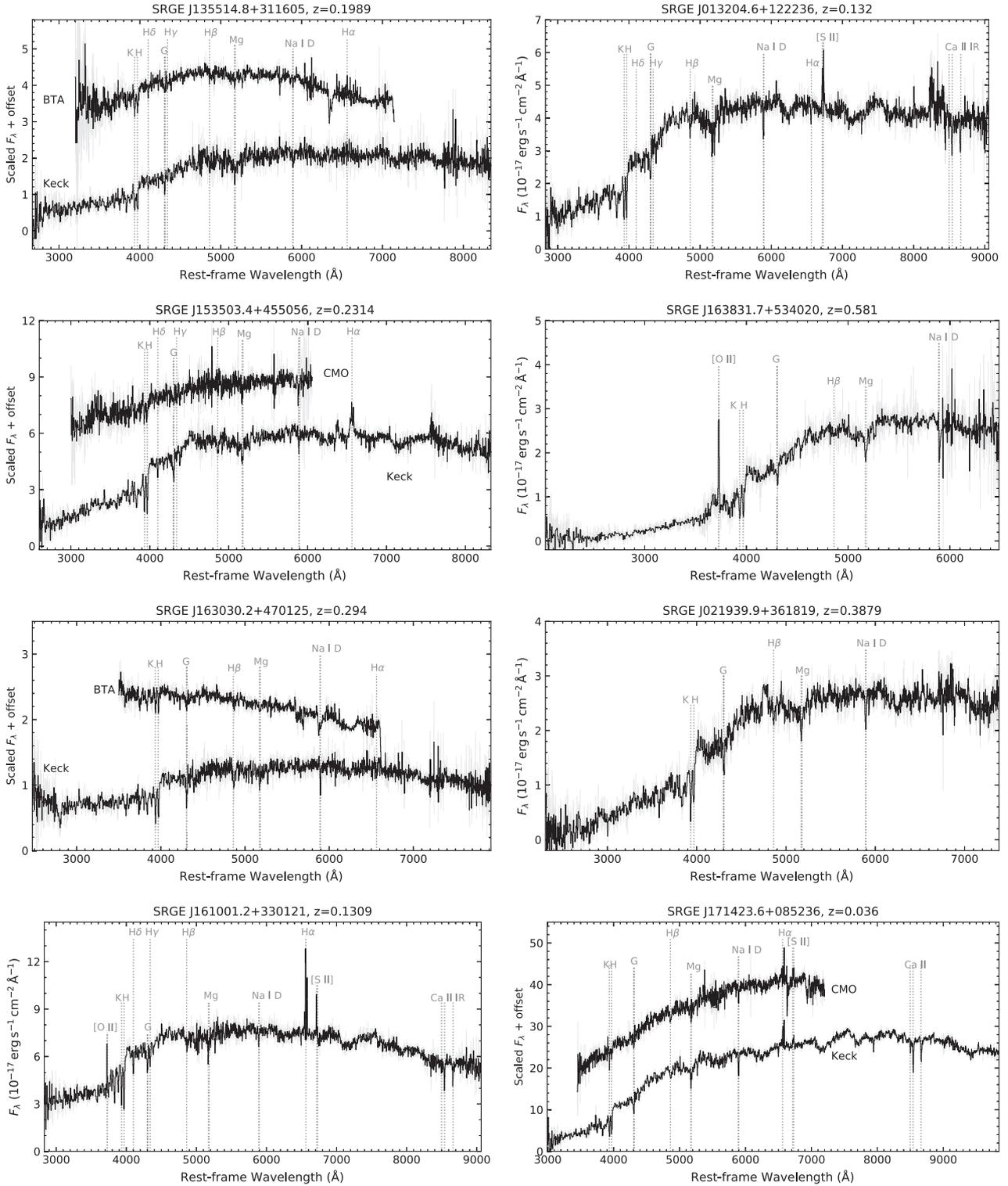


Figure 7. Optical spectra of SRG/eROSITA TDEs and/or their host galaxies.

based on the assumption of a standard accretion disc. As noted before, this assumption may fail for at least some of these events if eROSITA caught them in a super-Eddington accretion phase, when a slim rather than thin accretion disc would be expected. Furthermore, such estimates strongly depend on the adopted spin of the black hole. Bearing in mind these uncertainties, we plot in Fig. 10 the

inferred values of M_{BH} and λ_{Edd} for two extreme cases: $a_* = 0$ (a Schwarzschild black hole) and $a_* = 0.998$ (a maximally rotating Kerr black hole). In the case of a slowly rotating black hole, most of the TDEs would have been in a super-Eddington accretion phase at the epoch of their first detection by eROSITA, while in the $a_* = 0.998$ case, the deduced accretion rates are nearly critical. Regardless

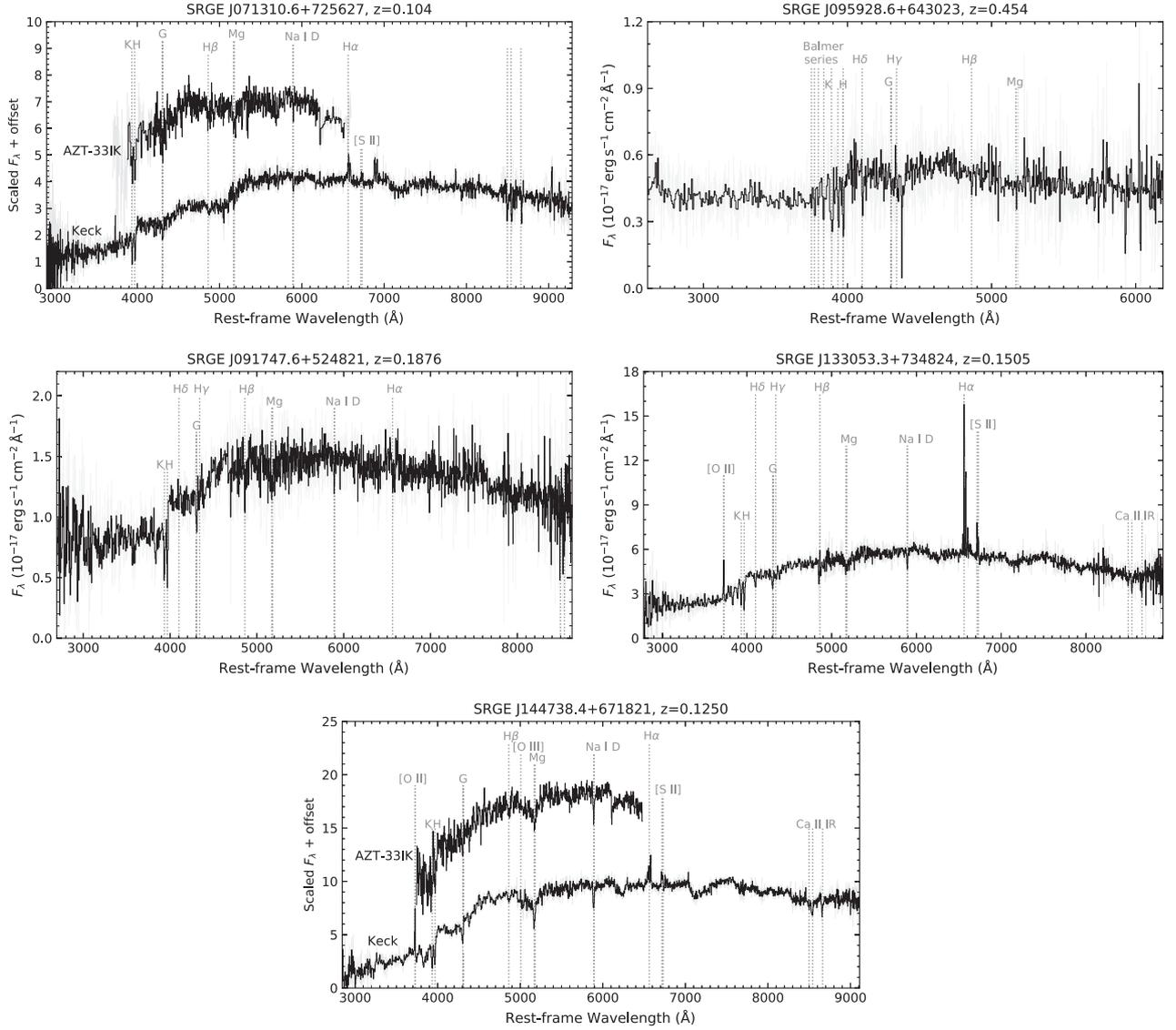


Figure 7. Continuation.

Table 9. Host properties (line fluxes).

Object (SRGE)	$f_{H\alpha}$	$f_{H\beta}$	$f_{[N\text{II}]6583}$	$f_{[S\text{II}]6717+6731}$	$f_{[O\text{II}]3726+3729}$	$f_{[O\text{III}]5007}$	$L_{[O\text{III}]5007}$	BPT class
J013204.6+122236	–	–	–	10.7 ± 2.4	–	3.5 ± 1.1	18.7 ± 6.0	–
J153503.4+455056	24.9 ± 1.2	–	14.2 ± 1.0	–	10.8 ± 2.7	4.4 ± 1.4	72.2 ± 23.5	–
J163831.7+534020	–	6.3 ± 0.8	–	–	39.6 ± 6.3	6.9 ± 1.1	987.3 ± 154.1	–
J161001.2+330121	30.4 ± 0.9	7.7 ± 1.1	19.1 ± 0.9	13.9 ± 2.2	24.7 ± 2.4	–	–	–
J171423.6+085236	42.3 ± 1.7	12.1 ± 1.3	38.5 ± 1.8	26.0 ± 3.5	24.7 ± 2.4	28.6 ± 1.5	12.6 ± 0.6	LINER
J071310.6+725627	11.8 ± 0.7	–	6.4 ± 0.7	4.3 ± 1.0	9.6 ± 2.5	–	–	–
J133053.3+734824	50.3 ± 0.9	10.4 ± 1.3	27.0 ± 0.8	15.5 ± 1.6	22.6 ± 2.0	3.9 ± 1.1	25.2 ± 6.8	Composite
J144738.4+671821	10.3 ± 1.0	–	19.8 ± 1.3	14.7 ± 2.4	44.4 ± 4.6	7.0 ± 2.4	29.9 ± 10.0	LINER/Seyfert

Notes. Observed line fluxes are given in units of $10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$. The luminosity of [O III]5007 (in units of $10^{38} \text{ erg s}^{-1}$) is corrected for Galactic extinction.

of the actual black hole spins, all the inferred black hole masses are consistent with being below the theoretical upper limits for non-spinning and rapidly spinning black holes in TDEs of $\sim 10^8$ and $\sim 7 \times 10^8 M_{\odot}$, respectively (Rees 1988; Kesden 2012).

5.2 Host galaxy properties

Following the procedures described by Mendel et al. (2014) and van Velzen et al. (2021b), we can estimate the stellar mass of the host galaxies of the SRG TDEs using pre-transient photometry from

Table 10. Host properties (equivalent width and *WISE* colours).

Object (SRGE)	EW(H α_{em})	EW(H δ_A)	W1–W2
J135514.8+311605	2.42 \pm 0.62	0.70	0.26 \pm 0.12
J013204.6+122236	0.11 \pm 0.31	–1.38	0.11 \pm 0.04
J153503.4+455056	5.88 \pm 0.22	0.46	0.29 \pm 0.06
J163831.7+534020	$\gtrsim 7$	–0.95	–0.04 \pm 0.05
J163030.2+470125	1.10 \pm 0.33	1.21	0.11 \pm 0.43
J021939.9+361819	0.52 \pm 0.37	–1.08	0.15 \pm 0.08
J161001.2+330121	6.71 \pm 0.26	2.60	0.10 \pm 0.04
J171423.6+085236	3.29 \pm 0.42	0.24	–0.05 \pm 0.04
J071310.6+725627	4.61 \pm 0.27	1.40	–0.01 \pm 0.12
J095928.6+643023	$\gtrsim 0$	3.74	–
J091747.6+524821	0.08 \pm 0.55	1.20	0.44 \pm 0.33
J133053.3+734824	12.11 \pm 0.46	2.73	0.16 \pm 0.08
J144738.4+671821	3.12 \pm 0.24	–1.91	0.12 \pm 0.03

Note. Equivalent widths are given in units of Å. The *WISE* colour W1–W2 is given in the Vega system.

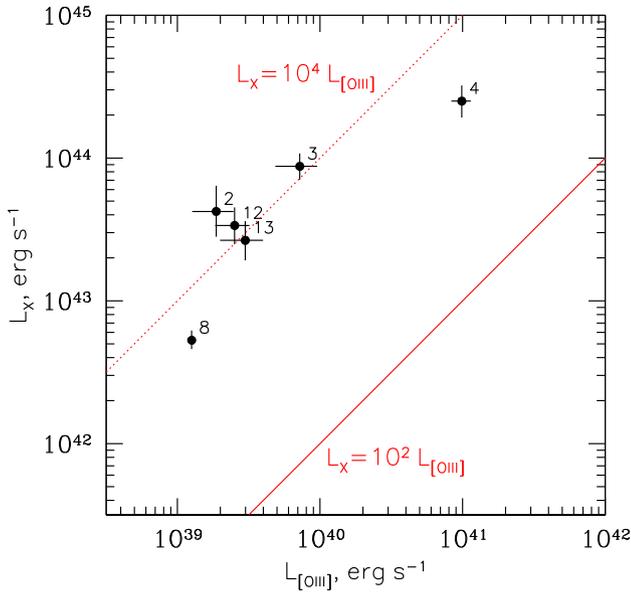


Figure 8. X-ray (0.2–6 keV) luminosities of the *SRG* TDEs (labelled by their internal numbers in Table 1) versus the [O III] $\lambda 5007$ luminosity of their host galaxies. Only objects exhibiting a significant [O III] $\lambda 5007$ emission line in the spectrum are shown. The solid line shows a typical relation between L_X and $L_{[OIII]}$ in AGNs (Heckman et al. 2005), while the dotted line indicates a factor of 100 larger $L_X/L_{[OIII]}$ ratio.

GALEX (Martin et al. 2005; Million et al. 2016), PS2, SDSS, the Two Micron All-Sky Survey (2MASS; Skrutskie et al. 2006), and the Near-Earth Object *WISE* Reactivation Mission (NEOWISE; Mainzer et al. 2014).

Briefly speaking, we employ a simple flexible stellar population synthesis (FSPS; Conroy & Wechsler 2009) model to fit the UV–MIR broad-band SEDs. The five free parameters are the total stellar mass (M_{gal}), the e-folding time of the star formation history (τ_{st}), the age of the stellar population, the metallicity (Z), and the Calzetti et al. (2000) dust model optical depth. We measure the Galactic extinction-corrected, rest-frame $u - r$ colour ($^{0.0}u - r$) from the synthetic photometry. The broad-band SEDs and fitted parameters are given in Appendix C.

Fig. 11 shows $^{0.0}u - r$ versus M_{gal} for the host galaxies of the TDEs discovered by *SRG/eROSITA* and discussed in this work, together with a comparison sample of $\approx 650\,000$ galaxies from SDSS (Mendel et al. 2014). For comparison, we also show a sample of optically selected TDEs (Hammerstein et al. 2021; van Velzen et al. 2021b). The majority of *SRG* TDE hosts reside in the green valley. This result is expected since previous studies on optically selected TDEs also show that green galaxies are overrepresented in TDE hosts (Law-Smith et al. 2017; Hammerstein et al. 2021). We note, however, that *SRG* TDEs are systematically hosted by galaxies with greater M_{gal} .

Moreover, it has been found that optically selected TDEs are preferentially hosted by rare quiescent post-starburst (QBS) galaxies or E + A galaxies characterized by strong H δ absorption and weak H α emission (French, Arcavi & Zabludoff 2016; Law-Smith et al. 2017; Hammerstein et al. 2021). In order to see if QBS galaxies are also overrepresented in X-ray TDEs, we show the distribution of *SRG* selected TDEs and optically selected TDEs on the H δ_A absorption index versus EW(H α_{em}) diagram in Fig. 12. In addition, we collect a sample of $\approx 516\,000$ SDSS galaxies from the MPA + JHU catalogs (Brinchmann et al. 2004) by requiring the H α EW error $0 < H_ALPHA_EQW_ERR < 4$, the H α continuum > 0 , the redshift $z > 0.01$, the median signal-to-noise ratio per pixel of the whole spectrum > 10 , and entries in the Mendel et al. (2014) catalogue. We construct two comparison samples with the distribution of M_{gal} similar to the *SRG* selected TDE hosts and the ZTF TDE hosts (see the blue and yellow contours in Fig. 12).

Fig. 12 shows that 0–8 per cent of *SRG* TDEs are hosted by QBS galaxies. In contrast, 42 per cent of ZTF TDEs are hosted by QBS galaxies. Compared with ZTF TDE hosts, *SRG* TDE hosts typically exhibit lower values of H δ_A , and there appears to be a dearth of host galaxies with strong H α emission. These differences might be partially accounted for by the greater M_{gal} of the *SRG* TDEs, as evidenced by the distributions shown in Fig. 12.

5.3 TDE occurrence rate

Although the first *SRG* sample of TDEs presented in this work is fairly small, comprising 13 objects, it is already well suited for drawing some inferences about the statistical properties of TDEs in the $z \approx 0.6$ Universe.

The detectability of a TDE in the *SRG/eROSITA* all-sky survey largely depends on the following physical parameters: z – TDE redshift, $L_{X,max}$ – X-ray luminosity (in the rest-frame 0.2–6 keV energy band) at the peak of TDE brightness, $t_{det} - t_0$ – time passed between TDE onset and its detection by *eROSITA*, T_{in} – characteristic temperature of the TDE X-ray spectrum assuming multiblackbody accretion disc emission, b_{ecl} – ecliptic latitude (since the *SRG* all-survey sensitivity increases towards the ecliptic poles), and $N_{H,Gal}$ – Galactic absorption in the TDE direction.

Since *SRG* passes each location in the sky every 6 months, only a fragmentary X-ray light curve can be obtained for a given TDE based on *eROSITA* data. Nevertheless, as we have seen in Section 3.2, the X-ray light curves obtained in this work suggest that most of the TDEs have been caught by *eROSITA* within 2 months after their onset. In addition, our X-ray spectral modelling indicates that during the discovery of these TDEs their Eddington ratios were all near or greater than 100 per cent (see Table 4). These facts together suggest that the X-ray luminosities measured by *eROSITA* during *eRASS2* should be close to the peak X-ray luminosities of these TDEs (i.e. $L_X \approx L_{X,max}$). We note in passing that it is hardly possible to predict

Table 11. TDE properties.

Object (SRGE)	Optical position		$N_{\text{H,Gal}}$ 10^{20} cm^{-2}	z	L_{X}^a ($10^{43} \text{ erg s}^{-1}$)		L_g^b $10^{43} \text{ erg s}^{-1}$	V_{max} Gpc^3
	RA	Dec.			eRASS2	eRASS3		
J135514.8+311605	208.812579	31.268121	1.21	0.1989 ± 0.0004	$5.8_{-1.4}^{+1.6}$	<0.3	<0.4	1.90
J013204.6+122236	23.018675	12.376562	4.80	0.132 ± 0.001	$4.2_{-1.4}^{+2.2}$	<0.4	<0.22	0.439
J153503.4+455056	233.763172	45.848598	1.22	0.2314 ± 0.0004	$8.8_{-1.7}^{+2.1}$	$3.1_{-0.9}^{+1.6}$	0.9–2.2	1.87
J163831.7+534020	249.633401	53.672931	2.93	0.581 ± 0.001	25_{-6}^{+7}	<1.8	<3.5	10.6
J163030.2+470125	247.626052	47.023730	1.62	0.294 ± 0.001	20_{-4}^{+5}	<1.0	0.9–2.4	4.16
J021939.9+361819	34.916264	36.305054	5.26	0.3879 ± 0.0002	25_{-8}^{+12}	15_{-3}^{+3}	<2.3	8.43
J161001.2+330121	242.505920	33.022416	1.56	0.1309 ± 0.0006	$1.2_{-0.3}^{+0.4}$	<0.15	<0.2	0.235
J171423.6+085236	258.598393	8.876918	5.39	0.036 ± 0.001	$0.53_{-0.07}^{+0.09}$	$0.03_{-0.01}^{+0.01}$	<0.07	0.0965
J071310.6+725627	108.293835	72.940751	3.53	0.104 ± 0.001	11_{-2}^{+3}	$1.9_{-0.8}^{+1.6}$	<0.15	1.06
J095928.6+643023	149.868660	64.506053	3.07	0.454 ± 0.001	89_{-32}^{+49}	87_{-35}^{+65}	1.2–2.6	13.8
J091747.6+524821	139.447492	52.805635	1.52	0.1876 ± 0.0002	48_{-16}^{+23}	<5	0.5–1.2	0.747
J133053.3+734824	202.720918	73.806739	1.59	0.1505 ± 0.0002	$3.4_{-0.9}^{+1.2}$	<0.2	<0.24	0.477
J144738.4+671821	221.912771	67.305094	0.88	0.1250 ± 0.0005	$2.7_{-0.7}^{+1.0}$	$3.2_{-0.6}^{+0.6}$	<0.6	0.179

Notes. ^aLuminosity in the rest-frame 0.2–6 keV energy range, corrected for Galactic absorption.

^bRest-frame g -band luminosity.

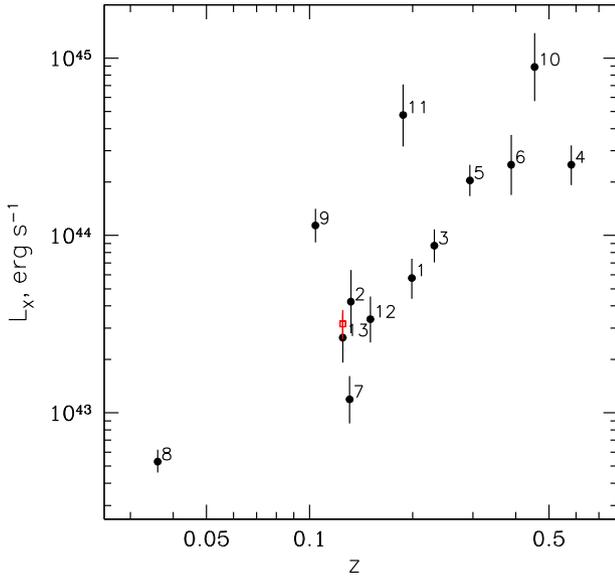


Figure 9. Intrinsic luminosity in the 0.2–6 keV energy band as a function of redshift for SRG TDEs (labelled by their number in Table 1). The luminosities are based on the eRASS2 measurements. For SRGE J144738.4+671821, we also show its X-ray luminosity during eRASS3 (red square), when it became brighter than in eRASS2.

the $t_{\text{det}} - t_0$ delay based on existing theoretical models, since they have too many parameters (see Rossi et al. 2021 for a recent review).

The detectability of TDEs in the SRG survey is strongly affected by their X-ray spectral hardness/softness. As we have seen in Section 3.3, the analysed eROSITA spectra can be described fairly well in terms of multiblackbody accretion disc emission, with their shape characterized by the single parameter T_{in} . Fig. 13 (upper panel) shows the expected eROSITA count rate in the 0.3–2.2 keV energy range (actually used for TDE detection) as a function of T_{in} for a fixed intrinsic flux of $10^{-12} \text{ erg s}^{-1} \text{ cm}^{-2}$ in the rest-frame 0.2–6 keV energy range. For nearby TDEs ($z \approx 0$), the sensitivity in the directions with low Galactic absorption does not vary by more

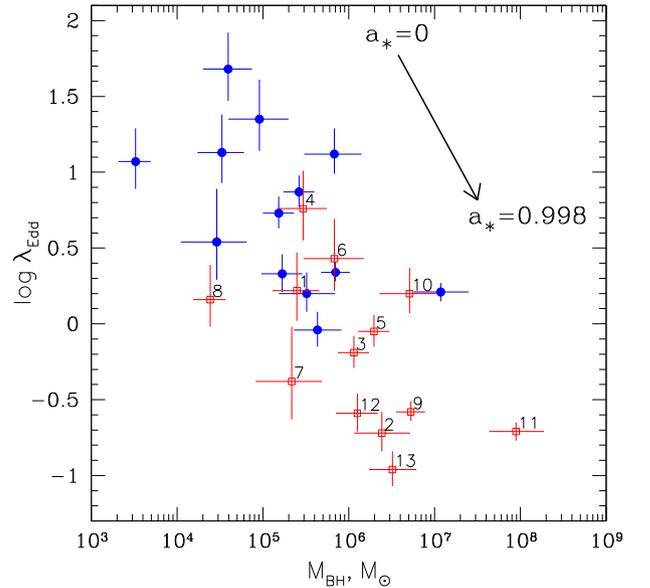


Figure 10. Black hole masses versus Eddington ratios for the SRG TDEs (labelled by the internal numbers in Table 1), estimated from the X-ray spectra assuming a standard accretion disc around a SMBH with $a_* = 0$ (blue points) and $a_* = 0.998$ (red points). The arrow illustrates the range of $(M_{\text{BH}}, \lambda_{\text{Edd}})$ values allowed for a given TDE depending on a_* .

than a factor of 2 for $0.1 \text{ keV} < kT_{\text{in}} < 1 \text{ keV}$, while it decreases dramatically below 0.1 keV. This effect is further strengthened if there is a significant column of cold gas ($N_{\text{H}} \gtrsim 5 \times 10^{20} \text{ cm}^{-2}$) in the direction of the source, and with redshift.

All but one (SRGE J091747.6+524821, with $kT_{\text{in}} \approx 0.05 \text{ keV}$) TDEs in our sample have $0.1 \text{ keV} \lesssim kT_{\text{in}} \lesssim 0.5 \text{ keV}$, i.e. fall within the temperature range favourable for detection at $z < 0.6$ (we recall that the most distant TDE in our sample is located at $z = 0.58$). While it is unlikely that we are missing TDEs with very hard spectra ($kT_{\text{in}} \gtrsim 1 \text{ keV}$), there is certainly a strong selection effect against TDEs with $kT_{\text{in}} \lesssim 0.1 \text{ keV}$. Within the standard TDE paradigm of nearly critical accretion on to an SMBH, such low temperatures correspond to black

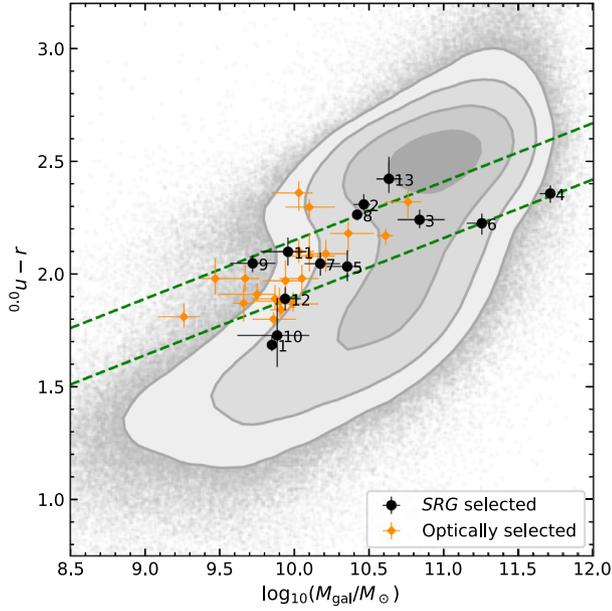


Figure 11. The Galactic extinction-corrected, synthetic rest-frame $u - r$ colour of TDE host galaxies. The 13 black data points are from this work (labelled by the internal numbers in Table 1). The 20 orange data points include 17 TDEs selected from the first 1.5 yr of ZTF (van Velzen et al. 2021b) and three events reported by Hammerstein et al. (2021) afterwards. The green valley is denoted by the dashed green lines, following the boundaries used by van Velzen et al. (2021b). The contours enclose a comparison sample of galaxies from SDSS DR7 (Mendel et al. 2014).

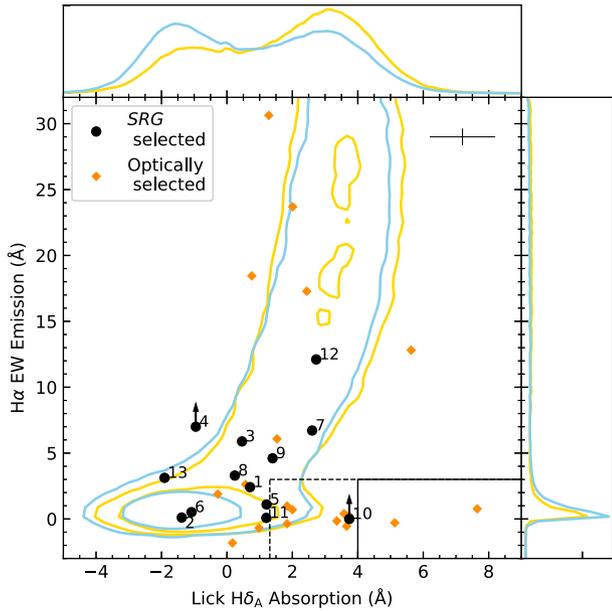


Figure 12. The $H\delta_A$ absorption index versus the $H\alpha$ EW emission of TDE host galaxies. The 13 black data points are from this work (labelled by their internal numbers in Table 1). The 19 orange data points are from Hammerstein et al. (2021). Median uncertainties are shown in the top right. Following the definition in Hammerstein et al. (2021), the solid black line marks the region of E + A galaxies ($H\delta_A - \sigma(H\delta_A) > 4.0$, $EW(H\alpha_{em}) < 3.0$), and the dashed black line marks the region of QBS galaxies ($H\delta_A > 1.31$, $EW(H\alpha_{em}) < 3.0$). We show two comparison samples with the distribution of M_{gal} similar to SRG selected TDE hosts (blue lines) and optically selected TDE hosts (yellow lines). For each comparison sample, the contours enclose 30 per cent and 85 per cent of the galaxies.

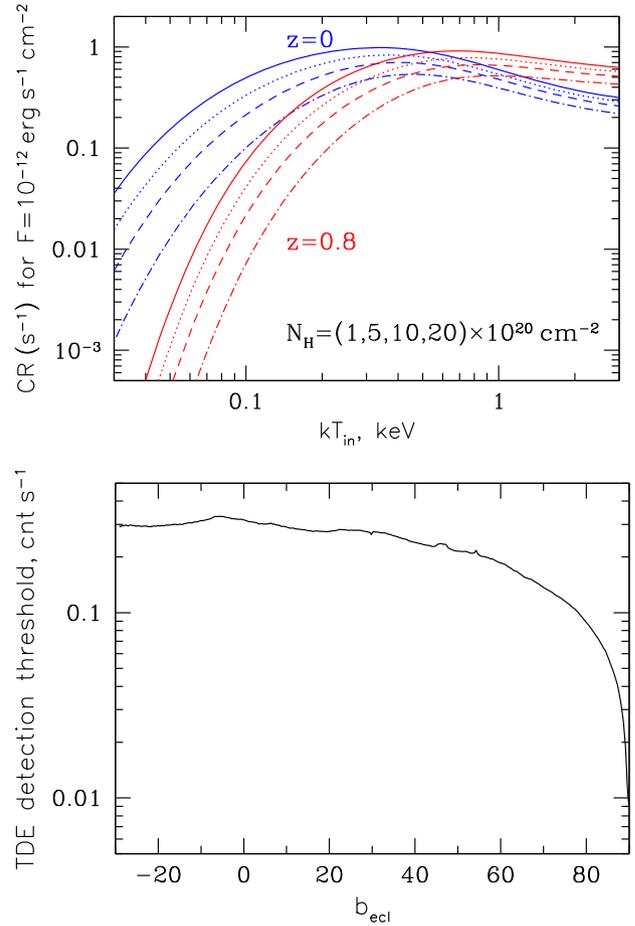


Figure 13. *Top:* eROSITA count rate in the 0.3–2.2 keV energy range as a function of the temperature of the multicolour accretion disc, T_{in} , for an intrinsic flux of 10^{-12} erg s^{-1} cm^{-2} in the 0.2–6 keV energy range, for different absorption columns (curves from top to bottom, $N_H = 10^{20}, 5 \times 10^{20}, 10^{21},$ and 2×10^{21} cm^{-2}). *Bottom:* Minimum count rate in the 0.3–2.2 keV energy band required for TDE detection as a function of ecliptic latitude, averaged over the $0^\circ < l < 180^\circ$ hemisphere.

holes with $M_{BH} \gtrsim 10^7 M_\odot$. Therefore, the current SRG/eROSITA TDE sample is well suited for estimating the rate of TDEs near SMBHs with $\lesssim 10^7 M_\odot$ in the $z \lesssim 0.6$ Universe, while our working energy range (0.3–2.2 keV) is too hard for systematic exploration of softer TDEs, presumably associated with more massive black holes.

We next need to know the minimum count rate, CR_{min} , required for detection of TDEs in the 0.3–2.2 keV band during eRASS2 according to the criteria outlined at the beginning of Section 2, namely that a candidate transient must be at least 10 times brighter during eRASS2 compared to the upper limit on its flux during eRASS1. This threshold depends primarily on the ecliptic latitude, since the exposure per point accumulated during the SRG all-sky survey is inversely proportional to $\cos b_{ecl}$. Fig. 13 (lower panel) shows the exact behaviour of the minimum count rate on b_{ecl} ; the details of this computation will be presented elsewhere. Count rates of at least $CR_{min} \sim 0.3$ and ~ 0.2 counts s^{-1} are needed to satisfy our TDE detection criterion at $b_{ecl} = 0^\circ$ and $b_{ecl} = 60^\circ$, respectively. Note that a typical (vignetting corrected) exposure during one eROSITA all-sky survey is ~ 120 s at $b_{ecl} = 0^\circ$, and ~ 240 s at $b_{ecl} = 60^\circ$, so that the quoted thresholds correspond to ~ 40 and ~ 50 counts for $b_{ecl} = 0^\circ$ and $b_{ecl} = 60^\circ$, respectively.

5.3.1 X-ray luminosity function

Having discussed the key properties of the SRG/eROSITA TDE sample and possible selection biases associated with it, we now proceed to evaluation of the TDE X-ray luminosity function (XLF). To this end, we may use the classical $1/V_{\max}$ method (Schmidt 1968).

The maximum observable volume V_{\max} for a given TDE in the studied sample depends on its intrinsic luminosity in the 0.2–6 keV energy band (L_X) and its intrinsic spectral shape, characterized by T_{in} . Given these quantities, the computation of V_{\max} is straightforward using the tabulated dependences CR/F (T_{in} , N_{H} , z) and CR_{min} (b_{ecl}) discussed above. Specifically, we should integrate the comoving volume over the $0^\circ < l < 180^\circ$ hemisphere (probed by this study) out to a luminosity distance D_{max} (or equivalently a redshift z_{max}) defined for each position in the sky so that $L_X/(4\pi D_{\text{max}}^2) \times CR/F(T_{\text{in}}, N_{\text{H,Gal}}, z_{\text{max}}) = CR_{\text{min}}(b_{\text{ecl}})$. The Galactic absorption column $N_{\text{H,Gal}}$ entering this equation varies across the sky, and we adopt it from (HI4PI Collaboration 2016).

The computed V_{\max} values for the studied TDEs are given in the last column of Table 11. We have not tried to estimate the corresponding uncertainties associated with the X-ray spectral analysis (i.e. the uncertainties in L_X and T_{in}) since these are less important compared to the scatter in the $1/V_{\max}$ values of individual objects used in the calculation of the XLF. To calculate the TDE XLF, we just need to sum the derived $1/V_{\max}$ values of individual TDEs in specified luminosity bins. To this end, we use five equal intervals in $\log L_X$ between 42.5 and 45. The uncertainty within a given bin is found as $\sqrt{\sum (1/V_{\max})_i^2}$, where the summation is done over the objects within that bin.

We next have to take into account that TDEs are transients. As discussed above, eROSITA appears to typically discover TDEs within 2 months after their onset. Adopting as a fiducial value $t_{\text{det}} - t_0 = 2$ months (ignoring subtleties associated with cosmological time dilation) and taking into account that the studied TDE sample was collected over a period of 6 months, we should multiply our $1/V_{\max}$ based estimates by $6/2 \times 12/6 = 6$ to evaluate the volume rate of TDEs per year.

We finally note that the current sample of 13 TDEs is not statistically complete. Indeed, as already mentioned, there are three additional TDEs discovered during eRASS2 based on the same criteria, which are discussed elsewhere (Gilfanov et al., in preparation) because of their pronounced activity in the optical band. To a first approximation, we can take this into account by multiplying the XLF by a factor of 16/13.

The computed TDE XLF is shown in Fig. 14. The largest uncertainty is associated with the lowest luminosity bin, $10^{42.5} \text{ erg s}^{-1} < \log L_X < 10^{43} \text{ erg s}^{-1}$, where there is just one, nearby TDE (SRGE J171423.6+085236 at $z = 0.036$) in the current SRG/eROSITA sample. There is a clear indication of the TDE volume rate declining with increasing X-ray luminosity. Fitting the XLF (using χ^2 minimization) by a power law, we obtain (see Fig. 14):

$$\frac{dN_{\text{TDE}}}{d \log L_X dV dt} = (1.4 \pm 0.8) \times 10^{-7} \left(\frac{L_X}{10^{43} \text{ erg s}^{-1}} \right)^{-0.6 \pm 0.2} \text{ Mpc}^{-3} \text{ yr}^{-1}. \quad (3)$$

If we exclude the $10^{42.5} \text{ erg s}^{-1} < \log L_X < 10^{43} \text{ erg s}^{-1}$ bin from this calculation, the slope of the power law becomes somewhat steeper, $\alpha = -0.8 \pm 0.3$, but consistent with the slope in equation (3).

According to theoretical predictions, TDEs should occur more frequently in lower mass galaxies with correspondingly lighter central black holes (Magorrian & Tremaine 1999; Wang & Merritt 2004; Kesden 2012; see Stone et al. 2020 for a recent review). This, together with the higher abundance of low-mass galaxies compared to high-mass ones, the approximately Eddington luminosities of TDEs,

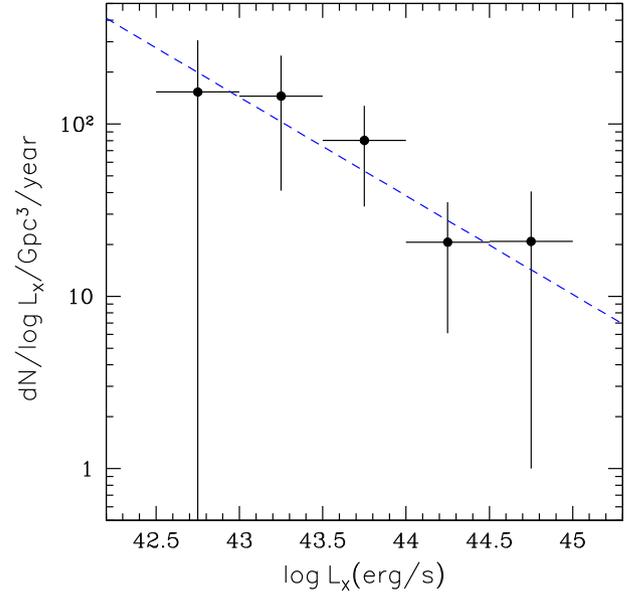


Figure 14. TDE X-ray (0.2–6 keV) luminosity function. The blue dashed line shows its best fit by a power law.

and the expectation that TDEs associated with smaller black holes should emit a larger fraction of their bolometric luminosity in X-rays (due to their hotter accretion discs), suggests that the X-ray TDE rate should increase within decreasing L_X . Apparently, we are starting to see this trend in the data of the SRG all-sky survey.

We note in this connection that the nearest and least X-ray luminous TDE in our sample, SRGE J171423.6+085236, has the hardest X-ray spectrum ($T_{\text{in}} \sim 0.5 \text{ keV}$). This suggests that it is also associated with the smallest black hole in the sample, with $M_{\text{BH}} \sim 10^4 M_{\odot}$, albeit with a large uncertainty in this estimate based on our simplistic X-ray spectral modelling (see Table 4). On the other hand, its host galaxy is fairly massive, with stellar mass $M_{\text{gal}} \approx 1.7 \times 10^{10} M_{\odot}$ (see Table C1).

The declining trend of the TDE rate with increasing X-ray luminosity inferred here is similar to the trend noticed before for optically UV selected TDEs. In that case, the measured slope of the luminosity function is $\alpha = -1.3 \pm 0.3$ (van Velzen 2018), somewhat steeper than inferred here for X-ray selected TDEs. However, this difference is only marginally significant.

5.3.2 Total rate

By integrating the XLF over the entire luminosity range probed here ($10^{42.5} \text{ erg s}^{-1} < \log L_X < 10^{45} \text{ erg s}^{-1}$), we can estimate the average TDE volumetric rate in the $z = 0-0.6$ Universe: $(2.1 \pm 1.0) \times 10^{-7} \text{ Mpc}^{-3} \text{ year}^{-1}$. Given the total galaxy volume density of $\sim 2 \times 10^{-2} \text{ Mpc}^{-3}$ (Bell et al. 2003), this translates to a rate of $R = (1.1 \pm 0.5) \times 10^{-5}$ TDEs per galaxy.

This estimated specific TDE rate is consistent with the earliest estimate of $R \sim 9 \times 10^{-6} \text{ yr}^{-1}$ per galaxy, based on three TDEs detected during the ROSAT all-sky survey and followed-up during ROSAT pointed observations (Donley et al. 2002). Another published estimate is much higher, $R \sim 2 \times 10^{-4} \text{ yr}^{-1}$ per galaxy (Esquej et al. 2008), but it is based on just two TDEs discovered during the XMM-Newton Slew Survey (Saxton et al. 2008) and undetected previously during the ROSAT all-sky survey. Yet another estimate was presented by Khabibullin & Sazonov (2014), who used a sample

of four candidate TDEs selected based on their brightness during the *ROSAT* all-sky survey and non-detection in subsequent *XMM-Newton* pointed observations: $R \sim 3 \times 10^{-5} \text{ yr}^{-1}$ per galaxy. In summary, all these previous estimates were based on extremely small samples of X-ray selected TDEs, and within their large (and poorly defined) uncertainties appear to be in line with the new result obtained here based on the first *SRG* TDE sample.

For comparison, the existing estimates of the total rate of optically UV selected TDEs are $\sim 10^{-4} \text{ yr}^{-1}$ per galaxy (van Velzen et al. 2020), which is significantly larger than our X-ray-based estimate, and is in better agreement with theoretical ‘loss cone’ predictions (Stone & Metzger 2016). This possibly indicates that the XLF shown in Fig. 14 continues to rise below $L_X \sim 10^{42.5} \text{ erg s}^{-1}$. However, the optical TDE rate, currently determined for $L_g > 10^{42.5} \text{ erg s}^{-1}$, may also increase if less luminous TDEs are included.

The lower volumetric rate of X-ray TDEs compared to that of optical TDEs may also imply that X-ray bright TDEs constitute a minority of all TDEs. The latter possibility would provide support to TDE models (e.g. Dai et al. 2018; Curd & Narayan 2019) that predict a strong dependence of the optical/X-ray brightness ratio on the viewing angle: namely, that we can only observe TDEs (during their luminous early phase) in X-rays from directions close to the axis of a thick accretion disc (formed from the debris of the disrupted star) with a powerful wind, while at larger inclination angles one can only observe the reprocessed optical-UV emission. In this connection, we also note that we have found no evidence of intrinsic absorption in the X-ray spectra of the 13 studied TDEs. However, since not only the viewing direction but also the black hole mass is expected to play a key role in diversity of TDE types (Mummery 2021), a more elaborate comparison between the TDE population properties inferred from this work and from optical studies should be done. In addition, slow or inefficient circularization of the stellar debris (Piran et al. 2015; Shiokawa et al. 2015) could also lead to a paucity of X-ray TDEs.

In future work, we plan to lower our X-ray detection threshold by a factor of ~ 2 , which should increase the TDE discovery rate by eROSITA by a factor of ~ 3 (assuming that the TDE XLF continues to grow to lower luminosities, so that the resulting eROSITA sample is dominated by low-redshift events). Given that our current sample of bright TDEs detected on one half of the sky over a period of half a year comprises 16 confirmed events, this implies that ~ 200 TDEs per year can be detected by eROSITA over the whole sky and a total of ~ 700 TDEs can be discovered by the end of the 4-yr *SRG* survey (excluding from this calculation the first of the eight planned sky scans). These numbers are consistent within an order of magnitude with the predictions done for the *SRG*/eROSITA all-sky survey prior to its beginning (Khabibullin et al. 2014; Jonker et al. 2020), which, however, strongly depend on the distribution of spins of the black holes associated with TDEs. Continued search during the *SRG* survey will allow us to further narrow down the rate of stellar disruptions in galactic nuclei and place tighter constraints on the underlying population of black holes.

5.4 Optical faintness

A salient feature of the TDEs discovered by *SRG*/eROSITA and studied here is their optical faintness. Fig. 15 shows the inferred optical versus X-ray luminosities (from Table 11) for our sample. For the four TDEs that have shown noticeable optical flares in addition to X-ray activity, the L_g/L_X ratio is constrained between 0.01 and 0.3. For the remaining TDEs, $L_g/L_X < 0.3$ and for most of them $L_g/L_X < 0.1$.

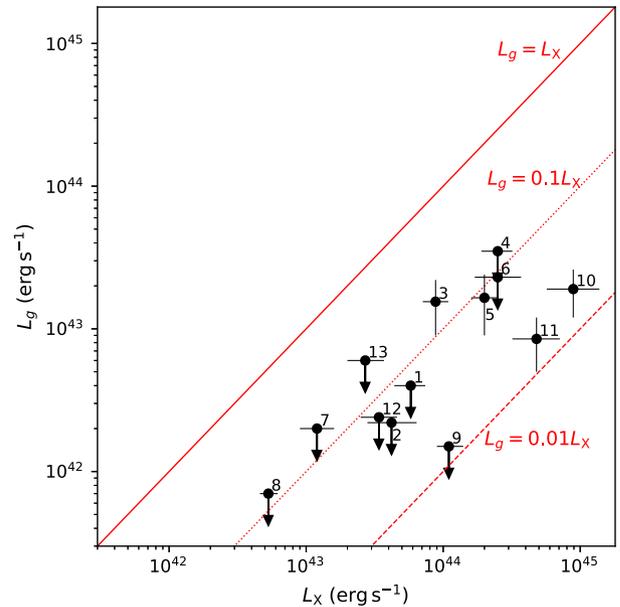


Figure 15. Rest-frame g-band luminosity versus rest-frame 0.2–6 keV luminosity for the *SRG* TDEs (labelled by the internal numbers in Table 1).

These constraints are conservative. Indeed, the possibility that the peak of TDE optical emission has been missed is already taken into account in our L_g estimates (see Section 4.2). On the other hand, as discussed before, our L_X estimates should be considered lower limits on the maximum X-ray luminosity. Therefore, the L_g/L_X ratios can actually be smaller than shown in Fig. 15.

The *SRG* TDEs studied here are similar to the very first TDEs that came to light thanks to the *ROSAT* X-ray observatory three decades ago, but drastically different from those discovered recently in the optical/UV, for which typically $L_g > L_X$ (e.g. Gezari 2021). Although it is certainly too early to draw firm conclusions on the underlying physical picture, it is possible that we are dealing here with the same orientation effect that was discussed above in relation to the observed TDE volumetric rate. Namely, that most of the TDEs detected by *SRG* are observed from small angles with respect to the accretion disc axis, whereas optical-UV surveys catch TDEs from quasi-random inclination angles, so that most of them prove to be X-ray faint.

We note, however, that apart from optically faint TDEs *SRG*/eROSITA has also discovered TDEs with prominent optical activity (Gilfanov et al. 2020, 2021). Several first events of this kind will be discussed by Gilfanov et al. (in preparation). Postponing further discussion of the relationship between optically bright and optically faint X-ray selected TDEs to that paper, we note that the former appear to constitute ~ 20 per cent of TDEs found during the *SRG* all-sky survey.

6 CONCLUSIONS

We have presented the first sample of TDEs discovered during the ongoing *SRG* all-sky survey. These 13 events were selected among the multitude of X-ray transients detected by eROSITA during its second scan of the sky (2020 June 10 to December 14) at a level exceeding at least tenfold the upper limit on the flux in the first scan and confirmed as TDEs by our optical follow-up observations. The most distant of these events (*SRG* J163831.7+534020) occurred at $z = 0.581$. Therefore, the *SRG* survey has already expanded the

horizon of TDE X-ray detectability by a factor of ~ 4 compared to the *ROSAT* all-sky survey, conducted 30 yr ago.

The properties of these events are similar to those of TDEs detected (in small numbers) by previous X-ray missions. Namely, the X-ray light curves, currently consisting of a flux measured by *eROSITA* during the second sky survey and a second measurement/upper limit obtained 6 months later, are in most cases consistent with a $t^{-5/3}$ decline that started shortly before the *eROSITA* discovery. Particularly interesting is the TDE SRGE J144738.4+671821, which continued to brighten after its discovery for at least another 6 months.

The early (*eRASS2*) X-ray spectra of these TDEs can be described by a multiblackbody accretion disc model with kT_{in} between ≈ 0.05 and 0.5 keV, which is consistent with nearly critical accretion on to black holes with masses between a few $\times 10^3$ and $\sim 10^8 M_{\odot}$ depending on their (unknown) spins. In reality, we may be dealing with supercritical accretion in most of these events. No evidence of intrinsic absorption is found in the X-ray spectra. Four TDEs remained sufficiently bright in *eRASS3* to allow us to analyse their spectra taken by *eROSITA* at this later epoch. In two cases, we observe a clear spectral hardening, possibly indicating the formation of an accretion disc corona.

Four TDEs have shown a brightening in their optical light curves, concurring with or preceding the X-ray outburst. One of these events (SRGE J163030.2+470125) also exhibited a blue excess in an optical spectrum taken two months after the *eROSITA* discovery, which disappeared in a spectrum obtained eight months later. The other nine TDEs show no signs of optical activity associated with stellar tidal disruption in existing photometric and spectroscopic data.

All of the TDEs (including those with optical flares) are optically faint in the sense that the estimated optical/X-ray luminosity ratio is less than 0.3 and in most cases $L_g/L_X < 0.1$. In this respect, this sample is drastically different from TDEs selected at optical-UV wavelengths, which typically have $L_g/L_X > 1$. However, apart from the events presented in this work, *SRG/eROSITA* has also discovered a few TDEs that demonstrate prominent activity in the optical-UV band (Gilfanov et al., in preparation).

The *SRG* TDEs are mostly hosted by galaxies located in the green valley on the colour–stellar mass diagram, similarly to optically selected TDEs. However, no excess of quiescent post-starburst galaxies or E + A galaxies is observed among the hosts of the *SRG* TDEs, unlike in some previous studies of optically selected TDEs. This may be partially accounted for by the fact that the host galaxies of the *SRG* TDEs are more massive on average.

We have constructed the X-ray luminosity function in the range from $10^{42.5}$ to 10^{45} erg s $^{-1}$, the first of its kind to our knowledge. It clearly shows that the TDE volumetric rate decreases with increasing X-ray luminosity. This trend can be described by a power law with a slope of $\alpha = -0.6 \pm 0.2$ ($dN/d \log L_X \propto L_X^{\alpha}$), or $\alpha = -0.8 \pm 0.3$ if we exclude from consideration the lowest luminosity bin of $10^{42.5}$ erg s $^{-1} < \log L_X < 10^{43}$ erg s $^{-1}$, which contains just one TDE. This behaviour is similar to the trend previously observed for optically selected TDEs, but in that case the decline is marginally steeper, with $\alpha = -1.3 \pm 0.3$.

The total rate of X-ray TDEs is estimated at $(1.1 \pm 0.5) \times 10^{-5}$ events per galaxy per year. This is an order of magnitude lower than recent estimates of optical TDEs, which possibly indicates that the TDE XLF continues to rise below $L_X \sim 10^{42.5}$ erg s $^{-1}$ and/or that X-ray bright events constitute a minority of TDEs. The latter possibility would provide support to the TDE models that predict a strong dependence on the viewing angle, namely, that TDEs can only be observed in X-rays from directions close to the axis of a thick accretion disc formed from the debris of the disrupted star.

The *SRG* all-sky survey is to continue until the end of 2023, opening up exciting prospects for TDE studies. In particular, we plan to lower our detection threshold for such events by a factor of ~ 2 , which should increase the *eROSITA* TDE discovery rate by a factor of ~ 3 . This implies that ~ 700 TDEs can be found by the end of the 4-yr *SRG* survey over the entire sky. Furthermore, regular (every 6 months) visits by *eROSITA* of previously discovered TDEs will provide valuable information on their long-term X-ray evolution.

To conclude, a continued search for TDEs during the *SRG* all-sky survey should allow us to narrow down the rate of stellar disruptions in galactic nuclei, tighten constraints on the properties of the underlying population of supermassive and, perhaps, intermediate-mass black holes, and shed light on the physics of near- and supercritical accretion on to such objects.

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⁵<http://ckp-rf.ru/ckp/3056/>

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DATA AVAILABILITY

X-ray data analysed in this article were used by permission of the Russian *SRG/eROSITA* consortium. The data will become publicly available as part of the corresponding *SRG/eROSITA* data release along with the appropriate calibration information. Optical data used in the article will be shared on reasonable request to the corresponding author.

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APPENDIX A: SHORT-TERM X-RAY LIGHT CURVES

During each *SRG* all-sky survey, *eROSITA* monitors any X-ray source in the sky for a period of at least 1 d. Each such visit consists of a series of ~ 40 s-long measurements taken at 4-h intervals. Therefore, apart from the long-term X-ray light curves of the TDEs presented in Fig. 1, we can also construct their short-term light curves during *eRASS2* and *eRASS3* (if a given transient remained detectable during *eRASS3*). These light curves are presented in Figs A1 and A2. They span between ~ 1 and ~ 7 d, depending on the ecliptic latitudes, b_{ecl} , of the TDEs (as a result of the *SRG* survey's strategy; Sunyaev et al. 2021).

None of the TDEs exhibit substantial X-ray variability on time-scales between 4 h and a few days. This may place interesting constraints on TDE models.

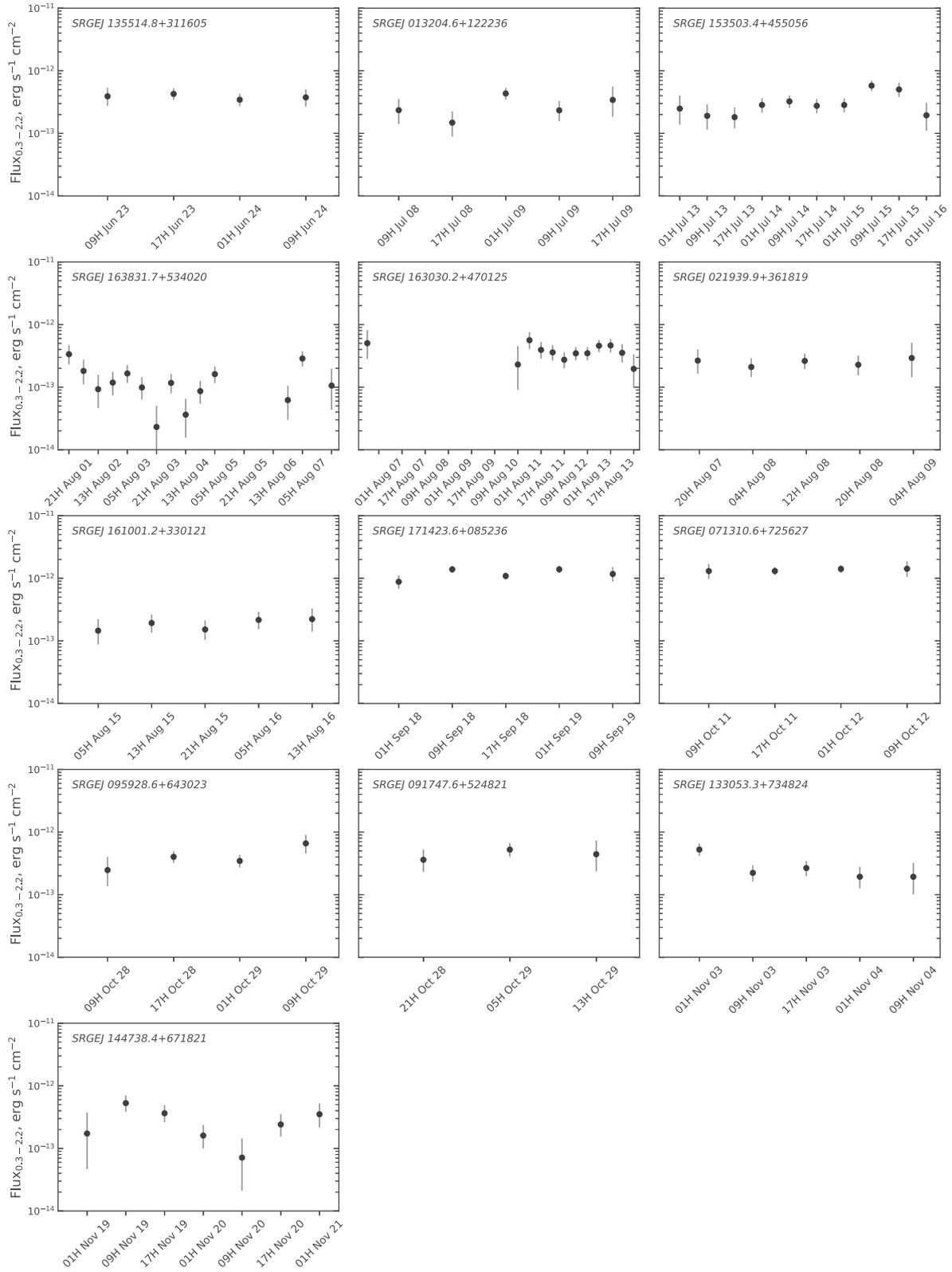


Figure A1. Short-term X-ray light curves of the TDEs obtained by eROSITA in the 0.3–2.2 keV energy range in 2020 (during eRASS2).

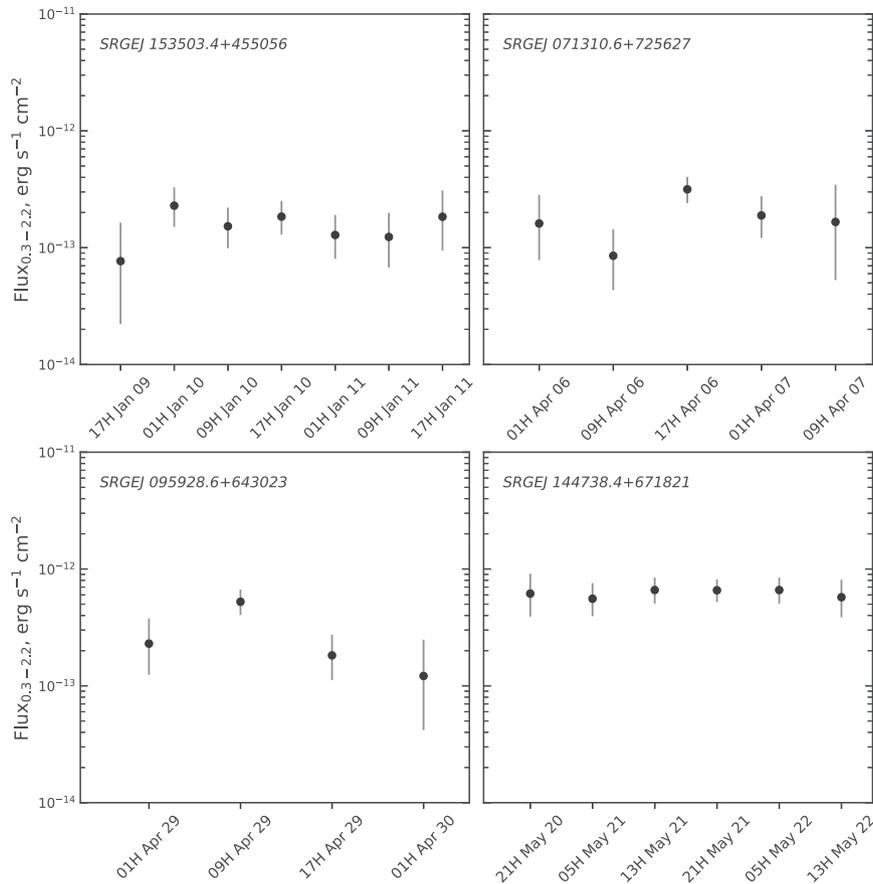


Figure A2. Short-term X-ray light curves of several TDEs obtained by eROSITA in the 0.3–2.2 keV energy range in 2021 (during eRASS3).

APPENDIX B: OPTICAL OBSERVATIONAL DETAILS AND DATA REDUCTION

All Keck-I/LRIS spectroscopic observations were conducted using the blue side grism 300/3400, the dichroic 560, the red side grating 400/8500, and a slit mask of 1 arcsec. This setup provides a spectral coverage of 3200–10 250 Å. All LRIS spectra were reduced and extracted using LPIPE (Perley 2019).

BTA/SCORPIO-2 spectroscopic observations were conducted with the VPHG940@600 grating, providing a spectral range of 3500–8500 Å. The spectrum of SRGEJ135514.8+311605 was obtained with a slit of 1 arcsec and that of SRGEJ163030.2+470125 with a slit of 2 arcsec. During both observations, the seeing was near 1.2 arcsec. SCORPIO-2 photometry was done in a 1024 × 1024 pixel imaging mode, corresponding to a binning of 2 × 2 (0.40 arcsec pixel scale).

Photometry at the CMO’s 2.5-m telescope was conducted using the NBI 4096 × 2048 × 2 CCD-photometer with a 0.155 arcsec pixel scale. Spectroscopy at the 2.5-m telescope was done using the TDS instrument. The spectra were obtained in a range of 3600–7500 Å with a 1 arcsec slit, at a spectral resolution of ∼1500. The spectrograph and data reduction are described in Potanin et al. (2020). Photometry at the CMO’s RC600 telescope was carried out using the Andor iKon-L camera with 2048 × 2048 pixels and a pixel scale of 0.67 arcsec.

AZT-33IK/ADAM spectra were obtained using a 2 arcsec slit and the VPHG600G grating (spectral range 3700–7340 Å, resolution 8.8 Å). Photometric observations at AZT-33IK were done using the

CCD photometer, consisting of a focal reducer and the Andor iKon-M 934 camera with 1024 × 1024 pixels, which provide a 6.3 arcmin × 6.3 arcmin field of view with a 0.372 arcsec pixel scale.

Photometric observations at RTT150 were conducted using the TFOSC instrument equipped with an ANDOR iKon-L DZ936N CCD 2048 × 2048 detector with a 0.327 arcsec pixel scale.

Images and spectroscopy from BTA, AZT-33IK, and RTT150 were processed using IRAF and our own software. Aperture photometry was done using the *apphot* task from the IRAF *digiphot* package. Measurement of source fluxes was done relative to nearby bright stars, with the aperture size for each observation series defined such that the total flux was obtained.

The obtained magnitudes were calibrated using the magnitudes of secondary photometric standards in the source field with a brightness comparable to or greater than the flux of the target object. We used point spread function magnitudes for the calibration stars from SDSS (Ahumada et al. 2020) if available, or from Pan-STARRS1 DR2 (Chambers et al. 2016) otherwise.

APPENDIX C: HOST GALAXY ANALYSIS

As described in Section 5.2, we fit the historical photometry of the TDE hosts following the approach used by van Velzen et al. (2021b). The FSPS (Conroy & Wechsler 2009) and Prospector (Johnson et al. 2021) packages are used to find the synthetic galaxy model that is the best description for the host galaxy photometry. Fig. C1 presents the host SEDs and the fitted models. The fitted parameters are given in Table C1.

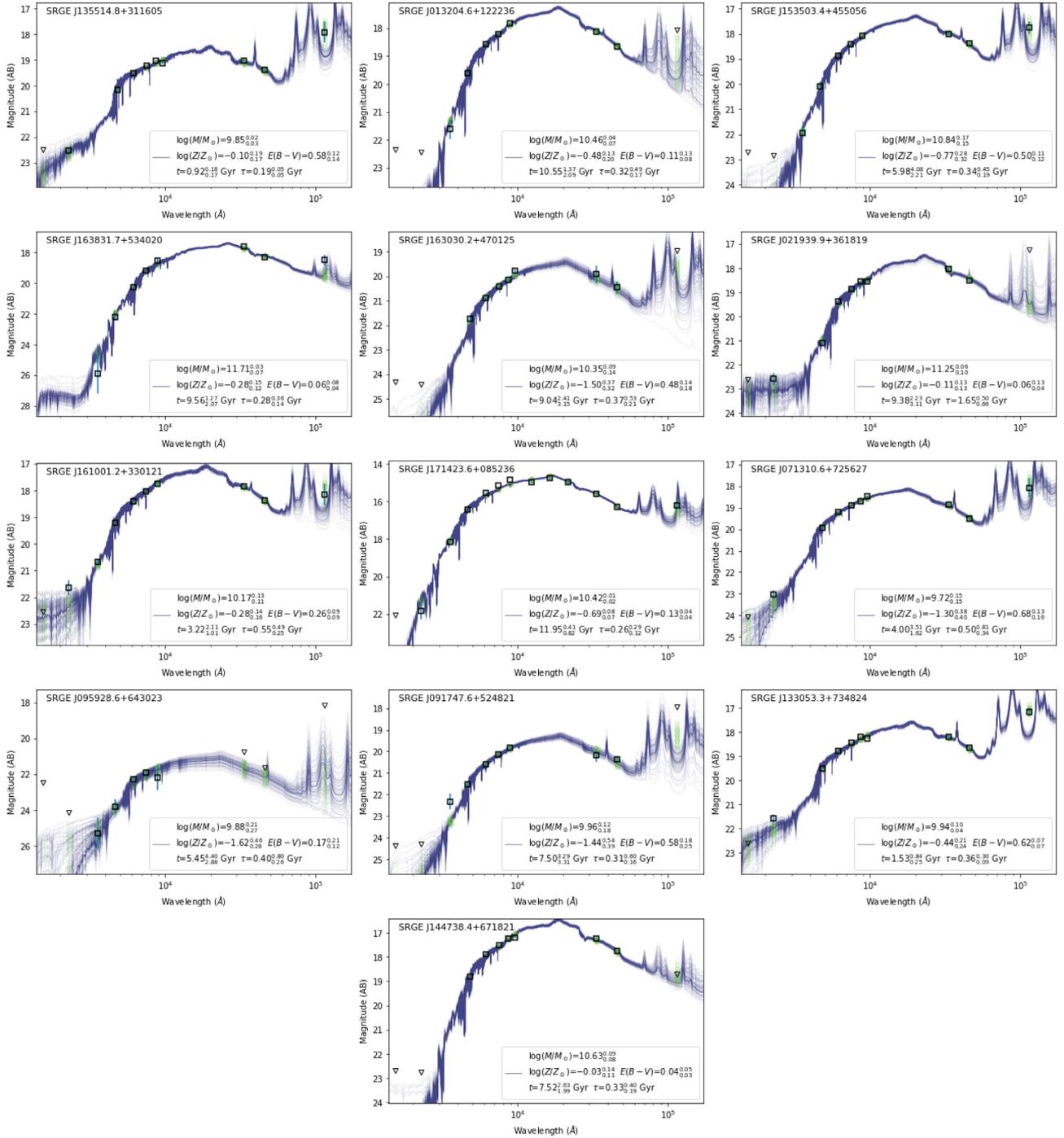


Figure C1. SED of TDE host galaxies. Squares (downward triangles) are observed photometric detections (upper limits) corrected by Galactic extinction. For each source, the purple lines show models of the 100 walkers in the Markov chain Monte Carlo (MCMC) sampler.

Table C1. Host properties inferred from SED fitting.

Object	M_{gal} $\log_{10} M_{\odot}$	${}^{0.0}u - r$	τ_{sfb} Gyr	Age Gyr	Z $\log_{10} Z_{\odot}$	Dust $E(B - V)$
SRGE J135514.8+311605	$9.85^{+0.02}_{-0.03}$	$1.69^{+0.04}_{-0.03}$	$0.19^{+0.05}_{-0.05}$	$0.92^{+0.18}_{-0.17}$	$-0.10^{+0.19}_{-0.17}$	$0.58^{+0.12}_{-0.14}$
SRGE J013204.6+122236	$10.46^{+0.04}_{-0.07}$	$2.31^{+0.05}_{-0.03}$	$0.32^{+0.49}_{-0.17}$	$10.55^{+1.37}_{-2.09}$	$-0.48^{+0.13}_{-0.20}$	$0.11^{+0.13}_{-0.08}$
SRGE J153503.4+455056	$10.84^{+0.17}_{-0.15}$	$2.24^{+0.04}_{-0.04}$	$0.34^{+0.45}_{-0.19}$	$5.98^{+4.08}_{-2.21}$	$-0.77^{+0.28}_{-0.32}$	$0.50^{+0.11}_{-0.12}$
SRGE J163831.7+534020	$11.71^{+0.03}_{-0.07}$	$2.36^{+0.04}_{-0.04}$	$0.28^{+0.36}_{-0.14}$	$9.56^{+1.27}_{-2.07}$	$-0.28^{+0.15}_{-0.12}$	$0.06^{+0.08}_{-0.04}$
SRGE J163030.2+470125	$10.35^{+0.09}_{-0.14}$	$2.03^{+0.07}_{-0.06}$	$0.37^{+0.53}_{-0.21}$	$9.04^{+2.41}_{-3.15}$	$-1.50^{+0.37}_{-0.32}$	$0.48^{+0.14}_{-0.18}$
SRGE J021939.9+361819	$11.25^{+0.06}_{-0.10}$	$2.23^{+0.04}_{-0.05}$	$1.65^{+0.50}_{-0.66}$	$9.38^{+2.23}_{-3.11}$	$-0.11^{+0.13}_{-0.13}$	$0.06^{+0.13}_{-0.04}$
SRGE J161001.2+330121	$10.17^{+0.13}_{-0.11}$	$2.05^{+0.05}_{-0.05}$	$0.55^{+0.49}_{-0.25}$	$3.22^{+2.11}_{-1.01}$	$-0.28^{+0.14}_{-0.16}$	$0.26^{+0.09}_{-0.09}$
SRGE J171423.6+085236	$10.42^{+0.01}_{-0.02}$	$2.26^{+0.02}_{-0.02}$	$0.26^{+0.29}_{-0.12}$	$11.95^{+0.41}_{-0.82}$	$-0.69^{+0.08}_{-0.07}$	$0.13^{+0.04}_{-0.04}$
SRGE J071310.6+725627	$9.72^{+0.15}_{-0.15}$	$2.05^{+0.04}_{-0.04}$	$0.50^{+0.81}_{-0.34}$	$4.00^{+3.51}_{-1.62}$	$-1.30^{+0.38}_{-0.40}$	$0.68^{+0.13}_{-0.16}$
SRGE J095928.6+643023	$9.88^{+0.21}_{-0.27}$	$1.73^{+0.17}_{-0.14}$	$0.40^{+0.80}_{-0.26}$	$5.45^{+4.40}_{-2.88}$	$-1.62^{+0.46}_{-0.28}$	$0.17^{+0.21}_{-0.12}$
SRGE J091747.6+524821	$9.96^{+0.12}_{-0.18}$	$2.10^{+0.06}_{-0.06}$	$0.31^{+0.60}_{-0.16}$	$7.50^{+3.29}_{-3.31}$	$-1.44^{+0.54}_{-0.39}$	$0.58^{+0.18}_{-0.25}$
SRGE J133053.3+734824	$9.94^{+0.10}_{-0.04}$	$1.89^{+0.05}_{-0.05}$	$0.36^{+0.30}_{-0.09}$	$1.53^{+0.84}_{-0.25}$	$-0.44^{+0.21}_{-0.24}$	$0.62^{+0.07}_{-0.07}$
SRGE J144738.4+671821	$10.63^{+0.09}_{-0.08}$	$2.42^{+0.10}_{-0.06}$	$0.33^{+0.40}_{-0.19}$	$7.52^{+2.63}_{-1.99}$	$-0.03^{+0.14}_{-0.11}$	$0.04^{+0.05}_{-0.03}$

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