

Folding energy landscape of cytochrome *cb*₅₆₂

Tetsunari Kimura, Jennifer C. Lee¹, Harry B. Gray², and Jay R. Winkler²

Beckman Institute, California Institute of Technology, Pasadena, CA 91125-7400

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Cytochrome *cb*₅₆₂ is a variant of an *Escherichia coli* four-helix bundle *b*-type heme protein in which the porphyrin prosthetic group is covalently ligated to the polypeptide near the terminus of helix 4. Studies from other laboratories have shown that the apoprotein folds rapidly without the formation of intermediates, whereas the holoprotein loses heme before native structure can be attained. Time-resolved fluorescence energy transfer (TRFET) measurements of cytochrome *cb*₅₆₂ refolding triggered using an ultrafast continuous-flow mixer (150 μ s dead time) reveal that heme attachment to the polypeptide does not interfere with rapid formation of the native structure. Analyses of the TRFET data produce distributions of Trp-59–heme distances in the protein before, during, and after refolding. Characterization of the moments and time evolution of these distributions provides compelling evidence for a refolding mechanism that does not involve significant populations of intermediates. These observations suggest that the cytochrome *b*₅₆₂ folding energy landscape is minimally frustrated and able to tolerate the introduction of substantial perturbations (i.e., the heme prosthetic group) without the formation of deep misfolded traps.

four-helix bundle | minimal frustration | protein folding | time-resolved fluorescence energy transfer | tryptophan

Energy landscape theory has delineated principles that underlie the conversion of disordered polypeptides into correctly folded functional proteins (1–8). A key element of this theory is the concept of minimal frustration that, in its qualitative formulation, predicts that the folding energy landscape is funneled toward the native structure and does not contain a large number of deep misfolded traps. This notion derives in part from the many experimental observations of proteins that rapidly fold to native structures without the apparent population of intermediates. Additional features of minimal frustration are the robustness of protein structures to mutation and the malleability of folding pathways (1, 3–5, 7).

The incorporation of prosthetic groups into protein structures introduces additional challenges for understanding folding and the maintenance of minimally frustrated folding pathways. The heme cofactors in *b*- and *c*-type cytochromes are a case in point. The apoprotein of mitochondrial cytochrome *c* (cyt *c*) is unstructured, demonstrating that heme is required to stabilize the native fold (9). The covalently bound porphyrin plays an integral role in cyt *c* folding, possibly as a hydrophobic nucleation site, but does not lead to nonnative clusters and misfolded traps (10). Many *b*-type cytochromes (e.g., cyt *b*₅ and cyt *b*₅₆₂), however, adopt native or near-native structures in the absence of their noncovalently bound porphyrins (11–23). On the basis of these observations, it has been suggested that folding precedes heme incorporation in the *b*-type cytochromes, whereas heme attachment is a prerequisite for folding in the *c*-type proteins. Because *b*-type cytochromes likely evolved to fold in the absence of heme, the conversion of a *b*-type cytochrome into a *c*-type protein might be expected to disrupt the folding landscape of the native protein.

We have examined a family of four-helix bundle cytochromes [cyt *b*₅₆₂ (24), cyt *cb*₅₆₂ (25, 26), cyt *c*₅₅₆ (25), and cyt *c*' (27–30)] in which refolding times differ by many orders of magnitude despite their strongly conserved structural topology (≈ 3 Å rmsd)

(25, 28, 31–33). In the *b*-type cytochromes (e.g., cyt *b*₅₆₂), the heme is attached to the polypeptide only through axial Fe ligation. The kinetics of heme dissociation from the cyt *b*₅₆₂ polypeptide ($k_{\text{diss}} \approx 2\text{--}7 \times 10^3 \text{ s}^{-1}$) compete with refolding dynamics, limiting the yield of the folding reaction in the reduced state and causing irreversible unfolding in the oxidized state (14, 24). To eliminate the complication of heme dissociation, we developed a protocol to overexpress a variant of cyt *b*₅₆₂ (R98C/Y101C) in which the α -carbons of the protoporphyrin vinyl groups form thioether linkages to two cysteines in a typical CXXCH *c*-type cytochrome-binding motif (e.g., cyt *c*, cyt *c*₅₅₆) (26). A K59W mutation, originally introduced to provide a fluorescent probe of folding, proved to increase yields of protein with properly attached hemes. We call this (K59W/R98C/Y101C) cyt *b*₅₆₂ variant cyt *cb*₅₆₂. We found that the presence of two *c*-type thioether linkages does not perturb the wild-type (cyt *b*₅₆₂) structure but does have a substantial impact on folding stability. The folding free-energy change for cyt *cb*₅₆₂, extrapolated to guanidine hydrochloride [GuHCl] = 0, is $-42 \pm 4 \text{ kJ mol}^{-1}$ (26); this value is substantially greater than that estimated for the wild-type (*b*₅₆₂) protein (-30 kJ mol^{-1}) (34) and three times greater than that of the apoprotein (-13 kJ mol^{-1}) (14).

The experimental characterization of conformational heterogeneity and structural changes during protein folding is particularly important to understand the molecular basis of the energy landscape and the detailed sequence of events that accompanies the transformation of an ensemble of denatured proteins into the native state. We have used time-resolved fluorescence energy transfer (TRFET) to estimate the distributions of distances between a fluorescent donor (D) and an energy acceptor (A) in an ensemble of protein molecules. By performing these measurements during protein folding, the time evolution of the polypeptide ensemble from heterogeneous unfolded to highly homogeneous native states can be determined. Our previous investigations of yeast cyt *c* (35–39) folding revealed the presence of both extended and collapsed conformations in the intermediate states formed within the mixing dead time, providing clear evidence that this protein does not fold by a simple two-state mechanism (U \leftrightarrow N).

Owing to the reversibility of cyt *cb*₅₆₂ denaturation, we were able to probe the kinetics of stopped-flow triggered cyt *cb*₅₆₂ refolding (26). The heme absorption spectrum reveals that formation of native cyt *cb*₅₆₂ is biphasic. The rate constant for the faster phase varies with denaturant concentration and has an extrapolated value of $4.2 \times 10^2 \text{ s}^{-1}$ at 0 M GuHCl. The rate of the slower phase is independent of denaturant concentration ($k_{\text{obs}} \approx 5 \text{ s}^{-1}$). When refolding is monitored by Trp-59 fluorescence, more than half of the signal amplitude is quenched during

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¹Present address: Laboratory of Molecular Biophysics, National Heart, Lung, and Blood Institute, National Institutes of Health, Bethesda, MD 20892.

²To whom correspondence may be addressed. E-mail: hbgray@caltech.edu or winklerj@caltech.edu.

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Data Fitting and Analysis. The resulting short- and long-time scale data were spliced together, and the combined traces were compressed logarithmically before fitting (70 points per decade). We confirmed that the compression does not alter the interpretation of data.

TRFET analysis involves the numerical inversion of a Laplace transform [$I(t) = \sum_k P(k) \exp(-kt)$] (50, 51). We have used two algorithms to invert our kinetics data with regularization methods that impose additional constraints on the properties of $P(k)$. The simplest constraint that applies to the TRFET is that $P(k) \geq 0$ ($\forall k$). We have fit the kinetics data by using a MATLAB (Mathworks) algorithm (LSQNONNEG) that minimizes the sum of the squared deviations (χ^2) between observed and calculated values of $I(t)$, subject to a nonnegativity constraint. It is our experience that LSQNONNEG produces the narrowest $P(k)$ distributions and smallest values of χ^2 with relatively few nonzero components. Information theory suggests that the least biased solution to this inversion problem minimizes χ^2 and maximizes the breadth of $P(k)$ (52). This regularization condition can be met by maximizing the Shannon–Jaynes entropy of the rate-constant distribution [$S = -\sum_k P(k) \ln[P(k)]$], implicitly requiring that $P(k) \geq 0$ ($\forall k$) (53). Maximum-entropy (ME) fitting produces stable and reproducible numerical inversions of the kinetics data. The balance between χ^2 minimization and entropy maximization is determined by graphical L curve analysis (54). This approach yields upper limits for the widths of $P(k)$ consistent with our experimental data. The $P(k)$ distributions from ME fitting are broader than those obtained with LSQNONNEG fitting but exhibit maxima in similar locations. A simple coordinate transformation using the Förster equation (Eq. 1)

$$k = k_0 \left(1 + \left(\frac{r_0}{r_{\text{DA}}} \right)^6 \right) \quad [1]$$

recasts the probability distribution of the decay rates, $P(k)$, obtained by LSQNONNEG or ME fitting as probability distributions over r_{DA} (55, 56). The

Förster critical length, r_0 , for the Trp-59–heme pair in cyt cb_{562} , is 34 Å under both native and unfolded conditions. The value of k_0 ($3.2 \times 10^8 \text{ s}^{-1}$) was obtained from luminescence decay measurements with NATA (10 μM) in the CTF mixer under various solvent conditions [50 mM NaOAc with and without GuHCl (0, 1.0, and 6.0 M) at pH 5.0]. At distances longer than $1.5r_0$, energy transfer quenching of **D** is not competitive with excited-state decay, and, at distances less than ≈ 10 Å, the Förster model does not reliably describe FET kinetics (57, 58); accordingly, our cyt cb_{562} FET kinetics measurements can provide information about **D–A** center-to-center distances only in the range $10 \leq r_{\text{DA}} \leq 44$ Å.

The $P(r_{\text{DA}})$ distributions are characterized by their moments: the first and second moments (M_n , $n = 1, 2$) (Eq. 2) correspond to the mean and mean-squared **D–A** distances in the ensemble; the second central moment or variance (V , Eq. 3) reflects the breadth of the distribution.

$$M_n = \langle r_{\text{DA}}^n \rangle = \frac{\sum P(r_{\text{DA}}) \cdot r_{\text{DA}}^n}{\sum P(r_{\text{DA}})} \quad [2]$$

$$V = M_2 - (M_1)^2. \quad [3]$$

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