

# Supplemental Material

## High pressure control of optical nonlinearity in the polar Weyl semimetal TaAs

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### I. SHG-RA DATA IN PARALLEL POLARIZATION GEOMETRY

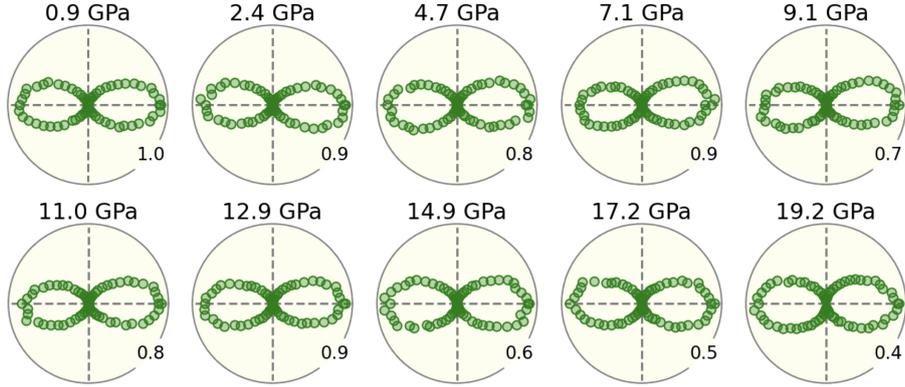


FIG. S1: SHG-RA patterns at select pressure in parallel polarization geometry measured upon compression.

The most general form of the electric-dipole SHG-RA intensity in parallel polarized geometry under normal incidence is given by  $I_{\parallel}(\varphi) \propto |B_1 \cos^3(\varphi) + B_2 \cos^2(\varphi) \sin(\varphi) + B_3 \cos(\varphi) \sin^2(\varphi) + B_4 \sin^3(\varphi)|^2$ , similar to the case for crossed polarized geometry described in the main text. From the (112) face of TaAs in the tetragonal  $4mm$  phase,  $B_2 = B_4 = 0$  by symmetry,  $B_1 = \gamma(2a^2\chi_{zxx} + 4a^2\chi_{xxz} + c^2\chi_{zzz})$  and  $B_3 = \gamma(4a^2 + 2c^2)\chi_{xxz} + (2a^2 + c^2)\chi_{zxx}$ , where  $\gamma = 1/\left(\frac{2a^3}{c} + ac\sqrt{2 + \frac{c^2}{a^2}}\right)$  and  $a$  and  $c$  are the TaAs lattice constants.

Figure S1 shows SHG-RA patterns at various pressures measured in parallel polarization geometry. For all pressures, the patterns are dominated by the  $\chi_{zzz}$  containing  $B_1$  term, which was previously shown to be an order of magnitude larger than both  $\chi_{zxx}$  and  $\chi_{xxz}$  [1]. Nonetheless, careful examination of the patterns shows that the nodes at  $90^\circ$  and  $270^\circ$  are lifted at high pressure, evidencing appearance of a  $B_4$  contribution that is only allowed in the hexagonal phase.

### II: DEPARTURE FROM $4mm$ POINT GROUP AT AMBIENT PRESSURE

For an ideal tetragonal crystal with point group  $4mm$ , the  $A_1$  term in  $I_{\perp}(\varphi)$  is forbidden by symmetry (see main text) and the pattern should exhibit nodes at  $0^\circ$ ,  $90^\circ$ ,  $180^\circ$  and  $270^\circ$ . Yet we observe a small but finite  $A_1$  component in our ambient pressure SHG-RA patterns taken outside the DAC, manifested as a lifting of the nodes at  $0^\circ$  and  $180^\circ$ . Our simulation results show that this is not due to misalignment, implying that the crystals may exhibit a slight departure from  $4mm$  symmetry.

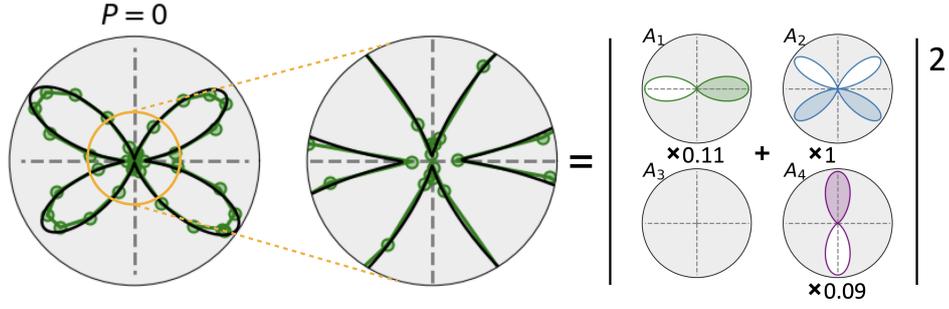


FIG. S2: Typical cross-polarized SHG-RA pattern from TaAs (112) acquired at ambient pressure outside the DAC (left). A zoom-in (middle) shows a lifting of the nodes at  $0^\circ$  and  $180^\circ$ , indicating a finite  $A_1$  term. Fits to the general form for  $I_\perp(\varphi)$  are overlaid as black lines. A decomposition of the fits into its  $A_1 \rightarrow A_4$  components is shown to the right. Filled versus empty lobes represent opposite signs of the associated trigonometric function. The fitted amplitude of the various terms is shown below, which are normalized to  $A_2$  for  $P < P_c$ .

### III: PRESSURE DEPENDENCE OF $A_1$ AND $A_4$

The SHG-RA patterns measured in crossed polarized geometry were fit to  $I_\perp(\varphi) \propto |A_1 \cos^3(\varphi) + A_2 \cos^2(\varphi) \sin(\varphi) + A_3 \cos(\varphi) \sin^2(\varphi) + A_4 \sin^3(\varphi)|^2$ . The fitted amplitudes of  $A_2$  and  $A_3$  are shown in the main text. Figure S3 shows the pressure dependence of the fitted amplitudes of  $A_1$  and  $A_4$ . Subtle kinks appear in  $|A_1|$  and  $|A_4|$  around the critical pressure  $P_c$ . However, as expected,  $|A_4|$  does not change as drastically as  $|A_2|$  because  $A_4$  does not contain  $\chi_{zzz}$ . Also, as expected,  $|A_1|$  does not exhibit a clear a kink at  $P_c$  as  $|A_3|$  because  $A_1$  is present in both the low and high pressure phases.

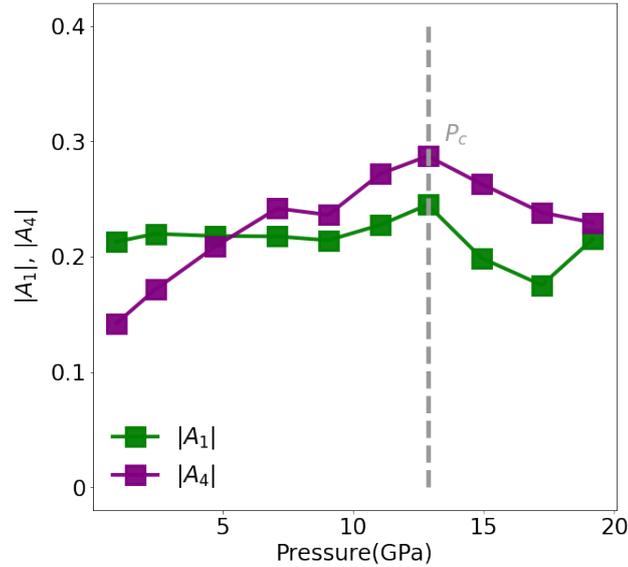


FIG. S3: Pressure dependence of the fitted amplitudes of  $A_1$  (green) and  $A_4$  (purple) normalized to the maximum value of  $A_2$  for  $P < P_c$

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[1] Wu, L. *et al.* Giant anisotropic nonlinear optical response in transition metal monpnictide Weyl semimetals. *Nature Physics* **13**, 350–355 (2017). URL <https://www.nature.com/articles/nphys3969>. Number: 4 Publisher: Nature Publishing Group.