

Constraining Type Ia supernova explosions and early flux excesses with the Zwicky Transient Factory

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ABSTRACT

In the new era of time-domain surveys, Type Ia supernovae are being caught sooner after explosion, which has exposed significant variation in their early light curves. Two driving factors for early-time evolution are the distribution of ^{56}Ni in the ejecta and the presence of flux excesses of various causes. We perform an analysis of the largest young SN Ia sample to date. We compare 115 SN Ia light curves from the Zwicky Transient Facility to the TURTLES model grid containing light curves of Chandrasekhar mass explosions with a range of ^{56}Ni masses, ^{56}Ni distributions, and explosion energies. We find that the majority of our observed light curves are well reproduced by Chandrasekhar mass explosion models with a preference for highly extended ^{56}Ni distributions. We identify six SNe Ia with an early-time flux excess in our *gr*-band data (four ‘blue’ and two ‘red’ flux excesses). We find an intrinsic rate of 18 ± 11 per cent of early flux excesses in SNe Ia at $z < 0.07$, based on three detected flux excesses out of 30 (10 per cent) observed SNe Ia with a simulated efficiency of 57 per cent. This is comparable to rates of flux excesses in the literature but also accounts for detection efficiencies. Two of these events are mostly consistent with circumstellar material interaction, while the other four have longer lifetimes in agreement with companion interaction and ^{56}Ni -clump models. We find a higher frequency of flux excesses in 91T/99aa-like events (44 ± 13 per cent).

Key words: surveys – supernovae: general.

1 INTRODUCTION

Despite the importance of Type Ia supernovae (SNe Ia) as cosmic distance indicators (Riess et al. 1998; Perlmutter et al. 1999; Scolnic et al. 2018), their explosion mechanisms and progenitor systems are still debated. There is a general consensus that SNe Ia are the consequence of the thermonuclear explosion of a white dwarf (WD) residing in a binary system, but the nature of the binary companion is uncertain (Maoz, Mannucci & Nelemans 2014; Livio & Mazzali 2018). There are two main proposed channels: a WD accreting mass from a non-degenerate star (SD), or two WD in a double-degenerate scenario (DD). Alternatively, theories can be categorized by the mass of the WD at the time of explosion: Chandrasekhar mass (M_{Ch}) and sub-Chandrasekhar mass ($\text{‘sub-}M_{\text{Ch}}\text{’}$), where the mass

of the WD is less than $1.4 M_{\odot}$ – see Hillebrandt et al. (2013), Maoz et al. (2014), Ruiter (2019), and Jha, Maguire & Sullivan (2019) for comprehensive reviews.

The specific explosion mechanism leading to SNe Ia also remains subject to debate. The most popular model is that of the deflagration, followed by the detonation of an M_{Ch} -WD (DDT) (e.g. Khokhlov 1991a, b), with predictions that are generally in good agreement with observed maximum light spectra and light curves (e.g. Seitzzahl et al. 2013). Plewa, Calder & Lamb (2004) suggest gravitationally confined detonations (GCD), where an initial off-centre deflagration ignites and expands to the surface. However, these models have a number of drawbacks, such as the overproduction of ^{56}Ni and highly polarized ejecta (Kasen & Plewa 2005), conflicting with spectropolarimetric observations of SNe Ia (Wang et al. 2008; Cikota et al. 2019). Explosions of sub-Chandrasekhar mass WDs have also been proposed. In particular, many investigations have focused on the double-detonation models (Fink, Hillebrandt & Röpke 2007; Fink et al. 2010;

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Kromer et al. 2010; Woosley & Kasen 2011; Sim et al. 2013; Shen & Bildsten 2014; Shen & Moore 2014; Polin, Nugent & Kasen 2019). In this scenario, a WD accretes a shell of He from a companion and this shell ignites and drives a shock wave into the WD, leading to the full ignition of the CO core. There are many more paths being studied that could lead to an SN Ia such as violent mergers of two WDs (Pakmor et al. 2012), triple collision models (Kushnir et al. 2013), rotating super-Chandrasekhar mass explosions (Di Stefano, Voss & W Claeys 2011), core degenerate explosions (CD, Soker 2013; Wang et al. 2017), or the explosions of WD merger remnants (Benz et al. 1990; Shen et al. 2012). There is now significant observational evidence that there may be more than one progenitor scenario/explosion mechanism that produces ‘normal’ SNe Ia, e.g. studies of early light-curve variations and flux excesses (Pskovskii 1984; Riess et al. 1999; Conley et al. 2006; Strovink, Strovink & Mark 2007; Hayden et al. 2010; Ganeshalingam, Li & Filippenko 2011; Firth et al. 2015; Miller et al. 2020b), ejecta masses (Stritzinger et al. 2006; Scalzo, Ruitter & Sim 2014), nucleosynthetic yields (Seitenzahl, Timmes & Magkotsios 2014; Maguire et al. 2018; Floers 2019), and the properties of SN Ia remnants (Martínez-Rodríguez et al. 2017).

Light curves of SNe Ia near maximum light are relatively homogeneous and do not provide the observer with sufficient information to unravel the pre-explosion conditions. However, light curves of SNe Ia obtained within the first days of explosion enable the observer to probe the outer layers of the ejecta and the immediate surroundings of the WD, providing essential information about the accretion/merger history, which are linked to progenitor scenarios and explosion mechanisms. A number of studies have investigated the early rise times of SNe Ia (where the early evolution is parametrized as a power law with exponent, α) to determine if it is consistent with the predictions of the fireball expansion model ($\alpha = 2$) and found mean power-law indices in the range of 1.80–2.40, and several SNe Ia were clearly inconsistent with $\alpha = 2$ (Conley et al. 2006; Hayden et al. 2010; Ganeshalingam et al. 2011; González-Gaitán et al. 2012; Firth et al. 2015; Papadogiannakis et al. 2019; Miller et al. 2020b).

The light curves of SNe Ia are energetically driven by the radioactive decay of ^{56}Ni , which is synthesized in the explosion (Arnett 1982; Pinto & Eastman 2000; Piro & Nakar 2013). The ^{56}Ni distribution has significant effects on the dark phase (the time that passes between the explosion epoch and time of first light; Piro & Nakar 2013) and the shape of the early rise (Gamezo et al. 2005; Piro & Nakar 2013; Mazzali et al. 2014; Piro & Morozova 2016; Magee et al. 2018). If ^{56}Ni extends to the outer regions of the ejecta, a shallower, earlier rise, and a shorter dark phase is expected. The full range of rise indices can be reproduced by simply varying the ^{56}Ni distributions (Piro & Nakar 2013; Magee et al. 2018, 2020). Dessart et al. (2014) argued that a pulsation delayed detonation of a WD can also produce light curves that are brighter and bluer at earlier times, synonymous to the effects of a highly mixed ejecta. However, early-time spectra and knowledge of the velocity gradients can be used to distinguish between these two scenarios (Dessart et al. 2014).

In a handful of cases, a flux excess in the early light curves of SNe Ia has been detected (commonly referred to as ‘bump’ or ‘flash’), e.g. SN 2012cg (Marion et al. 2016), SN 2014J (Goobar et al. 2015), iPTF14atg (Cao et al. 2015), SN 2016jhr (Jiang et al. 2017), SN 2017cbv (Hosseinzadeh et al. 2017), SN 2018oh (Dimitriadis et al. 2019; Li et al. 2019; Shappee et al. 2019), SN 2019yvq (Miller et al. 2020a), SN 2020hvf (Jiang et al. 2021), and SN 2021hpr (Lim, in preparation). Olling et al. (2015) analysed a sample of three SNe Ia with extremely high-cadence Kepler light curves and found no evidence of flux excesses at early times in these events. Fausnaugh et al. (2021) perform a similar study using Transiting Exoplanet Survey

Satellite (TESS) data and constrain the size of the companion to $<20 R_{\odot}$ for five SNe Ia, and $<4 R_{\odot}$ for two SNe Ia. Magee et al. (2020) (hereafter **M20**) investigated a literature sample of SNe Ia with early light-curve data and found that 22 per cent of objects required additional flux at early times to fit the light curve. Jiang et al. (2018) coined the term ‘Early-broad EExSNe Ia’, including objects with very broad early light curves in the class of excess flux objects (EExSNe). Under this broader definition, Jiang et al. (2018) found that all 91T/99aa-like events could be classified as EExSNe.

A number of different scenarios have been suggested to produce these flux excesses, such as interaction with a non-degenerate companion star (Kasen 2010), interaction with circumstellar material (CSM; Kromer et al. 2016; Piro & Morozova 2016), signatures of an He-shell detonation (Jiang et al. 2017; Noebauer et al. 2017; Polin et al. 2019; Magee et al. 2021, Ni et al. 2022), the presence of ^{56}Ni clumps in the outer ejecta (Dimitriadis et al. 2019; Shappee et al. 2019; Magee & Maguire 2020), or interaction with disk-originated matter (Levanon & Soker 2019). However, linking an observed flux excess to one particular scenario has proven difficult. The main reason for this is that while some of these models are able to reproduce the light curves with flux excesses, the accompanying model spectra are usually not in agreement with observations. Both the early and maximum light spectra can be significantly affected by He-shell detonations or ^{56}Ni clumps (e.g. Maeda et al. 2018; Polin et al. 2019; Magee & Maguire 2020; Miller et al. 2020a), and signatures expected from the interaction scenarios (e.g. H and He in nebular spectra) have been observed only in a very limited number of cases (Kollmeier et al. 2019; Prieto et al. 2020).

Our aim in this work is to extend the study of **M20** by analysing the ^{56}Ni distributions of a large, homogeneous sample of SNe Ia from the Zwicky Transient Facility (ZTF; Bellm et al. 2018; Graham et al. 2019; Masci et al. 2019; Dekany et al. 2020) to determine if the Chandrasekhar mass models of **M20**, with varying ^{56}Ni masses, ^{56}Ni distributions, shapes of the density profile, and kinetic energies, can reproduce the observed light curves. We also aim to constrain the fraction of objects in this sample that display a flux excess. The general properties, rise times, and colour evolution of the sample of SNe Ia in this sample are described in Yao et al. (2019), Miller et al. (2020b), and Bulla et al. (2020). In Section 2, we describe the ZTF SN Ia sample, as well as the model grid of **M20**. In Section 3, we discuss our methods for comparing the observed light curves to the model grid and identifying flux excess in the early light curves, as well as an efficiency analysis of the flux excess detection. In Section 4, we present the results of our light-curve analysis, and in Section 5, we discuss the implications of our results in terms of explosion mechanisms and the origin of flux excesses. We conclude our work in Section 6.

2 SAMPLE SELECTION AND MODEL GRID

Our aim in this work is to compare the 2018 ZTF SN Ia sample to a grid of 300 Chandrasekhar mass models presented in **M20**. The models are parametrized with different ^{56}Ni masses and ^{56}Ni distributions, kinetic energies, and density profiles. In Section 2.1, we introduce the sample of SN Ia light curves used in this analysis and discuss how we processed the light curves in order to compare them to the model light curves. In Section 2.2, we present a description of the **M20** model grid and in Section 2.3, we describe the cuts we applied to ensure that our sample has sufficient early-time data and consists solely of ‘normal’ SNe Ia (those that follow the width–luminosity relation; Pskovskii 1977, 1984; Phillips, Phillips & M. 1993, and can be used in cosmo-

logical analysis). Finally, in Section 2.4, we discuss our identification of the host galaxies for the sample and their general properties.

2.1 Data processing of ZTF light curves

The ZTF collaboration conducted a high-cadence extra-galactic survey during its first year, which obtained six epochs ($3g + 3r$) nightly over a sky coverage of 2500 deg^2 (Bellm et al. 2019; Dekany et al. 2020), resulting in 336 spectroscopically classified, high-cadence SNe Ia light curves (Graham et al. 2019).¹ The high-cadence and multiband photometry allows us to investigate the colour evolution, as well as the luminosity evolution of these objects from a very early stage. Yao et al. (2019) describe the general properties of a sub-sample of the 2018 ZTF data set, which consists of 127 SNe Ia with both g - and r -band observations at least 10 d before peak (obtained originally through ztfquery; Rigault 2018). All objects have a minimum of five observations in both bands before the time of peak in the B band.

In order to compare the ZTF light curves to the model grid of M20, we need to convert the raw counts to absolute flux. Following the prescription presented in section 3.5 in Yao et al. (2019), we convert the forced photometry counts to flux values for all 127 objects using the photometric zero-point of every image provided by the pipeline and correct the baseline by the offset values calculated in Yao et al. (2019). We apply an error floor of 2 per cent to the data to account for underestimated photometric uncertainties. The flux on each specific epoch has been stacked, with means weighted proportional to the size of the uncertainties. We also correct for time dilation, placing the light curves in the SN rest frame (for the rest of this paper, all quoted times have been corrected to the SN rest frame).

We apply the Milky Way extinction correction in the direction of each SN using dust maps from Schlafly & Finkbeiner (2011), assuming $R_V = 3.1$ and using the Cardelli, Clayton and Mathis extinction law (CCM, Cardelli et al. 1989). Bulla et al. (2020) performed SNooPy fitting (Burns et al. 2014) and derived host galaxy reddening values and K -corrections for 65 objects in the sample. The host extinction values derived by Bulla et al. (2020) range from 0.008 to 0.78 mag, with 10 objects having a host extinction greater than 0.3 mag. Following Yao et al. (2019) and Miller et al. (2020b), we do not correct for host extinction, and any highly reddened objects are excluded based on the colour (c) derived from fitting the light curves with the light-curve fitter SALT2 (see Section 2.3; Guy et al. 2007) since SALT2 does not explicitly fit for host reddening, but it is encapsulated in the c value. K -corrections at early times for SNe Ia are uncertain as they are derived by extrapolating from 15 d prior to maximum light to earlier epochs. Moreover, Bulla et al. (2020) found that the K -corrections for their subset of the ZTF SN Ia sample were small, with the objects at the highest redshifts having K -corrections of ≤ 0.1 mag. Therefore, no K -corrections have been applied in the following analysis.

2.2 Models

For a detailed description of the models used in this work, we direct the reader to Magee et al. (2018) and M20. Briefly, the model grid contains 300 light curves, which were produced using the one-dimensional, Monte Carlo radiative transfer code, TURTLS (Magee et al. 2018, M20). The code is designed to model the light curves

of thermonuclear SNe from explosion up to maximum light. Our particular suite of models was produced by varying the ejected ^{56}Ni mass, the ^{56}Ni distribution, the shape of the density profile, and the kinetic energy of the ejecta. We implement the Chandrasekhar mass model suite of M20 of 255 models with ^{56}Ni masses of 0.4, 0.6, and $0.8 M_\odot$, which was used to fit the light curves of 35 well-observed SNe Ia in M20. We supplement this model suite with 45 additional models with an Ni mass of 0.5 calculated as part of this work, in order to increase the resolution in the $0.4\text{--}0.6 M_\odot$ region because our initial analysis showed most objects to be best fit by models in this range.

The shape of the ^{56}Ni distribution as a function of the mass coordinate, m , is defined using the relation of Magee et al. (2018),

$$^{56}\text{Ni}(m) = \frac{1}{\exp(P[m - M_{\text{Ni}}]/M_\odot) + 1}, \quad (1)$$

where M_{Ni} is the total ^{56}Ni mass in M_\odot , and P parametrizes the shape of the ejecta (defined as s in Magee et al. 2018), with smaller values of P representing more extended ^{56}Ni distributions and larger values of P representing more compact distributions (^{56}Ni confined to the core). The model grid contains two density distributions: exponential (EXP) or double power law (DPL). There are four ejected ^{56}Ni masses (0.4, 0.5, 0.6, and $0.8 M_\odot$), nine kinetic energies ($0.5\text{--}2.18 \times 10^{51}$ erg), and five P values (spanning 3–100 in log space). The model names encapsulate all the information about the parameters (e.g. EXP_Ni0.4_KE0.65_p4.4).

It is important to note that the model composition is simplified to three ejecta zones: the innermost zone is composed of iron group elements (IGEs, assumed to be 100 per cent ^{56}Ni at the time of explosion), a middle zone of intermediate mass elements (IMEs), and an outer layer composed of $0.1 M_\odot$ of carbon and oxygen. The assumption of a pure radioactive ^{56}Ni composition in the inner zone at the time of explosion results in an overestimation of the brightness compared to a more physical composition. The ^{56}Ni masses derived should therefore be considered as lower limits. Because the models assume local thermal equilibrium (LTE), they predict only the light-curve evolution accurately up to approximately maximum light. Correspondingly, we use only light-curve data ranging from 0 to 20 d post-explosion for our fitting.

2.3 Sample cuts

This paper focuses on ‘normal’ SNe Ia to match the parameter space covered by the M20 model grid, which does not include the extremes of the SN Ia brightness distribution. Accordingly, we apply a number of cuts to our initial sample of 127 SNe Ia to remove peculiar objects. We have adopted the SALT2 parameters determined by Yao et al. (2019) and applied the cosmological cuts as quoted in Smith et al. (2020) ($-0.3 \leq c \leq 0.3$, $-3.0 \leq x_1 \leq 3.0$), reducing the sample down to 119 objects. These cuts removed one Ia-CSM (an SN Ia interacting with CSM, ZTF18aaykjei), three over-luminous ‘super-Chandrasekhar mass’/03fg-like objects (‘SC’, ZTF18abdvpnd, ZTF18aawpcel, ZTF18abhpgje), one 91T-like (ZTF18abealop), and three normal objects (ZTF18aasesgl, ZTF18abwdcdv, ZTF18abgmcmv), with spectroscopic classifications adopted from Yao et al. (2019).

For a fair comparison between data and models, we include only SNe Ia with absolute peak magnitudes covered by the M20 model grid, so any objects with an absolute peak g -band magnitude brighter than -19.95 mag or fainter than -18.35 mag are removed (see Fig. 1). This cut further excludes one 02cx-like (ZTF18abclfee), one

¹All public data are available on the GROWTH Marshall (Kasliwal et al. 2019, <http://skipper.caltech.edu:8080/cgi-bin/growth/marshal.cgi>).

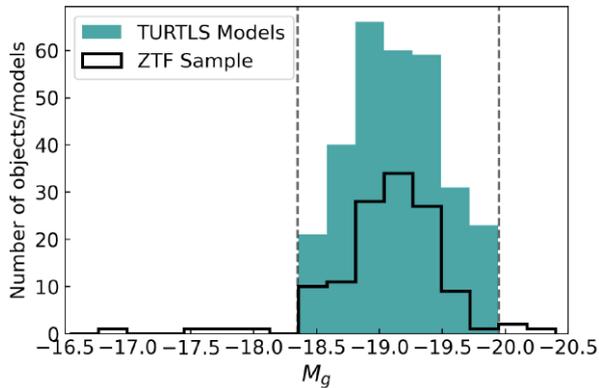


Figure 1. A plot showing the distribution of g -band magnitudes at maximum light covered by the TURTLS model grid, and those derived from SALT2 fits by Yao et al. (2019) for the observed 2018 ZTF SN Ia sample before making any of the cuts described in Section 2.3. The dashed lines indicate the absolute magnitude cuts we applied at $-19.95 \leq \text{peak } g \leq -18.35$.

Table 1. Summary of the cuts applied to the observed SN Ia sample.

Condition	No. of SNe Ia remaining
Starting sample	127
$-0.3 \leq c \leq 0.3$	122
$-3.0 \leq x_1 \leq 3.0$	119
$-19.95 \leq \text{peak } g \leq -18.35$	116
First detection \leq peak (g) – 14 d	115

normal (ZTF18aansqun), and one 99aa-like (ZTF18abjdjge) from our sample reducing the number to 116 SNe Ia.

M20 analysed the effects of limiting early-time data and concluded that in order to reliably constrain the ^{56}Ni distribution, a first $3\text{-}\sigma$ detection within ~ 3 d of the explosion epoch (approximately equivalent to 14 d prior to B -band peak) is required. To satisfy this requirement, we check that the SNe Ia in our sample have a first $3\text{-}\sigma$ detection in the g or r band at least 14 d before the time of the estimated g -band maximum derived from SALT2 fitting (in the SN rest frame). After applying these cuts, we are left with 115 SNe Ia in our final sample for investigation (see breakdown of cuts in Table 1). Our final sample still includes a number of more unusual SN Ia classes: nine 99aa-like SNe, two 91T-like SN, one 86G-like SN, and one ‘SC’/03fg-like SN.

2.4 Host galaxies

For each SN in our sample, the host galaxy was identified from PanSTARRS legacy imaging (Chambers et al. 2016) using the ‘Directional Light Radius’ method (DLR; Sullivan et al. 2006; Gupta et al. 2016). In detail, the closest galaxy (in units of DLR) of the SN position is considered the host galaxy of each event with a requirement of $\text{DLR} < 10$ (Gupta et al. 2016). SN with no host meeting this criterion (6 SNe Ia, 5.2 per cent) are considered hostless. For each host galaxy, $griz$ -band fluxes were determined from the PanSTARRS images using the SEP source extraction code (Barbary 2018), based on Source Extractor (Bertin & Arnouts 1996). To estimate the stellar mass and star formation rate of each host, we use the methodology described in Sullivan et al. (2010) and Smith et al. (2020). We use the PÉGASE.2 spectral synthesis code (Fioc & Rocca-Volmerange 1997) combined with a Kroupa (2001) initial mass function to calculate the spectral energy distribution (SED) of

a galaxy as a function of time using nine exponentially declining star formation histories (SFH) combined with seven foreground dust screens. To break the degeneracy between age and metallicity, for each SFH we generate 4000 models, each with a burst of star formation randomly superimposed between 1 and 10 Gyr after formation (Childress et al. 2013). Each model is then evaluated at 102 time-steps. For each host, we compare our $griz$ fluxes, after correcting for Milky Way extinction, to those of each model SED at the redshift of the SN to find the best-fitting template and stellar mass. Uncertainties on stellar mass are calculated using a Monte Carlo approach. The mean host galaxy stellar mass for our sample is $9.82 M_{\odot}$, and host masses and DLR values are given in Appendix A, Table A1.

3 METHODS

In this section, we discuss the fitting routine implemented to compare the model light curves to our data set of observed SN Ia light curves with early g - and r -band data. We describe the conditions for classifying objects as ‘well fit’ and compare the data to objects presented by M20 to ensure that the fitting routine is not affected by differences between the samples. The method for detecting flux excesses is described in Section 3.3, and the efficiency of this method is investigated in Section 4.3.1.

3.1 Comparing observed light curves to the model grid

We compare each observed SN Ia light curve to the grid of Chandrasekhar mass models of M20 in flux space using the χ^2 metric. Performing the fits in flux space has the benefit that we can include all the available data, including lower significance early data points. Following the methodology from M20, the first ‘detection’ is taken as the earliest 3σ detection in the g or r band. A rough date of maximum light is initially estimated as the mean of the epochs of peak magnitude in the g and r bands. We limit the range of data for fitting from 30 d before maximum light up to 20 d after first detection. We then fit simultaneously for the distance modulus, explosion date, and the best-fitting model.

We allow the distance modulus to vary up to 0.3 mag to account for the coarseness of the model grid, as well as potentially underestimated distance uncertainties. We iterate over the distance modulus range in 15 steps. As described in M20, it is assumed that the explosion could have occurred anytime between 30 d before maximum light, up to the time of the first $3\text{-}\sigma$ detection. The code iterates over this range in increments of 0.1 d. The χ^2 value is computed for each comparison of the model with the SN Ia light curve, for all distance modulus, and explosion date in the allowed ranges. Subsequently, we calculate χ_{red}^2 by dividing χ^2 by the degrees of freedom (DOF), estimated as the number of data points in each fit, minus the number of parameters fitted. We normalize the χ_{red}^2 values by dividing by the smallest χ_{red}^2 (such that the best-fitting model has a $\chi_{\text{norm}}^2 = 1.0$). Fits with χ_{norm}^2 values that fall within the $3\text{-}\sigma$ confidence region of $\chi_{\text{norm}}^2 = 1$, where the cutoff value is taken from the per cent point inverse function dependent on the DOF, are considered as potential matches. The $3\text{-}\sigma$ range is used as a measure for the uncertainty of the model parameters, distance moduli, and explosion epochs.

M20 consider an object consistent with the model if the mean residual is approximately zero, and the maximum value of the residual in each band is ≤ 1 mag. We base our well-fit criteria on the distribution of the χ^2 value, the number of unique fits within the $3\text{-}\sigma$ range (the number of unique matches corresponds to the

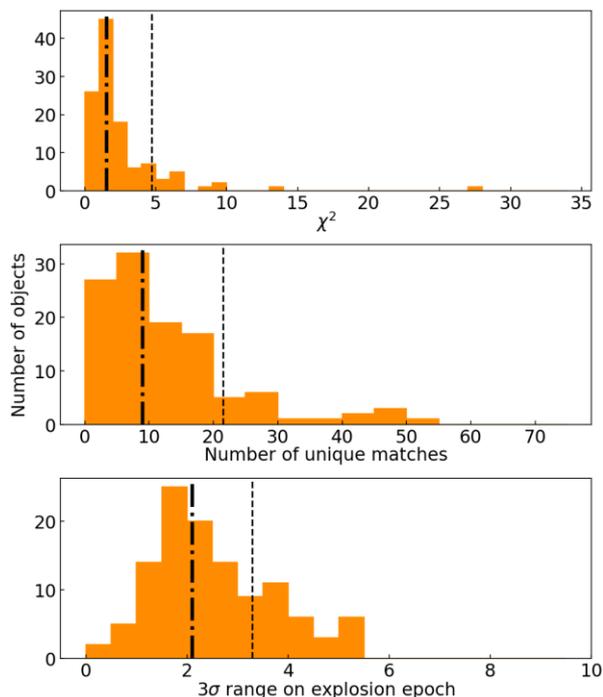


Figure 2. Histograms showing the distribution of χ^2 values, the number of unique fits, and the $3\text{-}\sigma$ range for the estimated explosion dates for each object. The vertical dot–dashed lines indicate the median of the distribution, and the dashed lines indicate one standard deviation from the median where we have placed the cut to separate well from badly fit objects.

number of models from the grid, which are matched to an object), and the $3\text{-}\sigma$ range on the explosion date. We find the median of each distribution (shown as the dash–dotted line in Fig. 2), and consider the upper limit as one standard deviation from the median (shown as the dashed line in Fig. 2). It should be noted that this study does not aim to constrain the intrinsic rate of M_{Ch} –mass explosions but rather it aims to test whether variation in the ^{56}Ni distribution is able to account for significant variation in the early light curves of normal SNe Ia. Therefore, the purpose of these cuts is to remove the tails of the distribution, leaving us with a normal distribution in these three parameters to describe the well-fit sample where these values vary randomly due to the uncertainties in the data. The limits are: a $\chi^2 \leq 4.7$, ≤ 22 unique potential matches, and a $3\text{-}\sigma$ explosion date range ≤ 3.3 d. The latter two conditions relate to the uniqueness of a fit, and how well the model parameters are constrained. Noise in the data means that we are unable to constrain the model, but we are not able to rule out a Chandrasekhar mass explosion with certainty. These parameters are also highly dependent on the resolution of the model grid, but since we are comparing fits to the same model grid, the relative values are still useful. When using this method to compare light curves to an expanded model grid, these limits should be adapted. Fits with high χ^2 values indicate that we are unable to fit the data with our models. Combined, these quality cuts remove 37 SNe Ia from the sample. This does not necessarily imply that all these SNe Ia originate from non-Chandrasekhar mass explosions but rather that we are unable to constrain model parameters from the fit. This could be due to poor data quality or intrinsic differences in the light curves compared to the models. See Section 4.2 for further discussion of these events. Our final ‘well-fit’ sample of 78 SNe Ia has a mean and standard deviation $\chi_{\text{red}}^2 = 1.6 \pm 1.0$, 8 ± 5 for the number of unique fits, and 1.9 ± 0.7 d for the $3\text{-}\sigma$ range of explosion dates.

3.2 Validation of the use of ZTF *gr*-band light curves

We analyse the potential effects of the differences between our sample and that presented by M20. We investigated whether the availability of only two bands (*g* and *r*) in the ZTF sample may affect our results, compared to those with multiband coverage in the M20 sample. We do this by selecting objects in M20 where the earliest data are in the *g* or *r* band but also have wider wavelength coverage, finding seven objects. We then refit these objects, limiting the bands to just *g* and *r* bands (Sloan *g* and *r* bands, which are comparable to ZTF bands). For five events (SN 2011fe, SN 2013gy, iPTF16abc, SN 2017cbv, SN 2017erp), the best-fitting model was not affected by reducing the number of bands available for fitting. Two events (iPTF13ebh, iPTF13dge) change from an original best fit of $P = 9.7$ to $P = 4.4$ when using only the *g*- and *r*-band data. However, for iPTF13ebh, only the best match changes, but the next best match is still included in the $3\text{-}\sigma$ range. In the case of iPTF13dge, we find that it has a large gap (4 d) between predicted explosion date and the epoch of the first data point. In this case, it is not possible to narrow down the explosion range sufficiently, and subsequently, the model parameters are badly constrained. Not all the light curves in our sample have a $3\text{-}\sigma$ detection prior to 3 d after explosion, but all have at least one $\leq 3\text{-}\sigma$ detection prior to this epoch, which we also include in the fitting. If there are insufficient early data to constrain the explosion epoch, the fit will not pass our quality cut on the $3\text{-}\sigma$ range of explosion dates. We conclude that *g*- and *r*-band data as in the ZTF sample are sufficient for constraining the ^{56}Ni distribution of SNe Ia when using the TURTLS models.

3.3 Detection of flux excess

SNe Ia with flux excesses in their early light curves are commonly presented in single object studies. In these cases, high-cadence observations at early times allow the bump to be clearly resolved (e.g. SN 2017cbv; Hosseinzadeh et al. 2017) or a data point is obtained during the decline of the bump (e.g. SN 2019yvq; Miller et al. 2020a), separating the rise of the bump from the ^{56}Ni -powered rise to the secondary peak. Jiang et al. (2018) were the first to systematically look for early flux excesses in SN Ia light curves. In their study, light curves with prominent flux excesses are identified by eye or their presence is taken from the literature, and early-broad EExSNe are identified through a comparison with the light curve of SN 2012cg. M20 search for flux excesses in the early light curves of SNe Ia in a more quantitative way by requiring that a data point produces a residual >1 mag when compared to the best matched model. We aim to search for flux excesses in a large sample of SNe Ia and want to use a quantitative reproducible method. To this end, we implement a similar method as presented in M20. However, we assess the residuals in flux rather than magnitude space so that all data points including upper limits can be analysed in a consistent fashion.

We base our parametrization of potential flux excesses in the data on the theoretical predictions from progenitor models. The additional power sources that are predicted to produce early flux excesses are expected to be blue (e.g. Kasen 2010; Polin et al. 2019). However, this additional luminosity must be combined with the underlying Ni-powered light curve that can make the early colour appear less blue (e.g. Piro & Morozova 2016). Therefore, we check for potential flux excesses in both the *g* and *r* bands. The longest flux excess in the *B* band in the interaction models of Kasen (2010) lasts for ~ 6 d and we search for flux excesses of up to this length. Specifically, we search

for positive residuals between the data and best-fitting model in the 6 d after the explosion epoch estimated from the best-fitting model, with an uncertainty on the explosion epoch corresponding to the $3\text{-}\sigma$ model range. This range also covers the time-scales of flux excesses seen in previous literature events (e.g. Hosseinzadeh et al. 2017; Dimitriadis et al. 2019; Shappee et al. 2019; Miller et al. 2020a). We require two consecutive epochs to have non-zero residuals in a single band in this range, where at least one has a residual of >2 per cent of the peak flux² and also require all models in the $3\text{-}\sigma$ range to have a non-zero residual outside their uncertainties at these epochs. Models of companion interaction (Kasen 2010), CSM interaction (Piro & Morozova 2016), and ^{56}Ni -clumps (Magee & Maguire 2020) all predict flux excesses that begin at approximately the same time as the explosion. The estimated explosion epoch from our fit is highly uncertain when a flux excess is present, but we rule out flux excesses that begin later than 1 d after the maximum explosion epoch in the $3\text{-}\sigma$ range. Some ZTF light curves show unphysical scatter that is most prominent around peak brightness, so again to be conservative in what we consider an early flux excess, we remove any objects with potential flux excesses at early times that also show a scatter of similar significance at later times (any days >10 d after explosion). Because this may remove events with true early flux excesses, for consistency we remove any SNe Ia in the full sample with similar levels of scatter prior to performing the flux excess rate calculation in Section 5.3.

We test the sample for contamination from false excesses through a search for negative flux excesses, performed by inverting the flux excess criteria. No negative flux excesses are found, indicating that the method can robustly detect true excesses.

We also considered an alternative method of removing the data at early epochs, refitting the light curve, and checking if the early data produce a positive residual relative to the best match. However, M20 emphasize the importance of the early data for reliably constraining the ^{56}Ni distribution, and we clearly notice the effect of the missing data. By leaving out the early epochs, we identify 70 per cent of objects as having a flux excess, but upon closer analysis, we find that the fits show a preference for models with more compact ^{56}Ni distributions. Consequently, the data points corresponding to the dark phase of the models are picked up as excess flux. Both methods have drawbacks, but we opt for the method described in detail above, as this method is less likely to produce false positives. It should be noted that the explosion parameters derived from a model fit, as well as the estimated explosion epoch, are not well constrained in the scenario where a flux excess is detected and these should not be used for further analysis.

4 RESULTS

Following the methodology presented in Section 3.1, we present the results of the analysis here. We first describe the objects that passed the quality cuts and appear to be well fit by Chandrasekhar mass explosion models and present the distributions of the model parameters in Section 4.1. This is followed by an analysis of objects

²This choice is motivated by the strength of early flux excesses seen in literature but the exact value is somewhat arbitrary. However, we consider it a conservative value that selects only the highest significance events. This method assumes that the brightness of a flux excess should scale with the overall brightness of the SN, which is not always the case, since the brightness of the flux excess in the case of companion interaction depends on external properties (e.g. companion separation, viewing angle) that are not inherently related to the exploding WD.

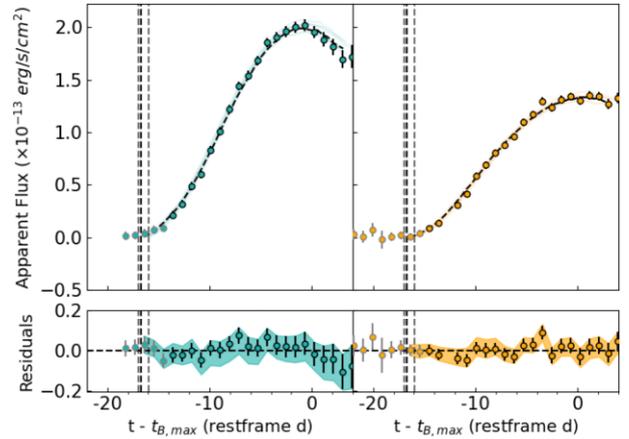


Figure 3. In the top panels, we show the flux g -band (left) and r -band (right) light curves of ZTF18abtnlk, which was well fit by models from the Chandrasekhar mass model grid. The light curve is plotted in the rest frame of the SNe, where $t_{B, \max}$ is taken from SALT2 fits performed by Yao et al. (2019). The models that fit within a $3\text{-}\sigma$ range are shown as the coloured shaded lines, and the best-fitting model is represented by the black dashed line. Any $\leq 3\sigma$ data points are included in the fitting routine but have a grey outline to highlight the amount of data that would be lost if we fitted in magnitude space where these points would be considered upper limits. The range of estimated explosion dates within $3\text{-}\sigma$ is shown by the grey dash-dotted lines, and the explosion date of the best-fitting model is indicated by the vertical black dash-dotted line. The residuals are also shown in flux in the same units as the top panel. The shaded region indicates the $3\text{-}\sigma$ region of the best-fitting models and also accounts for the uncertainty on the explosion epoch as well as the photometric uncertainty.

that have no good match in our model grid in Section 4.2. Lastly, in Section 4.3, we present our sample of SNe Ia with a potential flux excess, their light curves, host galaxy properties, and the details and results of a detection efficiency calculation.

4.1 SNe Ia well fit by Chandrasekhar mass explosions

We find that 78 out of 115 observed SN Ia light curves (67 per cent) can be well reproduced by the TURTLS Chandrasekhar mass model grid, implying that more than half of the variation observed in normal SN Ia light curves can be reproduced by simply varying the parameters (^{56}Ni mass, ^{56}Ni distribution, shape of the density profile, kinetic energy) in Chandrasekhar mass explosions. Tables A3 and A4 summarize the parameters of all SNe Ia in our sample, and the light curves are shown in Appendix A, Figs A1–A3. As discussed in Section 3.1, the conditions for being ‘well-fit’ are $\chi^2 \leq 4.7$, an uncertainty on the explosion epoch of ≤ 3.3 d, and ≤ 22 unique matches in the $3\text{-}\sigma$ range. In Fig. 3, we present a sample light curve that is very well reproduced by our models.

Fig. 4 shows the distribution of the ^{56}Ni mass, ^{56}Ni distribution in the ejecta, and kinetic energy of the well-fit and badly fit objects. The density distribution of our well-fit sample is dominated by exponential profiles; only 18 well-fit objects (23 per cent) have a double power-law density profile. The best-fitting ^{56}Ni masses are distributed between 0.4 and 0.6 M_{\odot} with a preference for lower masses. No objects were matched with an ejected ^{56}Ni mass of 0.8 M_{\odot} , although it should be noted that these values are lower limits. Moreover, the typical magnitude difference at peak between our models with different ^{56}Ni mass is ~ 0.5 mag and given that the distance modulus can vary by up to 0.3 mag, there is some overlap between the best-fitting values for our SN Ia sample. All well-fit

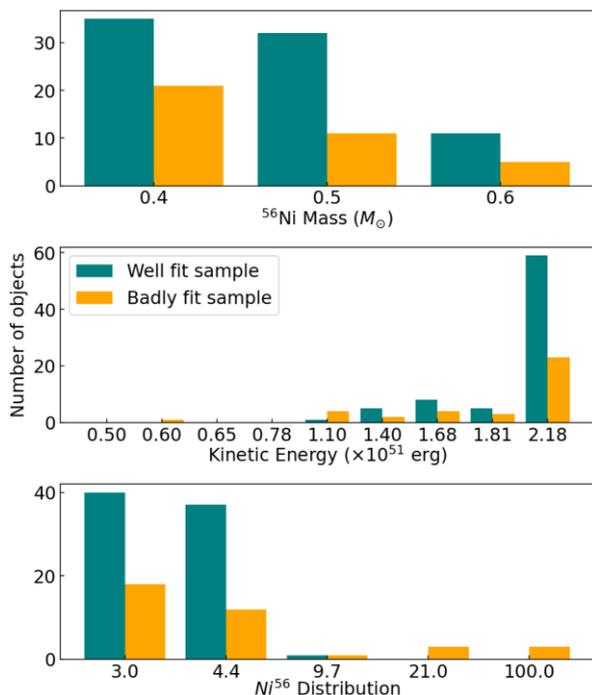


Figure 4. Histograms showing the distributions of ^{56}Ni mass (top), kinetic energy (middle), and ^{56}Ni distribution in the ejecta (low values indicate a very extended distribution) of the top match to each object in our well-fit sample consisting of 78 objects (teal), and the badly fit sample consisting of 37 objects (orange). It is clear that the ZTF SNe Ia sample shows a preference for highly energetic models with very extended ^{56}Ni distributions. We also find that the objects matched with compact ^{56}Ni distributions are all in the badly fit sample.

objects have more than one ^{56}Ni mass in the $3\text{-}\sigma$ range, and we therefore estimate the uncertainty on our ^{56}Ni mass estimates to be $\pm 0.1 M_{\odot}$. Models with highly extended ^{56}Ni distributions dominate the sample: 77 out of 78 well-fit objects are best matched by a model with $P = 3, 4.4$ (the two most extended ^{56}Ni distributions), with a single object being matched by $P = 9.7$. Fig. 4 shows that all SNe Ia matched by compact ^{56}Ni distributions are classified as badly fit. Finally, models with a high kinetic energy dominate the sample (82 per cent, have ejecta kinetic energies $>1.4 \times 10^{51}$ erg). The preference for models with higher kinetic energies is currently not understood, and we are unable to determine whether this is an artefact of the fitting routine or an intrinsic property of the sample. There is some degeneracy between the ^{56}Ni mass and the kinetic energy, which could be the cause of this skewed distribution, meaning these two parameters should be interpreted with caution. However, the ^{56}Ni distribution is the main driver of the shape of the rise and is not heavily affected by the other parameters.

We investigated the correlations between the observed light-curve properties (absolute peak g -band magnitude, M_g , and x_1 parameter from the SALT2 fitting performed by Yao et al. 2019) and the model ^{56}Ni masses and distributions for the 78 well-fit events. Table 2 shows a summary of the correlation coefficients. Unsurprisingly, M_g is positively correlated with the ^{56}Ni mass and x_1 value is correlated with M_g , in agreement with the well-known correlation between peak absolute magnitude and light-curve shape (Pskovskii 1977; Phillips et al. 1993; Hamuy et al. 1995; Riess, Press & Kirshner 1996). x_1 is also weakly correlated with ^{56}Ni mass. Our sample is heavily dominated by objects with very extended ^{56}Ni distributions in the

Table 2. The Pearson’s and Spearman’s correlation coefficients (r and ρ , respectively) and their respective p values for the model outputs (^{56}Ni mass and ^{56}Ni distribution, $\exp(P)$) compared to the light-curve observables (absolute peak g -band magnitude, M_g , and x_1 parameter from SALT2 fitting). We check correlations with $\exp(P)$ rather than P due to its formal definition in equation 1. We have highlighted in bold any significant correlations (those with correlation coefficients ≥ 0.50 and p values < 0.05).

Parameter 1	Parameter 2	r	p value	ρ	p value
^{56}Ni mass	$\exp(P)$	0.05	0.678	-0.09	0.415
^{56}Ni mass	x_1	0.40	3×10^{-4}	0.42	1×10^{-4}
^{56}Ni mass	M_g	0.70	2×10^{-12}	0.70	4×10^{-12}
$\exp(P)$	x_1	0.06	0.577	-0.07	0.519
$\exp(P)$	M_g	0.08	0.498	0.18	0.107
x_1	M_g	0.51	2×10^{-6}	0.52	9×10^{-7}

ejecta (very low P values), but due to the coarseness of the model grid in this parameter, it is difficult to extract any trends. However, we encourage a future analysis with an increased grid resolution for the ^{56}Ni distributions in order to further explore potential correlations with the shape of the rise.

4.2 Badly fit SNe Ia

For 37 out of 115 SNe Ia (32 per cent), the model fits did not pass the cuts defined in Section 3.1. Three objects (ZTF18abdfwur, ZTF18abpamut, and ZTF18abxxssh – see Fig. 5) fail to pass the cuts because they show a significant flux excesses at early times, which is inconsistent with any of the models. Further analysis of these light curves is presented in Section 4.3. ZTF18abdfzsk, ZTF18aayjvve, and ZTF18aaqqoqs also have a detected flux excess, but they pass the cuts and are included in the well-fit sample.

The other 34 excluded objects (See Appendix A, Fig. A4 for their light curves) fail the cuts for different reasons; eight fail because they have $\chi^2 \geq 4.7$, and seven fail only because the uncertainty on their explosion epoch is ≥ 3.3 d, with 19 failing on more than one criterion. It is important to note that for objects with noisy data, which were ruled out only due to their number of unique matches or large uncertainty on the explosion epoch (21 SNe Ia, 18 per cent), we are unable constrain their true explosion parameters, or conclusively determine whether they could be matched by Chandrasekhar mass explosions. We classify 13 light curves as ‘borderline’ by eye (ZTF18aaqcugm, ZTF18aatzygk, ZTF18aawjyww, ZTF18aaxdrjn, ZTF18abfhryc, ZTF18abjtgdg, ZTF18abjtger, ZTF18abkhcrj, ZTF18abkhwcl, ZTF18abkifng, ZTF18abkudjo, ZTF18abqjvyl, and ZTF18abrznrb). These objects have fits that generally look acceptable but marginally fail to pass the unique match cut of ≤ 22 , or uncertainties on the explosion epochs are > 3.3 d. As discussed in Section 3.1, the cuts on the light curves that were applied are based on the distributions of the values for the overall sample and are somewhat arbitrary but allow us to select the well-fit sample in a consistent and reproducible way.

Only 13 out of 115 (11 per cent), which have a $\chi^2 \geq 4.7$, could not be matched by our models. Interestingly, seven of these have SALT2 c value > 0.15 . Out of the 115 objects in our sample, only 10 have c value > 0.15 . One of the other three with $\chi^2 < 4.7$, ZTF18abcysdx, is ruled out because the explosion date range spans 5.4 d and it has 29 unique matches, implying that the parameters could not be adequately constrained. Of the two well-fitted objects with $c > 0.15$, ZTF18aaydmkh passes our cuts and falls into the well-fit sample, but extinction and light-curve shape appear to be an issue at late times,

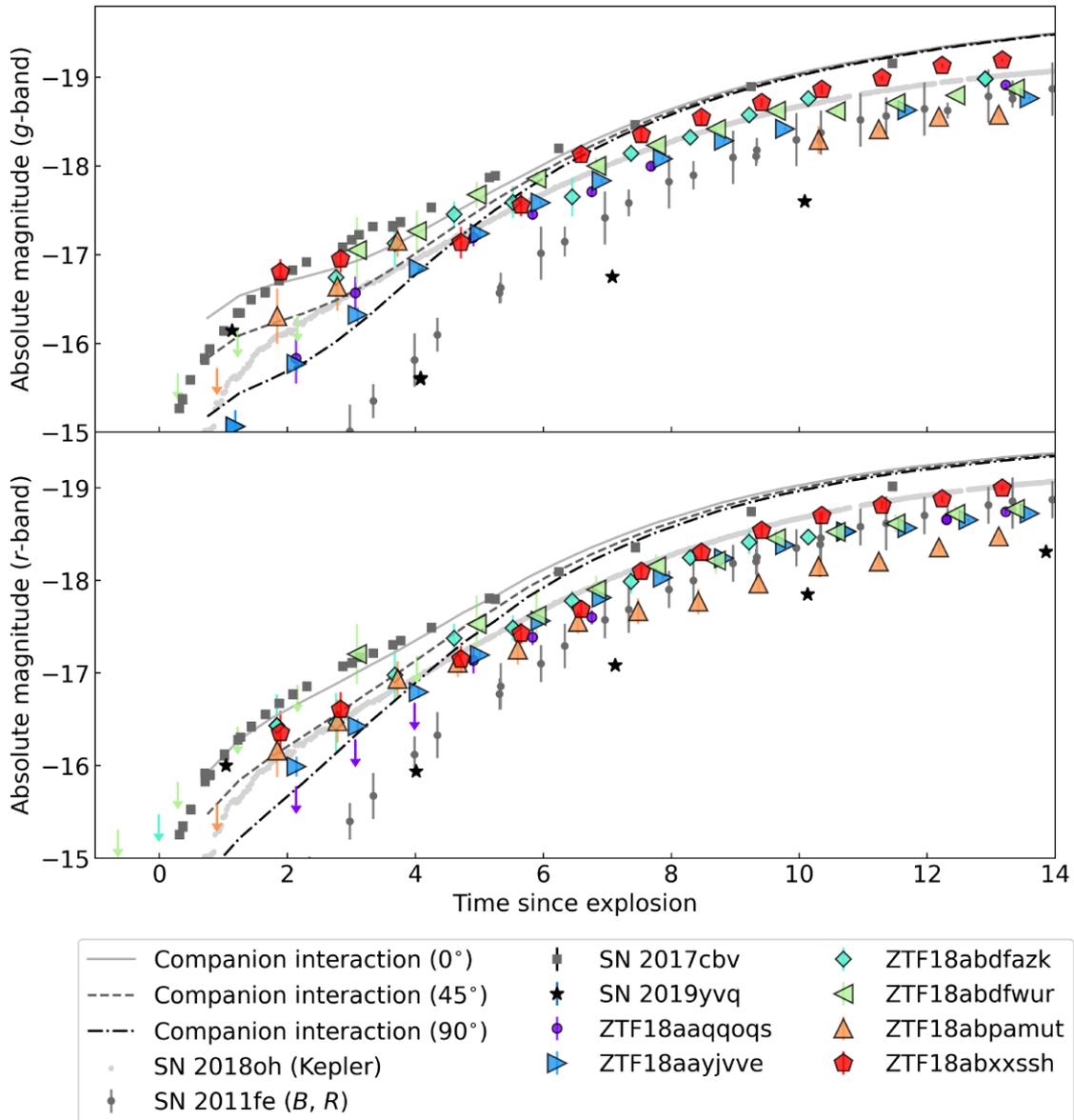


Figure 5. Plots showing the g - (top panel) and r -band (bottom panel) light curves of the six SNe Ia in our sample in magnitude space with a detected flux excess compared to three known SNe Ia with early flux excesses (SN 2017cbv, SN 2018oh, and SN 2019yvq) and the normal SN 2011fe, which shows no flux excess. The Kepler light curve of SN 2018oh is plotted in both panels because gr -band data were unavailable. For SN 2011fe, B -band and R -band data are shown in the upper and lower panels, respectively. Upper limits in magnitude space are shown as downward pointing arrows, although we note that all fitting is done in flux space as detailed in Section 3.3. The companion interaction models used for the efficiency analysis are also shown, where all three models have a separation of 2.00×10^{12} cm but the models have various viewing angles (0° , 45° , 90°) resulting in flux excesses of differing strengths. Photometric data, distance moduli, and explosion epochs for SN 2017cbv, SN 2018oh, SN 2019yvq, and SN 2011fe were taken from Hosseinzadeh et al. (2017), Dimitriadis et al. (2019), Shappee et al. (2019), Miller et al. (2020a), Nugent et al. (2011), and Bloom et al. (2012).

and it has $\chi^2 = 2.5$, which passes the cuts but is significantly higher than the mean of the sample. ZTF18abssdpi similarly passes the cuts, but with a range on the explosion dates of 2.9 d, it lies near the edge of the distribution. To test whether extinction could be the cause of the shape mismatch for these objects, we corrected all SNe Ia with $c > 0.15$ for extinction (values taken from Bulla et al. 2020) and re-ran the fitting routine. This did not improve the model fits, and we conclude that the TURTLS models are not suited to modelling very red SNe Ia, and any future analysis implementing TURTLS models should consider limiting their sample to objects with a SALT2 c value < 0.15 .

4.3 SNe Ia with early flux excesses

Using the method described in Section 3.3, we identify six objects as having an early (within 6 d of explosion) flux excess relative to our underlying explosion models. Their g - and r -band light curves are shown in Fig. 5. Three companion interaction models of Kasen (2010) that are implemented in a flux excess efficiency analysis in Section 4.3.1 are also shown, along with the light curves of three literature events with prominent early flux excesses, SNe 2017cbv, 2018oh, and 2019yvq (Hosseinzadeh et al. 2017; Shappee et al. 2019; Miller et al. 2020a), and the light curve of the normal SN 2011fe (Nugent et al. 2011; Bloom et al. 2012). Four

Table 3. Summary of the properties of SNe Ia with a detected flux excess.

Name	Spectral type ^a	Peak mag.	z	x_1	c excess (d)	Lifetime of flux excess ^b	Ratio of excess at 3 d ^c	$g-r$ colour
ZTF18abdfazk	Normal	−19.2	0.084	−0.28	−0.065	5.5 ± 0.7	$0.073^{+0.043}_{-0.021}$	-0.28 ± 0.43
ZTF18abdfwur	Normal	−18.9	0.070	−0.47	0.041	3.7 ± 0.7	$0.111^{+0.050}_{-0.061}$	0.15 ± 0.49
ZTF18abpamut	Normal	−18.6	0.064	0.83	0.103	6.6 ± 2.4	$0.069^{+0.037}_{-0.026}$	-0.16 ± 0.36
ZTF18abxxssh	91T-like ^d	−19.4	0.064	1.53	−0.017	2.3 ± 1.1	$0.067^{+0.055}_{-0.016}$	-0.35 ± 0.24
ZTF18aaqqoqs*	99aa-like	−19.3	0.082	1.22	−0.014	4.6 ± 0.7	$0.042^{+0.027}_{-0.0001}$	-0.30 ± 0.44
ZTF18aayjvve*	Normal	−18.8	0.0474	−1.48	0.060	2.4 ± 0.5	$0.022^{+0.012}_{-0.003}$	0.10 ± 0.11

^a Spectral classification from template fitting.

^b The ratio of the peak flux excess to the flux at maximum light. This ratio must exceed 0.02 in order for the SN Ia to be categorized as having a flux excess.

^c Measurement of the $g-r$ colour at 3 d post the estimated explosion epoch. This is not necessarily during the peak of the flux excess, but all SNe Ia with an identified flux excess have data at this epoch.

^d We have updated this classification from ‘normal SN Ia’ from Yao et al. (2019) based on the closer spectroscopic similarity to over-luminous 91T-like events.

* ZTF18aaqqoqs and ZTF18aayjvve have flux excesses that are less significant than those previously detected in the literature. Therefore, to be conservative, we do not include them in our intrinsic rate calculation. See Section 4.3.1 for further discussion.

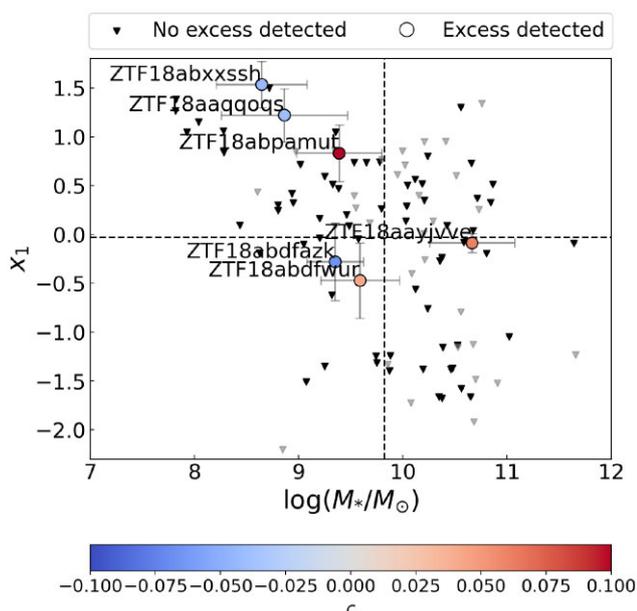


Figure 6. Plot of light-curve stretch, x_1 , as a function of the stellar mass of the host galaxy associated with each object. The colour indicates the c -value from the SALT2 fit. The black/grey triangles represent well-fit/badly fit objects for which no flux excess was detected. The vertical dashed line shows the mean host galaxy mass of our sample ($\log(M_*/M_\odot) = 9.82$) used here to show a division between low- and high-mass galaxies. The horizontal dashed line indicates the mean x_1 value of our sample (-0.03).

of our objects (ZTF18abdfazk, ZTF18abdfwur, ZTF18abpamut, and ZTF18abxxssh) have flux excesses with similar strengths to those of SN 2017cbv and 2018oh, while ZTF18aaqqoqs and ZTF18aayjvve have excesses that are less prominent. All the flux excesses detected in our sample occur later than that of SN 2018oh (although we note that the comparison light curve for SN 2018oh is in the Kepler band, which is roughly equivalent to a broad $g+r+i$ filter). None shows a flux excess that clearly declines and then rises again as was seen in SN 2019yvq but ZTF18abxxssh does show a plateau in its g -band light curve at the earliest epochs.

The properties (spectral type, peak magnitude, light-curve fit parameters, excess lifetime, significance of flux excess, $g-r$ colour

at 3 d post explosion) of the six SNe Ia showing early flux excesses are given in Table 3. We use the $g-r$ colour of the flux excesses at 3 d post the estimated explosion epoch as a way to determine if the early flux excess is intrinsically more prominent in ‘blue’ or ‘red’. We find that four of the early flux excesses have blue colours and two have red colours at 3 d post explosion. In Fig. 6, we show the SALT2 light-curve properties of x_1 and c as a function of the host galaxy mass for all SNe Ia in our sample with the objects displaying an early flux highlighted. Five out of the six SNe Ia with an identified early flux excess occur in a host galaxy with a mass lower than the mean of the sample, and three of the six events have x_1 values in the top 20 per cent of the sample. However, we are limited by the small number of SNe Ia in the flux excess sample, and it is not possible to draw any statistically significant conclusions. There is an even split in the maximum-light colours (c) with three events having red colours and three having blue colours.

All the objects with a detected flux excess have a 3- σ observation in the g or r band within 3 d of explosion. We show in Section 4.3.1 that the efficiency of detecting flux excesses drops off rapidly with increasing redshift, and the epoch of first detection scales with redshift. Early detections are therefore required to resolve the flux excess, but we also note that the presence of a flux excess will make it more likely that an SN Ia is detected earlier. Yao et al. (2019) identified two objects in the ZTF 2018 sample with an early excess, ZTF18aavrwhu and ZTF18abxxssh. ZTF18abxxssh was identified by our code (e.g. Fig. 5), while ZTF18aavrwhu is not picked up by our method because it can be well matched by Chandrasekhar mass explosion models in our grid. Bulla et al. (2020) highlighted a further four objects in their analysis of a subset of the ZTF 2018 sample, which have a rapid colour transition (‘red bumps’) at early times: ZTF18abcflnz, ZTF18abcrxoj, ZTF18abpaywm, and ZTF18abgxvra. These SN Ia do not show any form of flux excess when compared to our model grid but since they were identified based on their colour evolution, it is not surprising that their sample does not overlap with ours.

4.3.1 Flux excess detection efficiency

In order to determine the efficiency of our flux excess detection method, we use simulated light curves with early excesses based on the companion interaction models of Kasen (2010) and attempt

to recover them using the same detection method as for the data. We chose three models that reasonably resemble the range of flux excesses detected in our sample (Fig. 5) but note that there are differences in the lifetimes and strengths of the observed bumps as a function of time that are not completely captured. Our choice of models does not imply a preferred physical origin for the flux excesses but is purely to test the detection efficiency of early flux excess detection. The models were produced by Magee et al. (submitted) using the companion interaction formulations presented in Kasen (2010) combined with TURTLES SN Ia light curves (M20). The underlying Chandrasekhar mass model is chosen to have values resembling our well-fitting sample (^{56}Ni mass of $0.5 M_{\odot}$, kinetic energy of 1.68×10^{51} erg), and a ^{56}Ni distribution with a P value of 9.7 because more extended distributions result in a blended flux excess that would not be detected. All three companion interaction models have a radius of 2×10^{12} m between the WD and the companion, and we choose viewing angles of 0° , 45° , and 90° , resulting in a range of flux excesses of decreasing strength, which roughly covers the distribution of flux excesses found in the ZTF sample.

Next, we ensure that the model light curves resemble the ZTF data. We separate the observed ZTF light curves with $z \leq 0.1$ into three redshift bins of equal volume ($z \leq 0.07$, $0.07 < z \leq 0.087$, $0.087 < z \leq 0.1$). The bin sizes were chosen to ensure that each bin contains sufficient SNe Ia to sample from (the highest redshift bin has nine objects, which is the lowest number of SNe Ia across the bins). We draw photometric uncertainties from these redshift bins and apply these to the model light curves. The simulated light curves are initially set up at a 1-d cadence but to account for weather and other losses, we adjust each simulated light curve to have a cadence matched to a randomly drawn light curve from the observed sample. For ease of computation, the same number of data points are dropped in the g and r bands for an individual simulated event. We produce 15 000 simulated light curves, with equal contributions from each companion interaction model, spread evenly across the three redshift bins. We next apply our detection method to these simulated light curves in order to calculate the detection efficiencies for each redshift bin.

We obtain efficiencies of 68 per cent and 46 per cent in the lowest redshift bin for the strongest (0° viewing angle) and middle (45° viewing angle) model light curves, respectively. These detection efficiencies decline with increasing redshift as expected. For the weakest model light curve (90° viewing angle), we find very low-detection efficiencies (~ 1 per cent). This model was included to account for the two weakest flux excesses in our sample (ZTF18aaqqoqs and ZTF18aayjvve). However, the efficiency is too low to perform rate calculations for these objects. There are two possible explanations for why we observed these weak flux excesses even though the efficiencies deem this unlikely: (a) these two flux excesses are not real, or (b) the weakest flux excess model is not representative of our weakest flux excesses. Fig. 5 shows that these SNe Ia have a clear excess in comparison to SN 2011fe, and, therefore, we find the latter option more likely. This model light curve with the weakest early flux excess differs from the two weakest flux excess objects in both the excess duration and the peak luminosity from which the size of the flux excess is scaled. A further full parameter study, including a more detailed survey simulation (e.g. Feindt et al. 2019), is needed to accurately determine the efficiencies of these lower significance flux excesses but is beyond the scope of this work. Therefore, we exclude these two events and the model with the weakest flux from further discussion of the efficiencies.

Fig. 7 shows the combined g - and/or r -band efficiencies as a function of redshift for the two brightest models. We find the highest detection efficiency (57 ± 11 per cent) in the lowest redshift bin,

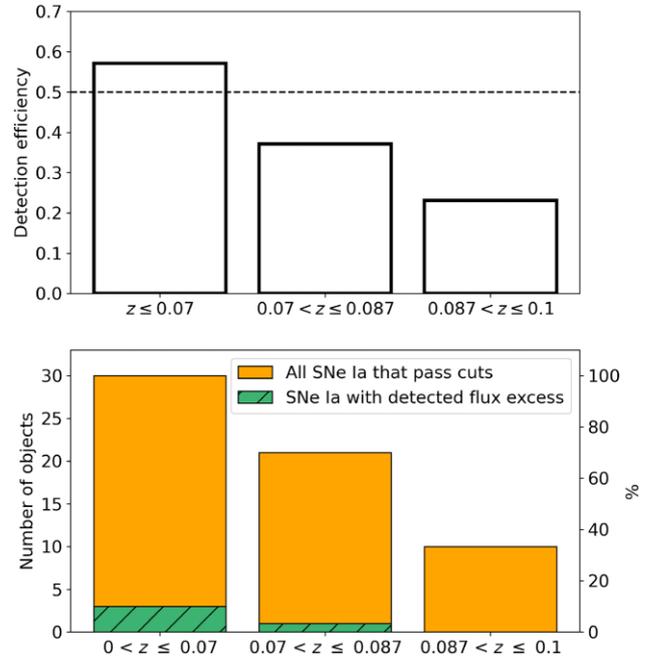


Figure 7. *Top:* Histogram showing the detection efficiency of the flux excesses for three redshift bins for detections in g and/or r band. The efficiencies are based on 10 000 simulated light curves produced by combining the formulations of Kasen (2010) for companion interaction with viewing angles 0° and 45° , with underlying Chandrasekhar mass model light curves (Magee et al., submitted) using typical uncertainties from the ZTF sample in each redshift bin. The efficiency drops below 50 per cent in the second redshift bin ($0.07 < z \leq 0.087$, black dashed line). These efficiencies represent only the strongest flux excesses in the sample and are based on the two strongest flux excess models. *Bottom:* Histogram showing the number of observed ZTF SNe Ia in our sample per redshift bin and the number of those that have a detected flux excess (as a percentage of each bin on the right axis).

with decreasing efficiency towards the highest redshift bin. We find that flux excesses are most frequently identified in the g band (50 per cent), and in half of these cases, we also detect the excess in the r band. The flux excess is detected solely in the r band in 10 per cent of cases.

5 DISCUSSION

In this section, we first discuss the results presented in Section 4.1 in the context of previous investigations of ^{56}Ni masses and ^{56}Ni distributions in Section 5.1. In Section 5.2, we discuss the implications of our results on potential explosion mechanisms of SNe Ia by comparing the light curves of popular explosion models to those produced by our best-fitting Chandrasekhar mass models. In Section 5.3, we discuss our estimate of the intrinsic rate of flux excesses and compare to literature estimates of the rates of early flux excesses. In Section 5.4, we present a summary of the properties of the sample of events with early excesses and discuss these in the context of previous observational and theoretical studies.

5.1 Constraints on ^{56}Ni masses and distributions from early light curves

We find that 67 per cent of SNe Ia in our sample are well matched by a Chandrasekhar mass model. This is similar to the 74 per cent found by M20 for a literature sample of 35 SNe Ia, in particular,

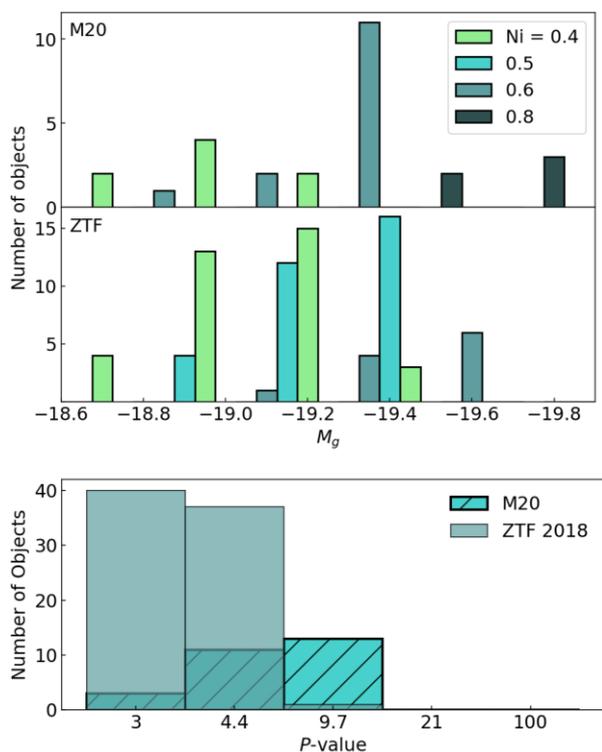


Figure 8. *Top:* The stacked distribution of absolute magnitudes (g band) of the sample in M20, colour coded by the fitted ejected ^{56}Ni mass. The panel below shows the distribution of peak magnitudes of the ZTF 2018 sample colour coded by fitted ^{56}Ni mass. This figure shows that there are four objects that pass our cuts that have a peak magnitude (~ -19.6 mag) that could correspond to a ^{56}Ni mass of $0.8 M_{\odot}$ when compared to the results from M20 but are matched by $0.6 M_{\odot}$ in our analysis. We note that there are models in the grid that reach -19.6 mag with $0.6 M_{\odot}$, and, moreover, the model grid used for fitting in M20 did not contain models with $^{56}\text{Ni} = 0.5 M_{\odot}$, rendering a direct comparison more complicated. *Bottom:* Histogram showing the distribution of P values of the well-matched objects in M20 compared to our sample. It is clear that the ZTF sample is skewed towards lower P values and therefore more extended ^{56}Ni distributions than the sample presented in M20. Nevertheless, both works agree that no normal SN Ia light curve can be reproduced by a model with a highly compact ^{56}Ni distribution ($P = 21, 100$).

when considering that for 18 percent of objects, we were not able to constrain the parameters but are unable to definitively rule out a Chandrasekhar mass model. We further find that ZTF SNe Ia are best matched by models with ^{56}Ni masses distributed between 0.4 and $0.6 M_{\odot}$ (Fig. 4), whereas the objects in M20 are distributed across the 0.4 – $0.8 M_{\odot}$ range, peaking at $0.6 M_{\odot}$. To understand this difference, we compare the peak magnitude distribution of both samples in Fig. 8. We find that the literature sample in M20 has three objects in the brightest magnitude bin, where the ZTF sample has zero. This is likely driven by selection effects: it is easier to identify flux excesses in brighter SNe Ia at earlier phases than fainter ones, and thus brighter objects get published more. Since the peak magnitude is correlated with the ^{56}Ni mass, it is not surprising that the M20 matches several SNe Ia with a ^{56}Ni mass of $0.8 M_{\odot}$. Moreover, the model grid implemented in M20 does not include models with $^{56}\text{Ni} = 0.5 M_{\odot}$, impeding a direct comparison between the distributions. Our ^{56}Ni mass range of 0.4 – $0.6 M_{\odot}$ is in general agreement with the large sample (337 SNe Ia) analysed in Scalzo et al. (2014), where they found that the majority of SNe Ia have ^{56}Ni masses of 0.3 – $0.7 M_{\odot}$.

Our sample of SNe Ia is dominated by model light curve with highly extended ^{56}Ni distributions ($P = 3, 4.4$), whereas the SNe Ia analysed by M20 also match models with more compact ^{56}Ni distributions ($P = 9.7$, Fig. 8). We cannot identify a clear reason for this difference in ^{56}Ni distributions between the literature sample and the ZTF sample. The ZTF SN Ia sample should be relatively unbiased in its selection, while the M20 sample is based on individual objects or small samples of objects that have been published independently. There is a clear difference in the peak absolute magnitudes and associated ^{56}Ni masses (Fig. 8) for the two samples. Since the ZTF sample is larger and lower bias, we take this to be a better representation of the SNe Ia population than the literature sample of M20 and conclude that the ^{56}Ni distributions of SNe Ia that are well matched by Chandrasekhar mass explosion models appear to be highly extended.

We perform a comparison between the rise times computed from the TURTLES models (measured from explosion to maximum light) and those computed by Miller et al. (2020b) (measured from first light to maximum light) and find these to be generally consistent. In approximately half of cases, the rise time computed from the TURTLES models is more than 1 d longer than those computed by Miller et al. (2020b), consistent with the existence of a dark phase. However, due to the breakdown of the LTE assumption around maximum light, the estimates of peak light from the TURTLES models have large uncertainties. The uncertainty on the explosion date and coarseness of the model grid further contributes to the uncertainty in the estimation of the rise time. Similarly, the rise time presented in Miller et al. (2020b) relies on a conversion from the peak in the g band to the B band, introducing non-negligible uncertainties. For these reasons, we do not present any further analysis on the duration of the dark phases.

Bulla et al. (2020) analysed the colour evolution of 65 SNe Ia from the ZTF 2018 data set that had a colour measurement within 5 d of first light. They found that the colour evolution of most objects was relatively flat and covered the full range of ^{56}Ni distributions presented by M20 implying that at least some degree of extended ^{56}Ni is required, but some objects can also be matched by the most compact distributions ($P = 21, 100$), which disagrees with both our findings and the results presented in M20. As noted by M20, a more realistic composition structure would shift colours to redder values meaning that caution must be taken when comparing the absolute colours of the models to observations.

5.2 Implications of extended ^{56}Ni distributions for SN Ia explosion mechanisms

In this section, we discuss the preference for highly extended ^{56}Ni distributions in the ZTF SNe Ia sample in the context of SNe Ia explosion models. In Fig. 9, we compare the light curves of models that were matched to SNe Ia in our sample to theoretical light curves for DDT (Seitenzahl et al. 2013), deflagration (‘DEF’ Fink et al. 2014), and GCD models (Seitenzahl et al. 2016). Fig. 9 shows that the DDT models produce light curves that are well matched by the TURTLES models, specifically those with 40–1600 ignition kernels (n_{40} – n_{1600}), where the n_{100} is generally taken to be representative of normal SNe Ia (Seitenzahl et al. 2013). The higher fraction of ^{56}Ni in the outer regions found in our objects also resembles the DDT models, although these models show a steeper decline in the ^{56}Ni mass fraction near the outer layers (see fig. 1 in M20), which results in a more rapid rise in the early light curve. By delaying the time between the detonation and deflagration, more mixing could bring the DDT models in better agreement with TURTLES models with the most extended ^{56}Ni distributions. As noted

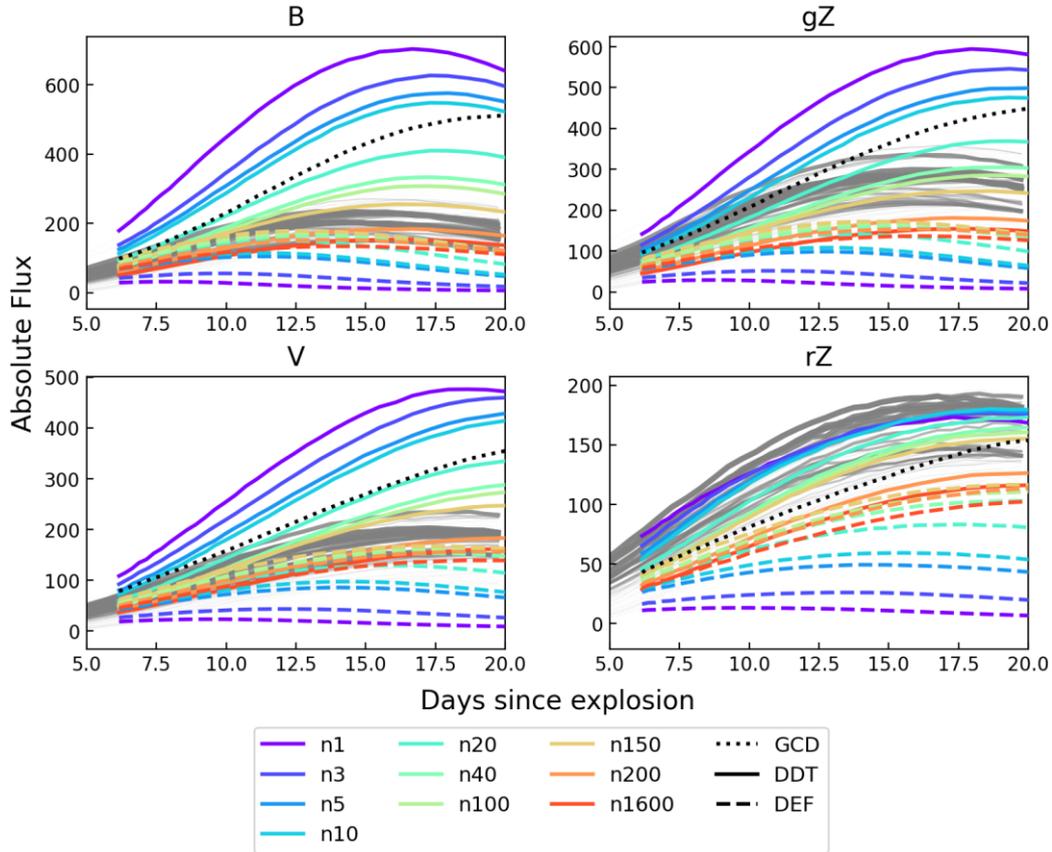


Figure 9. Plots showing the comparison of the TURTLS light curves to DDT (Seitenzahl et al. 2013), DEF (Fink et al. 2014), and GCD models (Seitenzahl et al. 2016). Only TURTLS models that were matched to objects in our well-fit sample are plotted (grey), where the line width indicating the number of times that particular model was matched by an SN Ia. The plots show the light curves in *B*, *V* (Bessel) and the *g*, *r* (ZTF) bands for ignition kernels ranging from $n = 1$ to $n = 1600$. While DDT models produce light curves that are well matched by our models, the DEF models are unable to reproduce the shape of the rise and the peak magnitude.

by M20, however, this could have a significant effect on the ejecta composition, and so further modelling is required to determine which configurations of the DDT models would best reproduce light curves from our model grid with the lowest P values.

It has been shown that DEF models struggle to reproduce the light-curve characteristics (peak magnitude and colour) and distinctive spectral features (Si, S, and Ca) seen in normal SNe Ia (Fink et al. 2014). We also find that the DEF models are unable to reproduce the shape of the rise and the peak magnitudes reached by the TURTLS models. DEF models are a better match to O2cx-like SNe (SNe Iax), a subclass easily distinguished from normal SNe Ia by their faint magnitudes and low ejecta velocities (Kromer et al. 2013; Fink et al. 2014; Leung & Nomoto 2020). Since the deflagration models often leave behind a bound remnant, the ejected ^{56}Ni masses are significantly lower ($0.035\text{--}0.315 M_{\odot}$; Fink et al. 2014), complicating a direct comparison of the absolute magnitudes reached between these models. None the less, the shape of the rising light curve has not been directly compared to a large sample of SNe Ia, and we show that deflagration models result in a flatter rise than the range set by the SNe Ia from the ZTF sample.

The GCD model describes the explosion of a near-Chandrasekhar mass WD, and the light-curve evolution of this model is also shown in Fig. 9. Although this model produces a ^{56}Ni distribution more comparable to TURTLS models with $P = 3$, GCD models tend to produce large amounts of ^{56}Ni ($0.74 M_{\odot}$; Seitenzahl et al. 2016), which is reflected in the high peak magnitudes of the GCD model

in Fig. 9. In addition to being over-luminous, the GCD models are expected to produce highly asymmetrical ejecta and since SNe Ia generally show little polarization indicative of a symmetric explosion geometry (Bulla, Sim & Kromer 2015), the GCD models are unlikely to be the dominating explosion mechanism for normal SNe Ia.

Scalzo et al. (2014) analysed the ejected mass distribution of SNe Ia and suggested that sub-Chandrasekhar mass WDs make a significant contribution to the SNe Ia population. To rule out sub-Chandrasekhar mass WDs for the SNe Ia in this sample, additional modelling would be required. Some models can be speculatively ruled out based on the distinct characteristics of their resulting light curves. Explosions of WD with a thick He-shell that burn to Fe-group elements produce a significant excess (Jiang et al. 2017; Noebauer et al. 2017; Polin et al. 2019), which would allow us to discern them from M_{Ch} explosions. Furthermore, some He-shell detonation mass models predict an early red colour evolution as a result of the line blanketing from the IGE produced in the He-shell (Nugent et al. 1997; Kromer et al. 2010; Wosley & Kasen 2011; Polin et al. 2019), and since the majority of objects in the ZTF sample show a relatively constant colour evolution, apart from six that are identified by Bulla et al. (2020) as having a ‘red’ bump, we speculate that a double detonation is unlikely to be the dominating explosion mechanism for our well-fit sample. However, some double-detonation models have also been shown to reproduce normal SNe Ia properties, without producing a significant excess at early times (Shen et al. 2018; Polin et al. 2019; Gronow et al. 2021;

Magee et al. 2021). In particular, double detonations with He-shells that do not burn to Fe-group elements cannot be excluded (Magee et al. 2021).

5.3 The intrinsic rate of flux excesses in SNe Ia

We combine the detection efficiencies presented in Section 4.3.1 with the observed rate of early flux excesses to estimate the intrinsic rate of events with flux excesses. As noted in Section 3.3, the observed rate is calculated by applying the same cuts on light-curve scatter to the full sample as to the flux excess sample. We use only the lowest redshift bin ($z \leq 0.07$), which contains three SNe Ia (out of 30 events) with a flux excess, to estimate the intrinsic rate of flux excesses. As discussed previously, ZTF18aaqqqs and ZTF18aayjvve are not included in the calculation of the intrinsic rate because we are unable to constrain the efficiency for these weaker flux excesses. We combine the g - and r -band flux excesses because all excesses are visible across both bands. The uncertainty on the intrinsic rate is derived by combining the uncertainty of the detection efficiency and the Poisson uncertainty on the number of detected flux excesses.

For ‘normal’ and 91T/99aa-like SNe Ia, we estimate an intrinsic rate of flux excesses of similar strength to those presented in the literature of 18 ± 11 per cent. This is consistent with the calculated intrinsic rate in the second redshift bin ($0.07 < z \leq 0.087$) of 12 ± 19 per cent, which contains one SN Ia with a flux excess out of 21 objects (observed rate of 5 per cent) in this bin. However, due to the efficiency estimate in this bin dropping below 50 per cent and only a single SN Ia with a detected flux excess, we quote the value from the lowest redshift bin as the intrinsic rate. We note that due to the simplified efficiency calculation and small number of flux excess SNe Ia in the sample, there are likely significant statistical and systematic uncertainties associated with this estimate of the intrinsic rate. A full survey simulation using, e.g. SIMSURVEY (Feindt et al. 2019) would be required for a more accurate estimate of the efficiency of flux excess detections, particularly, at the lower significance end.

Our calculated intrinsic rate of 18 ± 11 per cent of SNe Ia showing a flux excess is similar to that of M20, who found a rate of 22 per cent in a sample of 23 literature objects. Our rate is lower than Jiang et al. (2018), who found a rate of 35 per cent in a sample of 23 objects, although it should be noted that their analysis is focused on characterizing flux excesses rather constraining the rate of flux excesses. Since we have performed our analysis on the much larger unbiased year 1 ZTF SN Ia sample first presented in Yao et al. (2019), it is unsurprising that our rate is lower than previous estimates because of the untargeted nature of our sample. Moreover, the rates estimated by the above-mentioned studies represent absolute observed rates, whereas our measurement is an intrinsic rate. Olling et al. (2015) perform an in-depth study of the high-cadence Kepler light curves of 3 SNe Ia and find no evidence of any flux excesses. In total, there are four Kepler light curves of nearby SNe Ia (KSN 2012a, KSN 2011b, KSN 2011c, SN2018oh; Olling et al. 2015; Dimitriadis et al. 2019; Li et al. 2019; Shappee et al. 2019) with only one showing a flux excess (SN 2018oh), suggesting a rate of 25 per cent. However, we note the small number statistics for this sample.

5.4 Potential origins for early flux excesses in the ZTF SN Ia sample

In Section 4.3, we presented the SNe Ia in the ZTF 2018 sample for which we detected a flux excess and their general light curve and spectral properties. In Section 5.4.1, we present an overview of the properties of the early flux excesses and the relation to progenitor

scenarios, while in Section 5.4.2, we discuss their broader properties and the potential link in some cases to the 91T/99aa-like subclass and younger stellar populations.

5.4.1 Quantifying the diversity of early flux excesses

Fig. 10 shows the $g-r$ colours at 3 d post explosion and lifetimes of the excesses of the six SNe Ia with identified flux excesses. The duration of the bump is calculated assuming a start time of the bump as the median time between the last non-detection and the first non-zero detection of the residuals from the best-fitting model with an uncertainty on the start time of half the time between these points. The end time was calculated in a similar fashion. We also show the $g-r$ colours at 3 d post explosion and lifetimes of the excesses for the literature events, SNe 2017cbv and 2019yvq.

We compared the flux excess sample to models producing a flux excess (companion interaction, CSM interaction, and ^{56}Ni clumps; Kasen 2010; Piro & Morozova 2016; Maeda et al. 2018; Magee & Maguire 2020; Magee et al. 2021). The parameter space covered by the companion interaction and ^{56}Ni -clump models are overplotted in Fig. 10, along with the typical lifetime range for the CSM interaction models. For the companion interaction models, the interaction signature at early times is expected to produce blue emission since most of the emission from the shock is radiated in the UV (Kasen 2010). However, for a realistic comparison with observed events, it must be combined with an underlying ejecta model. Magee et al. (submitted) implemented the analytical formulae from Kasen (2010) to add a companion interaction component to an underlying Chandrasekhar mass model light curve. These models resulted in a $g-r$ range of -0.3 to -0.45 mag at 3 d post explosion and lifetimes of 2–7 d.

Piro & Morozova (2016) provided no direct $g-r$ colours for the CSM interaction models, but they predict a $B - V$ colour range of 0.1–0.4 mag during the interaction-powered flux excess. Assuming a blackbody-like spectrum at 3 d post explosion (near the end of the short-duration CSM interaction), we expect the $g-r$ colours to be 0.25-mag redder but again are complicated by the contribution of the underlying ejecta (see discussion in Piro & Morozova 2016). The lifetimes of flux excesses caused by CSM interaction are generally short, ranging from 1 to 3 d. The ^{56}Ni -clump models of Magee & Maguire (2020) also provide predictions for the early colour and duration of any flux excess. For this scenario, we exclude models that produce $g-r < -0.5$ mag near maximum light for being unrealistic matches for our sample with the remaining models predicting $g-r = -0.13$ – -0.34 at ~ 3 d post explosion and lifetimes of the flux excess of 3.6–7.0 d. We note that these models were produced to match SN 2017cbv and SN 2018oh, and therefore these models represent only a sub-section of the full parameter space. Double detonations models that burn to IGE in the shell predict a ‘red bump’ in the early light curve as the result of line blanketing (Polin et al. 2019; Bulla et al. 2020; Magee et al. 2021). However, we do not include these predictions in Fig. 10 because these models are also predicted to produce very red maximum light spectra that are not seen for these events. Double detonation models that burn to only IMEs look normal at maximum light but they are not predicted to have flux excesses at early times.

We have found that the early flux excesses have a variety of strengths, colours, and lifetimes (see Figs 5 and 10). The majority of the objects with a detected flux excess in our sample (apart from ZTF18abxxssh) do not show the distinct rise, followed by a drop/plateau as seen in, e.g. SN 2019yvq (Miller et al. 2020a), but

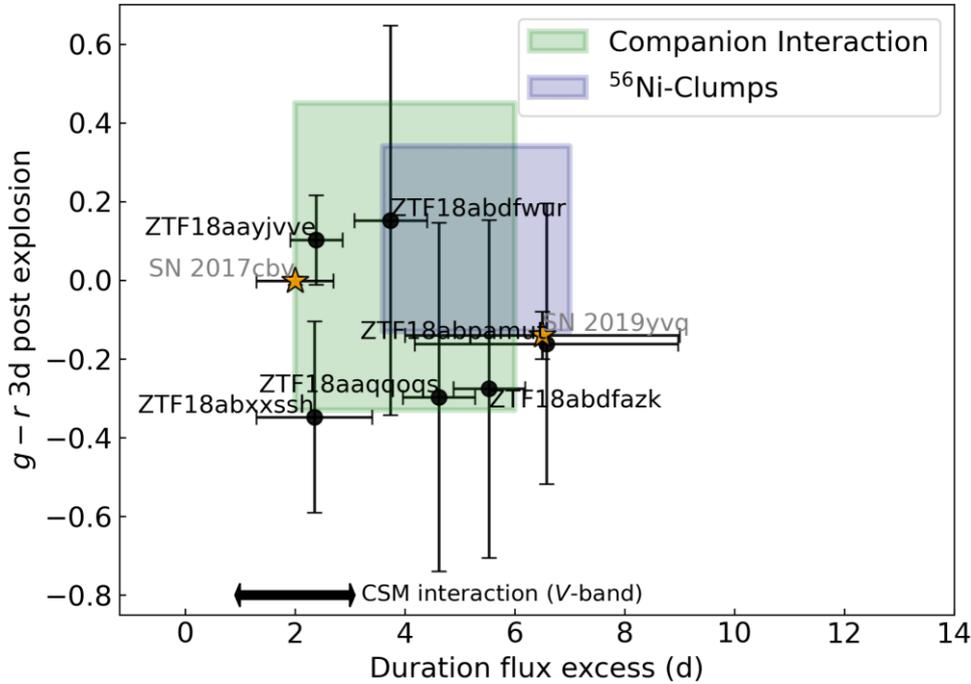


Figure 10. A plot showing the $g-r$ colour at 3 d post explosion of the SNe Ia with a detected flux excess, as a function of the duration of the flux excess. The shaded regions show the parameter space occupied by companion interaction (Kasen 2010; Magee et al. 2021) and ^{56}Ni -clump models (Magee & Maguire 2020). The range of lifetimes expected from CSM interaction models produced by Piro & Morozova (2016) is shown but no $g-r$ colour predictions are made. The $g-r$ colour and flux excess lifetimes for SN 2017cbv and SN 2019yvq are also plotted. The durations of the flux excesses for four of the events (ZTF18aaqqoqs, ZTF18abdfzfk, ZTF18abdfwur, ZTF18abpamut) appear to be longer than those predicted by the CSM interaction models of Piro & Morozova (2016).

four events (ZTF18abdfzfk, ZTF18abpamut, ZTF18aaqqoqs, and ZTF18aavrwhu) show flux excesses that are of a comparable strength to the literature objects, SN 2017cbv and SN 2018oh. ZTF18abxxssh shows the most prominent shoulder (Fig. 5) at early times of the sample. All of our events with early flux excesses rise earlier and more rapidly than predicted by the underlying Chandrasekhar mass models, producing a shoulder in the light curve and, therefore, likely need an external source of energy to produce them.

The peak magnitude of the excess can range between 2 and 11 per cent of the flux at maximum light (the lower limit is set by the detection method), have $g-r$ colours during the flux excess of -0.4 to 0.2 mag, and have lifetimes of the excesses of 2–7 d. For two of the events in our sample (ZTF18aayjvve, ZTF18abxxssh), the lifetimes of the bumps are short at ~ 2 d and are most consistent with the predictions of the companion (or CSM) interaction models (Kasen 2010; Piro & Morozova 2016) and are inconsistent with the ^{56}Ni clump model lifetime predictions. They fall in the region covered by the colours predicted for the companion interaction model (no $g-r$ colours are available for the CSM interaction models). For the other four events with flux excesses, the lifetimes are inconsistent with the CSM interaction predictions but overlap in the colour and lifetime space of the companion interaction and ^{56}Ni clumps models. Conclusively determining the cause of a flux excess based on the light curve without spectra obtained during the excess is difficult, although we are able to rule out CSM interaction as a cause for excesses lasting more than 3 d, corresponding to four events in our sample (ZTF18abdfzfk, ZTF18abdfwur, ZTF18abpamut, and ZTF18aaqqoqs).

We find that three (ZTF18abdfzfk, ZTF18abxxssh, and ZTF18aaqqoqs) of the four SNe Ia displaying a blue colour at 3 d post explosion also have a negative c value at maximum light, following

the previous results for SNe 2017cbv and 2018oh that had bluer than average colours (Hosseinzadeh et al. 2017; Dimitriadis et al. 2019; Li et al. 2019; Shappee et al. 2019). ZTF18abxxssh and ZTF18aaqqoqs also have very broad light curves and maximum light spectra consistent with overluminous SNe Ia (see Section 5.4.2 for further comparison with such events). The two events (ZTF18abdfwur, ZTF18aayjvve) with red early colours also show red colours at maximum light but not extreme values. An interesting conclusion of our model and data comparisons of the colour of the early flux excesses is that although the physical processes involved in the production of the flux excesses may produce hot blue emission, this needs to be combined with the underlying SN ejecta component that can result in an overall redder colour. As shown in Fig. 10, the parameter space of companion interaction and ^{56}Ni clump model predictions with an Chandrasekhar mass model ejecta cover relatively red, as well as blue, early light colours. This suggests that studies should be careful to not exclude early flux excess scenarios that predict blue colour solely because the observed colour is redder. Spectra are key for distinguishing between the origins of flux excesses, and rapid spectroscopy should be a focus of future flux excess studies.

5.4.2 A connection to young stellar populations?

The sample size is very small but as demonstrated in Section 4.3, the SNe Ia with flux excesses in our sample have a preference for lower stellar mass hosts (Fig. 6). This suggests that these objects originate from younger stellar populations (Rigault et al. 2020), which is likely more consistent with the SD scenario, or the CD scenario. Alternatively, since lower mass galaxies tend to host brighter SNe Ia (e.g. those with broader light curves; Sullivan et al. 2010), it is easier to detect brighter SNe Ia at early times. However, our method

measures the relative flux excess between the data and the model, and assuming that the luminosity of the flux excess does not scale with the luminosity of the SNe explosion, it should be easier to identify a bump in fainter SNe Ia, or one with its ^{56}Ni distribution more constrained towards the centre (high P value). This is not evident from our small sample, which is populated by both faint and bright SNe Ia, making it unclear whether brighter, more energetic, SNe Ia explosions produce more prominent bumps.

Jiang et al. (2018) performed a study of the light curves of SNe Ia from the literature with early flux excesses and categorized them based on the shapes and evolution of their early light curves. They found that all three 91T-like events in their sample have strong flux excess in their early light lasting over a week and the three 99aa-like events in their sample have either broad early light curves or flux excesses lasting up to a week. Of our six events with early flux excesses, we identify one object as a 99aa-like SN Ia (ZTF18aaqqoqs) and one as a 91T-like SN Ia (ZTF18abxxssh). In keeping with the typical characteristics of these classes, these two events have some of the broadest light curves (high x_1 values) of the sample and are found in lower stellar mass hosts than average. Their early bumps are also the two bluest of the sample and they also show blue colours at maximum light. However, the lifetimes of their early excesses differ with ZTF18aaqqoqs being significantly longer (4.6 ± 0.7 d) than the flux excess lifetime of ZTF18abxxssh (2.3 ± 1.1 d). Therefore, it is unclear if they have the same origin for their early flux excesses but their overlap in other properties is of note.

We also wish to determine if we can confirm the result of Jiang et al. (2018) that all 91T/99aa-like SNe Ia have some form of excess at early times. To do this, we focus on the early light curves of SNe Ia in our full sample that are classified as 99aa/91T-like but were not identified by our method as having an early flux excess. We find 10 such events. Seven of these (ZTF18aaumeys, ZTF18aaxakh, ZTF18abixjey, ZTF18abkljlv, ZTF18abmmkaz, ZTF18abpmpo, and ZTF18abrznb) have insufficient early data to detect a flux excess or the uncertainties in the early data are too large. One event (ZTF18abauprj) does not show a flux excess but is not well matched to any model at peak, which may affect the fit at early times. Two events (ZTF18aaytovs, ZTF18abfwuwn) have sufficient early-time data to search for an early excess but show no excess suggesting that not all 99aa/91T-like events in an unbiased sample show an early excess. We can estimate a rough rate of 91T/99aa-like events showing early flux excesses using the observed rates and our detection efficiencies. The lowest redshift bin ($0 < z < 0.07$) of our sample contains four 99aa/91T-like SNe Ia (including our reclassification of ZTF18abxxssh). We identify one flux excess in this lowest redshift bin (ZTF18abxxssh), setting a lower limit on the observed rate of flux excesses for 91T/99aa-like SNe Ia at 25 per cent. The intrinsic rate derived from the four 91T/99aa-like SNe Ia in the first redshift bin is 44 ± 13 per cent when considering the detection efficiency (57 per cent) in this bin. This rate should be considered with caution since there is only one 91T/99aa-like SN Ia with a detected flux excess in this redshift bin. We do not confirm the results found by Jiang et al. (2018) that 100 per cent of 91T/99aa-like SNe Ia have some form of excess, but the frequency does appear to exceed that of normal SNe Ia and we encourage future studies with larger samples of these sub-classes to investigate this in greater detail.

6 CONCLUSIONS

We fit 115 SNe Ia light curves from the ZTF 2018 sample (Yao et al. 2019; Miller et al. 2020b) with a grid of Chandrasekhar mass explosions (Magee et al. 2018, 2020) with varying degrees of ^{56}Ni

mixing and find that approximately 67 per cent of objects can be fit by a model from our grid with ^{56}Ni masses in the range of 0.4–0.6 M_{\odot} . We are unable to adequately constrain model parameters for 19 per cent of the sample, and only 11 per cent could not be matched by a Chandrasekhar mass model. We find that no light curves can be fit by the most compact ^{56}Ni distributions ($P = 21, 100$), confirming the result found by M20, although our sample shows a stronger preference for the most highly extended distribution (lowest P values). The best-fitting models from our grid are found to be generally consistent with the DDT of Seitzzahl et al. (2013), although the steep decline towards the outer layers is not matched by the TURTLES models, and further explosion modelling is needed to explore this.

We performed a search for the presence of flux excesses in the early light curves of the sample and find six objects that display a flux excess. We calculate detection efficiencies from a sample of 10 000 simulated light curves in the g and r bands. Based on three SNe Ia with a flux excess out of 30 SNe Ia (detected rate of 10 per cent) in the lowest redshift bin ($z \leq 0.07$), and a detection efficiency of 57 per cent in this redshift bin, we find an intrinsic rate of strong flux excesses in SNe Ia of 18 ± 11 per cent. This is consistent with the intrinsic rate derived from the $0.07 < z \leq 0.087$ redshift bin of 12 ± 19 per cent, which contained one SN Ia with a flux excess out of 21 SNe Ia (5 per cent). We analysed the $g-r$ colours and lifetimes of the flux excesses and found that these are generally consistent with the ranges predicted by both interaction (companion or CSM) and ^{56}Ni clump models, although longer lasting excesses are unlikely to be caused by CSM interaction. SNe Ia displaying a flux excesses tend to occur in lower mass galaxies. We could be missing flux excesses because the SNe Ia are not detected early enough, which highlights the importance of finding infant SNe within 3 d of explosion. Moreover, flux excesses are predicted to peak in X-rays, and a strong signal is expected in the UV/NUV. The various models predict some of the signal to leak into the optical bands, enabling us to detect the flux excesses. However, with more early UV/NUV data, the detection efficiencies could be significantly improved, particularly since the contribution from the underlying SN Ia ejecta will be weaker in the UV/NUV allowing a cleaner detection.

We encourage future searches for early flux excesses with 1-d cadence surveys (e.g. ZTF-II) and UV/NUV surveys. This will allow us to detect SNe Ia flux excesses at a higher efficiency, and catching them earlier will enable us to trigger spectroscopic follow-up. Rapid spectroscopic follow-up is crucial and could provide us with the information needed to discern between the different origins of the flux excesses (e.g. interaction with a companion or CSM, He-shell detonations, or ^{56}Ni clumps in the outer ejecta). Further analysis of the properties of flux excesses provides a promising method to unveiling the progenitors and explosion scenarios of SNe Ia.

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DATA AVAILABILITY

The TURTLS model grid implemented in this study is available at <https://github.com/MarkMageeAstro/turtles-Light-curves> and is described in detail in Magee et al. (2018). All the ZTF light curves from the sample are described in Yao et al. (2019) and any public data are available on the ZTF GROWTH Marshall (Kasliwal et al. 2019). The forced photometry light curves will be shared on reasonable request to the corresponding author. The simulated light curves produced for the efficiency analysis are available at https://github.com/deckersm/ZTF2018_Deckers_Paper.

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APPENDIX A

The g - and r -band light curves and the matched TURTLES models are shown in Figs A1–A4. We present the host galaxy data in Tables A1 and A2. A summary of the TURTLES fits, as well as the SALT2 parameters derived in Yao et al. (2019), is presented in Tables A3 and A4.

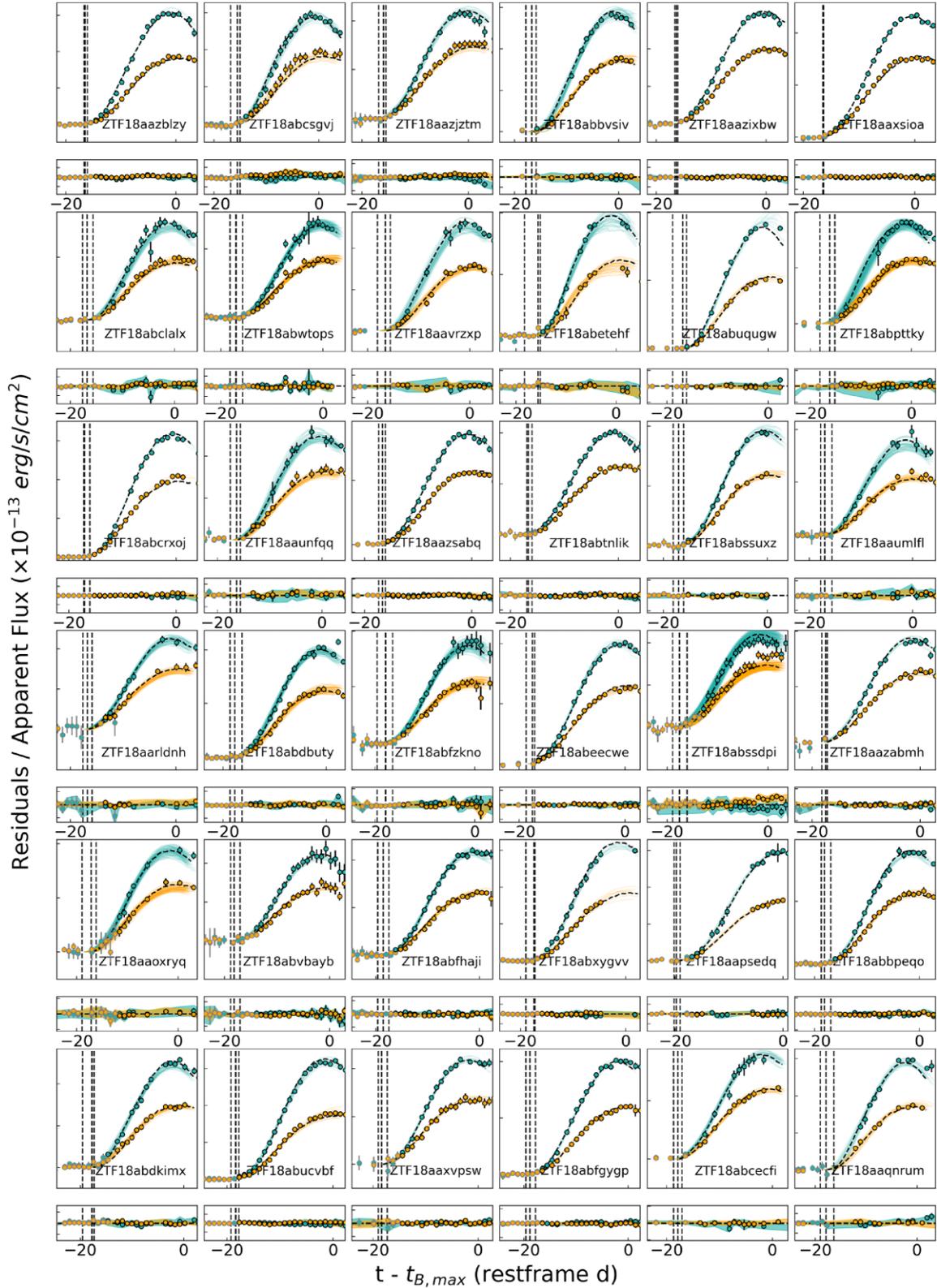


Figure A1. The light curves of the well-fit SNe Ia in the sample, ordered alphabetically.

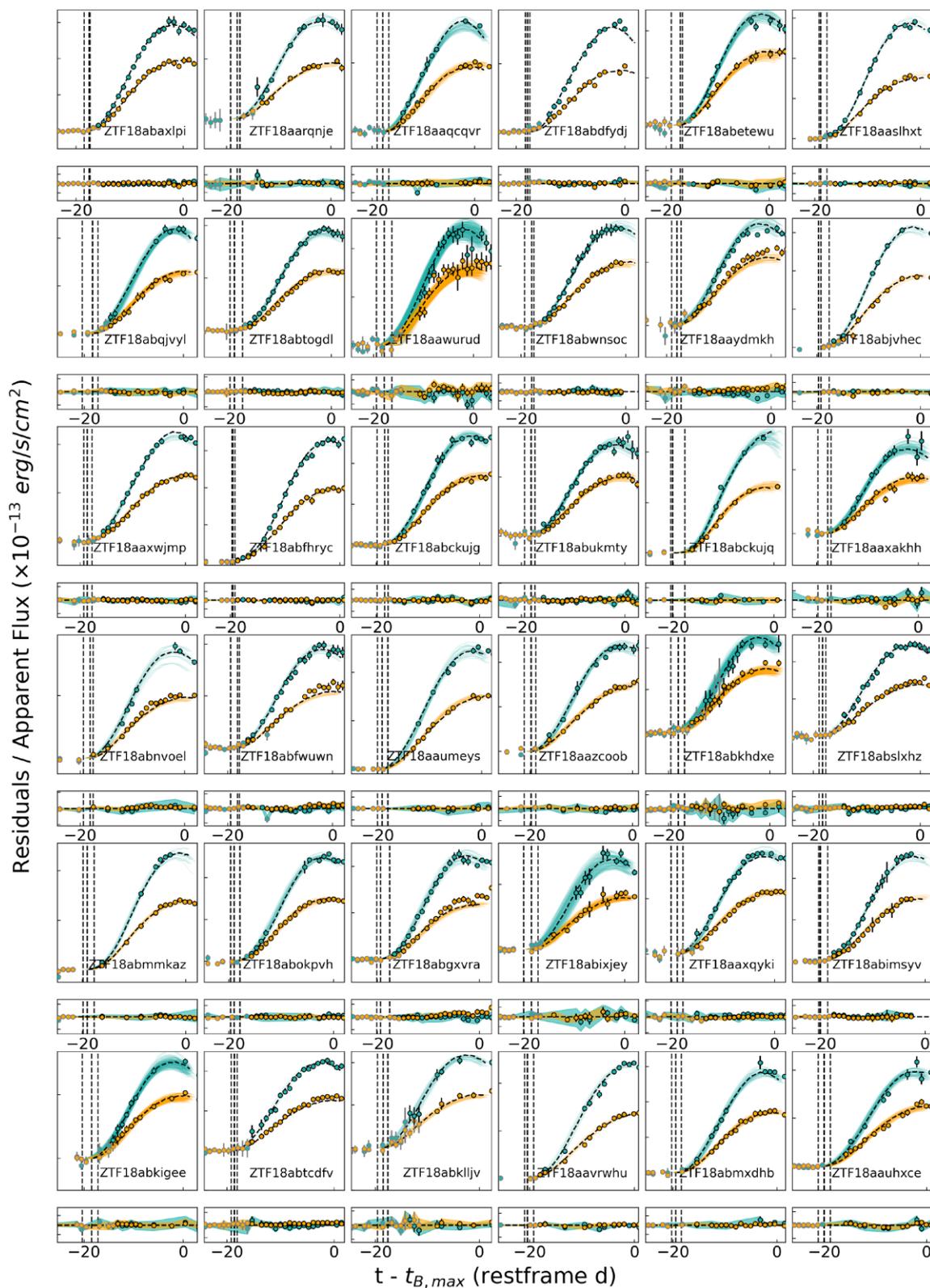


Figure A2. Continuation of Fig. A1.

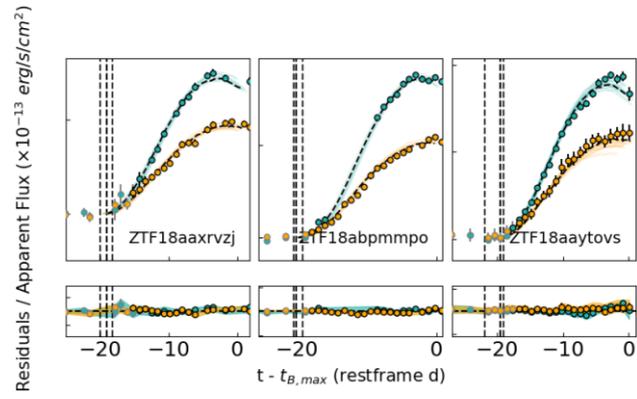


Figure A3. Continuation of Fig. A1.

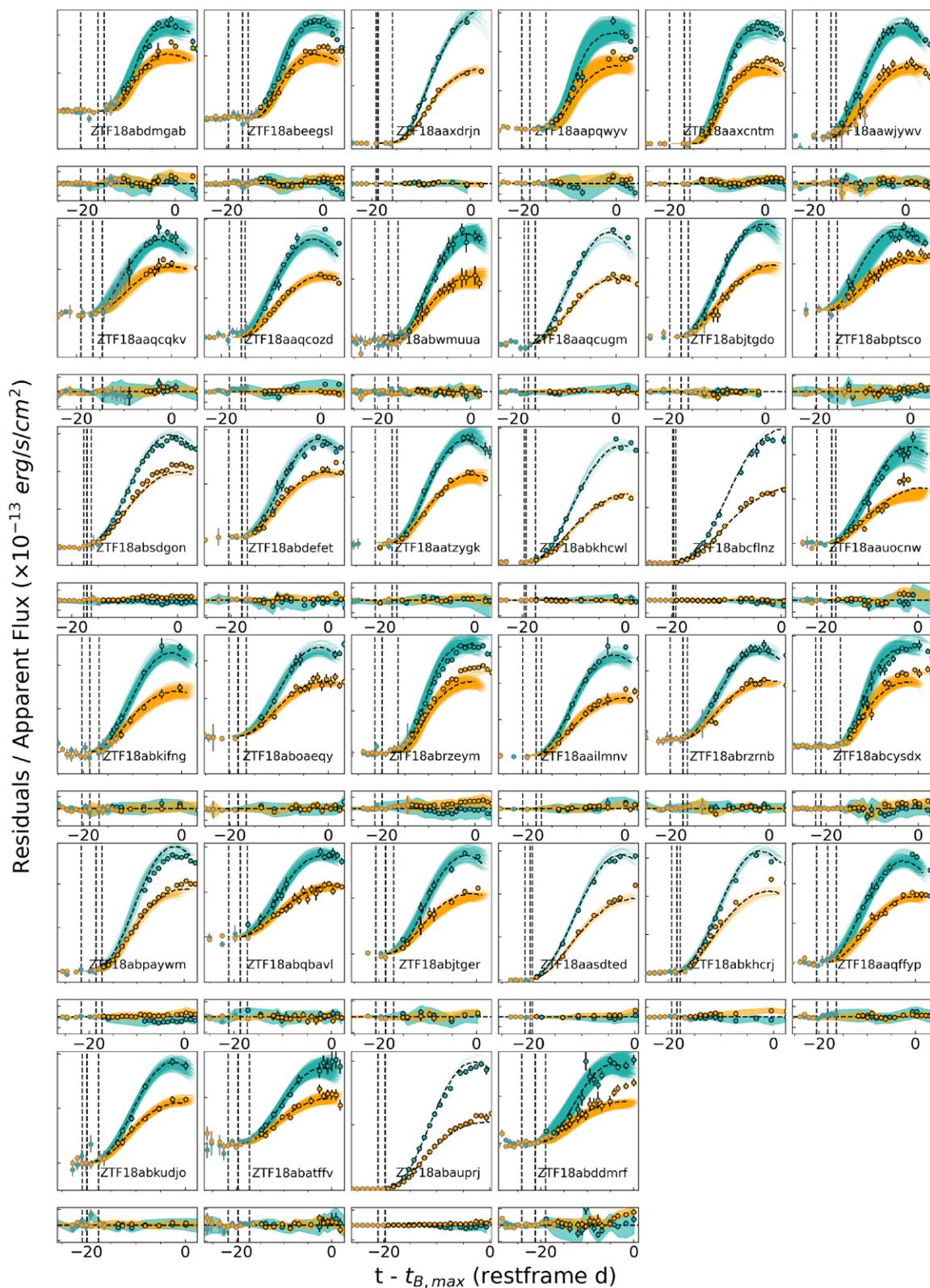


Figure A4. The light curves of the badly fit SNe Ia in the sample, ordered alphabetically.

Table A1. Values of the mean host galaxies stellar mass and uncertainties for the sample, as well as the DLR calculated for between each object and its identified host galaxy.

Object	Host galaxy stellar mass	DLR
ZTF18aailmnv	10.150 ± 0.319	0.925
ZTF18aaoxryq	10.374 ± 0.351	0.126
ZTF18aapqwyy	10.077 ± 0.339	0.161
ZTF18aapsedq	11.640 ± 0.383	1.913
ZTF18aaqcozd	11.657 ± 0.355	2.851
ZTF18aaqcqkv	9.854 ± 0.503	0.509
ZTF18aaqcqvr	9.457 ± 0.348	1.190
ZTF18aaqcugm	10.671 ± 0.403	5.718
ZTF18aaqffyp	9.342 ± 0.298	1.810
ZTF18aaqnrnm	8.433 ± 0.762	4.221
ZTF18aaqqoqs	8.863 ± 0.605	1.908
ZTF18aarldnh	11.018 ± 0.270	1.590
ZTF18aarqnje	9.202 ± 0.546	3.391
ZTF18aasdted	8.969 ± 0.385	0.652
ZTF18aaslhtx	10.034 ± 0.186	1.778
ZTF18aatzygk	10.201 ± 0.307	0.182
ZTF18aauhxce	10.555 ± 0.227	3.445
ZTF18aaumeys	9.013 ± 0.338	1.877
ZTF18aaumlfl	10.524 ± 0.348	2.626
ZTF18aaunfqq	9.742 ± 0.397	0.773
ZTF18aauoconw	10.288 ± 0.403	1.522
ZTF18aavrwhu	–	1.573
ZTF18aavrzxp	10.460 ± 0.422	1.150
ZTF18aawjyvw	10.697 ± 0.353	0.524
ZTF18aawurud	10.842 ± 0.324	0.395
ZTF18aaxakhh	10.184 ± 0.334	0.774
ZTF18aaxcntm	10.904 ± 0.401	0.800
ZTF18aaxdrjn	10.679 ± 0.326	2.211
ZTF18aaxqyki	–	27.734
ZTF18aaxrvzj	7.820 ± 1.540	1.550
ZTF18aaxsioa	9.073 ± 0.253	1.277
ZTF18aaxvpsw	10.671 ± 0.391	0.241
ZTF18aaxwjmp	8.933 ± 0.564	0.262
ZTF18aaydmkh	10.204 ± 0.321	0.378
ZTF18aayjvve	10.663 ± 0.408	2.362
ZTF18aaytovs	10.779 ± 0.348	0.936
ZTF18aazabmh	–	36.070
ZTF18aazblzy	8.638 ± 0.671	16.070
ZTF18aazcoob	10.654 ± 0.361	2.661
ZTF18aazixbw	10.561 ± 0.326	1.421
ZTF18aazjztm	10.345 ± 0.347	0.669
ZTF18aazsabq	9.876 ± 0.367	0.786
ZTF18abatffv	10.409 ± 0.499	2.577
ZTF18abaupri	10.758 ± 0.333	3.989
ZTF18abaxlpi	10.026 ± 0.310	0.887
ZTF18abbpeqo	10.585 ± 0.308	1.204
ZTF18abbvsiv	10.647 ± 0.341	1.303
ZTF18abcecfi	10.421 ± 0.366	1.641
ZTF18abcflnz	9.682 ± 0.323	3.049
ZTF18abcckujg	10.041 ± 0.225	2.529
ZTF18abcckujq	10.863 ± 0.467	2.731
ZTF18abclalx	9.869 ± 0.463	1.559
ZTF18abcrxoj	–	1.376
ZTF18abcsqvj	10.373 ± 0.377	0.931
ZTF18abcysdx	10.510 ± 0.371	1.505
ZTF18abdbuty	10.237 ± 0.348	0.317
ZTF18abdmdrf	9.197 ± 0.946	3.000
ZTF18abdfazk	9.350 ± 0.271	1.446

Table A2. Table A1 continued.

Object	Host galaxy stellar mass	DLR
ZTF18abdfwur	9.589 ± 0.376	0.137
ZTF18abdfydj	8.801 ± 0.536	1.140
ZTF18abdtkimx	9.571 ± 0.387	0.764
ZTF18abdmgab	10.919 ± 0.325	17.923
ZTF18abeecwe	10.118 ± 0.334	0.918
ZTF18abeegsl	8.847 ± 0.664	2.966
ZTF18abetehf	10.473 ± 0.393	1.609
ZTF18abetewu	9.793 ± 0.350	0.299
ZTF18abfgygp	9.482 ± 0.504	0.601
ZTF18abfhaji	8.623 ± 0.798	3.552
ZTF18abfhrzc	9.382 ± 0.073	3.567
ZTF18abfwuwn	9.250 ± 0.591	3.794
ZTF18abfzkno	9.316 ± 0.523	0.113
ZTF18abgxvra	10.238 ± 0.410	7.141
ZTF18abimsyv	9.353 ± 0.455	1.792
ZTF18abixjey	8.281 ± 1.186	1.301
ZTF18abjstcm	10.941 ± 0.367	0.109
ZTF18abjtgdg	10.555 ± 0.346	0.823
ZTF18abjtger	9.798 ± 0.538	1.110
ZTF18abjvhcj	10.711 ± 0.410	1.413
ZTF18abkherc	9.993 ± 0.351	0.183
ZTF18abkhcwl	9.351 ± 0.338	1.137
ZTF18abkhdxe	9.530 ± 0.718	0.924
ZTF18abkifng	10.728 ± 0.353	0.337
ZTF18abkigee	7.925 ± 1.389	0.360
ZTF18abklljv	8.038 ± 1.666	6.409
ZTF18abkudjo	10.204 ± 0.390	1.770
ZTF18abmmkaz	9.776 ± 0.380	1.304
ZTF18abmxdhh	7.819 ± 1.133	0.702
ZTF18abnvoel	10.114 ± 0.344	0.584
ZTF18aboaeqy	9.549 ± 0.565	1.043
ZTF18abokpvh	7.235 ± 1.516	12.803
ZTF18abpamut	9.388 ± 0.410	0.958
ZTF18abpaywm	9.946 ± 0.429	1.727
ZTF18abpmpo	8.715 ± 0.550	1.496
ZTF18abptsco	10.084 ± 0.458	0.305
ZTF18abpttky	9.750 ± 0.377	0.283
ZTF18abqbavl	10.019 ± 0.448	2.496
ZTF18abqjvyl	8.803 ± 0.629	0.566
ZTF18abrzeym	9.532 ± 0.366	0.409
ZTF18abrzrnb	8.602 ± 1.119	2.519
ZTF18absdgon	10.210 ± 0.403	0.922
ZTF18abslxhz	9.649 ± 0.575	0.673
ZTF18absdpi	10.355 ± 0.373	0.162
ZTF18abssexz	–	0.218
ZTF18abtcdfv	8.277 ± 1.478	2.419
ZTF18abtnlik	10.382 ± 0.366	3.543
ZTF18abtogdl	8.947 ± 0.483	0.315
ZTF18abucvbf	9.200 ± 0.331	1.151
ZTF18abukmty	9.326 ± 0.474	0.433
ZTF18abuqugw	9.248 ± 0.380	0.244
ZTF18abvbayb	10.801 ± 0.377	3.175
ZTF18abwmuua	10.527 ± 0.383	0.254
ZTF18abwnsoc	8.124 ± 1.481	13.050
ZTF18abwtpos	10.192 ± 0.419	4.662
ZTF18abxssh	8.645 ± 0.436	5.987
ZTF18abxygvv	9.048 ± 0.483	0.332

Table A3. A table summarizing the properties of SNe Ia in our sample. We have included the SNe Ia, which could not be well fit by our models, but in these cases, the model parameters are not indicative of the true SN parameters.

Name ^a	z^b	Model ^c	Exp. date ^d	μ	x_1	c	χ^2	# Unique fits
ZTF18aailmnv	0.080	EXP_Ni0.6_KE2.18.P3	2276211.2 ± 3.6	37.71 ± 0.4	0.4	−0.06	2.3	26
ZTF18aaoxryq	0.0940	EXP_Ni0.5_KE2.18.P3	2247096.7 ± 2.7	37.95 ± 0.3	−0.2	0.02	1.1	10
ZTF18aapqwyyv	0.0560	EXP_Ni0.4_KE0.78.P21	2327864.7 ± 5.0	36.85 ± 0.5	−1.7	0.17	8.0	40
ZTF18aapsedq	0.0650	EXP_Ni0.4_KE1.10.P4.4	2308194.3 ± 1.1	37.16 ± 0.3	−0.1	−0.09	0.6	4
ZTF18aaqcozd	0.0732	DPL_Ni0.4_KE2.18.P3	2290994.3 ± 3.0	37.46 ± 0.4	−1.2	−0.1	6.3	19
ZTF18aaqcqkv	0.1174	DPL_Ni0.4_KE2.18.P3	2194847.7 ± 4.2	38.64 ± 0.6	−1.3	−0.04	2.2	14
ZTF18aaqcqvr	0.070	DPL_Ni0.4_KE2.18.P4.4	2297418.5 ± 2.1	37.24 ± 0.3	0.2	−0.02	2.8	16
ZTF18aaqcugm	0.066	DPL_Ni0.4_KE2.18.P4.4	2319092.1 ± 2.0	36.95 ± 0.3	−1.1	−0.09	7.0	5
ZTF18aaqffyp	0.070	DPL_Ni0.4_KE1.81.P3	2297416.1 ± 4.0	37.36 ± 0.4	0.9	0.06	3.9	25
ZTF18aaqnrum	0.066	DPL_Ni0.4_KE2.18.P4.4	2306038.3 ± 2.5	37.35 ± 0.3	0.1	−0.03	3.3	6
ZTF18aaqqoqs	0.082	EXP_Ni0.5_KE1.10.P4.4	2271941.5 ± 1.2	37.46 ± 0.0	1.2	−0.01	0.9	2
ZTF18aarldnh	0.1077	EXP_Ni0.4_KE2.18.P3	2219188.3 ± 1.7	38.22 ± 0.4	−1.0	0.07	0.7	8
ZTF18aarqnje	0.117	EXP_Ni0.6_KE2.18.P3	2200752.6 ± 1.8	38.29 ± 0.0	0.2	−0.18	1.0	4
ZTF18aasdted	0.0182	EXP_Ni0.4_KE1.40.P3	2414424.0 ± 1.3	34.64 ± 0.2	0.8	0.16	13.9	7
ZTF18aaslhxt	0.0550	EXP_Ni0.6_KE2.18.P4.4	2319098.2 ± 1.4	36.98 ± 0.1	0.3	−0.11	1.7	3
ZTF18aatzygk	0.077	EXP_Ni0.5_KE2.18.P3	2282492.9 ± 4.6	37.75 ± 0.5	0.0	0.05	2.3	46
ZTF18aauhxce	0.0831	EXP_Ni0.4_KE1.10.P3	2269684.1 ± 2.3	37.54 ± 0.4	1.3	0.12	1.8	9
ZTF18aaumeyys	0.0365	EXP_Ni0.5_KE1.10.P3	2369304.1 ± 2.0	35.85 ± 0.4	0.7	−0.02	3.6	15
ZTF18aaumlfl	0.0874	DPL_Ni0.4_KE2.18.P4.4	2260727.9 ± 2.0	37.87 ± 0.3	−1.1	−0.09	1.5	9
ZTF18aauinfqq	0.0711	EXP_Ni0.5_KE2.18.P3	2295092.5 ± 1.8	37.65 ± 0.4	−1.2	0.01	1.3	10
ZTF18aauocnw	0.102	EXP_Ni0.6_KE2.18.P3	2230717.4 ± 3.8	37.97 ± 0.2	0.1	−0.05	2.7	49
ZTF18aavrwhu	0.0620	EXP_Ni0.5_KE0.78.P4.4	2314718.7 ± 1.6	36.83 ± 0.1	1.2	−0.02	2.2	3
ZTF18aavrzxp	0.069	DPL_Ni0.4_KE2.18.P4.4	2299582.2 ± 2.3	37.45 ± 0.3	−1.4	−0.07	2.7	7
ZTF18aawjyvv	0.0509	DPL_Ni0.4_KE2.18.P3	2363707.4 ± 3.8	36.68 ± 0.6	−1.5	−0.01	2.4	23
ZTF18aawurud	0.0531	EXP_Ni0.4_KE2.18.P3	2341246.2 ± 3.0	36.64 ± 0.6	0.3	0.14	1.2	16
ZTF18aaxakhh	0.117	EXP_Ni0.5_KE2.18.P3	2200759.5 ± 2.5	38.33 ± 0.2	0.5	−0.08	1.3	9
ZTF18aaxcntm	0.0269	EXP_Ni0.4_KE1.68.P100	2394119.0 ± 4.2	35.51 ± 0.4	−1.5	0.1	27.0	26
ZTF18aaxdrjn	0.0340	EXP_Ni0.5_KE0.50.P9.7	2377657.9 ± 3.3	35.56 ± 0.6	−1.9	−0.08	4.1	32
ZTF18aaxqyki	0.1003	EXP_Ni0.5_KE1.40.P3	2234787.1 ± 2.2	37.97 ± 0.2	0.9	−0.05	1.2	10
ZTF18aaxrvzj	0.114	EXP_Ni0.5_KE1.81.P3	2206700.8 ± 1.8	38.23 ± 0.0	1.4	−0.01	0.9	6
ZTF18aaxsioa	0.0315	EXP_Ni0.4_KE2.18.P3	2386669.7 ± 0.2	35.69 ± 0.0	−1.5	0.05	2.3	1
ZTF18aaxvpsw	0.0916	EXP_Ni0.5_KE2.18.P4.4	2251985.8 ± 1.7	37.85 ± 0.1	0.0	−0.01	0.6	6
ZTF18aaxwjmp	0.084	DPL_Ni0.4_KE1.81.P4.4	2267775.6 ± 1.5	37.56 ± 0.3	0.4	−0.01	0.9	11
ZTF18aaydmkh	0.077	EXP_Ni0.4_KE2.18.P3	2282515.7 ± 0.78	37.83 ± 0.1	0.3	0.19	2.8	4
ZTF18aayjvve	0.0474	EXP_Ni0.4_KE1.68.P3	2347316.1 ± 0.4	36.43 ± 0.1	−0.1	0.06	2.1	4
ZTF18aaytovs	0.0746	EXP_Ni0.5_KE1.81.P3	2276179.5 ± 2.8	37.25 ± 0.1	2.0	0.12	1.4	10
ZTF18aazabmh	0.0746	EXP_Ni0.4_KE2.18.P3	2291027.2 ± 1.0	37.67 ± 0.1	−0.3	0.11	1.4	3
ZTF18aazblzy	0.0653	DPL_Ni0.4_KE2.18.P4.4	2319127.3 ± 0.6	37.29 ± 0.0	−1.7	−0.07	1.7	1
ZTF18aazcoob	0.0845	EXP_Ni0.5_KE1.10.P3	2259443.2 ± 2.2	37.62 ± 0.1	0.7	0.06	2.2	5
ZTF18aazixbw	0.0594	EXP_Ni0.4_KE2.18.P3	2320440.9 ± 0.5	37.16 ± 0.0	−1.6	0.05	1.3	1
ZTF18aazjztm	0.0721	EXP_Ni0.5_KE2.18.P3	2293166.6 ± 1.3	37.77 ± 0.3	−1.7	0.09	1.3	4
ZTF18aazsabq	0.060	EXP_Ni0.4_KE2.18.P3	2319129.2 ± 1.2	37.01 ± 0.3	−1.2	0.03	1.5	3
ZTF18abatffv	0.143	EXP_Ni0.5_KE1.40.P4.4	2150720.4 ± 4.4	38.85 ± 0.3	1.0	−0.01	1.2	46
ZTF18abauprj	0.0242	EXP_Ni0.6_KE2.18.P4.4	2391325.3 ± 1.4	34.89 ± 0.2	1.3	−0.01	6.9	4
ZTF18abaxlpi	0.0642	EXP_Ni0.5_KE2.18.P3	2309111.6 ± 0.7	37.46 ± 0.3	0.1	0.06	0.7	3
ZTF18abbpeqo	0.0667	EXP_Ni0.5_KE2.18.P4.4	2304394.4 ± 1.8	37.29 ± 0.3	−0.1	0.02	0.8	7
ZTF18abbvsiv	0.051	DPL_Ni0.4_KE2.18.P4.4	2338991.4 ± 1.9	36.6 ± 0.3	−1.7	0.01	2.3	17
ZTF18abcecfi	0.079	EXP_Ni0.4_KE2.18.P3	2278297.1 ± 1.4	37.72 ± 0.3	0.1	0.09	1.1	6
ZTF18abcflnz	0.0273	EXP_Ni0.5_KE0.78.P4.4	2400670.4 ± 0.5	35.07 ± 0.3	0.1	−0.03	5.8	4
ZTF18abcujukg	0.075	EXP_Ni0.5_KE1.40.P3	2286775.8 ± 2.1	37.52 ± 0.3	0.5	0.03	1.1	14
ZTF18abcujukq	0.0638	EXP_Ni0.5_KE1.40.P9.7	2310786.4 ± 2.6	36.89 ± 0.3	0.5	−0.02	2.5	18
ZTF18abclalx	0.105	DPL_Ni0.4_KE2.18.P4.4	2224690.0 ± 2.0	38.03 ± 0.3	−1.4	−0.08	1.7	6
ZTF18abcxroj	0.0309	EXP_Ni0.5_KE2.18.P4.4	2384604.8 ± 1.2	35.69 ± 0.3	−1.3	0.04	1.7	3
ZTF18abcsgvj	0.060	DPL_Ni0.4_KE2.18.P3	2319138.9 ± 1.7	37.31 ± 0.2	−1.7	0.04	1.4	4
ZTF18abcysdx	0.066	EXP_Ni0.4_KE0.78.P21	2306083.8 ± 5.3	37.22 ± 0.5	0.6	0.16	3.7	29
ZTF18abdbuty	0.059	EXP_Ni0.5_KE2.18.P4.4	2321331.2 ± 2.3	37.18 ± 0.4	−0.8	0.01	1.5	16
ZTF18abdmdrf	0.163	EXP_Ni0.6_KE2.18.P4.4	2113743.2 ± 5.0	39.07 ± 0.3	2.5	0.08	2.7	50

^a The names of SNe Ia with a detected flux excess are highlighted in bold.^c File name of model, encapsulating all the model parameters (⁵⁶Ni mass, P value, kinetic energy, and density profile).^d Estimated explosion date from fit in the rest frame of the SN.^b x_1 values, c colours, and redshifts (z) are taken from Yao et al. (2019). z is noted to four decimal places if the value is taken from NED, or measured from the galaxy spectrum obtained by Yao et al. (2019), or measured from an SN spectrum where the host H α line can be identified. If z is taken from an SNID fit, it is noted to three decimal places.

Table A4. Continuation of Table A3.

Name	z	Model	Exp. date	μ	x_1	c	χ^2	# Unique fits
ZTF18abdefet	0.074	EXP_Ni0.4_KE2.18_P3	2288909.8 ± 2.3	37.7 ± 0.1	−0.1	0.14	3.5	6
ZTF18abdfazk	0.084	DPL_Ni0.4_KE2.18_P4.4	2267794.3 ± 1.8	37.56 ± 0.4	−0.3	−0.06	1.1	10
ZTF18abdfwur	0.070	DPL_Ni0.4_KE1.81_P3	2297469.8 ± 3.4	37.36 ± 0.4	−0.5	0.04	1.3	24
ZTF18abdfydj	0.076	DPL_Ni0.4_KE2.18_P4.4	2284658.7 ± 0.8	37.29 ± 0.0	0.2	−0.04	1.8	1
ZTF18abdкимx	0.077	DPL_Ni0.4_KE2.18_P4.4	2282540.1 ± 2.1	37.49 ± 0.4	0.0	−0.04	1.6	10
ZTF18abdmgab	0.0803	EXP_Ni0.4_KE1.68_P100	2276198.2 ± 5.0	37.97 ± 0.4	−2.3	0.12	4.6	28
ZTF18abeecwe	0.0393	EXP_Ni0.4_KE1.81_P3	2363751.7 ± 1.6	35.84 ± 0.4	−0.6	0.02	1.8	11
ZTF18abeegsl	0.072	EXP_Ni0.4_KE1.68_P100	2293187.5 ± 4.0	37.64 ± 0.4	−2.2	0.16	9.9	19
ZTF18abetehf	0.0649	DPL_Ni0.4_KE2.18_P3	2308423.1 ± 3.0	36.97 ± 0.3	−1.4	−0.14	2.2	15
ZTF18abetewu	0.077	EXP_Ni0.5_KE2.18_P3	2282544.3 ± 2.0	37.7 ± 0.5	0.3	0.11	1.0	10
ZTF18abfgygp	0.064	DPL_Ni0.4_KE1.81_P4.4	2310434.7 ± 1.9	36.94 ± 0.4	0.1	−0.07	1.0	6
ZTF18abfhaji	0.084	EXP_Ni0.5_KE1.68_P4.4	2267805.7 ± 2.0	37.82 ± 0.3	−0.2	−0.04	0.8	14
ZTF18abfhryc	0.0323	EXP_Ni0.4_KE1.81_P4.4	2386703.0 ± 0.6	35.36 ± 0.3	0.5	0.01	4.7	6
ZTF18abfwuwn	0.109	EXP_Ni0.6_KE2.18_P3	2216686.0 ± 1.6	38.12 ± 0.0	0.6	−0.03	1.9	6
ZTF18abfzkno	0.100	EXP_Ni0.5_KE2.18_P4.4	2234821.0 ± 2.9	38.09 ± 0.3	−0.6	−0.07	0.8	16
ZTF18abgxvra	0.104	EXP_Ni0.6_KE2.18_P4.4	2226728.3 ± 2.4	38.01 ± 0.0	0.8	−0.04	4.1	8
ZTF18abimsyv	0.088	EXP_Ni0.5_KE1.40_P4.4	2259479.0 ± 1.6	37.63 ± 0.1	1.0	−0.03	0.9	6
ZTF18abixjey	0.1218	EXP_Ni0.6_KE2.18_P3	2194920.6 ± 2.8	38.47 ± 0.1	0.8	−0.07	1.3	12
ZTF18abjgdo	0.0741	EXP_Ni0.5_KE1.40_P3	2289086.7 ± 3.6	37.44 ± 0.4	−0.8	0.0	0.6	27
ZTF18abjtger	0.107	EXP_Ni0.5_KE2.18_P4.4	2220701.0 ± 3.4	38.08 ± 0.2	0.8	−0.01	2.6	17
ZTF18abjvhec	0.0570	EXP_Ni0.5_KE1.40_P4.4	2325972.6 ± 2.0	36.63 ± 0.2	0.4	−0.04	1.3	5
ZTF18abkhcrj	0.0383	EXP_Ni0.4_KE2.18_P3	2367803.3 ± 1.5	36.34 ± 0.1	0.9	0.24	6.7	4
ZTF18abkhwcl	0.0317	EXP_Ni0.5_KE1.10_P4.4	2382929.2 ± 2.2	35.49 ± 0.4	0.1	−0.08	5.7	12
ZTF18abkhdxe	0.104	EXP_Ni0.4_KE2.18_P3	2226737.4 ± 2.6	38.23 ± 0.5	0.7	0.12	1.0	10
ZTF18abkifng	0.0880	EXP_Ni0.5_KE2.18_P4.4	2259489.0 ± 3.3	37.8 ± 0.4	0.3	−0.04	0.9	23
ZTF18abkigee	0.0936	EXP_Ni0.5_KE1.68_P3	2234838.4 ± 2.9	38.15 ± 0.5	1.0	−0.03	1.1	20
ZTF18abkljvo	0.141	EXP_Ni0.6_KE2.18_P3	2154530.4 ± 1.9	38.73 ± 0.0	1.2	−0.1	0.5	3
ZTF18abkudjo	0.0921	EXP_Ni0.4_KE1.81_P4.4	2250882.8 ± 3.3	37.73 ± 0.4	0.9	0.02	0.9	24
ZTF18abmmkaz	0.063	EXP_Ni0.5_KE1.40_P4.4	2312638.7 ± 1.9	36.95 ± 0.3	0.7	−0.07	1.8	7
ZTF18abmxdbh	0.070	DPL_Ni0.4_KE1.81_P4.4	2297515.6 ± 1.9	37.31 ± 0.5	1.3	0.01	1.0	16
ZTF18abnvoel	0.083	EXP_Ni0.6_KE2.18_P3	2269937.5 ± 1.7	37.49 ± 0.0	0.6	−0.01	4.6	4
ZTF18aboaeqy	0.129	EXP_Ni0.5_KE2.18_P4.4	2177447.1 ± 3.4	38.52 ± 0.1	0.3	−0.03	1.4	10
ZTF18abokpvh	0.081	EXP_Ni0.4_KE2.18_P4.4	2274140.1 ± 1.8	37.43 ± 0.4	0.8	−0.05	1.1	9
ZTF18abpamut	0.064	EXP_Ni0.4_KE1.10_P3	2310477.5 ± 3.3	37.28 ± 0.5	0.8	0.1	1.8	7
ZTF18abpaywm	0.040	EXP_Ni0.4_KE1.81_P3	2363797.4 ± 4.3	36.44 ± 0.5	0.6	0.23	6.9	9
ZTF18abpmpo	0.076	EXP_Ni0.5_KE1.40_P4.4	2284709.7 ± 1.2	37.29 ± 0.0	1.5	−0.01	2.1	3
ZTF18abptsco ^a	0.12	DPL_Ni0.4_KE2.18_P3	2194951.2 ± 4.2	38.6 ± 0.4	−0.4	0.07	1.4	15
ZTF18abpttky	0.084	DPL_Ni0.4_KE2.18_P3	2267849.2 ± 2.8	37.82 ± 0.3	−1.3	0.02	1.7	15
ZTF18abqbavl	0.1392	EXP_Ni0.5_KE1.40_P3	2156439.7 ± 3.6	38.7 ± 0.2	0.7	−0.03	0.9	18
ZTF18abqjvyl	0.0835	DPL_Ni0.4_KE1.68_P4.4	2260133.4 ± 3.1	37.59 ± 0.5	0.3	−0.06	0.6	17
ZTF18abrzeym	0.055	EXP_Ni0.4_KE0.78_P21	2330192.7 ± 4.9	36.81 ± 0.5	0.4	0.23	9.9	37
ZTF18abrznb	0.120	EXP_Ni0.5_KE2.18_P3	2194959.4 ± 3.3	38.35 ± 0.1	0.4	−0.02	1.4	9
ZTF18absdgon	0.0620	EXP_Ni0.4_KE2.18_P3	2255374.0 ± 1.4	37.17 ± 0.3	−0.3	0.17	5.0	3
ZTF18abslxhz	0.134	EXP_Ni0.6_KE2.18_P3	2167861.9 ± 1.2	38.61 ± 0.0	0.7	−0.08	1.3	2
ZTF18abssdpi	0.103	EXP_Ni0.4_KE2.18_P3	2228791.2 ± 2.9	38.25 ± 0.4	−0.3	0.2	2.0	9
ZTF18absxuz	0.0649	EXP_Ni0.5_KE2.18_P4.4	2308536.0 ± 2.0	37.23 ± 0.3	−1.1	0.02	0.7	9
ZTF18abtcdfv	0.140	EXP_Ni0.6_KE2.18_P3	2156453.8 ± 1.1	38.71 ± 0.0	1.1	−0.05	1.2	1
ZTF18abtnlik	0.084	EXP_Ni0.4_KE2.18_P3	2267861.3 ± 0.9	37.6 ± 0.3	−1.2	0.06	0.9	3
ZTF18abtogdl	0.100	EXP_Ni0.5_KE1.40_P4.4	2234873.4 ± 2.5	37.92 ± 0.2	0.3	−0.06	1.1	10
ZTF18abucvbf	0.0549	EXP_Ni0.6_KE2.18_P4.4	2341301.2 ± 1.5	36.85 ± 0.3	0.0	−0.07	4.3	9
ZTF18abukmty	0.104	EXP_Ni0.5_KE1.81_P3	2226778.4 ± 2.0	38.14 ± 0.2	0.5	0.0	1.3	8
ZTF18abuqugw	0.0313	DPL_Ni0.4_KE2.18_P4.4	2384357.8 ± 2.4	35.38 ± 0.3	−1.3	−0.12	4.7	11
ZTF18abvbyab	0.132	EXP_Ni0.5_KE2.18_P4.4	2171699.8 ± 1.8	38.57 ± 0.0	−0.2	−0.05	0.5	5
ZTF18abwmuua	0.0643	EXP_Ni0.4_KE2.18_P4.4	2309976.4 ± 4.8	37.08 ± 0.6	−1.2	−0.01	0.5	81
ZTF18abwnsoc	0.099	EXP_Ni0.5_KE1.68_P3	2236916.6 ± 1.8	38.07 ± 0.3	0.3	0.01	0.4	10
ZTF18abwtops	0.101	EXP_Ni0.5_KE2.18_P4.4	2232851.4 ± 2.5	38.37 ± 0.5	−1.4	−0.05	0.8	19
ZTF18abxssh	0.064	EXP_Ni0.6_KE1.81_P4.4	2310503.7 ± 3.7	37.57 ± 0.4	1.5	−0.02	3.2	44
ZTF18abxygvv	0.079	EXP_Ni0.6_KE2.18_P3	2278384.0 ± 1.6	37.55 ± 0.3	−0.1	−0.04	0.6	7

^a The redshift of ZTF18abptsco was reported by ATel 12052 (Gomez et al. 2018) and was shown with two decimal places.This paper has been typeset from a \LaTeX file prepared by the author.