

# Folding energy landscape of cytochrome *cb*<sub>562</sub>

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Cytochrome *cb*<sub>562</sub> is a variant of an *Escherichia coli* four-helix bundle *b*-type heme protein in which the porphyrin prosthetic group is covalently ligated to the polypeptide near the terminus of helix 4. Studies from other laboratories have shown that the apoprotein folds rapidly without the formation of intermediates, whereas the holoprotein loses heme before native structure can be attained. Time-resolved fluorescence energy transfer (TRFET) measurements of cytochrome *cb*<sub>562</sub> refolding triggered using an ultrafast continuous-flow mixer (150  $\mu$ s dead time) reveal that heme attachment to the polypeptide does not interfere with rapid formation of the native structure. Analyses of the TRFET data produce distributions of Trp-59–heme distances in the protein before, during, and after refolding. Characterization of the moments and time evolution of these distributions provides compelling evidence for a refolding mechanism that does not involve significant populations of intermediates. These observations suggest that the cytochrome *b*<sub>562</sub> folding energy landscape is minimally frustrated and able to tolerate the introduction of substantial perturbations (i.e., the heme prosthetic group) without the formation of deep misfolded traps.

four-helix bundle | minimal frustration | protein folding | time-resolved fluorescence energy transfer | tryptophan

Energy landscape theory has delineated principles that underlie the conversion of disordered polypeptides into correctly folded functional proteins (1–8). A key element of this theory is the concept of minimal frustration that, in its qualitative formulation, predicts that the folding energy landscape is funneled toward the native structure and does not contain a large number of deep misfolded traps. This notion derives in part from the many experimental observations of proteins that rapidly fold to native structures without the apparent population of intermediates. Additional features of minimal frustration are the robustness of protein structures to mutation and the malleability of folding pathways (1, 3–5, 7).

The incorporation of prosthetic groups into protein structures introduces additional challenges for understanding folding and the maintenance of minimally frustrated folding pathways. The heme cofactors in *b*- and *c*-type cytochromes are a case in point. The apoprotein of mitochondrial cytochrome *c* (cyt *c*) is unstructured, demonstrating that heme is required to stabilize the native fold (9). The covalently bound porphyrin plays an integral role in cyt *c* folding, possibly as a hydrophobic nucleation site, but does not lead to nonnative clusters and misfolded traps (10). Many *b*-type cytochromes (e.g., cyt *b*<sub>5</sub> and cyt *b*<sub>562</sub>), however, adopt native or near-native structures in the absence of their noncovalently bound porphyrins (11–23). On the basis of these observations, it has been suggested that folding precedes heme incorporation in the *b*-type cytochromes, whereas heme attachment is a prerequisite for folding in the *c*-type proteins. Because *b*-type cytochromes likely evolved to fold in the absence of heme, the conversion of a *b*-type cytochrome into a *c*-type protein might be expected to disrupt the folding landscape of the native protein.

We have examined a family of four-helix bundle cytochromes [cyt *b*<sub>562</sub> (24), cyt *cb*<sub>562</sub> (25, 26), cyt *c*<sub>556</sub> (25), and cyt *c*' (27–30)] in which refolding times differ by many orders of magnitude despite their strongly conserved structural topology ( $\approx 3$  Å rmsd)

(25, 28, 31–33). In the *b*-type cytochromes (e.g., cyt *b*<sub>562</sub>), the heme is attached to the polypeptide only through axial Fe ligation. The kinetics of heme dissociation from the cyt *b*<sub>562</sub> polypeptide ( $k_{\text{diss}} \approx 2\text{--}7 \times 10^3 \text{ s}^{-1}$ ) compete with refolding dynamics, limiting the yield of the folding reaction in the reduced state and causing irreversible unfolding in the oxidized state (14, 24). To eliminate the complication of heme dissociation, we developed a protocol to overexpress a variant of cyt *b*<sub>562</sub> (R98C/Y101C) in which the  $\alpha$ -carbons of the protoporphyrin vinyl groups form thioether linkages to two cysteines in a typical CXXCH *c*-type cytochrome-binding motif (e.g., cyt *c*, cyt *c*<sub>556</sub>) (26). A K59W mutation, originally introduced to provide a fluorescent probe of folding, proved to increase yields of protein with properly attached hemes. We call this (K59W/R98C/Y101C) cyt *b*<sub>562</sub> variant cyt *cb*<sub>562</sub>. We found that the presence of two *c*-type thioether linkages does not perturb the wild-type (cyt *b*<sub>562</sub>) structure but does have a substantial impact on folding stability. The folding free-energy change for cyt *cb*<sub>562</sub>, extrapolated to guanidine hydrochloride [GuHCl] = 0, is  $-42 \pm 4 \text{ kJ mol}^{-1}$  (26); this value is substantially greater than that estimated for the wild-type (*b*<sub>562</sub>) protein ( $-30 \text{ kJ mol}^{-1}$ ) (34) and three times greater than that of the apoprotein ( $-13 \text{ kJ mol}^{-1}$ ) (14).

The experimental characterization of conformational heterogeneity and structural changes during protein folding is particularly important to understand the molecular basis of the energy landscape and the detailed sequence of events that accompanies the transformation of an ensemble of denatured proteins into the native state. We have used time-resolved fluorescence energy transfer (TRFET) to estimate the distributions of distances between a fluorescent donor (D) and an energy acceptor (A) in an ensemble of protein molecules. By performing these measurements during protein folding, the time evolution of the polypeptide ensemble from heterogeneous unfolded to highly homogeneous native states can be determined. Our previous investigations of yeast cyt *c* (35–39) folding revealed the presence of both extended and collapsed conformations in the intermediate states formed within the mixing dead time, providing clear evidence that this protein does not fold by a simple two-state mechanism (U  $\leftrightarrow$  N).

Owing to the reversibility of cyt *cb*<sub>562</sub> denaturation, we were able to probe the kinetics of stopped-flow triggered cyt *cb*<sub>562</sub> refolding (26). The heme absorption spectrum reveals that formation of native cyt *cb*<sub>562</sub> is biphasic. The rate constant for the faster phase varies with denaturant concentration and has an extrapolated value of  $4.2 \times 10^2 \text{ s}^{-1}$  at 0 M GuHCl. The rate of the slower phase is independent of denaturant concentration ( $k_{\text{obs}} \approx 5 \text{ s}^{-1}$ ). When refolding is monitored by Trp-59 fluorescence, more than half of the signal amplitude is quenched during

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the stopped-flow mixing dead time ( $\approx 5$  ms); the residual signal amplitude decays exponentially with a denaturant-dependent rate constant that corresponds to the faster phase observed using heme absorption spectroscopy.

To clarify the missing early stage of the folding reaction, we have used a continuous-flow (CTF) mixer ( $\tau_{\text{mixing}} < 150 \mu\text{s}$ ) (40, 41) to trigger the folding reaction. The mixer has been coupled to a picosecond streak camera with a fiberoptic bundle that allows luminescence decays for 25 different folding times to be acquired simultaneously (29). We measured TRFET between a Trp residue (Trp-59) and the heme to monitor refolding of cyt *cb*<sub>562</sub>. Analysis of the TRFET data produced distributions of Trp-59–heme center-to-center distances  $[P(r_{\text{DA}})]$  that provide detailed structural insights into the early events in cyt *cb*<sub>562</sub> folding.

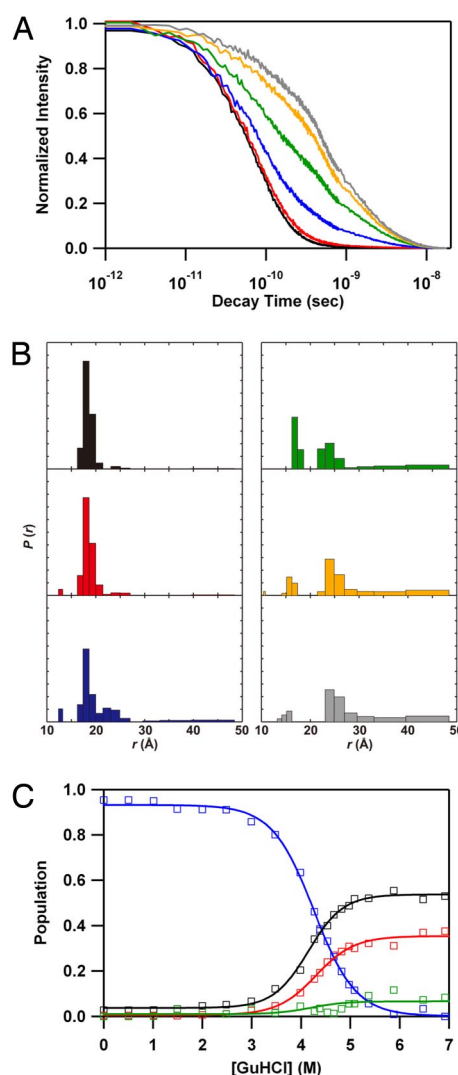
## Results

**Moment Analyses.** The properties of heterogeneous systems such as denatured proteins are most conveniently described in terms of distribution functions. TRFET measurements provide the data necessary to generate estimated distributions of distances  $[P(r_{\text{DA}})]$  between energy donors and acceptors in a polypeptide ensemble. We characterize these distributions by their moments (see *Methods*) to quantify the average properties and the heterogeneity of the protein under both equilibrium and transient conditions.

**GuHCl-Induced Unfolding.** We first characterized the equilibrium unfolding transition of cyt *cb*<sub>562</sub> by measuring fluorescence energy transfer kinetics between Trp-59 (**D**) and the heme group (**A**). In folded cyt *cb*<sub>562</sub> (Fig. 1; 0 M GuHCl), the Trp-59 fluorescence decay kinetics are consistent with a narrow distribution of Trp-59–heme distances centered at 19.4 Å (the center of the distribution is given by the first moment,  $M_1$ ; see *Methods*), a value that is in good agreement with that determined from the X-ray crystal structure (18.9 Å; Protein Data Bank ID code 2bc5).

More complex behavior is evident after the addition of GuHCl. The Trp-59 fluorescence decay kinetics become slower and nonexponential, and the mean **D–A** distance moves to larger values as [GuHCl] increases (Fig. 1 *A* and *B*). The TRFET kinetics of the denatured protein are indicative of a heterogeneous collection of compact and extended conformations. Near the midpoint of the GuHCl titration ( $[\text{GuHCl}]_{1/2} = 4.2 \text{ M}$ ), substantial populations of polypeptides in three different **D–A** distance ranges are apparent: shorter distances than folded state ( $r_{\text{DA}} \leq 16 \text{ \AA}$ ; *S* component), intermediate distances ( $22 \leq r_{\text{DA}} \leq 30 \text{ \AA}$ ; *I* component) and extended distances ( $r_{\text{DA}} \geq 30 \text{ \AA}$ ; *E* component). Finally, at GuHCl concentrations  $>5.5 \text{ M}$ , the population of proteins with native **D–A** distances is diminished in favor of *S* ( $\approx 10\%$ ), *I* ( $\approx 40\%$ ), and *E* components ( $\approx 50\%$ ) as shown in Fig. 1C. The changes in the populations of all components show cooperative unfolding of cyt *cb*<sub>562</sub> with the transition midpoint of 4.2 M and *m* value of  $7.5 \text{ kJ mol}^{-1} \text{ M}^{-1}$ .

The TRFET data do not allow us to say with confidence that the three components identified in the  $P(r_{\text{DA}})$  distributions represent distinct minima in the energy landscape of the denatured protein. It is possible that the true  $P(r_{\text{DA}})$  function is a broad, asymmetric single-mode distribution and that the trimodal functions extracted from the fluorescence data are simply a consequence of the uncertainties inherent in the data analysis. Although, the maximum-entropy method used to extract the  $P(r_{\text{DA}})$  distributions tends to give maximally broadened functions, there is no guarantee that a general single-mode function cannot provide adequate fits to the data. Neither random polymer distribution functions (e.g., freely jointed chain, worm-like chain) nor a simple Gaussian distribution, however, ade-

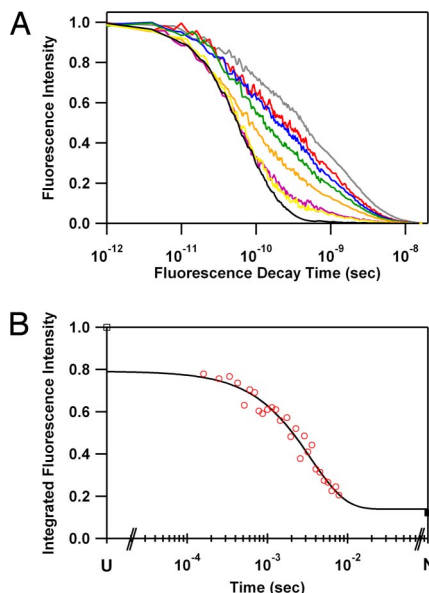


**Fig. 1.** GuHCl-induced unfolding of *cyt cb<sub>562</sub>* (pH 5.0, 50 mM NaOAc) probed by using Trp fluorescence decay. (A) Normalized Trp-59 fluorescence decay curves of *cyt cb<sub>562</sub>*. For clarity, only 6 of the 19 observed decays are displayed; 0 M (black), 2.0 M (red), 4.0 M (blue), 4.5 M (green), 5.4 M (yellow), and 6.9 M (gray). (B) Distributions of Trp-59–heme center-to-center distances [ $r_{DA}$ ] extracted from maximum entropy analyses of the TRFET data. The area of each bar reflects the probability amplitude over the corresponding distance range. (C) Population changes of each component as a function of the GuHCl concentration (native, blue; short, green; intermediate, red; and extended, black). Solid lines are two-state fitting curves with  $[GuHCl]_{1/2} = 4.2$  M and  $m = 7.5$  kJ mol<sup>-1</sup> M<sup>-1</sup>.

quately modeled the experimental data (Figs. S1 and S2 and *SI Appendix*).

**Refolding Kinetics.** In prior studies using a conventional stopped-flow mixer, we were able to capture only the final stages of *cyt cb<sub>562</sub>* refolding. The 150- $\mu$ s dead time of our CTF mixer provides insights into the early refolding events in this four-helix bundle. The evolution of the Trp-59 fluorescence decay kinetics as the protein refolds (Fig. 24) reveals a gradual transition from the denatured to the folded state. Integration of the individual Trp-59 fluorescence decay traces yields a quantity proportional to the fluorescence quantum yield ( $\Phi_{\text{fl}}$ ). Plots of integrated fluorescence intensity as a function of folding time reveal an exponential process with a time constant of  $\approx 4$  ms (Fig. 2B).

Immediately after the CTF mixing dead time, we observe a

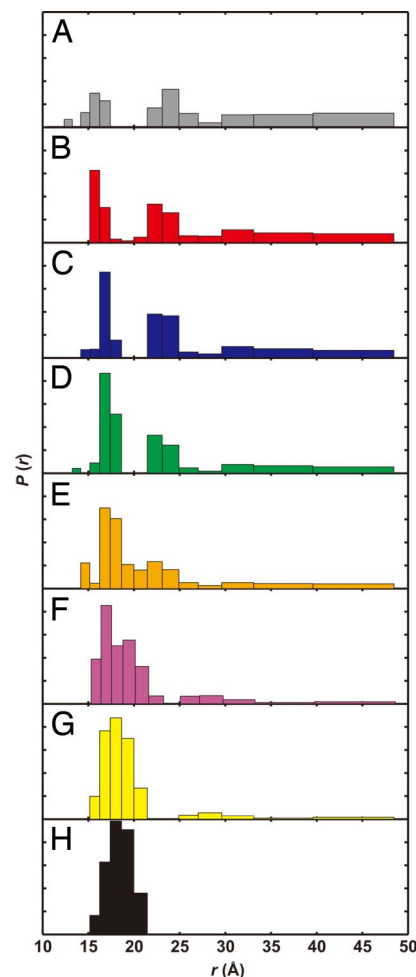


**Fig. 2.** Folding kinetics of *cyt cb*<sub>562</sub> triggered by GuHCl jump from 6.0 M to 1.0 M in a CTF mixer (pH 5.0, 18 °C) probed by Trp-59 fluorescence. (A) TRFET data (8 of the 112 measured decays are displayed): 0 (unfolded in 6.0 M GuHCl, gray), 150  $\mu$ s (red), 430  $\mu$ s (blue), 1.25 ms (green), 3.41 ms (orange), 6.50 ms (purple), 7.70 ms (yellow), and >30 min (native state in 1.0 M GuHCl, black) after the initiation of folding reaction. (B) Integrated Trp-59 fluorescence intensity ( $\Phi_{\text{H}}$ ) as a function of folding time. The 112 original data points are logarithmically smoothed for clarity. Solid line is a fit of the full dataset to a single exponential function with a rate constant of  $2.6 (\pm 0.2) \times 10^2 \text{ s}^{-1}$ .

21% decrease in Trp-59 integrated fluorescence intensity (Fig. 2B). This reduction in Trp-59 fluorescence is likely the result of a reduction in the average Trp-59–heme distance because the change is greater than can be accounted for on the basis of the change in solvent conditions. Maximum entropy analysis of the Trp-59 fluorescence decay kinetics measured 150  $\mu$ s after mixing reveals very modest changes in the Trp-59–heme distance distribution compared with that of the denatured protein (Fig. 3A and B). The presence of substantial populations of extended polypeptides indicates that the protein has not undergone a large-scale collapse upon dilution of the denaturant; the first moment of the 150- $\mu$ s distribution reveals a 3-Å reduction in the average Trp-59–heme distance (Fig. 4A and Fig. S3). These burst-phase changes are consistent with a response of the polypeptide to the change in solvent conditions upon denaturant dilution. It is clear that 150  $\mu$ s after dilution of denaturant, the *cyt cb*<sub>562</sub> ensemble is highly heterogeneous, with a considerable range of Trp-59–heme distances.

Trp-59–heme TRFET measurements at folding times between 150  $\mu$ s and 12 ms allowed us to follow the evolution of the *cyt cb*<sub>562</sub> ensemble (Fig. 3). After the modest collapse at 150  $\mu$ s is a major drop in Trp-59 fluorescence intensity (Fig. 2B) with a rate constant of  $240 \text{ s}^{-1}$ . This value is consistent (within experimental error) with the results from our earlier stopped-flow studies in which the folding was probed by using Trp-59 fluorescence intensity and heme UV/visible absorption measurements (26). The stopped-flow burst phase corresponded to a  $\approx 70\%$  decrease of the total fluorescence intensity (26), but with the CTF mixer we resolve the entire kinetics phase.

The time dependences of the moments of the  $P(r_{\text{DA}})$  distributions (see *Methods* and *SI Appendix*) enable us to characterize the time evolution of the protein ensemble. The time courses of the first and second moments of Trp-59–heme  $P(r_{\text{DA}})$  can be described by single exponential functions with nearly identical rate constants ( $260$  and  $240 \text{ s}^{-1}$ , respectively) (Fig. 4A and B and



**Fig. 3.** Distributions of  $P(r_{\text{DA}})$  for refolding of *cyt cb*<sub>562</sub>. The  $P(r_{\text{DA}})$  distributions are extracted from maximum entropy analyses of the FET kinetics data. Folding was triggered by GuHCl jump from 6.0 M to 1.0 M in a CTF mixer (pH 5.0, 18 °C). For clarity, only 8 of the 112 observed decays are displayed: (A) unfolded in 6.0 M GuHCl; (B) 150  $\mu$ s after the initiation of folding reaction; (C) 430  $\mu$ s; (D) 1.25 ms; (E) 3.41 ms; (F) 6.50 ms; (G) 7.70 ms; (H) >30 min (native state in 1.0 M GuHCl).

Fig. S3). The time dependence of the  $P(r_{\text{DA}})$  variance is biphasic (Fig. 4C and Fig. S3): a fast phase ( $k_{\text{fast}} = 510 \text{ s}^{-1}$ ) corresponding to a slight increase in  $V$  is followed by a slower large-amplitude reduction in  $V$  ( $k_{\text{slow}} = 260 \text{ s}^{-1}$ ). This behavior is precisely that expected for a folding process in which just the unfolded ( $U$ ) and native ( $N$ ) distributions are present during refolding: the moments evolve with a rate constant equal to the sum of the folding and unfolding rates ( $k_{\text{obs}} = k_u + k_f$ ), whereas the time dependence of the variance is biexponential with rate constants equal to  $k_{\text{obs}}$  and  $2 \times k_{\text{obs}}$  (*SI Appendix*).

## Discussion

Characterization of the conformational properties of the unfolded protein is a necessary precursor to any investigation of refolding kinetics. The *cyt cb*<sub>562</sub> Trp-59–heme distance distribution under strongly denaturing conditions (6.9 M GuHCl) is poorly described by random polymer models (freely jointed chain, wormlike chain) and simple Gaussian distribution functions (Figs. S1 and S2 and *SI Appendix*). Intrachain interactions in the unfolded state produce higher populations of proteins with Trp-59 closer to the heme than expected for a random polymer. We can exclude the possibility that the misligation of His-63 to





**Data Fitting and Analysis.** The resulting short- and long-time scale data were spliced together, and the combined traces were compressed logarithmically before fitting (70 points per decade). We confirmed that the compression does not alter the interpretation of data.

TRFET analysis involves the numerical inversion of a Laplace transform [ $I(t) = \sum_k P(k) \exp(-kt)$ ] (50, 51). We have used two algorithms to invert our kinetics data with regularization methods that impose additional constraints on the properties of  $P(k)$ . The simplest constraint that applies to the TRFET is that  $P(k) \geq 0$  ( $\forall k$ ). We have fit the kinetics data by using a MATLAB (Mathworks) algorithm (LSQNONNEG) that minimizes the sum of the squared deviations ( $\chi^2$ ) between observed and calculated values of  $I(t)$ , subject to a nonnegativity constraint. It is our experience that LSQNONNEG produces the narrowest  $P(k)$  distributions and smallest values of  $\chi^2$  with relatively few nonzero components. Information theory suggests that the least biased solution to this inversion problem minimizes  $\chi^2$  and maximizes the breadth of  $P(k)$  (52). This regularization condition can be met by maximizing the Shannon–Jaynes entropy of the rate-constant distribution [ $S = -\sum_k P(k) \ln[P(k)]$ ], implicitly requiring that  $P(k) \geq 0$  ( $\forall k$ ) (53). Maximum-entropy (ME) fitting produces stable and reproducible numerical inversions of the kinetics data. The balance between  $\chi^2$  minimization and entropy maximization is determined by graphical L curve analysis (54). This approach yields upper limits for the widths of  $P(k)$  consistent with our experimental data. The  $P(k)$  distributions from ME fitting are broader than those obtained with LSQNONNEG fitting but exhibit maxima in similar locations. A simple coordinate transformation using the Förster equation (Eq. 1)

$$k = k_0 \left( 1 + \left( \frac{r_0}{r_{DA}} \right)^6 \right) \quad [1]$$

recasts the probability distribution of the decay rates,  $P(k)$ , obtained by LSQNONNEG or ME fitting as probability distributions over  $r_{DA}$  (55, 56). The

Förster critical length,  $r_0$ , for the Trp-59–heme pair in cyt  $cb_{562}$ , is 34 Å under both native and unfolded conditions. The value of  $k_0$  ( $3.2 \times 10^8 \text{ s}^{-1}$ ) was obtained from luminescence decay measurements with NATA (10  $\mu\text{M}$ ) in the CTF mixer under various solvent conditions [50 mM NaOAc with and without GuHCl (0, 1.0, and 6.0 M) at pH 5.0]. At distances longer than  $1.5r_0$ , energy transfer quenching of **D** is not competitive with excited-state decay, and, at distances less than  $\approx 10$  Å, the Förster model does not reliably describe FET kinetics (57, 58); accordingly, our cyt  $cb_{562}$  FET kinetics measurements can provide information about **D**–**A** center-to-center distances only in the range  $10 \leq r_{DA} \leq 44$  Å.

The  $P(r_{DA})$  distributions are characterized by their moments: the first and second moments ( $M_n$ ,  $n = 1, 2$ ) (Eq. 2) correspond to the mean and mean-squared **D**–**A** distances in the ensemble; the second central moment or variance ( $V$ , Eq. 3) reflects the breadth of the distribution.

$$M_n = \langle r_{DA}^n \rangle = \frac{\sum P(r_{DA}) \cdot r_{DA}^n}{\sum P(r_{DA})} \quad [2]$$

$$V = M_2 - (M_1)^2 \quad [3]$$

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