

HYPERBOLIC CONE-MANIFOLDS, SHORT GEODESICS, AND SCHWARZIAN DERIVATIVES

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With his hyperbolic Dehn surgery theorem and later the orbifold theorem, Thurston demonstrated the power of using hyperbolic cone-manifolds to understand complete, non-singular hyperbolic 3-manifolds. Hodgson and Kerckhoff introduced analytic techniques to the study of cone-manifolds that they have used to prove deep results about finite volume hyperbolic 3-manifolds. In this paper we use Hodgson and Kerckhoff's techniques to study infinite volume hyperbolic 3-manifolds. The results we will develop have many applications: the Bers density conjecture, the density of cusps on the boundary of quasiconformal deformation spaces, and for constructing type preserving sequences of Kleinian groups.

The simplest example of the problem we will study is the following: Let M be a hyperbolic 3-manifold and c a simple closed geodesic in M . Then the topological manifold $M \setminus c$ also has a complete hyperbolic metric which we call \hat{M} . How does the geometry of M compare to that of \hat{M} ? Before attempting to answer such a question, we need to note that if M has infinite volume, the hyperbolic structure will not be unique. If we do not make further restrictions on the choice of \hat{M} , then there is no reason to expect that M and \hat{M} will be geometrically close. If M is convex co-compact, there is a natural choice to make for \hat{M} . Namely M is compactified by a conformal structure X . We then choose \hat{M} to be the unique geometrically finite hyperbolic structure on $M \setminus c$ with conformal boundary X .

We can now return to our question: How do the geometry of M and \hat{M} compare? We will quantify this question in two ways. We will measure the length of geodesics in M and \hat{M} and we will measure the geometry of the ends of M and \hat{M} by bounding the distance between the projective structures on their boundaries. What we will see is that the change in geometry is bounded by the length of the geodesic c in the original manifold M .

Results of this type were first obtained by McMullen [Mc], in the case of a quasifuchsian manifold, where the geodesic c is also short on a component of the conformal boundary. This work has been extended to arbitrary geometrically finite manifolds by Canary, Culler, Hersensky and Shalen [CCHS]. Their techniques are entirely different from ours and one goal of this paper is to give new proofs of their estimates. These estimates are a key step in proving the density of cusps in the boundary of quasiconformal deformation spaces of hyperbolic 3-manifolds.

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Our original motivation was another application, the density conjecture. In [Br2] we construct a family of quasifuchsian cone-manifolds, with cone-angle 4π , approximating any singly degenerate hyperbolic manifold with arbitrarily short geodesics. The estimates developed here imply that the quasifuchsian cone-manifolds are geometrically close to smooth quasifuchsian manifolds which will approximate the degenerate manifold. This is a special case of the Bers density conjecture.

In a joint work with Brock [BB] we extend the results of this paper. In particular we show that there is a diffeomorphism from $M \setminus c$ to \hat{M} that is bi-Lipschitz outside of a tubular neighborhood of c where the bi-Lipschitz constant is bounded by a constant depending only on the length of the c . Using this result we are able to prove the density conjecture for all freely indecomposable Kleinian groups without parabolics.

Another question is the following: Assume Γ is a Kleinian group and Γ_i is a sequence of geometrically finite Kleinian groups such that $\Gamma_i \rightarrow \Gamma$, algebraically. Does there exist a type preserving sequence Γ'_i of geometrically finite groups also converging to Γ ? Here type preserving means that if elements γ_i converge to γ , then γ is parabolic if and only if the γ_i are parabolic. In a joint work with Brock, Evans and Souto [BBES] we show that the answer to the question is yes and the type preserving sequence can be constructed by pinching the short geodesics in the Γ_i to cusps. The estimates developed here and extended in [BB] can then be used to show that the new sequence has the same limit. This question is important because in many cases the work of Anderson and Canary [AC1, AC2] implies that type preserving sequences are strong. Work of Evans [Ev], expanding on work of Canary and Minsky [CM], then implies that the limit is tame.

The starting point for this paper is the local parameterization of hyperbolic cone-manifolds developed by Hodgson and Kerckhoff for finite volume manifolds in [HK1] and [HK3] and extended to geometrically finite cone-manifolds in [Br1]. These local results tell us that we can make a small decrease in the cone-angle. To decrease the cone-angle to zero, we need to ensure that there are no degenerations. There are three possible types of degenerations that need to be avoided. First, we need to make sure that the cone-singularity does not develop a point of self intersection. Second we must show that there is a lower bound for the injectivity radius of the cone-manifold. Finally the geometry of the geometrically finite ends must be controlled. The first two problems are taken care of by work of Hodgson and Kerckhoff. It is the last problem that is the main work of this paper.

We now outline the remainder of the paper.

In §1 we make enough definitions to state our main theorems.

In §2 we review background material on projective structures, hyperbolic cone-manifolds and deformations thereof.

In §3 we describe some of Hodgson and Kerckhoff's results on tubes in cone-manifolds. In particular, they derive estimates on the radius of embedded tubes about the cone-singularity. We also use these ideas to show that embedded hyperbolic half spaces are disjoint from these tubular neighborhoods of the cone-singularity.

In §4 we use the analytic deformation theory of cone-manifolds to control the length of geodesics as the cone-angle decreases (Theorem 1.4). Again the key estimates are work of Hodgson and Kerckhoff.

§5 is the heart of the paper. In this section we show that the L^2 -norm of the cone-manifold deformation bounds the change in projective structure (Theorem 1.3). We first do this for a hyperbolic half space and then use the work from §3 to find a large embedded half space in each geometrically finite end which allows us to globally bound the deformation of the entire projective structure.

The final step is to understand the geometric limit of a sequence of cone-manifolds. Our approach is essentially the same as in [HK4] although a bit of extra care is needed to take care of the geometrically finite ends. Under certain conditions we then show that given a geometrically finite cone-manifold, there is a one-parameter family of cone-manifolds decreasing the cone-angle to zero (Theorem 1.2).

Although all the results of this paper hold for manifolds with rank two cusps, they are, as is often the case, an annoyance and distraction from the main line of argument. For this reason we defer the discussion of rank two cusps until §7 where we show how the L^2 -norm bounds control the shape of the cusp. We then outline how this result can be used to finish the proof of the main theorems for manifolds with rank two cusps.

In §8 we derive our versions of the estimates of McMullen and Canary et al. mentioned above.

In the appendix we recount mean value theorems for harmonic vector and strain fields that Hodgson and Kerckhoff proved in an early version of [HK4].

Acknowledgments. As should be obvious to the reader of the above outline, none of the results in this paper would be possible without the important work of Hodgson and Kerckhoff on hyperbolic cone-manifolds. Their analysis of harmonic forms in a neighborhood of the cone-singularity underlies all the estimates derived in this paper.

I would also like to thank Jeff Brock who originally suggested to me that cone-manifolds could be used to measure the effect of drilling short geodesics in hyperbolic 3-manifolds and the referee whose detailed comments have vastly improved the exposition of the paper.

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1. GEOMETRICALLY FINITE HYPERBOLIC CONE-MANIFOLDS

Before we state the main theorems, we need to define the main object of study, geometrically finite hyperbolic cone-manifolds. Let N be a compact, differentiable, 3-manifold with boundary, let \mathcal{C} be a collection of simple closed curves in the interior of N and let M be the interior of $N \setminus \mathcal{C}$. A *hyperbolic cone-metric* is a complete metric g on the interior of N that restricts to a Riemannian metric with constant sectional curvature equal to -1 on M ; i.e., a hyperbolic metric. At all components c of \mathcal{C} the metric will be singular; in cylindrical coordinates (r, θ, z) the metric will locally have the form

$$dr^2 + \sinh^2 r d\theta^2 + \cosh^2 r dz^2$$

where θ is measured modulo the cone-angle α . In this coordinate system the singular locus will be identified with the z -axis. Note that the cone-angle will be constant along each component of the cone-singularity.

A *complex projective structure* on a surface is an atlas of charts to $\widehat{\mathbb{C}}$ with transition maps Möbius transformations. A hyperbolic cone-metric g is *geometrically*

finite without rank one cusps if it extends to a projective structure on the non-toroidal components $\partial_0 N$ of ∂N . To state this more precisely, we recall that hyperbolic 3-space \mathbb{H}^3 is naturally compactified by the Riemann sphere $\widehat{\mathbb{C}}$. Then g is geometrically finite if each $p \in \partial_0 N$ has a neighborhood V in N and a map $\phi : V \rightarrow \mathbb{H}^3 \cup \widehat{\mathbb{C}}$ such that ϕ restricted to $V \cap \text{int } N$ is an isometry into \mathbb{H}^3 and ϕ restricted to $V \cap \partial_0 N$ defines an atlas for a projective structure on $\partial_0 N$. As we will not discuss rank one cusps in this paper (except briefly in §8), we simply refer to such metrics as geometrically finite.

A projective structure on a surface S also determines a conformal structure on S . Moreover, for a fixed conformal structure there will be many projective structures. We will often need to distinguish between the *projective boundary* and *conformal boundary* of a geometrically finite hyperbolic manifold.

If ∂N contains a torus T , the behavior near infinity is different. A neighborhood of T in M will be foliated by Euclidean tori of a fixed conformal class with area decreasing exponentially as the tori exit the end. Such a neighborhood is a *rank two cusp*. More explicitly every rank two cusp is the quotient of a subspace of \mathbb{H}^3 (in the upper half space model) of the form $\{(z, t) | t \geq t_0\}$ by parabolic isometries $z \mapsto z + 1$ and $z \mapsto z + \tau$ where $\text{Im } \tau > 0$. Note that τ is the Teichmüller parameter of the tori that foliate the cusp.

As we mentioned in the introduction, we will postpone discussion of rank two cusps whenever possible. However they cannot be completely avoided because as the cone-angle limits to zero, a tubular neighborhood of the cone-singularity will limit geometrically to a cusp. For this reason it is natural to think of a rank two cusp as a cone-singularity with cone-angle zero. This is of particular importance for the local parameterization for cone-manifolds that we are about to state.

Let $\mathcal{GF}(N, \mathcal{C})$ be the space of geometrically finite hyperbolic cone-metrics on the interior of N with cone-singularity \mathcal{C} and assume $\mathcal{GF}(N, \mathcal{C})$ has the compact- C^∞ topology. Then $GF(N, \mathcal{C})$ will be $\mathcal{GF}(N, \mathcal{C})$ modulo diffeomorphisms which are isotopic to the identity and that fix each component of \mathcal{C} . An equivalence class of metrics $[g] \in GF(N, \mathcal{C})$ assigns to each component of \mathcal{C} a cone-angle and to $\partial_0 N$ a conformal structure in the Teichmüller space $T(\partial_0 N)$. If \mathcal{C} has n components, there is a map

$$\Psi : GF(N, \mathcal{C}) \rightarrow [0, \infty)^n \times T(\partial_0 N)$$

with $\Psi([g]) = (\alpha_1, \dots, \alpha_n, X)$ where α_i is the cone-angle of the i th component of \mathcal{C} and X is the conformal boundary. Note that as we discussed in the previous paragraph, we have allowed the possibility for the cone-angle to be zero in which case the cone-singularity becomes a rank two cusp. We then have the following local parameterization:

Theorem 1.1 ([HK1], [HK3], [Br1]). *Let $[g] \in GF(N, \mathcal{C})$ be a geometrically finite hyperbolic cone-metric. Suppose the tube radius of the cone-singularity is $\geq \sinh^{-1} \left(\frac{1}{\sqrt{2}} \right)$. Then Ψ is a local homeomorphism at $[g]$.*

Remark. Hodgson and Kerckhoff first proved Theorem 1.1 for finite volume cone-manifolds with all cone-angles $\leq 2\pi$ without the restriction on the tube radius ([HK1]). More recently they have announced that the parameterization holds with the tube radius condition we give here ([HK2, HK3]). (When the cone-angle is zero, the tube radius is infinite and the result holds.) The parameterization was extended to geometrically finite cone-manifolds in [Br1].

Although it is not necessary, we simplify our notation by assuming that all the cone-angles are equal to a single cone-angle α . By Theorem 1.1 there is a neighborhood V of $\Psi(g)$ where Ψ is invertible. Let $[g_t] = \Psi^{-1}(V \cap (t, \dots, t, X))$ with $[g] = [g_\alpha]$.

We set notation that we will use throughout the remainder of the paper. Let $M_t = (M, g_t)$ be the one-parameter family of hyperbolic cone-manifolds coming from the cone-metrics $[g_t]$. Although the conformal boundary is a fixed conformal structure X , the projective boundary will change. Let Σ_t denote the projective boundary of M_t . We label the connected components of the conformal boundary X^1, \dots, X^k and the corresponding components of the projective boundary $\Sigma_t^1, \dots, \Sigma_t^k$.

If γ is a simple closed curve in M , then $\mathcal{L}_\gamma(g) = L_\gamma(g) + i\Theta_\gamma(g)$ is the complex length of the geodesic representative (if it exists) of γ in (M, g) where $L_\gamma(g)$ is the length of γ and $\Theta_\gamma(g)$ is the twisting. Note that $\Theta_\gamma(g)$ is defined modulo the cone-angle if γ is a component of the cone-singularity and modulo 2π if γ is a smooth geodesic. For the one-parameter family of metrics g_t we simplify notation and write the complex length $\mathcal{L}_\gamma(t)$. We also simplify notation by setting $L_C(t) = \sum_{c \in \mathcal{C}} L_c(t)$.

We now state our first main result:

Theorem 1.2. *Let $M_\alpha \in GF(N, \mathcal{C})$ be a geometrically finite hyperbolic cone-metric with all cone-angles α . Suppose the tube radius of the cone-singularity is $\geq \sinh^{-1} \sqrt{2}$. Then there exists an ℓ_0 depending only on α such that if $L_c(\alpha) \leq \ell_0$ for all $c \in \mathcal{C}$, then the one-parameter family of cone-manifolds $M_t \in GF(N, \mathcal{C})$ is defined for all $t \leq \alpha$.*

Remark. The value $\sinh^{-1} \sqrt{2}$ is arbitrary. We have chosen it because it simplifies some of the other constants that appear in the paper (see Theorem 3.5).

In our next result we control the geometry of the geometrically finite ends as the cone-angle decreases. In particular we will measure the distance between the projective boundaries of M_α and M_t . This distance is defined in the next section. We also note that $\|\Sigma_\alpha^i\|_F$ is the distance between Σ_α^i and the unique fuchsian projective structure with conformal structure X^i . This is also defined in the next section.

Theorem 1.3. *There exists a C depending only on α , the injectivity radius of the unique hyperbolic metric on X^i and $\|\Sigma_\alpha^i\|_F$ such that*

$$d(\Sigma_\alpha^i, \Sigma_t^i) \leq CL_C(\alpha)$$

for all $t \leq \alpha$.

We can also control the complex length of geodesics in M_t .

Theorem 1.4. *For each $L > 0$ there exists an $\epsilon > 0$ and $A > 0$ such that if γ is a simple closed curve in M with $L_\gamma(\alpha) \leq L$ and $L_c(\alpha) \leq \epsilon$ for all $c \in \mathcal{C}$, then*

$$e^{-AL_C(\alpha)} L_\gamma(\alpha) \leq L_\gamma(t) \leq e^{AL_C(\alpha)} L_\gamma(\alpha)$$

and

$$(1 - AL_C(\alpha))\Theta_\gamma(\alpha) \leq \Theta_\gamma(t) \leq (1 + AL_C(\alpha))\Theta_\gamma(\alpha)$$

for all $t \leq \alpha$.

We note that Theorems 1.3 and 1.4 are proved before Theorem 1.2 at least for all $t \leq \alpha$ where M_t is known to exist. In fact these results are used in the proof of Theorem 1.2.

2. DEFORMING $(PSL_2\mathbb{C}, X)$ -STRUCTURES

Although there is a very general theory of (G, X) -structures, for simplicity we will restrict ourselves to $(PSL_2\mathbb{C}, X)$ -structures here. In fact for our purposes X will either be $\widehat{\mathbb{C}}$ or \mathbb{H}^3 . Then a $(PSL_2\mathbb{C}, X)$ -structure on a manifold is an atlas of charts to X with transition maps in $PSL_2\mathbb{C}$; i.e., either a 2-dimensional projective or a 3-dimensional hyperbolic structure. We use the $(PSL_2\mathbb{C}, X)$ notation simply so that we can develop together the common elements of the deformation theory of projective and hyperbolic structures.

An equivalent way to define a $(PSL_2\mathbb{C}, X)$ -structure is through a developing map and holonomy representation. A *developing map* D is a local diffeomorphism from the universal cover \tilde{M} to X that commutes with a holonomy representation $\rho : \pi_1(M) \rightarrow PSL_2\mathbb{C}$. That is

$$(2.1) \quad D(\gamma(x)) = \rho(\gamma)(D(x))$$

for all $\gamma \in \pi_1(M)$ and $x \in \tilde{M}$.

Let $\mathcal{D}(M)$ be the space of developing maps for M which we topologize with the compact- C^∞ topology. We also define an equivalence relation for developing maps. We say $D_1 \sim D_2$ if there exists a diffeomorphism $\psi : M \rightarrow M$ isotopic to the identity and element $\alpha \in PSL_2\mathbb{C}$ such that $D_1 = \alpha \circ D_2 \circ \tilde{\psi}$ where $\tilde{\psi}$ is the unique lift of ψ that is equivariantly isotopic to the identity. Let $D(M)$ be the quotient space $\mathcal{D}(M)/\sim$.

To study one-parameter families of $(PSL_2\mathbb{C}, X)$ -structures we need to make a definition about vector fields on X and M . We say a vector field v on X is *geometric* if the homeomorphisms in the one-parameter flow it defines are elements of $PSL_2\mathbb{C}$. As is well known, the space of geometric vector fields is the Lie algebra $sl_2\mathbb{C}$. Geometric vector fields are analytic in the sense that any geometric vector field is determined uniquely by its germ at a single point. If M has a $(PSL_2\mathbb{C}, X)$ -structure, then a vector field v on M is geometric if for every chart ϕ , ϕ_*v is geometric.

A one-parameter family M_t of $(PSL_2\mathbb{C}, X)$ -structures on M can be defined through a one-parameter family of developing maps D_t and holonomy representations ρ_t . By taking the derivative of D_t , we can define a family of vector fields v_t on the universal covers \tilde{M}_t . More precisely, if x is a point in \tilde{M} , then $D_t(x)$ is a smooth path in X . The derivative $D'_t(x)$ will be a tangent vector to the path at $D_t(x)$. We pull back this tangent vector to $T_x\tilde{M}$ via D_t to define the vector field v_t at x .

Although these vector fields are defined on the differentiable manifold \tilde{M} , the vector field v_t has a special automorphic property in the $(PSL_2\mathbb{C}, X)$ -structure on M_t . Explicitly, by differentiating (2.1), we see that for each $\gamma \in \pi_1(M)$ the vector field $v - \gamma_*v$ is a geometric on \tilde{M}_t . We say that any vector field that satisfies this relationship is *automorphic*. To see how an automorphic vector field describes the infinitesimal change in geometry, we need to discuss projective structures and hyperbolic structures separately.

2.1. Projective structures. A $(PSL_2\mathbb{C}, \widehat{\mathbb{C}})$ -structure on a surface S is usually called a projective structure. As we noted when we first defined projective structures, a projective structure also defines a conformal structure on S and a fixed conformal structure X will have many distinct projective structures. We let $P(X)$ denote the space of projective structures on S with conformal structure X . $P(X)$ inherits a topology as a subspace of the space of developing maps $D(S)$. We will only be interested in projective structure deformations contained in $P(X)$, i.e., deformations preserving a given conformal structure X . This greatly simplifies the theory.

The objects that distinguish different projective structures in $P(X)$ are holomorphic quadratic differentials. In a local, conformal chart for X a holomorphic quadratic differential Φ has the form $\Phi(z) = \phi(z)dz^2$ where ϕ is a holomorphic function. On this local chart a conformal metric is determined by a positive, real valued function σ and has the form $\sigma(z)^2 dzd\bar{z}$. On a Riemann surface there is at most one complete, conformal hyperbolic metric and this will always be the metric we will use. In particular if σ is the hyperbolic metric on X , then $\|\Phi(z)\|_X = \sigma^{-2}|\phi(z)|$ is a well-defined function on S which we define to be the point-wise norm of Φ with respect to the σ metric. When it is clear which conformal structure is determining the metric, we will drop the subscript and write the norm $\|\Phi(z)\|$.

Our first construction of a holomorphic quadratic differential will come from a conformal vector field v on a projective structure Σ . In a local chart, v has the form $v(z) = g(z)\frac{\partial}{\partial z}$. Since v is conformal, g is a holomorphic function. The *Schwarzian derivative* Sv of v is a quadratic differential defined in a local projective chart by

$$Sv(z) = g_{zzz}(z)dz^2.$$

Note that this will only be a well-defined quadratic differential if the derivative of g is taken in projective charts. If it is taken in an arbitrary conformal chart, the equation will not define a quadratic differential. For projective structures a geometric vector field is usually called *projective*. As is well known, v will be projective if and only if $g(z)$ is a quadratic polynomial in z . In particular, $Sv \equiv 0$ if and only if v is a projective vector field. The flow of a projective field preserves the projective structure so the Schwarzian measures the infinitesimal change in projective structure.

Note that there are no global conformal vector fields on a closed Riemann surface of genus > 1 . The conformal vector fields we will be interested in are automorphic vector fields v on the universal cover $\tilde{\Sigma}$ of our projective structure Σ . Then Sv will be a holomorphic quadratic differential on $\tilde{\Sigma}$. However, by the automorphic property, Sv will be equivariant and will descend to a quadratic differential on Σ .

The second holomorphic quadratic differential we will construct will measure the distance between two projective structures Σ_0 and Σ_1 in $P(X)$. There is an obvious map f between Σ_0 and Σ_1 , namely the unique conformal map. It is the existence of this map that simplifies the deformation theory of projective structures in $P(X)$. The Schwarzian derivative of f , defined using projective charts for Σ_0 and Σ_1 , is the quadratic differential

$$(2.2) \quad Sf = \left[\left(\frac{f_{zz}}{f_z} \right)_z - \frac{1}{2} \left(\frac{f_{zz}}{f_z} \right)^2 \right] dz^2.$$

Again we must use projective charts for this equation to give a well-defined quadratic differential.

Conversely given any projective structure $\Sigma_0 \in P(X)$ and any holomorphic quadratic differential Φ on X , there exists a unique projective structure Σ_1 in $P(X)$ such that $\Phi = Sf$. In particular, after fixing Σ_0 as a basepoint, there is a canonical isomorphism from $P(X)$ to $Q(X)$, the space of holomorphic quadratic differentials on X . The space $Q(X)$ is a finite dimensional vector space. The L^∞ -norm on quadratic differentials makes $Q(X)$ a normed vector space. The identification of $P(X)$ with $Q(X)$ gives $P(X)$ a euclidean metric which will not depend on the choice of basepoint. In this metric the distance between two projective structures Σ_0 and Σ_1 will be $d(\Sigma_0, \Sigma_1) = \|Sf\|_\infty$.

A projective structure is *fuchsian* if its developing map is a homeomorphism onto a round disk in $\widehat{\mathbb{C}}$. The uniformization theorem implies that there is a unique fuchsian projective structure in $P(X)$. As we will frequently need to know the distance between an arbitrary projective structure $\Sigma \in P(X)$ and the unique fuchsian element Σ_F , we define $\|\Sigma\|_F = d(\Sigma, \Sigma_F)$.

A vector space is its own tangent space so the derivative of a smooth path Σ_t in $P(X) = Q(X)$ will also be a smooth path of holomorphic quadratic differentials Φ_t on X . To see this more explicitly, we let \tilde{X} be the universal cover of the conformal structure X and choose conformal developing maps $D_t : \tilde{X} \rightarrow \widehat{\mathbb{C}}$. The vector fields v_t obtained by differentiating D_t will be conformal and automorphic on $\tilde{\Sigma}_t$; therefore Sv_t will be a holomorphic quadratic differential on X . By noting that on suitably chosen local charts $f_t = D_t \circ D_0^{-1}$ and differentiating (2.2), we see that Sv_t is exactly Φ_t . This implies the following proposition:

Proposition 2.1. *The length of a smooth path, Σ_t with $a < t < b$, in $P(X)$ is*

$$\int_a^b \|\Phi_t\|_\infty dt.$$

2.2. Hyperbolic structures. For a family of hyperbolic structures there is no obvious choice of maps between the structures that plays the role of the conformal map in the deformation theory of projective structures. However, recent work has shown that there is a canonical choice of an automorphic vector field describing an infinitesimal deformation of the hyperbolic structure. We describe this in more detail below. Our review will be brief. See [HK1] for more details.

Assume M has a fixed hyperbolic metric coming from a developing map D with holonomy ρ . In hyperbolic space a geometric vector field is an infinitesimal isometry or *Killing field*. Let E and \tilde{E} be the bundles over M and \tilde{M} , respectively, of germs of Killing fields. For \tilde{M} the developing map identifies germs of Killing fields at a point in \tilde{M} with a Killing field on \mathbb{H}^3 so \tilde{E} has a global product structure, i.e. $\tilde{E} \cong \tilde{M} \times sl_2\mathbb{C}$. Then E is the quotient of \tilde{E} by the action of $\pi_1(M)$ where the action on the first factor is by deck transformations and on the second by the holonomy representation. The product structure on \tilde{E} defines a flat connection which descends to the quotient E . This flat connection has a covariant derivative d which we use to define the deRham cohomology groups $H^i(M; E)$. As we shall see next, an automorphic vector field on \tilde{M} determines a cohomology class in $H^1(M; E)$.

Given a vector field v on M (or \tilde{M}), the *canonical lift* s of v is the section of E (or \tilde{E}) determined by the relationship that $s(p)$ is the unique Killing field that agrees with v at p and whose curl agrees with the curl of v at p . If a section s of \tilde{E} is the canonical lift of an automorphic vector field on \tilde{M} , then $s - \gamma_*s$ will be a

constant section. Therefore $\omega = ds$ will descend to a closed E -valued 1-form on M which determines a cohomology class in $H^1(M; E)$.

If we let s_t be the canonical lift of v_t , then $\omega_t = ds_t$ is a family of E_t -valued 1-forms. Furthermore if the paths D_t and D'_t are equivalent in $D(M)$, the ω_t and ω'_t will be cohomologous. Therefore the derivative of a path in $D(M)$ is a one-parameter family of cohomology classes in $H^1(M_t; E_t)$. This cohomology class plays the role of the holomorphic quadratic differential in the study of projective structures.

For many calculations we will take advantage of the complex structure on E . In particular, the Lie algebra $sl_2\mathbb{C}$ has a complex structure that can be interpreted geometrically. If v is a Killing field on \mathbb{H}^3 , then $\text{curl } v$ is also a Killing field and $\text{curl curl } v = -v$. Therefore taking the curl of v is equivalent to multiplying by ι . For a section s of E this leads us to define ιs by the relationship $\iota s(p) = \text{curl}(s(p))$. We make a similar definition for E -valued n -forms. A section s of E is *real* if the Killing field $\iota s(p)$ is zero at p , while s is *imaginary* if $s(p)$ is zero at p . Every section s has a unique decomposition into a real section $\text{Re } s$ and an imaginary section $\text{Im } s$. A real section determines a vector field by the formula $v(p) = (s(p))(p)$ and vice versa. If s is an imaginary section, then ιs is a real section so the formula $v(p) = (\iota s(p))(p)$ also identifies each imaginary section with a vector field. Returning to a general section s , we have a map $s \mapsto (\text{Re } s, \text{Im } s)$. If we view both $\text{Re } s$ and $\text{Im } s$ as vector fields, this defines an isomorphism $E \rightarrow TM \oplus TM$. The canonical lift of a vector field v is then $(v, -\text{curl } v)$ under this isomorphism.

This identification of E with $TM \oplus TM$ gives E a natural metric; we simply use the hyperbolic metric on each copy of TM . This metric defines an isomorphism from E to the dual bundle E^* . For an E -valued k -form α we let α^\sharp be the image of α in E^* under this isomorphism, while if α is an E^* -valued k -form, we let α^\flat be the image of α under the inverse of the isomorphism. The bundle E^* has a flat connection d^* and we define $\partial\alpha = (d^*\alpha^\sharp)^\flat$. The formal adjoint for d defined on k -forms is $\delta = (-1)^k * \partial*$ where $*$ is the Hodge star operator. We also define the Laplacian $\Delta = d\delta + \delta d$.

In §1 of [HK1] there are explicit formulas for d and δ in local coordinates in terms of the Riemannian connection ∇ and algebraic operators. Let $\{e_i\}$ be an orthonormal frame field with dual co-frame field $\{\omega_i\}$. Then

$$(2.3) \quad d = \sum_i \omega^i \wedge (\nabla_{e_i} + \text{ad}(E_i))$$

and

$$(2.4) \quad \delta = \sum_j i(e_j)(\nabla_{e_j} - \text{ad}(E_j)).$$

Here $i(\)$ is the interior product on forms. The operator $\text{ad}(E_i)$ takes a Killing field Y to the Killing field $[E_i, Y]$. We also decompose d and δ into the real and imaginary parts. Namely we let $D = \text{Re } d$, $T = \text{Im } d$, $D^* = \text{Re } \delta$ and $T^* = \text{Im } \delta$. Formulas for D , T , D^* and T^* follow easily from (2.3) and (2.4) since ∇_{e_i} is a real operator and $\text{ad}(E_i)$ is an imaginary operator. In particular T and T^* are algebraic operators and therefore are easy to calculate. It is also worth noting that $\partial = D - T$. That is, the flat connection on E^* is the “conjugate” of the flat connection on E .

We make three more definitions that will be useful later. Let

$$\|\alpha\|^2 = \alpha \wedge *\alpha^\sharp.$$

Strictly speaking, $\|\alpha\|^2$ is a real valued 3-form while $\ast(\alpha \wedge \ast\alpha^\sharp)$ is a function. We will abuse notation and use $\|\alpha\|^2$ to refer to both the 3-form and the function. It will be clear from the context which meaning is correct. A vector field v is *harmonic* if $\Delta s = 0$ where s is the canonical lift of v . A closed E -valued 1-form ω on M is a *Hodge form* if it is the projection from \tilde{M} to M of ds , where s is the canonical lift of a harmonic, automorphic, divergence free vector field on \tilde{M} . By the work in §2 of [HK1] ω is a Hodge form if and only if ω is closed and co-closed and the real and imaginary parts of ω are symmetric and traceless vector valued 1-forms.

There is a nice formula for the L^2 -norm of a Hodge form:

Theorem 2.2 ([HK1]). *Let ω be an E -valued Hodge form on a compact hyperbolic 3-manifold M with boundary. Then*

$$2 \int_M \|\omega\|^2 = \int_{\partial M} \iota\omega \wedge \omega^\sharp$$

where ∂M is oriented with inward pointing normal.

2.3. Extending deformations to the projective boundary. Let M be a hyperbolic 3-manifold with projective boundary Σ . Together M and Σ form a differentiable 3-manifold with boundary so if v is a vector field on M , we can discuss its continuous extension to Σ and vice versa. We will always want the extended vector field to be tangent to the boundary. We will use this notion to discuss extending E -valued 1-forms on M to holomorphic quadratic differentials on Σ . Essentially an E -valued 1-form extends continuously to a holomorphic quadratic differential if the vector field that generates the 1-form extends continuously to a vector field that generates the quadratic differential.

Our vector fields will in general be automorphic vector fields on the universal cover. However, the extension property is a local one so we will work in a neighborhood V in $\mathbb{H}^3 \cup \widehat{\mathbb{C}}$. We assume that $V \cap \widehat{\mathbb{C}}$ is an open disk and that $V \cap \mathbb{H}^3$ is a topological ball. Let ω be a closed E -valued 1-form on $V \cap \mathbb{H}^3$ and let Φ be a holomorphic quadratic differential on $V \cap \widehat{\mathbb{C}}$. Since V is simply connected, $\omega = ds$ where s is an E -valued section on V and $\Phi = Sv_\infty$ where v_∞ is a conformal vector field on ∂V . Neither s nor v_∞ are uniquely determined; we can add a constant section to s and a projective vector field to v_∞ . We say that ω extends continuously to Φ if s and v_∞ can be chosen such that $\operatorname{Re} s$ extends continuously to v_∞ and $-\operatorname{Im} s$ extends continuously to v_∞ .

Returning to our hyperbolic 3-manifold M with projective boundary Σ , the E -valued 1-form ω extends continuously to the holomorphic quadratic differential Φ if it does so in a neighborhood of every point on Σ . In general, an E -valued 1-form is *conformal at infinity* if it is cohomologous to a 1-form that extends continuously to a holomorphic quadratic differential on Σ .

Recall the one-parameter family of cone-manifolds M_t that we defined in the previous section. The derivative of this path will be a path of cohomology classes $[\omega_t]$ in $H^1(M_t; E_t)$. The derivative of the projective boundary will be a path of holomorphic quadratic differentials Φ_t in $P(X)$. The following Hodge theorem is Theorem 4.3 in [Br1] plus Theorem 9.5 in the appendix of this paper.

Theorem 2.3. *The cohomology class $[\omega_t]$ is represented by a Hodge form ω_t which extends continuously to Φ_t on Σ_t . Furthermore the Hodge form will be in L^2 outside a neighborhood of the singular locus.*

Theorem 2.2 tells us how to calculate the L^2 -norm of a Hodge form on a compact hyperbolic 3-manifold with boundary. We will need to calculate the L^2 -norm of conformal Hodge forms on geometrically finite manifolds.

Theorem 2.4. *Let M be a hyperbolic 3-manifold with boundary such that the union of M with its projective boundary is compact. If ω is an L^2 Hodge form that extends continuously to a holomorphic quadratic differential on the projective boundary, then*

$$2 \int_M \|\omega\|^2 = \int_{\partial M} \omega \wedge \omega^\sharp$$

where ∂M is oriented with inward pointing normal.

Note that ∂M is the hyperbolic boundary of M and does not include the projective boundary Σ . That is, M is a smooth manifold with boundary. Its interior has a hyperbolic metric that extends to a smooth metric on ∂M . In the previous theorem the disjoint union $M \cup \Sigma$ is a smooth, compact manifold with boundary and the hyperbolic metric naturally extends to a projective structure on Σ .

3. TUBES AND HALF SPACES

In this section we make a digression from studying families of cone-manifolds to prove some results about a single hyperbolic cone-manifold $M_\alpha = (M, g)$ with cone-metric g and all cone-angles α . Our goal is to find a constant ℓ_0 such that if the length of the cone-singularity is less than ℓ_0 , it will have a “large” tubular neighborhood and this neighborhood will be disjoint from any embedded half space in the geometrically infinite ends. We will prove a succession of results, each producing its own constant. At the end of the section we will simply take the minimum of these constants to find a single constant which will be used throughout the rest of the paper.

We first review an estimate of Hodgson and Kerckhoff on the size of embedded tubes about the cone-locus. These should be thought of as cone-manifold versions of the Margulis lemma with explicit constants. The main difficulty is of course the cone-singularity. As a first stab at proving these results, one might hope to smooth the cone-metric and then apply the Margulis lemma for manifolds with pinched negative curvature. There are two problems with this approach. If the cone-angles are $< 2\pi$, then it may not be possible to smooth the metric to a negatively curved metric. On the other hand if there are cone-angles $> 2\pi$, we can always smooth the metric to one that is negatively curved; however we cannot bound the amount of negative curvature required even if all the cone-angles are bounded. To get around both of these problems, we assume a priori that the cone-singularity has a tubular neighborhood of definite size. In practice this is not much of a restriction since in most applications the hyperbolic cone-manifolds arise from smooth hyperbolic structures where the standard Margulis lemma applies.

Proposition 3.1. *Given an $\alpha > 0$ and an $R > 0$, there exists an $\ell_1 > 0$ such that the following holds. Let (M, g) be a hyperbolic cone-manifold with all cone-angles less than α and let γ be a closed non-singular geodesic in (M, g) . Suppose the tube radius of the cone-singularity is $\geq R$. If $L_\gamma(g) \leq \ell_1$ and $L_c(g) \leq \ell_1$ for all $c \in \mathcal{C}$, then γ has an embedded tube of radius R which is disjoint from the R -neighborhood of the cone singularity.*

Proof. We can construct a complete Riemannian metric h on M such that h agrees with g outside the $R/2$ -neighborhood of the cone-singularity and such that h has pinched negative sectional curvature (see [Ko2], Theorem 1.2.1). If $L_\gamma(g) \leq R/2$, then γ will be disjoint from the R neighborhood of \mathcal{C} and therefore γ will also be a geodesic in the h metric with $L_\gamma(g) = L_\gamma(h)$. Furthermore there is a universal bound on the sectional curvature of h depending only on our choice of tube radius $R/2$.

Let M_h^ϵ be the ϵ -thin part of M for the h metric. That is, M_h^ϵ is the subset of M consisting of those points whose injectivity radius is $< \epsilon$. By the Margulis lemma ([BGS]) there is an ϵ_0 , depending only on the curvature bounds, such that each component of $M_h^{\epsilon_0}$ has virtually nilpotent fundamental group. Since M is an orientable, hyperbolizable 3-manifold, the only possible virtually nilpotent subgroups of M are \mathbb{Z} and $\mathbb{Z} \oplus \mathbb{Z}$, and the second case will only occur at peripheral tori.

Let c be a component of \mathcal{C} and let V_c be the R -neighborhood of c in the g metric. Choose δ_1 such that if $L_c(g) \leq \delta_1$, then V_c will be contained in the ϵ_0 -thin part of the g metric. This is the one place where our choice depends on α for as the cone-angle increases, δ_1 will decrease. The h metric will decrease the injectivity radius in V_c so V_c will also be contained in $M_h^{\epsilon_0}$.

Next choose δ_2 such that if $L_\gamma(g) \leq \delta_2$, then the R -neighborhood V_γ of γ in the h metric is also contained in $M_h^{\epsilon_0}$. Note that δ_2 will only depend on the curvature bounds and not on the cone-angle. Any component of $M_h^{\epsilon_0}$ with fundamental group $\mathbb{Z} \oplus \mathbb{Z}$ does not contain any closed geodesics; therefore the component of $M_h^{\epsilon_0}$ containing γ must be a solid torus with fundamental group \mathbb{Z} . This implies that V_γ is a solid torus and is disjoint from V_c . Since V_γ is disjoint from V_c , the g and h metrics agree on V_γ so γ has an embedded tubular neighborhood of radius R in the original metric if $L_\gamma(g) \leq \ell_1$ where $\ell_1 = \min\{R/2, \delta_1, \delta_2\}$. \square

Combining this proposition with Theorem 4.4 in [HK4] we have

Proposition 3.2. *Let (M, g) be a hyperbolic cone-manifold with all cone-angles less than α and γ a closed geodesic in M_α . Suppose the tube radius of the cone-singularity is $\geq R$, $L_c(g) \leq \ell_1$ for each $c \in \mathcal{C}$ and $L_\gamma(g) \leq \ell_1$. Then for each $c \in \mathcal{C}$, c and the geodesic γ have disjoint tubular neighborhoods such that the area of the boundary tori is $\geq 1.6978 \frac{\sinh^2 R}{\cosh(2R)}$.*

We next prove similar results for hyperbolic half spaces.

Lemma 3.3. *Let D be an embedded round disk in the projective boundary of a hyperbolic cone-manifold (M, g) . Then there is an embedded hyperbolic half space H in M whose projective boundary is D .*

Proof. In general a hyperbolic half space H is bounded by a hyperbolic plane P and its projective boundary, a round disk D . The half space is foliated by planes P_d of constant curvature where P_d is the set of points a distance d from P . If D is an embedded round disk on the boundary of a geometrically finite cone-manifold, then the planes P_d will be embedded for large d . Let d' be the inf of all such d . If $d' > 0$, then the metric closure of $\bigcup_{d>d'} P_d$ will have strictly concave boundary so $M' = M \setminus \bigcup_{d>d'} P_d$ will have strictly convex boundary $P_{d'}$. This implies that $P_{d'}$ is embedded; hence we must have $d' = 0$. The only way P_0 cannot be embedded is for it to intersect an element c of the cone-singularity. In this case c must be tangent to

P_0 and therefore be contained in P_0 . Since P_0 does not contain a closed geodesic, it cannot intersect c . Therefore P_0 and hence all of H must be embedded. \square

Next we see that these embedded half spaces do not intersect the tubular neighborhood of the cone-singularity if it is sufficiently short.

Proposition 3.4. *Given an $\alpha > 0$ and an $R > 0$, there exists an $\ell_2 > 0$ such that the following holds. Let (M, g) be a hyperbolic cone-manifold with all cone-angles less than α . Suppose that the tube radius of the cone-singularity is $\geq R$ and that $L_c(g) \leq \ell_2$ for all $c \in \mathcal{C}$. Then each $c \in \mathcal{C}$ has a tubular neighborhood V disjoint from any embedded half space H with $\text{area}(\partial V) \geq 1.6978 \frac{\sinh^2 R}{\cosh(2R)}$.*

Proof. We first show that ℓ_2 can be chosen such that if $L_c(g) \leq \ell_2$ for all $c \in \mathcal{C}$, then H does not intersect the R -neighborhood of the cone-singularity. We need the following simple geometric fact. If U is a tube of radius $R' > R$ and H intersects the R -neighborhood of the core curve of the tube, then the area of the intersection of H with ∂U is bounded below by a function $A(R')$ with $A(R') \rightarrow \infty$ as $R' \rightarrow \infty$. Note that $\text{area}(\partial U) > A(R')$.

Choose R' such that $A(R') = 1.6978 \frac{\sinh R}{\cosh(2R)}$ and choose δ such that

$$\alpha \delta \sinh R' \cosh R' = 1.6978 \frac{\sinh R}{\cosh(2R)}.$$

Then let $\ell_2 = \min\{\ell_1, \delta\}$. If $L_c(g) \leq \ell_2$ for all $c \in \mathcal{C}$, then by Proposition 3.2 the R' -neighborhood of each $c \in \mathcal{C}$ will be an embedded tube U . Furthermore

$$\text{area}(\partial U) = \alpha L_c(g) \sinh R' \cosh R' \leq 1.6978 \frac{\sinh R}{\cosh(2R)}.$$

If H intersect the R -neighborhood of c , then

$$\text{area}(\partial U) > A(R') = 1.6978 \frac{\sinh R}{\cosh(2R)}.$$

This contradiction implies that H does not intersect the R -neighborhood.

The rest of the proof is similar to the proof of Theorem 4.4 in [HK4] and we will only sketch it. Define R_m to be the maximal radius such that the tube U_m of radius R_m about c is embedded and disjoint from H . We can assume that $T = \partial U_m$ intersects the hyperbolic plane $P = \partial H$ for otherwise we can simply apply Proposition 3.2. Except for possible self-tangencies T will be embedded and P will be tangent to T at a point p . Let B be the ball of radius R contained in H and tangent to T at p .

We now lift to the universal cover. Since M is hyperbolizable, any component (again ignoring self-tangencies) \tilde{T} of the pre-image of T in \tilde{M} will be a topological plane and the stabilizer of \tilde{T} will be a $\mathbb{Z} \oplus \mathbb{Z}$ subgroup Γ_T of $\pi_1(M)$. Let \tilde{p} be a point in the pre-image of p contained in \tilde{T} and let \tilde{B} be the component of the pre-image B that is tangent to \tilde{T} at \tilde{p} . Let C be the orthogonal projection of \tilde{B} onto \tilde{T} . One needs to be careful here to make sure that the cone-singularity does not interfere with this orthogonal projection. It is at this point that we refer to Hodgson and Kerckhoff. In particular they show that C is well defined and disjoint from its translates under the action of Γ_T . This implies that $\text{area}(T) \geq \text{area}(C)$. Hodgson and Kerckhoff also show $\text{area}(C) \geq 1.6978 \frac{\sinh^2 R}{\cosh(2R)}$. Therefore $\text{area}(T) \geq 1.6978 \frac{\sinh^2 R}{\cosh(2R)}$ as desired. \square

We now define two constants determined by a cone-metric g . For $c \in \mathcal{C}$ let R_g^c be chosen such that

$$\alpha L_c(g) \sinh R_g^c \cosh R_g^c = \frac{1}{2}.$$

Similarly if γ is a closed geodesic, let R_g^γ be chosen such that

$$2\pi L_\gamma(g) \sinh R_g^\gamma \cosh R_g^\gamma = \frac{1}{2}.$$

Let U_g^c and U_g^γ be the R_g^c and R_g^γ neighborhoods of c and γ , respectively, and let $U_g^c = \bigcup_{c \in \mathcal{C}} U_t^c$. Note that the area of both ∂U_g^c and ∂U_g^γ is $\frac{1}{2}$. The value of $\frac{1}{2}$ is essentially arbitrary. We could have chosen any constant less than $\lim_{R \rightarrow \infty} 1.6978 \frac{\sinh^2 R}{\cosh(2R)} = \frac{1.6978}{2}$.

In our next result we summarize the work of this section for a fixed choice of minimal tube radius. Our choice, although essentially arbitrary, will simplify some of the constants in the rest of the paper.

Theorem 3.5. *Given an $\alpha > 0$, there exists an $\ell_0 > 0$ such that the following holds. Let (M, g) be a hyperbolic cone-manifold with all cone-angles less than α and assume that the tube radius of the cone-singularity is greater than $\sinh^{-1} \sqrt{2}$. If $L_c(g) \leq \ell_0$ for all $c \in \mathcal{C}$, then the U_g^c are embedded tubular neighborhoods, each pairwise disjoint and disjoint from any embedded half space and $R_g^c \geq \sinh^{-1} \sqrt{2}$. Furthermore if γ is a closed geodesic with $L_\gamma(g) \leq \ell_0$, then U_g^γ is also an embedded tubular neighborhood disjoint from the U_g^c .*

Proof. Let ℓ_1 and ℓ_2 be the constants given by Propositions 3.2 and 3.4, respectively, with $R = \sinh^{-1} \sqrt{2}$. Choose ℓ_3 such that

$$\alpha \ell_3 \sinh R \cosh R = \alpha \ell_3 (\sqrt{2})(\sqrt{3}) = \frac{1}{2}.$$

Then $\ell_0 = \min\{\ell_1, \ell_2, \ell_3\}$ is the desired constant. □

For each point p in the projective boundary we will also need to estimate the size of the largest embedded round disk containing p . Here size will be measured by comparing the hyperbolic metric on the round disk to the hyperbolic metric on the entire projective boundary. By the Schwarz lemma the metric will always be bigger on the disk so we want to find a disk where we can bound the ratio of the two metrics. This bound will depend both on the injectivity radius of the hyperbolic metric and on the deviation of the projective boundary from a fuchsian projective structure.

We begin with a simple lemma about hyperbolic geometry.

Lemma 3.6. *Let D be a round disk in $\widehat{\mathbb{C}}$ with hyperbolic metric σ . Let γ be an isometry of \mathbb{H}^2 with translation length $\geq \ell$. For every $z \in D$ there exists a round disk $D' \subset D$ such that $D' \cap \gamma(D') = \emptyset$ and $\sigma'(z) = \coth(\ell/4)\sigma(z)$, where σ' is the hyperbolic metric for D' .*

Proof. Without loss of generality we assume that $D = \Delta$, the unit disk in \mathbb{C} , and $z = 0$. Then let D' be the euclidean disk of radius $\tanh(\ell/4)$ centered at 0. The hyperbolic diameter of D' in the σ metric is $\ell/2$ so $D' \cap \gamma(D') = \emptyset$. Finally

$$\sigma(z) = \frac{2}{1-|z|^2} \text{ while } \sigma'(z) = \frac{2 \tanh(\ell/4)}{\tanh^2(\ell/4)-|z|^2} \text{ so}$$

$$\sigma'(0) = \coth(\ell/4)\sigma(0)$$

as desired. □

Next, we use the previous lemma to estimate the size of embedded round disks in a projective structure.

Proposition 3.7. *Let Σ be a projective structure, σ the hyperbolic metric on Σ and κ the injectivity radius of σ . Then every $z \in \Sigma$ is contained in an embedded round disk D in Σ such that*

$$\sigma_D(z) < \sigma(z) \coth(\kappa/2) \sqrt{1 + 2\|\Sigma\|_F}$$

where σ_D is the hyperbolic metric on D .

Proof. Let X be the conformal structure induced by Σ and let Σ_F be the unique fuchsian structure in $P(X)$. Then $\tilde{\Sigma}_F$ is projectively isomorphic to D and the group of deck transformations Γ satisfies the conditions of Lemma 3.6. In particular for each $z \in \tilde{\Sigma}_0$ there is a round disk D' containing z such that for each $\gamma \in \Gamma$, $D \cap \gamma(D) = \emptyset$ and

$$(3.1) \quad \sigma_{D'}(z) = \coth(\kappa/2)\sigma(z)$$

where $\sigma_{D'}$ is the hyperbolic metric on D' . Therefore D' descends to an embedded round disk in Σ_F .

Let f be the unique conformal map from Σ_F to Σ . Then $f(D')$ will be an embedded topological disk but not a round disk in the projective structure Σ . However by Theorem 4.2 in [And] there exists a round disk in $D \subset f(D')$ in Σ with $z \in D$ such that

$$\sigma_D(z) \leq \sigma_{D'}(z) \sqrt{1 + 2\|Sf\|_{D',\infty}}.$$

By the Schwarz Lemma $\sigma_{D'} > \sigma$ so $\|Sf\|_{D',\infty} < \|Sf\|_{\Sigma,\infty} = \|\Sigma\|_F$. Combining the two inequalities with (3.1) gives

$$\sigma_D(z) < \sigma(z) \coth(\kappa/2) \sqrt{1 + 2\|\Sigma\|_F}. \quad \square$$

4. BOUNDING THE LENGTH OF GEODESICS

We can now return to investigating the one-parameter family of hyperbolic cone-manifolds M_t . As we shall see, the estimates we derive are simple consequences of the work of Hodgson and Kerckhoff.

Recall that the derivative of the path M_t is a cohomology class in $H^1(M_t; E_t)$ represented by a Hodge form ω_t . Throughout this section we assume that the one-parameter family is defined for t in a half-open interval $(\alpha', \alpha]$ and that at the starting structure $L_c(\alpha) \leq \ell_0$ and $R_\alpha^c > \sinh^{-1} \sqrt{2}$ for all $c \in \mathcal{C}$. Here ℓ_0 is the constant given in Theorem 3.5.

In our first result we show that the tube radius does not decay and we bound the length of the cone-singularity.

Proposition 4.1. *For all $c \in \mathcal{C}$, $L_c(t) \leq L_c(\alpha) \leq \ell_0$, $R_t^c > \sinh^{-1} \sqrt{2}$ and*

$$(4.1) \quad \frac{tL_c(\alpha)}{\alpha + 2L_c(\alpha)(\alpha^2 - t^2)} \leq L_c(t) \leq \frac{tL_c(\alpha)}{\alpha - 2L_c(\alpha)(\alpha^2 - t^2)}$$

if $t \in (\alpha', \alpha]$.

Proof. Before we begin, note that the expression on the far right of (4.1) is positive since $L_c(\alpha) \leq \ell_0$ and by the proof of Theorem 3.5, $\ell_0 \leq \frac{1}{2\sqrt{6}\alpha} < \frac{1}{2\alpha}$.

By assumption $L_c(\alpha) \leq \ell_0$ and $R_\alpha^c > \sinh^{-1} \sqrt{2}$. We will show that these two properties hold for all $t \in (\alpha', \alpha]$. If the tube radius condition does not hold, by the continuity of R_t^c , there exists a largest value $T < \alpha$ such that $R_T^c = \sinh^{-1} \sqrt{2}$. We will work by contradiction and show that such a T cannot exist.

To do so, we show that $L_c(t) \leq \ell_0$ for $t \in [T, \alpha]$, also working by contradiction. If this does not hold, there exists a T' with $L_c(T') = \ell_0$, $\frac{dL_c}{dt}(T') < 0$ and $L_c(t) \leq \ell_0$ if $t \in [T', \alpha]$. Theorem 3.5 implies that the tubular neighborhood of c of radius R_t^c is embedded for $t \in [T, \alpha]$ and therefore by an estimate of Hodgson and Kerckhoff (Theorem 2.7 of [HK4]) we have

$$(4.2) \quad \frac{L_c(t)}{t} \left(1 - \frac{1}{\sinh^2 R_t^c} \right) \leq \frac{dL_c}{dt}(t) \leq \frac{L_c(t)}{t} \left(1 + \frac{1}{\sinh^2 R_t^c} \right).$$

Since $\frac{1}{\sinh^2 R_t^c} \leq \frac{1}{2}$, the left-hand side of (4.2) implies $\frac{dL_c}{dt}(T') > 0$. This contradiction implies that $L_c(t) \leq \ell_0$ for all t between α and T . Since $TL_c(T) \sinh R_T^c \cosh R_T^c = \frac{1}{2}$, our choice of ℓ_0 implies that $R_T^c > \sinh^{-1} \sqrt{2}$. Again, this is a contradiction and hence $R_t^c > \sinh^{-1} \sqrt{2}$ for all $t \in (\alpha', \alpha]$. The previous argument also shows that $\frac{dL_c}{dt}(t) > 0$ and therefore $L_c(t) \leq L_c(\alpha) \leq \ell_0$ for all $t \in (\alpha', \alpha]$.

Next we combine the inequality

$$\frac{1}{\sinh^2 R_t^c} \leq \frac{2}{\sinh R_t^c \cosh R_t^c}$$

and the equality

$$tL_t(c) \sinh R_t^c \cosh R_t^c = \frac{1}{2}$$

with (4.2) to obtain

$$(4.3) \quad \frac{L_c(t)}{t} (1 - 4tL_c(t)) \leq \frac{dL_c}{dt}(t) \leq \frac{L_c(t)}{t} (1 + 4tL_c(t)).$$

To prove (4.1), we need to integrate this inequality. To do so we make the substitution $y(t) = \frac{L_c(t)}{t}$. Then the first inequality of (4.3) becomes

$$y + t \frac{dy}{dt} \geq y(1 - 4t^2y).$$

Rearranging and integrating, we get

$$\begin{aligned} \int_T^\alpha \frac{1}{y^2} \frac{dy}{dt} dt &\geq - \int_T^\alpha 4t dt, \\ -\frac{1}{y(\alpha)} + \frac{1}{y(T)} &\geq \frac{4(T^2 - \alpha^2)}{2}, \\ \frac{T}{L_c(T)} &\geq 2(T^2 - \alpha^2) + \frac{\alpha}{L_c(\alpha)}. \end{aligned}$$

As we remarked at the start of the proof, the right-hand side of the final inequality is positive since $L_c(\alpha) < \frac{1}{2\alpha}$. This final inequality is equivalent to the second inequality of (4.1). The other inequality is derived similarly. \square

Recall that in §3 we made a standard choice of tubular neighborhoods of the singular locus in M_t . We labeled the union of these neighborhoods U_t^c (see the

paragraph proceeding Theorem 3.5). The L^2 -norm of ω_t will be infinite on all of M_t . However if we let $M'_t = M_t \setminus U_t^c$, then the L^2 -norm will be bounded on M'_t .

Proposition 4.2.

$$\int_{M'_t} \|\omega_t\|^2 \leq L_C(t)^2.$$

Futhermore for any A and R there exists a K such that if for each $c \in \mathcal{C}$, V^c is a tubular neighborhood of c with the $\text{area}(\partial V^c) \geq A$ and the radius of V^c greater than R , then

$$\int_{M_t \setminus V^c} \|\omega_t\|^2 \leq K^2 L_C(t)^2$$

where $V^c = \bigcup_{c \in \mathcal{C}} V^c$.

Proof. By Proposition 4.1, $R_t^c > \sinh^{-1} \sqrt{2}$ for all $c \in \mathcal{C}$ and $t \in (\alpha', \alpha]$. By Theorem 3.5 the tube U_t^c of radius R_t^c about c is embedded. Using Theorem 2.4 along with the work of Hodgson and Kerckhoff (see (17) on p. 14 of [HK4]), we see that

$$\int_{M'_t} \|\omega_t\|^2 \leq \sum_{c \in \mathcal{C}} \frac{\cosh R_t^c \text{area}(\partial U_t^c)}{t^2 \sinh^3 R_t^c (2 \cosh^2 R_t^c + 1)}.$$

From the area formula for the tube boundary and our definition of R_t^c , $\text{area}(\partial U_t^c) = t L_t(c) \sinh R_t^c \cosh R_t^c = \frac{1}{2}$. Using this and also the fact that $\sinh R_t^c > \sqrt{2}$, we have

$$\begin{aligned} \frac{\cosh R_t^c \text{area}(\partial U_t^c)}{t^2 \sinh^3 R_t^c (2 \cosh^2 R_t^c + 1)} &= \frac{2 L_C(t)^2 \cosh^3 R_t^c}{\sinh R_t^c (2 \cosh^2 R_t^c + 1)} \\ &\leq L_C(t)^2 \end{aligned}$$

and therefore

$$\int_{M'_t} \|\omega_t\|^2 \leq \sum_{c \in \mathcal{C}} L_C(t)^2 \leq L_C(t)^2.$$

The proof of the more general inequality is essentially the same. □

Remark. In [HK4] instead of using the cone-angle as the parameter for the family of the hyperbolic cone-manifolds, they use the cone-angle squared. This accounts for the difference in the constants in their paper and the constants in this paper.

We can also bound the length of short curves that are not part of the cone-singularity:

Proposition 4.3. *If γ is a simple closed curve in M with $L_\gamma(T) \leq e^{-4\alpha\ell_0} \ell_0$ for some $T \in (\alpha', \alpha]$, then*

$$(4.4) \quad e^{-4\alpha L_C(\alpha)} L_\gamma(\alpha) \leq L_\gamma(t) \leq e^{4\alpha L_C(\alpha)} L_\gamma(\alpha)$$

and

$$(4.5) \quad (1 - 4\ell_0 L_C(\alpha)) \Theta_\gamma(\alpha) \leq \Theta_\gamma(t) \leq (1 + 4\ell_0 L_C(\alpha)) \Theta_\gamma(\alpha)$$

for all $t \in (\alpha', \alpha]$.

Proof. By Proposition 4.1, $R_t^c > \sinh^{-1} \sqrt{2}$ for all $c \in \mathcal{C}$ and $t \in (\alpha', \alpha]$. Therefore by Theorem 3.5, if $L_\gamma(t) \leq \ell_0$, the tube U_t^γ is embedded and contained in M_t' .

We need to show that the L^2 -norm of ω_t on U_t^γ is bounded by the derivative $\mathcal{L}'_\gamma(t)$. The essential estimates again come from §2 of [HK4]. In particular they decompose a Hodge form in a tubular neighborhood of a component of the singular locus as the sum of certain Hodge forms of standard type and a correction term. To apply this decomposition to our situation, we view γ as a component of the singular locus with cone-angle 2π . According to this decomposition, in the complement of γ on U_t^γ

$$\omega_t = C_m \omega_m + C_l \omega_l + \omega_c$$

where ω_m and ω_l are radially symmetric Hodge forms with C_m and C_l complex constants while ω_c represents a trivial cohomology class on $U_t^\gamma \setminus \gamma$.

The constants C_m and C_l are determined by the derivatives of the complex lengths of the meridian and longitude (recall that we are viewing γ as a component of the singular locus). The complex derivative of the meridian is zero since the cone-angle is fixed at 2π . By Lemma 2.1 in [HK4], $C_m = 0$ and $C_l = \frac{\mathcal{L}'_\gamma(t)}{2L_\gamma(t)}$. That is,

$$\omega_t = \frac{\mathcal{L}'_\gamma(t)}{2L_\gamma(t)} \omega_l + \omega_c.$$

For an explicit description of ω_l see p. 9 of [HK4]. For our purposes we only need the following two facts. First, by Lemma 2.5 in [HK4]

$$\int_{U_t^\gamma} \|\omega_t\|^2 = \left(\frac{|\mathcal{L}'_\gamma(t)|}{2L_\gamma(t)} \right)^2 \int_{U_t^\gamma} \|\omega_l\|^2 + \int_{U_t^\gamma} \|\omega_c\|^2$$

and second, by the formulas on p. 14 of [HK4]

$$\int_{U_t^\gamma} \|\omega_l\|^2 = \frac{\sinh R_t^\gamma}{\cosh R_t^\gamma} \left(2 + \frac{1}{\cosh^2 R_t^\gamma} \right) \text{area}(\partial U_t^\gamma).$$

Together this implies

$$\int_{U_t^\gamma} \|\omega_t\|^2 \geq \frac{\sinh R_t^\gamma}{\cosh R_t^\gamma} \left(2 + \frac{1}{\cosh^2 R_t^\gamma} \right) \left(\frac{|\mathcal{L}'_\gamma(t)|}{2L_\gamma(t)} \right)^2 \text{area}(\partial U_t^\gamma).$$

Since $\sinh R_t^\gamma > \sqrt{2}$ and $\text{area}(\partial U_t^\gamma) = \frac{1}{2}$,

$$\frac{\sinh R_t^\gamma}{\cosh R_t^\gamma} \left(2 + \frac{1}{\cosh^2 R_t^\gamma} \right) \text{area}(\partial U_t^\gamma) \geq \frac{\sqrt{2}}{\sqrt{3}} \left(2 + \frac{1}{3} \right) \left(\frac{1}{2} \right) > \frac{1}{4}.$$

We also know that

$$\int_{M_t'} \|\omega_t\|^2 \geq \int_{U_t^\gamma} \|\omega_t\|^2.$$

By Proposition 4.2

$$L_C(t)^2 \geq \int_{M_t'} \|\omega_t\|^2.$$

By combining these four inequalities, we have

$$L_C(t)^2 \geq \left(\frac{|\mathcal{L}'_\gamma(t)|}{4L_\gamma(t)} \right)^2$$

which rearranges to give

$$(4.6) \quad |\mathcal{L}'_\gamma(t)| \leq 4L_\gamma(t)L_C(t)$$

if $L_\gamma(t) \leq \ell_0$.

Next we show that if $L_\gamma(T) \leq e^{-4\alpha\ell_0}\ell_0$, then $L_\gamma(t) \leq \ell_0$ for all $t \in (\alpha', \alpha]$. Let I be the largest interval containing T such that $L_\gamma(t) \leq \ell_0$ if $t \in I$. We will show that I is an open and closed subset of $(\alpha', \alpha]$. By the continuity of the $L_\gamma(t)$, I is closed. Furthermore if T' is the right endpoint of I , then either $L_\gamma(T') = \ell_0$ or $T' = \alpha$. Since $|L'_\gamma(t)| \leq |\mathcal{L}'_\gamma(t)|$ and $L_C(t) \leq L_C(\alpha)$, (4.6) becomes

$$(4.7) \quad |L'_\gamma(t)| \leq 4L_\gamma(t)L_C(\alpha)$$

if $t \in I$. By integrating (4.7) from T to T' , we get

$$(4.8) \quad e^{-4|T'-T|L_C(\alpha)}L_\gamma(T) \leq L_\gamma(T') \leq e^{4|T'-T|L_C(\alpha)}L_\gamma(T).$$

Since $|T' - T| \leq \alpha$ with equality only holding if $T' = \alpha$ and $T = 0$, the right-hand side of (4.8) implies that either $L_\gamma(T') < \ell_0$, which contradicts the definition of I , or $T' = \alpha$. A similar argument shows that the left endpoint of I is α' and $I = (\alpha', \alpha]$. Therefore we can integrate (4.7) to get (4.4).

To prove (4.5), we note that (4.6) implies that

$$|\Theta'_\gamma(t)| \leq 4L_\gamma(t)L_C(t) \leq 4\ell_0L_C(\alpha).$$

Integrating this inequality gives (4.5). □

Recall that the *injectivity radius* at a point x in a Riemannian manifold is the radius of the largest embedded ball centered at x or equivalently half the length of the shortest geodesic arc with both endpoints at x . We define $\text{inj}_x(t)$ to be the injectivity radius at x for the g_t metric on M . Recall that $M'_t = M \setminus U_t^C$ is the complement of the tubular neighborhoods of the singular locus. We then have the following proposition.

Proposition 4.4. *The injectivity radius $\text{inj}_x(t)$ is bounded away from zero on M'_t for all $t \leq \alpha$.*

Proof. Note that the function $\text{inj}_x(t)$ measures the injectivity of x in all of M but we are only showing that it is bounded below on M'_t . Clearly as $t \rightarrow 0$ there will be points $x_t \in U_t^C$ such that $\text{inj}_{x_t}(t) \rightarrow 0$.

If the injectivity radius is not bounded below, there are two possibilities.

First, there could be a simple closed curve γ in M such that $L_\gamma(t) \rightarrow 0$. This is not possible by (4.4).

The second possibility is that there are points $x_t \in \partial U_t^C$ for some $c \in \mathcal{C}$ such that $\text{inj}_{x_t}(t) \rightarrow 0$ as $t \rightarrow 0$. The tori ∂U_t^C have an induced Euclidean metric and it is not hard to see that the hyperbolic injectivity radius will decay to zero if and only if the Euclidean injectivity radius decays to zero.

The Euclidean metric on ∂U_t^C can be constructed by gluing together (possibly with a twist) the boundary components of a Euclidean cylinder of height $L_c(t) \cosh R_t^c$ and radius $t \sinh R_t^c$. The area of this cylinder $tL_c(t) \sinh R_t^c \cosh R_t^c$ is always $\frac{1}{2}$ so the injectivity radius will be bounded below if and only if the height $H(t)$ is bounded above and below. By (4.1), $\frac{L_c(t)}{t}$ is bounded above and below. Since $R_t^c \geq \sinh^{-1} \sqrt{2}$, $\frac{\cosh R_t^c}{\sinh R_t^c}$ is also bounded above and below and therefore so is $\frac{L_c(t) \cosh R_t^c}{t \sinh R_t^c}$. Finally by multiplying the numerator and denominator by

$L_c(t) \cosh R_t^c$, we have that

$$\frac{(L_c(t) \cosh R_t^c)^2}{tL_c(t) \sinh R_t^c \cosh R_t^c} = 4H(t)^2$$

is also bounded above and below, as desired. □

We can also bound the length of arbitrary geodesics although the estimates are slightly different.

Lemma 4.5. *Let (M, g) be a hyperbolic cone-manifold. For each $c \in \mathcal{C}$ assume that U_g^c is an embedded tubular neighborhood with $R_g^c > \sinh^{-1} \sqrt{2}$ and assume that $L_c(g) \leq \ell_0$. For each $L > 0$ there exists an $\epsilon > 0$ such that if γ is a closed geodesic in M with $\epsilon < L_\gamma(g) < L$, then γ is disjoint from the M_g^ϵ , the ϵ -thin part of (M, g) .*

Proof. Given our upper and lower bounds for $L_c(g)$ and R_g^c , respectively, there exists a ‘‘Margulis constant’’ ϵ_0 depending only on α such that $M_g^{\epsilon_0}$ consists of tubes and rank two cusps. Furthermore for any $K > 0$ we can choose an $\epsilon(K) < \epsilon_0$ such that the distance between $\partial M_g^{\epsilon_0}$ and $\partial M_g^{\epsilon(K)}$ is greater than K . The number $\epsilon(K)$ will only depend on α and K . Let $\epsilon = \epsilon(L/2)$. Therefore if γ intersects M_g^ϵ , it will be entirely contained in a component of $M_g^{\epsilon_0}$. The only closed geodesic in a component of $M_g^{\epsilon_0}$ will be a core curve of one of the tubes; hence $L_\gamma(g) \leq \epsilon$. □

Lemma 4.6. *Let γ be a closed, non-singular geodesic in M_t such that $\|\omega_t(p)\| \leq K$ for all $p \in \gamma$. Then*

$$|\mathcal{L}'_\gamma(t)| \leq \sqrt{\frac{2}{3}} K L_\gamma(t).$$

Proof. Let M_t^γ be the cover of M_t associated to γ . For small values of R , γ will have an embedded tubular neighborhood $U(R)$ of radius R in M_t^γ . The E -valued 1-form ω_t lifts to M_t^γ and let $K(R)$ be an upper bound for $\|\omega_t(p)\|$ for all $p \in U(R)$. Then $K(R)$ is continuous and $K(0) = K$. As we noted in the proof of Proposition 4.3, Hodgson and Kerckhoff show that

$$\int_{U(R)} \|\omega_t\|^2 \geq \frac{\sinh R}{\cosh R} \left(2 + \frac{1}{\cosh^2 R} \right) \left(\frac{|\mathcal{L}'_\gamma(t)|}{2L_\gamma(t)} \right)^2 \text{area}(\partial U(R)).$$

We also know that

$$\int_{U(R)} \|\omega_t\|^2 \leq K(R)^2 \int_{U(R)} dV = K(R)^2 \pi L_\gamma(t) \sinh^2 R$$

and

$$\text{area}(\partial U(R)) = 2\pi L_\gamma(t) \sinh R \cosh R.$$

Together this implies that

$$K(R)^2 \geq 2 \left(2 + \frac{1}{\cosh^2 R} \right) \left(\frac{|\mathcal{L}'_\gamma(t)|}{2L_\gamma(t)} \right)^2.$$

Taking the limit of both sides as $R \rightarrow 0$ and rearranging terms gives the desired inequality. □

We are now ready to prove one of our main theorems, stated in the introduction.

Theorem 1.4. *For each $L > 0$ there exists an $\epsilon > 0$ and an $A > 0$ such that if γ is a simple closed curve in M with $L_\gamma(\alpha) \leq L$ and $L_c(\alpha) \leq \epsilon$ for all $c \in \mathcal{C}$, then*

$$e^{-AL_c(\alpha)}L_\gamma(\alpha) \leq L_\gamma(t) \leq e^{AL_c(\alpha)}L_\gamma(\alpha)$$

and

$$(1 - AL_c(\alpha))\Theta_\gamma(\alpha) \leq \Theta_\gamma(t) \leq (1 + AL_c(\alpha))\Theta_\gamma(\alpha)$$

for all $t \leq \alpha$.

Proof. We begin by noting that if $L_\gamma(t) \leq e^{-4\alpha\ell_0}\ell_0$ for any $t \leq \alpha$, then the theorem follows from Proposition 4.3 with $\epsilon = \ell_0$ and $A = 4L_c(\alpha)$. Therefore we will assume for the remainder of the proof that $L_\gamma(\alpha) \leq L$ but $L_\gamma(t) > e^{-4\alpha\ell_0}\ell_0$ for all $t \leq \alpha$.

By Lemma 4.5 there exists a $\delta > 0$ such that if $\delta < L_\gamma(t) \leq 2L$, then the geodesic representative γ_t of γ in M_t is disjoint from M_t^δ . We assume that $\delta \leq e^{-4\alpha\ell_0}\ell_0$.

We need to bound the pointwise norm of ω_t on the geodesic γ_t . To do so, we bound the L^2 -norm of ω_t using Proposition 4.2 and then apply the mean value theorem developed in the appendix.

Choose $\epsilon_1 > 0$ such that if $L_c(t) < \epsilon_1$, then $R_t^c > 3L$. Then if $L_\gamma(t) < 2L$, γ_t will not intersect the radius $R_t^c/3$ tube about c . In fact any ball B_δ centered at a point p on γ_t will not intersect this tube. Furthermore, since $R_t^c > \sinh^{-1}\sqrt{2}$, the area of the boundary of this tube will be universally bounded below and therefore by Proposition 4.2 there exists a K_1 such that

$$\int_{B_\delta} \|\omega_t\|^2 \leq (K_1 L_c(\alpha))^2.$$

Let

$$K = \frac{3\sqrt{2 \operatorname{vol}(B_\delta)}K_1}{2\pi(\cosh(\delta) \sin(\sqrt{2}\delta) - \sqrt{2} \sinh(\delta) \cos(\sqrt{2}\delta))}.$$

Then the norm bound and Theorem 9.9 imply that

$$\|\omega_t(p)\| \leq K L_c(t)$$

and therefore

$$|\mathcal{L}'_\gamma(t)| \leq \sqrt{\frac{2}{3}} K L_c(\alpha) L_\gamma(t)$$

if $\delta < L_\gamma(t) \leq 2L$. Next we choose ϵ_2 such that

$$e^{\sqrt{\frac{2}{3}}K\epsilon_2} \leq 2$$

and we let $\epsilon = \min\{\epsilon_1, \epsilon_2\}$ and $A = \sqrt{\frac{2}{3}}K$. The rest of the argument is a repeat of the proof of Proposition 4.3. In particular we can first show that $L_\gamma(t) \leq 2L$ for all $t \leq \alpha$ and then integrate the derivative bound to get the final estimate. \square

5. BOUNDING THE SCHWARZIAN DERIVATIVE

Recall that Σ_t^i is the projective boundary of M_t corresponding to a component of ∂N . The Hodge forms ω_t extend to a holomorphic quadratic differential Φ_t^i on Σ_t^i . We now use our bound on the L^2 -norm of ω_t to bound the L^∞ -norm of Φ_t^i . To do so, we first need a local result: given a bound on the L^2 -norm of a Hodge form on a half space H , we bound the L^∞ -norm of the quadratic differential on the projective boundary, a round disk D .

We begin with a lemma from complex analysis.

Lemma 5.1. *Let Φ be a holomorphic quadratic differential on a round disk D . Then*

$$\|\Phi\|_2 \geq 2\sqrt{\frac{\pi}{3}}\|\Phi(z)\|$$

for all $z \in D$.

Proof. Without loss of generality we assume that D is the unit disk Δ in \mathbb{C} and $z = 0$. The hyperbolic metric on Δ is $\sigma(z) = \frac{2}{1-|z|^2}$ and the area form in polar coordinates is $\sigma^2 r dr d\theta$. Then $\Phi = \phi dz^2$ where ϕ is a holomorphic function on Δ . Let $\phi(z) = \sum_{n=0}^{\infty} a_n z^n$ be the Taylor series for ϕ so $\|\Phi(0)\| = \frac{|a_0|}{4}$. We then calculate

$$\begin{aligned} \|\Phi\|_2^2 &= \int_{\Delta} |\phi|^2 \sigma^{-4} dA \\ &= \sum_{n,m} \int_0^1 \int_0^{2\pi} a_n \overline{a_m} r^{n+m} e^{i(n-m)\theta} \sigma^{-2} r d\theta dr \\ &= \frac{\pi}{2} \sum_n |a_n|^2 \int_0^1 (1-r^2)^2 r^{2n+1} dr \\ &\geq \frac{\pi}{2} |a_0|^2 \int_0^1 (1-r^2)^2 r dr \\ &= \frac{\pi}{12} |a_0|^2 \\ &= \frac{4\pi}{3} \|\Phi(0)\|^2. \end{aligned}$$

□

Next we construct an extension of a holomorphic quadratic differential Φ on D to an E -valued 1-form ω_Φ on H . We will use a non-standard model for \mathbb{H}^3 . Namely, let $\mathbb{H}^3 = D \times \mathbb{R}$ with the Riemannian metric $\sigma^2 \cosh^2 t dx^2 + \sigma^2 \cosh^2 t dy^2 + dt^2$ where σ is the hyperbolic metric on D . Then $H = D \times [0, \infty)$ is then a half space in \mathbb{H}^3 whose projective boundary is naturally identified with D . We also fix an orthonormal framing of \mathbb{H}^3 by letting $\omega^1 = \sigma \cosh t dx$, $\omega^2 = \sigma \cosh t dy$ and $\omega^3 = dt$. Let e_1, e_2 and e_3 be the vector fields dual to the ω^i . Define E_i to be the lift of e_i to $\text{Re } E$.

The holomorphic quadratic differential Φ is the Schwarzian derivative of a conformal vector field $v = f \frac{\partial}{\partial z}$ where f is a holomorphic function with $\Phi = f_{zzz} dz^2$. To extend Φ , we first use v to construct a section of E . At each point $w = x + iy$ in D let

$$(5.1) \quad s_\infty(w) = \left[f(w) + f_z(w)(z-w) + \frac{f_{zz}(w)}{2}(z-w)^2 \right] \frac{\partial}{\partial z}$$

be the projective vector field that best approximates v at w . A projective vector field extends to a Killing field in \mathbb{H}^3 so the function $s(w, t) = s_\infty(w)$ is an E -valued section on H . The vector fields $\text{Re } s$ and $-\text{Im } s$ on H will extend continuously to v and w on D so $\omega_\Phi = ds$ extends continuously to Φ . We will show that ω_Φ minimizes the L^2 -norm among all Hodge forms that extend continuously to Φ .

To do any calculations, we need to understand how a projective vector field on $\widehat{\mathbb{C}}$ extends continuously to a Killing field on \mathbb{H}^3 . It turns out that all projective fields that appear in our calculations are parabolic. The parabolic vector field $v = \frac{\partial}{\partial z}$

has a very simple extension if we use the upper half space model of \mathbb{H}^3 . That is, if $\mathbb{H}^3 = \{(x, y, t) | t > 0\}$, then v extends to the Killing $\frac{\partial}{\partial x}$ (here $z = x + iy$).

To extend the parabolic vector field $\frac{(z-w)^2}{2} \frac{\partial}{\partial z}$ to \mathbb{H}^3 in our unusual coordinate system, we need to perform a change of coordinates. Let $r_w = \{w\} \times \mathbb{R}$ be a geodesic in \mathbb{H}^3 . The next lemma describes the extension of $\frac{(z-w)^2}{2} \frac{\partial}{\partial z}$ to a Killing field on the geodesic r_w .

Lemma 5.2. *The E -valued section $-\sigma^{-1}e^{-t}(E_1 + \iota E_2)$ evaluated on the geodesic r_w is a Killing field which extends to the projective vector field $\frac{(z-w)^2}{2} \frac{\partial}{\partial z}$ on D .*

Proof. Let v be the Killing field in \mathbb{H}^3 that continuously extends to $\frac{(z-w)^2}{2}$ on D . Then v restricted to r_w will be $-\sigma^{-1}e^{-t}e_1$. We also know that $\text{curl } v$ is the Killing field that continuously extends to the projective vector field $\frac{i(z-w)^2}{2} \frac{\partial}{\partial z}$ on D . From this we see that $\text{curl } v$ restricted to r_w is the vector field $\sigma^{-1}e^{-t}e_2$.

The section $-\sigma^{-1}e^{-t}(E_1 + \iota E_2)$ evaluated at any point p is the unique Killing field w with $w(p) = -\sigma^{-1}e^{-t}e_1$ and $(\text{curl } w)(p) = \sigma^{-1}e^{-t}e_2$. Therefore for $p \in r_w$, $w = v$, proving the lemma. \square

Next we calculate the pointwise and L^2 -norms of ω_Φ .

Lemma 5.3. *If Φ is a holomorphic quadratic differential on D , then ω_Φ is a Hodge form. Furthermore the pointwise norm of ω_Φ is*

$$(5.2) \quad \|\omega_\Phi(w, t)\| = 2e^{-t} \text{sech } t \|\Phi(w)\|$$

while the L^2 -norm on H is

$$(5.3) \quad \int_H \|\omega_\Phi\|^2 = 2\|\Phi\|_2^2.$$

Proof. Recall that $\Phi = \phi dz^2 = f_{zz} \frac{\partial}{\partial z}$ where f and hence ϕ are holomorphic functions. Since $s(w, t) = s_\infty(w)$,

$$\omega_\Phi(w, t) = ds(w, t) = ds_\infty(w) = \left[\phi(w) \frac{(z-w)^2}{2} \frac{\partial}{\partial z} \right] dw.$$

Note that the last equality is obtained by differentiating (5.1) with respect to w . We need to rewrite this expression in terms of E_i and ω^i . We first note that $dw = dx + i dy = \sigma^{-1} \text{sech } t (\omega^1 + i \omega^2)$. From Lemma 5.2 we know that

$$\frac{(z-w)^2}{2} \frac{\partial}{\partial z} = -\sigma^{-1} e^{-t} (E_1 + \iota E_2).$$

Together this implies that

$$(5.4) \quad \omega_\Phi(w, t) = -\sigma^{-2} e^{-t} \text{sech } t \phi(w) (E_1 + \iota E_2) (\omega^1 + i \omega^2).$$

Next we show that ω_Φ is a Hodge representative. First we know that ω_Φ is closed since $d\omega_\Phi = d(ds_\Phi) = 0$. To see that ω_Φ is co-closed, we need to calculate $\delta\omega_\Phi$. This can be done using the formula for δ given in §2. Finally, to see that s is the canonical lift of a divergence free vector field, we note that $\omega_\Phi = ds$ is symmetric and traceless. The result then follows from the work in §2 of [HK1].

To calculate the pointwise norm of ω_Φ on H , we note that

$$\|(E_1 + \iota E_2)(\omega^1 + i \omega^2)\|^2 = 4\omega^1 \wedge \omega^2 \wedge \omega^3$$

from which it follows that

$$\begin{aligned} \|\omega_\Phi(w, t)\|^2 &= 4\sigma^{-4}e^{-2t} \operatorname{sech}^2 t |\phi(w)|^2 \omega^1 \wedge \omega^2 \wedge \omega^3 \\ &= 4e^{-2t} \operatorname{sech}^2 t \|\Phi(w)\|^2 \omega^1 \wedge \omega^2 \wedge \omega^3. \end{aligned}$$

Next we calculate the L^2 -norm:

$$\begin{aligned} \int_H \|\omega_\Phi(w, t)\|^2 &= \int_H 4e^{-2t} \operatorname{sech}^2 t \|\Phi(w)\|^2 \omega^1 \wedge \omega^2 \wedge \omega^3 \\ &= \left(\int_0^\infty 4e^{-2t} dt \right) \left(\int_D \|\Phi(w)\|^2 \sigma^2 dx dy \right) \\ &= 2\|\Phi\|_2^2. \end{aligned}$$

□

We can now show that ω_Φ minimizes the L^2 -norm.

Theorem 5.4. *Let ω be a Hodge form on H that extends continuously to a holomorphic quadratic differential Φ on D and assume that Φ extends to a neighborhood of D in $\hat{\mathbb{C}}$. Then*

$$\int_H \|\omega\|^2 \geq \int_H \|\omega_\Phi\|^2 \geq \frac{8\pi}{3} \|\Phi(z)\|^2$$

for all $z \in D$.

Proof. Let p be a point on P and H_t the set of points in H that are within a distance t from p . Then the boundary of H_t consists of a disk P_t in the plane P and a hemisphere S_t . By Theorem 2.2

$$2 \int_{H_t} \|\omega\|^2 = \int_{P_t \cup S_t} \iota\omega \wedge \omega^\sharp.$$

We first examine the integral over S_t . By definition of the Hodge star operator $\iota\omega \wedge \omega^\sharp = *_\partial(\iota\omega \wedge \omega^\sharp)dA$ where $*_\partial$ is the star operator for S_t and dA the area form, while $\omega \wedge *\omega^\sharp = *(\omega \wedge *\omega^\sharp)dV$ where dV is the volume form. Both $*_\partial(\iota\omega \wedge \omega^\sharp)$ and $*(\omega \wedge *\omega^\sharp)$ are real valued functions with $|*_\partial(\iota\omega \wedge \omega^\sharp)| \leq *(\omega \wedge *\omega^\sharp)$. Since ω is in L^2 on H and $dV = dAdt$, we have

$$\int_0^\infty \int_{S_t} *(\omega \wedge *\omega^\sharp)dAdt = \int_H \omega \wedge *\omega^\sharp \leq \infty$$

and therefore

$$\lim_{t \rightarrow \infty} \int_{S_t} *(\omega \wedge *\omega^\sharp)dA = 0.$$

Since $|*_\partial(\iota\omega \wedge \omega^\sharp)| \leq *(\omega \wedge *\omega^\sharp)$, we have

$$\lim_{t \rightarrow \infty} \int_{S_t} |*_\partial(\iota\omega \wedge \omega^\sharp)|dA = 0$$

and

$$\int_H \|\omega\|^2 = \lim_{t \rightarrow \infty} \int_{P_t} \iota\omega \wedge \omega^\sharp.$$

We now calculate the integral over the disk P_t . To do so, we decompose the Hodge form ω as the sum of the model Hodge form ω_Φ and a correction term $d\tau$

where τ is the canonical lift of a vector field w such that both w and $\text{curl } w$ extend to the zero vector field on D . Therefore

$$\begin{aligned} \int_{P_t} \omega \wedge \omega^\sharp &= \int_{P_t} \iota(\omega_\Phi + d\tau) \wedge (\omega_\Phi + d\tau)^\sharp \\ &= \int_{P_t} \omega_\Phi \wedge \omega_\Phi^\sharp \\ &\quad + 2 \int_{P_t} \iota d\tau \wedge \omega_\Phi^\sharp \\ &\quad + \int_{P_t} \iota d\tau \wedge d\tau^\sharp. \end{aligned}$$

By our previous remarks we know that

$$\lim_{t \rightarrow \infty} \int_{P_t} \omega_\Phi \wedge \omega_\Phi^\sharp = \int_H \|\omega_\Phi\|^2$$

and

$$\lim_{t \rightarrow \infty} \int_{P_t} \iota d\tau \wedge d\tau^\sharp = \int_H \|d\tau\|^2$$

so we need to show that

$$\lim_{t \rightarrow \infty} \int_{P_t} \iota d\tau \wedge \omega^\sharp = 0.$$

Recall that d decomposes into its real and imaginary parts, D and T , respectively. Then $\partial = D - T$ is the covariant derivative of the flat connection on E^* conjugated by the isomorphism between E and E^* . If α and β are E -valued forms with α a k -form, then

$$d(\alpha \wedge \beta^\sharp) = d\alpha \wedge \beta^\sharp + (-1)^k \alpha \wedge (\partial\beta)^\sharp.$$

(On the left-hand side of this equation $\alpha \wedge \beta^\sharp$ is a real valued form and d is the covariant derivative for real forms.) Therefore

$$\int_{P_t} \iota d\tau \wedge \omega_\Phi^\sharp = - \int_{P_t} \iota \tau \wedge (\partial\omega_\Phi)^\sharp + \int_{\partial P_t} \iota \tau \wedge \omega_\Phi^\sharp.$$

To calculate $\partial\omega_\Phi$, we recall that $d\omega_\Phi = 0$ and $\partial - d = -2T$ (see §2.2) so

$$\begin{aligned} \partial\omega_\Phi &= -2T\omega_\Phi \\ &= -\sigma^{-2} e^{-t} \text{sech } t\phi (TE_1 + \iota TE_2) \wedge (\omega^1 + \omega^2) \\ &= -\sigma^{-2} e^{-t} \text{sech } t\phi ((E_1 + \iota E_2)\omega^1 \wedge \omega^3 + (\iota E_1 - E_2)\omega^2 \wedge \omega^3). \end{aligned}$$

Since there are no $\omega^1 \wedge \omega^2$ terms in $\partial\omega_\Phi$,

$$\int_{P_t} \iota \tau \wedge (\partial\omega_\Phi)^\sharp = 0.$$

To finish the proof, we need to calculate the boundary term. To do so, we use the Poincare disk model for P with p the center of the disk. In particular we reset our coordinates identifying the unit disk Δ in \mathbb{C} with the hyperbolic plane P . The conformal factor of the hyperbolic metric on $\Delta = P$ is $\sigma(z) = \frac{1}{1-|z|^2}$. In these coordinates (5.2) becomes

$$\|\omega_\Phi(z)\|^2 = 4\|\Phi(z)\|^2 = 4\sigma(z)^{-4}|\phi(z)|^2$$

where ϕ is a holomorphic function on a neighborhood of Δ . In particular, $|\phi|$ is bounded on Δ .

We also know that τ is the canonical lift of a vector field w such that w and $\text{curl } w$ limit to zero on ∂P . Therefore the Euclidean length of these vector fields must decay to zero at ∂P ; i.e., $\sigma^{-1}\|w\|$ and $\sigma^{-1}\|\text{curl } w\|$ extend to the zero function on ∂P , where $\|w\|$ and $\|\text{curl } w\|$ are hyperbolic lengths. So although $\|\tau\|^2 = *(\tau \wedge *\tau^\sharp) = \|w\|^2 + \|\text{curl } w\|^2$ may not decay to zero, there is a function g on P that does extend continuously to zero on ∂P with $\|\tau\| = \sigma g$.

Putting this all together we have

$$\begin{aligned} \int_{\partial P_t} \iota\tau \wedge \omega_\Phi^\sharp &\leq \int_{\partial P_t} \|\tau\| \|\omega_\Phi\| dR \\ &\leq \int_{\partial P_t} 2\sigma^{-1}g|\phi| dR \\ &\leq \int_0^{2\pi} 2g|\phi| d\theta \end{aligned}$$

and since $g|\phi|$ limits to zero on ∂P , we have

$$\lim_{t \rightarrow \infty} \int_{\partial P_t} \iota\tau \wedge \omega_\Phi^\sharp = 0.$$

Therefore

$$\int_H \|\omega\|^2 = \int_H \|\omega_\Phi\|^2 + \int_H \|d\tau\|^2 \geq 2\|\Phi\|_2^2 \geq \frac{8\pi}{3}\|\Phi(z)\|$$

for all $z \in D$. □

With this local result in place we can now bound the norm of Φ_t^i . The projective structures Σ_t^i have a fixed conformal structure X^i . Let κ be the injectivity radius for the hyperbolic metric on X^i .

Theorem 5.5. *Assume that the tube radius of \mathcal{C} in M_α is greater than $\sinh^{-1} \sqrt{2}$ and that $L_{\mathcal{C}}(\alpha) \leq \ell_0$. Then*

$$(5.5) \quad L_{\mathcal{C}}(t) \geq 2\sqrt{\frac{2\pi}{3} \frac{\tanh^2(\kappa/2)}{1+2\|\Sigma_t\|_F}} \|\Phi_t^i\|_\infty.$$

Proof. By Proposition 3.7, for each $z \in \Sigma_t$ there exists an embedded round disk D with hyperbolic metric σ_D such that

$$\sigma_D(z) \leq \sigma(z) \coth(\kappa/2) \sqrt{1+2\|\Sigma_t\|_F}.$$

Therefore if we compare the norm of a holomorphic quadratic differential Φ in the σ_D -metric to the norm in the σ -metric, we get

$$\|\Phi(z)\|_{\sigma_D} \geq \frac{\tanh^2(\kappa/2)}{1+2\|\Sigma_t\|_F} \|\Phi(z)\|_\sigma.$$

Let H be the half space bounded by D on Σ_t . By Theorem 3.5, H is disjoint from $U_t^{\mathcal{C}}$ and therefore

$$L_{\mathcal{C}}(t)^2 \geq \int_{M_t} \|\omega_t\|^2 \geq \int_H \|\omega_t\|^2.$$

By Theorem 5.4

$$\int_H \|\omega_t\|^2 \geq \frac{8\pi}{3} \|\Phi_t(z)\|_{\sigma_D}^2.$$

Combing these three formulas, we have

$$\begin{aligned} L_C(t)^2 &\geq \int_H \|\omega_t\|^2 \\ &\geq \frac{8\pi}{3} \|\Phi_t(z)\|_{\sigma_D}^2 \\ &\geq \frac{8\pi}{3} \frac{\tanh^4(\kappa/2)}{(1+2\|\Sigma_t\|_F)^2} \|\Phi_t(z)\|_{\sigma}^2 \end{aligned}$$

which implies the result. \square

Next we use the bound on $\|\Phi_t^i\|_{\infty}$ to bound the distance between projective structures Σ_{α}^i and Σ_t^i . This is another of the main results mentioned in the introduction.

Theorem 1.3. *There exists a C depending only on α , the injectivity radius of the unique hyperbolic metric on X^i and $\|\Sigma_{\alpha}^i\|_F$ such that*

$$d(\Sigma_{\alpha}^i, \Sigma_t^i) \leq CL_C(\alpha)$$

for all $t \leq \alpha$.

Proof. We will integrate (5.5). Let $\sigma(T)$ be the length, in $P(X)$, of the path of projective structures Σ_t with $t \in [T, \alpha]$. Since Σ_t is a smooth path in $P(X)$, $\sigma(t)$ will be a smooth function and by definition $\|\Sigma_t\|_F \leq \|\Sigma_{\alpha}\|_F + \sigma(t)$. By Proposition 2.1, $-\frac{d\sigma}{dt}(t) = \|\Phi_t\|_{\infty}$. By Proposition 4.1, $L_C(\alpha) \geq L_C(t)$ for all $t \leq \alpha$. Then by (5.5)

$$-\frac{d\sigma}{1/2 + \|\Sigma_{\alpha}\|_F + \sigma} \leq KL_C(\alpha)dt$$

where $K = 2\sqrt{\frac{2\pi}{3}} \coth^2(\kappa/2)$. Integrating both sides, we have

$$\begin{aligned} \int_T^{\alpha} -\frac{d\sigma}{1/2 + \|\Sigma_{\alpha}\|_F + \sigma} &\leq \int_T^{\alpha} KL_C(\alpha)dt \\ \log\left(\frac{1/2 + \|\Sigma_{\alpha}\|_F + \sigma(T)}{1/2 + \|\Sigma_{\alpha}\|_F + \sigma(\alpha)}\right) &\leq (\alpha - T)KL_C(\alpha) \leq \alpha KL_C(\alpha) \\ 1 + \frac{\sigma(T)}{1/2 + \|\Sigma_{\alpha}\|_F} &\leq e^{\alpha KL_C(\alpha)} \\ \sigma(T) &\leq (1/2 + \|\Sigma_{\alpha}\|_F) \left(e^{\alpha KL_C(\alpha)} - 1\right). \end{aligned}$$

There exists a C' depending only on α and K (and hence κ) such that $e^{\alpha KL_C(\alpha)} - 1 \leq C' L_C(\alpha)$. Therefore

$$\sigma(t) \leq C' L_C(\alpha)(1/2 + \|\Sigma_{\alpha}\|_F).$$

Since $d(\Sigma_{\alpha}, \Sigma_t) \leq \sigma(t)$, we have

$$d(\Sigma_{\alpha}, \Sigma_t) \leq CL_C(\alpha)$$

where $C = C'(1/2 + \|\Sigma_{\alpha}\|_F)$ depends only on κ , α and $\|\Sigma_{\alpha}\|_F$. \square

Corollary 5.6. *The projective structures Σ_t^i converge to a projective structure $\Sigma_{\alpha'}^i$ as $t \rightarrow \alpha'$.*

6. GEOMETRIC LIMITS

We know that the projective boundary of M_t converges at α' . Now we need to show that the entire cone-manifold converges. We will need to examine geometric limits of hyperbolic cone-manifolds. Our approach will follow that of [HK4].

If X and Y are metric spaces and $f : X \rightarrow Y$ is a map, define $\text{lip}(f)$ to be the infimum of all K such that $d_Y(f(x_1), f(x_2)) \leq K d_X(x_1, x_2)$. The *bi-Lipschitz distance* between X and Y is $\text{bilip}(X, Y) = \inf\{|\log \text{lip}(f)| + |\log \text{lip}(f^{-1})| : f : X \rightarrow Y \text{ is bi-Lipschitz}\}$. This defines the *bi-Lipschitz topology* on the set of metric spaces.

To show that a sequence of compact hyperbolic manifolds with boundary converges, we need to control three quantities: the principal curvatures of the boundary, the injectivity radius and the width of collar neighborhoods of the boundary. If M is a hyperbolic manifold, we let $\text{inj}_M = \inf\{\text{inj}_x \mid x \in M\}$. We define $\partial M(t)$ to be those points in M whose distance from ∂M is less than t . Then $\text{width}(\partial M) = \inf\{t \mid \partial M(t) \text{ is a collar of } M \text{ for all } t' < t\}$. The geometric convergence theorem that follows is essentially due to Kodani [Kod1] (see the remarks on p. 20 of [HK4]):

Theorem 6.1. *Let λ^- , λ^+ , i_0 and W be real constants with $\lambda^+ \geq \lambda^-$ and $i_0, W > 0$. Let M_n be a sequence of hyperbolic manifolds with boundary such that the principal curvatures of ∂M_n are contained in the interval $[\lambda^-, \lambda^+]$, $\text{inj}_{M_n} \geq i_0$ and $\text{width}(\partial M_n) \geq W$. Then there exists a hyperbolic 3-manifold with boundary M_∞ and a subsequence $\{n_k\}$ such that $M_{n_k} \rightarrow M_\infty$ in the bi-Lipschitz topology. Furthermore if all the M_n are diffeomorphic and have bounded volume (or diameter), then M_∞ is diffeomorphic to M_n .*

We will apply this result to suitably chosen compact submanifolds of our hyperbolic cone-manifolds M_t .

6.1. Geometric limit of geometrically finite ends. To construct these compact submanifolds, we remove a neighborhood of the geometrically finite ends. To do this, we need to understand how the projective boundary determines the hyperbolic geometry of the geometrically finite ends. This information comes from work of Epstein ([Ep]) and Anderson ([And]) which we will review here.

We will use the same coordinates for \mathbb{H}^3 as we did in §5. In particular, let U be the upper half space of \mathbb{C} with hyperbolic metric σ . Then $\mathbb{H}^3 = U \times \mathbb{R}$ with metric

$$\sigma^2 \cosh^2 t dx^2 + \sigma^2 \sinh^2 t dy^2 + dt^2.$$

Let P_d be the set of points of the form (z, d) in \mathbb{H}^3 . Then P_0 is a hyperbolic plane and P_d is a constant curvature plane a (signed) distance d from P_0 .

Let $\psi : U \rightarrow \widehat{\mathbb{C}}$ be a conformal, locally univalent map and let $\Phi = S\psi$ be its Schwarzian derivative. The *osculating Möbius transformation* $M_{\psi(z)}$ is the unique Möbius transformation whose two jet agrees with ψ at z . We define $\Psi : \mathbb{H}^3 \rightarrow \mathbb{H}^3$ by

$$\Psi(z, d) = M_{\psi(z)}(z, d).$$

Note that Ψ extends continuously to ψ on U .

The following two results can be found in §3 of [And].

Proposition 6.2. *Let $p = (z, d)$ be a point in \mathbb{H}^3 . There exist an orthonormal basis $\{e_1, e_2, e_3\}$ for $T\mathbb{H}_p^3$ with e_1 and e_2 spanning the plane normal to P_d and an*

orthonormal basis for $T\mathbb{H}_{\Psi(p)}^3$ such that $d\Psi$ at p in these coordinates is

$$\begin{pmatrix} 1 + \frac{\|\Phi(z)\|}{4e^d \cosh d} & 0 & 0 \\ 0 & 1 - \frac{\|\Phi(z)\|}{4e^d \cosh d} & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

In particular if $4e^d \cosh d > \|\Phi\|_\infty$, then Ψ is an orientation preserving local diffeomorphism at p .

Remark. For each d the nearest point retraction π_d defines a natural map from U to P_d . In [And] what is actually calculated is the derivative of the composition $\Psi \circ \pi_d$. Since the derivative of π_d is easy to calculate, we can translate the work in [And] to the above proposition.¹

When Ψ is an immersion, there are also formulas for the curvature of the image surface $\Psi(P_d)$.

Proposition 6.3. *Let $p = (z, d)$ be a point on P_d and let $k_1 = -\frac{\|\Phi(z)\|}{\|\Phi(z)\|-1}$ and $k_2 = -\frac{\|\Phi(z)\|}{\|\Phi(z)\|+1}$. Then the principal curvatures (if defined) of $\Psi(P_d)$ at $\Psi(p)$ are*

$$\frac{\sinh d + k_i \cosh d}{\cosh d + k_i \sinh d}$$

for $i = 1, 2$.

Let Σ be a projective structure on a surface S with conformal structure X and let Σ_F be the fuchsian projective structure with conformal structure X . Then there is a representation $\rho_F : \pi_1(S) \rightarrow PSL_2\mathbb{R}$ such $\Sigma_F = U/\rho_F(\pi_1(S))$. Identifying U with the universal cover \tilde{S} and the deck transformations with $\rho_F(\pi_1(S))$, there is a conformal developing map $\psi : U \rightarrow \hat{\mathbb{C}}$ for Σ . In particular, Σ has a holonomy representation ρ and $\psi \circ \rho_F(\gamma) = \rho(\gamma) \circ \psi$ for all $\gamma \in \pi_1(S)$. As above ψ extends to a map $\Psi : \mathbb{H}^3 \rightarrow \mathbb{H}^3$. It is clear from the definition that this construction is natural. That is, $\Psi \circ \rho_F(\gamma) = \rho(\gamma) \circ \Psi$ for all $\gamma \in \pi_1(S)$.

The group $\rho_F(\pi_1(S))$ also acts on \mathbb{H}^3 with quotient homeomorphic to $S \times \mathbb{R}$. We can therefore view \mathbb{H}^3 as the universal cover of $S \times \mathbb{R}$, identifying $\tilde{S} \times \{d\}$ with P_d in \mathbb{H}^3 . Then Ψ is a map from $\tilde{S} \times \mathbb{R}$ to \mathbb{H}^3 . Restricted to $\tilde{S} \times [d, \infty)$ where $e^d > \sqrt{2\|\Sigma\|_F + 1}$, Ψ is a diffeomorphism. Let $\mathcal{E}(\Sigma, d)$ be the hyperbolic structure on $S \times [d, \infty)$ defined by this developing map. The hyperbolic structure $\mathcal{E}(\Sigma, d)$ extends to the projective structure Σ on $S \times \{\infty\}$ so $\mathcal{E}(\Sigma, d)$ is a geometrically finite end with projective boundary Σ . The plane P_t descends to surfaces $\mathcal{S}(\Sigma, t)$ that foliate $\mathcal{E}(\Sigma, d)$.

Proposition 6.4. *The surfaces $\mathcal{S}(\Sigma, t)$ are convex in $\mathcal{E}(\Sigma, d)$ and are strictly convex if $t > 0$.*

Proof. This is a direct consequence of Proposition 6.3. □

The foliation of $\mathcal{E}(\Sigma, d)$ by convex surfaces implies that $\mathcal{E}(\Sigma, d)$ embeds in a hyperbolic cone-manifold with projective boundary Σ . More precisely:

Proposition 6.5. *If Σ is a component of the projective boundary of a hyperbolic cone-manifold M , then $\mathcal{E}(\Sigma, d)$ embeds in M if $d > 0$.*

¹There is an error in the calculation of the eigenvalues on p. 35 of [And]. They should be $\frac{1}{2}(1 + 1/t) + \|Sf(0)\|/4t$ and $\frac{1}{2}(1 + 1/t) - \|Sf(0)\|/4t$ not $\frac{1}{2}(1 + 1/t) + \|Sf(0)\|$ and $\frac{1}{2}(1 + 1/t) - \|Sf(0)\|$.

Proof. The proof is the same as Lemma 3.3. □

Although $\mathcal{E}(\Sigma, d)$ embeds, it may intersect the tubes U_t^c . We need to show that d can be chosen large enough so that this does not happen. To do this, we will use an alternative construction of the surfaces $\mathcal{S}(\Sigma, d)$ as an envelope of horospheres.

For each $p \in \mathbb{H}^3$ the identification of the unit sphere in $T_p\mathbb{H}^3$ with the ideal boundary $\widehat{\mathbb{C}}$ of \mathbb{H}^3 determines a *visual measure* μ_p of $\widehat{\mathbb{C}}$. Given a conformal metric σ on a domain $\Omega \subseteq \widehat{\mathbb{C}}$ and a point $z \in \widehat{\mathbb{C}}$, the set of points in \mathbb{H}^3 whose visual measure equals σ at z is a horosphere $\mathcal{H}_{z,\sigma}$. This horosphere also has the following property: A plane P in \mathbb{H}^3 limits to a round circle in $\widehat{\mathbb{C}}$. This circle bounds two disks and we assume one of these disks D contains z . Then hyperbolic metric σ_D on D will agree with σ at z if and only if P is tangent to $\mathcal{H}_{z,\sigma}$.

The envelope of this family of horospheres $\mathcal{H}_{z,\sigma}$ is a surface in \mathbb{H}^3 . A similar construction works in a geometrically finite end. In fact if σ is the hyperbolic metric on Σ and $\sigma_d = e^d\sigma$, the envelope of the family of horospheres \mathcal{H}_{z,σ_d} is the surface $\mathcal{S}(\Sigma, d)$.

Theorem 6.6. *Let κ be the injectivity radius of the hyperbolic metric on Σ . If $e^d > \coth(\kappa/2)\sqrt{1 + 2\|\Sigma\|_F}$, then $\mathcal{E}(\Sigma, d)$ is embedded and disjoint from U_t^c .*

Proof. Since the envelope of the family of horospheres \mathcal{H}_{z,σ_d} is the surface $\mathcal{S}(\Sigma, d)$, the union of the horoballs bounded by \mathcal{H}_{z,σ_d} is the entire end $\mathcal{E}(\Sigma, d)$. By Proposition 3.7 there exists an embedded round disk D' with hyperbolic metric $\sigma_{D'}$ such that

$$\sigma_{D'}(z) < \sigma(z) \coth(\kappa/2)\sqrt{1 + 2\|\Sigma\|_F}.$$

By our choice of d there exists a round disk $D \subset D'$ such that $\sigma_D(z) = \sigma_d(z)$. The round disk D is the projective boundary of an embedded half space H which will have hyperbolic boundary a plane P . The plane P will be tangent to \mathcal{H}_{z,σ_d} so the horoball bounded by \mathcal{H}_{z,σ_d} will be contained in H . By Theorem 3.5, H is disjoint from U_t^c and therefore the horoball bounded by \mathcal{H}_{z,σ_d} is disjoint from U_t^c , proving the theorem. □

For each Σ the map Ψ canonically identifies $\mathcal{E}(\Sigma, d)$ as a Riemannian metric on a fixed copy of $S \times [d, \infty)$. We then have the following proposition:

Proposition 6.7. *Let Σ_t be a sequence of projective structures which converge to Σ_∞ in $P(X)$ and assume that $e^d > \sqrt{2\|\Sigma_t\|_F + 1}$ for all t . Then the metrics $\mathcal{E}(\Sigma_t, d)$ converge to $\mathcal{E}(\Sigma, \infty)$ in the compact- C^∞ topology on metrics on $S \times [d, \infty)$.*

Proof. This is a simple consequence of the fact that the maps $\tilde{\Psi}_t$ depend continuously on the projective structures Σ_t . In particular, if the Σ_t converge to Σ_∞ , then $\tilde{\Psi}_t$ can be chosen to converge to $\tilde{\Psi}_\infty$ in the C^∞ -topology on maps from $\tilde{S} \times [d, \infty)$ to \mathbb{H}^3 which implies that the metrics converge. □

On a compact manifold convergence in the compact- C^∞ topology of a sequence of metrics implies that the associated metric spaces converge in the bi-Lipschitz topology. Since $S \times [d, \infty)$ is non-compact, we only get bi-Lipschitz convergence on compact submanifolds such as collars $S \times [d, d']$. This will be enough for our applications.

6.2. The Schläfli formula. We will need to bound the volume of the complement of the geometrically finite ends in the cone-manifolds M_t . To do so, we will use the generalized Schläfli formula of Rivin and Schlenker. Although their formula applies in much greater generality, we will stick to the case of a 3-manifold with boundary M and a smooth family of hyperbolic cone-metrics g_t where t is the cone-angle. Let $I(t)$ and $II(t)$ be the first and second fundamental forms for ∂M in the g_t metric and $H(t)$ the mean curvature. Finally let $V(t)$ be the volume of M in the g_t metric. The generalized Schläfli formula is then

Theorem 6.8 (Rivin and Schlenker [RS]).

$$-\frac{1}{3}V'(t) = \int_{\partial M} \left(H'(t) + \frac{1}{2}\langle I'(t), II(t) \rangle \right) dA + L_C(t).$$

We will use the following simple corollary of this result:

Corollary 6.9. *Assume that the metrics g_t are defined for $a < t \leq b$ and that as $t \rightarrow a$, g_t converges in the C^∞ -topology on a collar neighborhood of ∂M . Furthermore assume that $L_C(t)$ is bounded. Then $V(t)$ is bounded.*

Proof. Since g_t converges, the quantities $H'(t)$ and $I'(t)$ and $II(t)$ are all bounded and therefore the integral

$$\int_{\partial M} (H' + \langle I', II \rangle) dA$$

will be bounded. Since L_C is bounded, Theorem 6.8 implies that V' is bounded which in turn implies that

$$V(T) = \int_1^T V' dt + V(1)$$

is bounded. □

Remark. This result could also have been proven using the standard Schläfli formula for manifolds with polyhedral boundary. One simply needs to construct a polyhedral approximation for the smooth boundary.

6.3. Geometric limits of cone-manifolds. By Corollary 5.6 we know that the projective structures Σ_t^i converge to a projective structure $\Sigma_{\alpha'}^i$, as $t \rightarrow \alpha'$. Therefore there exists a $d > 0$ such that $e^{d-1} > \sqrt{2\|\Sigma_t^i\|_F + 1}$ for all $t \in (\alpha', \alpha]$ and $i = 1, \dots, n$. We then let \mathcal{M}_t be the closure of $M'_t \setminus \left(\bigcup_{i=1}^n \mathcal{E}(\Sigma_t^i, d) \right)$. The boundary of \mathcal{M}_t consist of the boundary ∂U_t^c of the tubes and the boundary $\mathcal{S}(\Sigma_t^i, d)$ of the geometrically finite ends.

Theorem 6.10. *There exists a sequence $\{t_n\}$ in $(\alpha', \alpha]$ with $t_n \rightarrow \alpha'$ such that \mathcal{M}_{t_n} converges in the bi-Lipschitz topology to a hyperbolic cone-manifold $\mathcal{M}_{\alpha'}$ homeomorphic to \mathcal{M}_t .*

Proof. We need to see that the conditions of Theorem 6.1 hold for the family \mathcal{M}_t . As we have already noted, the norms $\|\Sigma_t^i\|_F$ are bounded since the projective structures converge. Proposition 6.3 then implies that the principal curvatures of $\mathcal{S}(\Sigma_t^i, d)$ are bounded above and below. Since the radii of the tubes U_t^c are greater than $\sinh^{-1} \sqrt{2}$, the principal curvatures of ∂U_t^c are bounded between $\coth(\sinh^{-1} \sqrt{2})$ and $\tanh(\sinh^{-1} \sqrt{2})$. We bounded the injectivity radius in Proposition 4.4. Our choice of d guarantees that $\text{width}(\mathcal{S}(\Sigma_t^i, d)) \geq 1$. To see that the

boundary components ∂U_t^c have a definite width, we note that the tubes U_t^c are not maximal. In fact from Proposition 3.2 we see that there are disjoint embedded tubes whose torus boundary has area 0.51. The difference between these tubes and the U_t^c will be a collar of definite width for all t .

This is enough to obtain a limiting hyperbolic manifold $\mathcal{M}_{\alpha'}$. To see that $\mathcal{M}_{\alpha'}$ is homeomorphic to the \mathcal{M}_t , we need to bound the volume of the \mathcal{M}_t . If we consider the \mathcal{M}_t as a family of metrics on a fixed manifold, Proposition 6.7 implies that we can choose these metrics such that they converge on a neighborhood of each $\mathcal{S}(\Sigma_t^i, d)$. We then apply Corollary 6.9 to $M_t \setminus \left(\bigcup_{i=1}^n \mathcal{E}(\Sigma_t^i, d) \right)$ to bound the volume. □

Theorem 6.11. *There exists an $M_{\alpha'} \in GF(N, \mathcal{C})$ such that $M_t \rightarrow M_{\alpha'}$.*

Proof. The boundary of $\mathcal{M}_{\alpha'}$ consists of tori and higher genus surfaces. On a collar of the higher genus ends, the manifolds \mathcal{M}_{t_n} converge to a collar of the geometrically finite ends $\mathcal{E}(\Sigma_{\alpha'}^i, d-1)$. Therefore we can glue the ends $\mathcal{E}(\Sigma_{\alpha'}^i, d-1)$ to the higher genus boundary components. It is shown in §3 of [HK4] that the metric can be extended to a cone-singularity (or cusp if $\alpha' = 0$) at the torus components of $\partial \mathcal{M}_{\alpha'}$ with cone-angle α' . In fact they show more than this. For large n there are bi-Lipschitz diffeomorphisms f_n from \mathcal{M}_{t_n} to $\mathcal{M}_{\alpha'}$. For a fixed $c \in \mathcal{C}$, f_n maps $\partial U_{t_n}^c$ to a fixed component $\partial U_{\alpha'}^c$ of $\partial \mathcal{M}_{\alpha'}$. The *meridian* of $\partial U_{t_n}^c$ is the unique homotopy class a of non-trivial simple closed curve on $\partial U_{t_n}^c$ that bounds a disk in $U_{t_n}^c$. Hodgson and Kerckhoff further show that f_n maps the meridians to a fixed homotopy class on $\partial U_{\alpha'}^c$, and this homotopy class is a meridian of the cone-singularity in the extended structure.

The extended manifold $M_{\alpha'}$ is then a geometrically finite cone-metric on a pair $(\hat{N}, \hat{\mathcal{C}})$. Since the maps f_n take meridians to meridians, the f_n extend to homeomorphisms from (N, \mathcal{C}) to $(\hat{N}, \hat{\mathcal{C}})$. These extensions of f_n can also be chosen to be conformal maps from the conformal boundaries of M_{t_n} to $M_{\alpha'}$. By Theorem 1.1 in an open interval about α' there exists a one-parameter family of cone-manifolds \hat{M}_t with cone-angle t , conformal boundary X and $\hat{M}_{\alpha'} = M_{\alpha'}$. We need to show that each f_n is homotopic to an isometry from M_{t_n} to \hat{M}_{t_n} .

Let ρ_t and $\hat{\rho}_t$ be the holonomy representations of M_t and \hat{M}_t , respectively. (Note that they are representations of $\pi_1(N \setminus \mathcal{C}) \cong \pi_1(\hat{N} \setminus \hat{\mathcal{C}})$ not of $\pi_1(N) \cong \pi_1(\hat{N})$.) Convergence in the bi-Lipschitz topology implies that the representations $(f_n)_* \rho_{t_n}$ converge to $\hat{\rho}_{\alpha'}$. By Theorem 5.7 in [Br1] the space of conjugacy classes of representations is locally parameterized by the complex length of the meridians and the conformal boundary. This is a stronger version of Theorem 1.1 which allows representations where the holonomy of the meridians is not elliptic. It implies that $\hat{\rho}_{t_n} = (f_n)_* \rho_{t_n}$ for large n . By Theorem 1.7.1 of [CEG] on \mathcal{M}_{t_n} the maps f_n will be homotopic to an isometric embedding of \mathcal{M}_{t_n} in \hat{M}_{t_n} . Since \mathcal{M}_{t_n} extends to a geometrically finite cone-manifold, in a unique way this implies that f_n is homotopic to an isometry from all of M_{t_n} to \hat{M}_{t_n} .

To finish the proof, we choose a fixed large value of n and use the map f_n to pull back metrics in $GF(\hat{N}, \hat{\mathcal{C}})$ to metrics in $GF(N, \mathcal{C})$. Under this identification $M_{t_n} = \hat{M}_{t_n}$. Theorem 1.1 then implies that $M_t = \hat{M}_t$ wherever both are defined. Therefore $M_t \rightarrow M_{\alpha'}$ as desired. □

We are now ready to prove Theorem 1.2.

Theorem 1.2. *Let $M_\alpha \in GF(N, \mathcal{C})$ be a geometrically finite hyperbolic cone-metric with cone-angle α . Suppose the tube radius of the cone-singularity is $\geq \sinh^{-1} \sqrt{2}$. Then there exists an ℓ_0 depending only on α such that if $L_c(\alpha) \leq \ell_0$ for all $c \in \mathcal{C}$, then the one-parameter family of cone-manifolds $M_t \in GF(N, \mathcal{C})$ is defined for all $t \leq \alpha$.*

Proof. By Theorem 1.1 the interval for which the family M_t is defined is open in $[0, \alpha]$. By Theorem 6.11 this interval is also closed. Therefore M_t is defined for all $t \in [0, \alpha]$. \square

7. RANK TWO CUSPS

In this section we show how bounds on the L^2 -norm control the shape of a rank two cusp. Recall that a rank two cusp is the quotient of a horoball, centered at infinity in the upper half space model, by parabolic elements $z \mapsto z+1$ and $z \mapsto z+\tau$ with $\text{Im } \tau > 0$. The horoball is foliated by horospheres which are horizontal planes in \mathbb{H}^3 . The quotients of these planes are tori which foliate the rank two cusp. Each tori will be conformally equivalent with τ , the Teichmüller parameter of the tori. To normalize the cusp, we choose the horoball so that in the quotient the boundary torus has area $\frac{1}{2}$.

We will use similar notation for cusps as we do for short geodesics and their tubular neighborhoods. In particular if γ is a torus component of ∂N , then U_t^γ will be the associated rank two cusp. We also let $\mathcal{L}_\gamma(t)$ be the Teichmüller parameter of the cusp.

We do not know, a priori, that the cusps U_t^γ are embedded. The proof of this is essentially the same as Theorem 3.5. In particular we have

Proposition 7.1. *If the tube radius of the cone-singularity is greater than $\sinh^{-1} \sqrt{2}$ and if $L_c(t) \leq \ell_0$ for all $c \in \mathcal{C}$, then the cusps U_t^γ are embedded and disjoint from the tubes U_t^c for all $c \in \mathcal{C}$.*

Next we need to control the derivative of the Teichmüller parameter as t varies. Note that $\mathcal{L}_\gamma(t)$ is a point in the Teichmüller space of a torus which is canonically identified with \mathbb{H}^2 so we will measure the derivative $\mathcal{L}'_\gamma(t)$ in the hyperbolic metric. We then have

Theorem 7.2.

$$\mathcal{L}_c(\alpha)^2 \geq |\mathcal{L}'_\gamma(t)|^2.$$

Proof. The proof is similar to the proof Theorem 5.4.

On the rank two cusp the Hodge form ω_t is a sum of a model deformation ω_t^m and a correction term ω_t^c . The model is constructed in §3.7 of [Br1]. The term ω_t^c is trivial in cohomology and therefore exact, so $\omega_t^c = d\psi_t$ where ψ_t is an E -valued section on U_t^γ .

Recall that to calculate the L^2 -norm of ω_t on U_γ^t we need to integrate $\imath\omega_t \wedge \omega_t^\sharp$ over the boundary torus ∂U_γ^t . As in the proof of Theorem 5.4 we expand $\imath\omega_t \wedge \omega_t^\sharp$ to get

$$\imath\omega_t \wedge \omega_t^\sharp = \imath\omega_t^m \wedge (\omega_t^m)^\sharp + 2\imath\omega_t^c \wedge (\omega_t^m)^\sharp + \imath\omega_t^c \wedge (\omega_t^c)^\sharp.$$

Using the integration by parts argument from Theorem 5.4, we have

$$\int_{\partial U_\gamma^t} \omega_t^c \wedge (\omega_t^m)^\sharp = - \int_{\partial U_\gamma^t} \psi_t \wedge (\partial \omega_t^m)^\sharp.$$

From the explicit form of ω_t^m we see that $\partial \omega_t^m$ has no terms tangent to ∂U_t^γ so

$$\int_{\partial U_t^\gamma} \psi_t \wedge (\partial \omega_t^m)^\sharp = 0$$

which implies that

$$\int_{U_\gamma^t} \|\omega_t\|^2 = \int_{U_\gamma^t} \|\omega_t^m\|^2 + \int_{U_\gamma^t} \|\omega_t^c\|^2.$$

We can also calculate the L^2 -norm

$$\int_{U_\gamma^t} \|\omega_t^m\|^2 = |\mathcal{L}'_\gamma(t)|^2.$$

By Propositions 4.2 and 7.1 we know that

$$\mathcal{L}_C(\alpha)^2 \geq \int_{M_t} \|\omega_t\|^2 \geq \int_{U_t^\gamma} \|\omega_t\|^2.$$

Combining this last inequality with the previous two equalities completes the proof. □

As an immediate corollary we have

Corollary 7.3. *The length of the path $\mathcal{L}_\gamma(t)$ with $t \in (\alpha', \alpha]$ is bounded and therefore $\mathcal{L}_\gamma(t)$ converges to $\mathcal{L}_\gamma(\alpha')$ as $t \rightarrow \alpha'$, where $\mathcal{L}_\gamma(\alpha')$ is a complex number with $\text{Im } \mathcal{L}_\gamma(\alpha') > 0$.*

The proofs of Theorems 1.3 and 1.4 do not change when rank two cusps are allowed. To prove Theorem 1.2, we need to modify the geometric limit argument. When there are rank two cusps, we define \mathcal{M}_t to be the complement of geometrically finite ends, the tubes about the cone-singularity and the rank two cusps. Then just as in Proposition 4.4, Corollary 7.3 implies that the injectivity radius of \mathcal{M}_t is bounded below. Once we have this, the geometric limit argument and the proof of Theorem 1.2 follow as before.

8. APPLICATIONS TO KLEINIAN GROUPS

As mentioned in the introduction, estimates similar to the ones in this paper were first proven by McMullen as part of his proof of the density of cusps on the boundary of a Bers slice [Mc]. His methods were entirely different than ours and in this section we will compare the two results and show how our estimates can give a new approach to the density of cusps. Later Canary, Culler, Hersonsky and Shalen extended McMullen's estimates to prove the density of cusps in a more general setting. We will also discuss this generalization.

A *Bers slice*, $T(X)$, is an embedding of the Teichmüller space $T(S)$ in the space of projective structures, $P(X)$, on a fixed conformal structure X . In particular for each $Y \in T(S)$ there is a unique projective structure $\Sigma_Y \in P(X)$ such that the developing map for Σ_Y is a homeomorphism onto a Jordan domain in $\widehat{\mathbb{C}}$. The holonomy for Σ_Y acts as deck transformations on the complementary Jordan domain uniformizing Y . Most interestingly, $T(X)$ is a bounded domain in $P(X)$ so

its closure $U(X)$ is a compactification of Teichmüller space. Bers began a systematic study of this compactification in [Bers]. In this paper he made a number of conjectures, one of which is the density of cusps which we are now discussing.

Projective structures on the boundary $\partial T(X) = U(X) \setminus T(X)$ will have discrete and faithful holonomy. Hence the image of the holonomy is a Kleinian group whose quotient hyperbolic 3-manifold will be homotopy equivalent to S . If the image of the holonomy contains a parabolic, then the projective structure is a *cuspidal*. The following theorem was conjectured by Bers:

Theorem 8.1 (McMullen [Mc]). *Cusps are dense in $\partial T(X)$.*

The parabolics in a cusp correspond to a collection \mathcal{C} of homotopically distinct, disjoint and essential simple closed curves on S . For each such \mathcal{C} there is a projective structure for whose holonomy each curve in \mathcal{C} is parabolic, and if \mathcal{C} is maximal, there is a unique such projective structure $\Sigma_{\mathcal{C}}$. In this case $\Sigma_{\mathcal{C}}$ is a *maximal cusp*.

The first step in McMullen’s proof of the density of cusps is to find a dense set of $\Sigma \in \partial T(X)$ with the following property: There exists a sequence of $Y_i \in T(S)$ and maximal curve collections \mathcal{C}_i such that $L_{\mathcal{C}_i}(Y_i) \rightarrow 0$ and $\Sigma_{Y_i} \rightarrow \Sigma$ as $i \rightarrow \infty$. The construction of the Y_i is a fairly straightforward argument involving the compactness of k -quasiconformal maps and Sullivan rigidity (see p. 221 of [Mc]). The proof is then completed with the following estimate which is a combination of Theorem 1.2 and Corollary 1.3 in [Mc]:

Theorem 8.2 (McMullen). *Given a Bers slice $T(X)$, there exists a $k > 0$ such that if \mathcal{C} is a maximal collection of simple closed curves on S and $Y \in T(S)$ with $L_{\mathcal{C}}(Y) \leq 1/2$, then the following hold:*

- (1) *There exists a $Y' \in T(S)$ with $L_{\mathcal{C}}(Y') \leq L_{\mathcal{C}}(Y)/2$ and*

$$d(\Sigma_Y, \Sigma_{Y'}) \leq k(L_{\mathcal{C}}(Y) \log(1/L_{\mathcal{C}}(Y)))^2.$$

- (2) *We can bound the distance between Σ_Y and the maximal cusp $\Sigma_{\mathcal{C}}$:*

$$d(\Sigma_Y, \Sigma_{\mathcal{C}}) \leq k(L_{\mathcal{C}}(Y) \log(1/L_{\mathcal{C}}(Y)))^2.$$

Note that (2) follows from (1) since for any sequence $Z_i \in T(S)$ with $L_{\mathcal{C}}(Z_i) \rightarrow 0$ we have $\Sigma_{Z_i} \rightarrow \Sigma_{\mathcal{C}}$. Applying (2) to each element of the sequence Y_i , we have $d(\Sigma_{Y_i}, \Sigma_{\mathcal{C}_i}) \leq k(L_{\mathcal{C}_i}(Y_i) \log(1/L_{\mathcal{C}_i}(Y_i)))^2$. Therefore $d(\Sigma_{Y_i}, \Sigma_{\mathcal{C}_i}) \rightarrow 0$ and $\Sigma_{\mathcal{C}_i} \rightarrow \Sigma_{\mathcal{C}}$ as $i \rightarrow \infty$. Since a dense set of $\Sigma \in \partial T(X)$ is approximated by maximal cusps, maximal cusps are dense in $\partial T(X)$. This is, of course, stronger than Bers’ conjecture that cusps are dense.

Theorem 8.2 bounds the effect of a quasiconformal deformation of a quasifuchsian manifold. One obtains a bound on the distance between the quasifuchsian manifold and the cusp by taking a limit of quasiconformal deformations. From Theorem 1.3 we can obtain similar estimates using hyperbolic cone-manifolds to interpolate between manifolds.

Before describing these estimates, we need to reset our notation. Let M be a complete, smooth hyperbolic 3-manifold and \mathcal{C} a collection of disjoint simple closed curves on the conformal boundary ∂M of M . Following [CCHS], the collection \mathcal{C} is *pinchable* if the curves in \mathcal{C} are homotopically distinct in M and each $c \in \mathcal{C}$ represents a non-trivial primitive element of $\pi_1(M)$. Then $L_c(M)$ will be the length of the geodesic representative of c in M and $L_c(\partial M)$ will be the length of the geodesic representative of c for the hyperbolic metric on the conformal boundary.

We need the following preliminary result which is a combination of theorems of Canary (Theorem 5.1 in [Can]) and Otal (Theorem 3 in [Ot]).

Theorem 8.3. *Let \mathcal{C} be a collection of pinchable curves on ∂M . There exists an $\epsilon_0 > 0$ such that if $L_c(\partial M) \leq \epsilon_0$ for each $c \in \mathcal{C}$, then the geodesic representative c^* of c in M is isotopic to c on ∂M . Furthermore this isotopy is disjoint from the geodesic representatives of the other curves in \mathcal{C} .*

With this result in hand we can now apply Theorem 1.3 to a manifold with short curves on its boundary. The next result is our version of McMullen's estimate (Theorem 8.2).

Theorem 8.4. *Assume M is a smooth geometrically finite hyperbolic 3-manifold without rank one cusps and that \mathcal{C} is a pinchable collection of curves on ∂M . There exists an $\ell'_0 > 0$ such that if $L_c(\partial M) \leq \ell'_0$ for all $c \in \mathcal{C}$, the following hold:*

- (1) *There exists a smooth, geometrically finite hyperbolic structure \hat{M} homeomorphic to M with each curve c pinched to a rank one cusp.*
- (2) *The components of the conformal boundaries of M and \hat{M} that are disjoint from \mathcal{C} are conformally equivalent.*
- (3) *If X is a component of the conformal boundary disjoint from \mathcal{C} and Σ and $\hat{\Sigma}$ are the projective boundaries on X for M and \hat{M} , respectively, then there exists a C depending only on the injectivity radius of the hyperbolic metric on X and $\|\Sigma\|_F$ such that*

$$d(\Sigma, \hat{\Sigma}) \leq CL_{\mathcal{C}}(\partial M).$$

Furthermore if X is incompressible, then C depends only on the injectivity radius of X and not on $\|\Sigma\|_F$.

Proof. Let $\ell'_0 = \min\{\ell_0, \epsilon_0\}$. By Theorem 1.2 there exists a one-parameter family of cone-manifolds M_t for $t \in [0, 2\pi]$ with cone-singularity \mathcal{C}^* such that $M_{2\pi} = M$. By Theorem 8.3, the manifold M has a compact core M' which is disjoint from \mathcal{C}^* . Furthermore each c^* will be isotopic to a curve on the boundary of M' . Recall that M_0 is a complete hyperbolic structure on the topological manifold $M \setminus \mathcal{C}$. Then the cover \hat{M} of M_0 associated to the compact submanifold M' is the hyperbolic manifold satisfying (1) and (2).

Finally, the first statement in (3) holds by Theorem 1.3. For the second statement we note that if X is incompressible, then $\|\Sigma\|_F \leq 3/2$ by Nehari's Theorem and therefore C only depends on the injectivity radius of X . \square

If M is quasifuchsian, then (3) of Theorem 8.4 is essentially the same as (2) of Theorem 8.2 although the bound in Theorem 8.4 is weaker.

We also remark that the only new part of Theorem 8.4 is (3). The existence of the manifold \hat{M} given by (1) and (2) is a well-known consequence of Thurston's hyperbolization theorem. In particular, if every component of the conformal boundary contains a curve in \mathcal{C} , then (3) is vacuous and Theorem 8.4 gives no new information. Furthermore, even if the conformal boundary contains a component X disjoint from \mathcal{C} , unlike in the quasifuchsian case, the projective structure on X may not determine the hyperbolic structures on M and \hat{M} . For these reasons it is often more useful to bound the change in complex length of closed geodesics as short curves are pinched on the conformal boundary. This is exactly what is done in the next theorem. Note that $\Theta_\gamma(M)$ is the imaginary part of the complex length.

Theorem 8.5. *For each $L > 0$ there exists an $\epsilon > 0$ and an $A > 0$ such that if γ is a closed geodesic in M with $L_\gamma(M) \leq L$ and $L_c(\partial M) \leq \epsilon$ for all $c \in \mathcal{C}$, then*

$$e^{-AL_c(\partial M)} L_\gamma(\hat{M}) \leq L_\gamma(M) \leq e^{AL_c(\partial M)} L_\gamma(\hat{M})$$

and

$$(1 - AL_c(M))\Theta_\gamma(\hat{M}) \leq \Theta_\gamma(M) \leq (1 + AL_c(M))\Theta_\gamma(\hat{M}).$$

Proof. This follows immediately using the construction of \hat{M} in the previous result and Theorem 1.4. □

This estimate should be compared to Proposition 5.1 of [CCHS] and Theorem 9.1 of [CH]. These results are essentially the same as Theorem 8.5 although they require both an upper and lower bound on the length of γ . Their proof is a generalization of McMullen’s methods.

The motivation for these results in [CCHS] and [CH] was to generalize the density of cusps. A Bers slice is the simplest example of a quasiconformal deformation space. For an arbitrary quasiconformal deformation space maximal cusps will not be dense on the boundary. The best one can hope for is for geometrically finite manifolds to be dense. This is shown in [CH]. Although this result of Canary and Hersensky follows from Theorem 8.5, in the next theorem we will restrict ourselves to the class of quasiconformal deformation spaces where maximal cusps are dense.

Theorem 8.6 ([CCHS]). *Let M be a geometrically finite hyperbolic 3-manifold without rank one cusps and assume that the conformal boundary of M has exactly one connected component. Then maximal cusps are dense on the boundary of the space of quasiconformal deformations of M .*

We remark that although a quasifuchsian manifold’s conformal boundary has two components, a Bers slice behaves like the quasiconformal deformation space of a manifold with one boundary component. This is because the quasiconformal deformations in a Bers slice are supported on only one of the boundary components while the other component is fixed.

The proof Theorem 8.6 follows the same outline as the density of cusps on the boundary of a Bers slice. In the final step we replace the main estimate of [CCHS] with Theorem 8.5. Let $MP(M)$ be the space of quasiconformal deformations of M . ($MP(M)$ is the space of *minimally parabolic* and geometrically finite manifolds homeomorphic to M .) Let \hat{M} be a hyperbolic manifold in $\partial MP(M)$ whose limit set is $\hat{\mathcal{C}}$. When the conformal boundary of M has exactly one component, such \hat{M} are dense in $\partial MP(M)$. (The proof of this fact is more complicated than it is for a Bers slice.) We can then find a sequence \mathcal{C}_i of maximal collections of pinchable curves on ∂M and $M_i \in MP(M)$ such that $M_i \rightarrow \hat{M}$ and $L_{\mathcal{C}_i}(\partial M_i) \rightarrow 0$. (Again, the construction of the manifolds M_i is more involved than it is for a Bers slice. There are topological considerations in finding the pinchable curves \mathcal{C}_i that do not occur for quasifuchsian manifolds.) Let $M_{\mathcal{C}_i}$ by the hyperbolic manifolds in $\partial MP(M)$ where the curves \mathcal{C}_i have become rank one cusps. Since $M_i \rightarrow \hat{M}$ for every simple closed curve γ we have $\mathcal{L}_\gamma(M_i) \rightarrow \mathcal{L}_\gamma(\hat{M})$. Theorem 8.5 then implies that $\mathcal{L}_\gamma(M_{\mathcal{C}_i}) \rightarrow \mathcal{L}_\gamma(\hat{M})$ from which it follows that $M_{\mathcal{C}_i} \rightarrow \hat{M}$. Once again we have shown that a dense set of points in $\partial MP(M)$ is limits of cusps and therefore cusps are dense in $\partial MP(M)$.

9. APPENDIX: MEAN VALUE INEQUALITIES

In this appendix we prove mean value inequalities for harmonic vector and strain fields. For strain fields this inequality was proved by Hodgson and Kerckhoff in an early version of [HK4]. Their proof is not in the current version of [HK4] as they have found a simpler proof of their main results which does not require the inequality. With their permission we recount the result here.

We also prove a mean value inequality for vector fields v where $\|\Delta v\|$ is bounded. The proof is essentially the same as for strain fields, if not simpler. We will start with the vector field inequality along with an application to Hodge forms on geometrically finite ends.

Using the identification of the tangent bundle with the real part of E , a vector valued k -form ω can be identified as a real E -valued k -form. Then $\Delta\omega$ will be the Laplacian associated to the bundle E . For a function u , Δu will be the standard Laplacian on functions. It will be clear from the context which Laplacian we are using.

Lemma 9.1. *If $\|\Delta v\| \leq b$, then $-(\Delta\|v\| + 2\|v\|) \geq -b$.*

Proof. For vector fields we have the Weitzenböck formula,

$$\Delta v = \nabla^* \nabla v + 2v$$

where ∇ is the Riemannian connection and ∇^* its adjoint (see §2 of [HK1]).

Let

$$\nabla_{XY}^2 = \nabla_X \nabla_Y - \nabla_{\nabla_X Y}.$$

Then $\nabla^* \nabla = -\sum_i \nabla_{e_i}^2$, where $\{e_1, e_2, e_3\}$ is an orthonormal frame field on M . For functions, $\Delta = \nabla^* \nabla$. Using this formula, we see that for any tensor S on a Riemannian manifold

$$\Delta\|S\|^2 = 2\langle \nabla^* \nabla S, S \rangle - 2\|\nabla S\|^2.$$

Combining this formula with the Weitzenböck formula for v , we have

$$(9.1) \quad \Delta\|v\|^2 = 2\langle \Delta v, v \rangle - 4\|v\|^2 - 2\|\nabla v\|^2.$$

Let $u = \|v\|^2$. Applying the product formula for the Laplacian to u times itself gives

$$\Delta(u^2) = 2u\Delta u - 2\|\nabla u\|^2.$$

Combining this formula with (9.1), we get

$$2u\Delta u - 2\|\nabla u\|^2 = 2\langle \Delta v, v \rangle - 4u^2 - 2\|\nabla v\|^2$$

or

$$u\Delta u + 2u^2 = \langle \Delta v, v \rangle + \|\nabla u\|^2 - \|\nabla v\|^2.$$

We also know that $\|\nabla v\| \geq \|\nabla u\| = \|\nabla(\|v\|)\|$ and $-\langle \Delta v, v \rangle \geq -\|\Delta v\|\|v\| \geq -bu$. Therefore

$$-u(\Delta u + 2u) \geq -bu.$$

If $u \neq 0$, we have

$$-(\Delta u + 2u) \geq -b.$$

On the other hand if $u = 0$, u has a local minimum so $-\Delta u \geq 0$ and the inequality still holds. \square

Define the operator L by $L = -(\Delta + 2)$. Let B_r be a ball of radius r centered at a point p . We first need a fundamental solution for L , that is, a radially symmetric, smooth function $v(r)$ on $B_R \setminus p$ such that $Lv = 0$ and

$$\int_{B_R} vL\phi dV = \lim_{\epsilon \rightarrow 0} \int_{B_R \setminus B_\epsilon} vL\phi dV = \phi(p)$$

for all smooth functions ϕ with support in the interior of B_R .

Lemma 9.2. *The function*

$$v(r) = \frac{-\cosh(\sqrt{3}r) + \coth(\sqrt{3}R) \sinh(\sqrt{3}r)}{4\pi \sinh r}$$

is a fundamental solution for L . Furthermore, $v(R) = 0$, $v(r) \leq 0$ for $0 < r < R$ and

$$v'(R) = \frac{\sqrt{3}}{4\pi \sinh(R) \sinh(\sqrt{3}R)}.$$

Proof. For any radial function

$$-\Delta f = \frac{\partial^2 f}{dr^2} + 2 \coth r \frac{df}{dr}.$$

Using this formula, it is easy to check that $Lv = 0$.

The operator L is self-adjoint so on any compact manifold M with boundary ∂M , L satisfies Green's identity:

$$\int_M fLg dV = \int_M gLf dV + \int_{\partial M} \left(f \frac{\partial g}{\partial n} - g \frac{\partial f}{\partial n} \right) dA$$

where f and g are smooth functions on M and $\frac{\partial}{\partial n}$ is the derivative in the direction of the outward normal.

Applying Green's identity to v and a test function ϕ , we have

$$\begin{aligned} \int_{B_R \setminus B_\epsilon} vL\phi dV &= \int_{B_R \setminus B_\epsilon} \phi Lv dV + \int_{\partial B_R} \left(v \frac{\partial \phi}{\partial n} - \phi \frac{\partial v}{\partial r} \right) dA \\ &\quad - \int_{\partial B_\epsilon} \left(v \frac{\partial \phi}{\partial n} - \phi \frac{\partial v}{\partial r} \right) dA \\ &= \int_{\partial B_\epsilon} v' \phi dA - \int_{\partial B_\epsilon} v \frac{\partial \phi}{\partial n} dA. \end{aligned}$$

Clearly, $v(\epsilon) \sim -\frac{1}{4\pi \sinh R}$ for ϵ near 0 and it is easy to check that $v'(\epsilon) \sim \frac{1}{4\pi \sinh^2 r}$. We also know that $\text{area}(\partial B_\epsilon) = 4\pi \sinh^2(\epsilon)$ and therefore

$$\lim_{\epsilon \rightarrow 0} \int_{\partial B_\epsilon} v' \phi dA = \phi(p)$$

and

$$\lim_{\epsilon \rightarrow 0} \int_{\partial B_\epsilon} v \frac{\partial \phi}{\partial n} dA = 0.$$

Therefore

$$\int_{B_R} vL\phi dA = \phi(p)$$

and v is our fundamental solution.

The other properties of v are a straightforward calculation. □

Lemma 9.3. *Let u be a smooth function on B_R such that $Lu \geq -b$. Then*

$$u(p) \leq \frac{1}{\sqrt{\text{vol}(B_R)}} \sqrt{\int_{B_R} u^2 dV} + b/2.$$

Proof. We first apply Green's identity on $B_{r'} \setminus B_\epsilon$:

$$(9.2) \quad \begin{aligned} \int_{B_{r'} \setminus B_\epsilon} vLudV &= \int_{B_{r'} \setminus B_\epsilon} uLv dV + \int_{\partial B_{r'}} \left(v \frac{\partial u}{\partial r} - u \frac{\partial v}{\partial r} \right) dA \\ &\quad - \int_{\partial B_\epsilon} \left(v \frac{\partial u}{\partial r} - u \frac{\partial v}{\partial r} \right) dA \end{aligned}$$

where v is the fundamental solution on the ball $B_{r'}$. Recall that $Lv = 0$, $v(r') = 0$ and

$$\lim_{\epsilon \rightarrow 0} - \int_{\partial B_\epsilon} \left(v \frac{\partial u}{\partial r} - u \frac{\partial v}{\partial r} \right) dA = u(p).$$

Therefore after taking the limit of (9.2) as $\epsilon \rightarrow 0$ and rearranging terms, we have

$$u(p) = \int_{B_{r'}} vLudV + \int_{\partial B_{r'}} uv' dA = \int_{B_{r'}} vLudV + v'(R) \int_{\partial B_{r'}} u dA.$$

By letting $u \equiv -1/2$ and solving the above equation for $\int_{B_{r'}} vLudV = \int_{B_{r'}} v dV$, we have

$$0 \geq \int_{B_{r'}} v dV = \frac{\sqrt{3}}{2} \frac{\sinh r'}{\sinh(\sqrt{3}r')} - \frac{1}{2} \geq -\frac{1}{2}.$$

Now if u is any smooth function on $B_{r'}$ with $Lu \geq -b$, we have

$$\int_{B_{r'}} vLudV \leq b/2.$$

Therefore

$$u(p) \leq b/2 + \frac{\sqrt{3}}{4\pi \sinh(r') \sinh(\sqrt{3}r')} \int_{\partial B} u dA.$$

Rearranging, we have

$$\frac{4\pi}{\sqrt{3}} (u(p) - b/2) (\sinh(\sqrt{3}r') \sinh r') \leq \int_{\partial B_{r'}} u dA.$$

Next we integrate both sides from 0 to R :

$$\frac{4\pi}{\sqrt{3}} (u(p) - b/2) \int_0^R (\sinh(\sqrt{3}r') \sinh r') dr \leq \int_0^R \left(\int_{\partial B_{r'}} u dA \right) dr' = \int_{B_R} u dV.$$

Since $\sinh r' \leq \sinh(\sqrt{3}r')/\sqrt{3}$, we have

$$\frac{4\pi}{\sqrt{3}} \int_0^R (\sinh(\sqrt{3}r') \sinh r') dr' \leq 4\pi \int_0^R \sinh^2 r' dr' = \text{vol}(B_R).$$

By Holder's inequality

$$\int_{B_R} u dV \leq \sqrt{\text{vol}(B_R)} \sqrt{\int_{B_R} u^2 dV}.$$

Therefore

$$u(p) \leq \frac{1}{\sqrt{\text{vol}(B_R)}} \sqrt{\int_{B_R} u^2 dV} + b/2.$$

□

The three previous lemmas easily lead to the following theorem:

Theorem 9.4. *Let v be a vector field on B_R and assume $\|\Delta v\| < b$. Then*

$$\|v(p)\| \leq \frac{1}{\sqrt{\text{vol}(B_R)}} \sqrt{\int_{B_R} \|v\|^2 dV} + b/2.$$

By Theorem 4.3 of [Br1] every cohomology class in $H^1(M; E)$ that extends to a conformal deformation Φ of the projective boundary is represented by a Hodge form ω . However, it is not shown that ω extends continuously to Φ . We show this now.

Theorem 9.5. *The Hodge form ω extends continuously to the holomorphic quadratic differential Φ .*

Proof. We need to recall some of the proof of Theorem 4.3 in [Br1]. In the proof $\omega = \omega_m + \omega_c$ where ω_m is a model deformation that extends continuously to Φ and the correction term ω_c is a trivial deformation. Then $\omega_m = ds_m$ where s_m is the canonical lift of an automorphic vector field on \tilde{M} while $\omega_c = ds_c$ where s_c is the canonical lift of a vector field on M . The model w_m is an automorphic vector field that is in a standard form on the geometrically finite ends and near the cone-singularity. From this standard form we know that the norms $\|\Delta w_m\|$ and $\|\Delta(\text{curl } w_m)\|$ decay to zero at the projective boundary.

Since we know that ω_m extends continuously to Φ , to prove the theorem we need to show that w_c and $\text{curl } w_c$ extend continuously to the zero vector field. From the proof of Theorem 4.3 in [Br1] we know that both w_c and $\text{curl } w_c$ are in L^2 . We also know that $\Delta w = \Delta(\text{curl } w) = 0$ so $\Delta w_c = -\Delta w_m$ and $\Delta(\text{curl } w_c) = -\Delta(\text{curl } w_m)$. Therefore the norms $\|\Delta w_c\|$ and $\|\Delta(\text{curl } w_c)\|$ decay to zero at the projective boundary.

We now apply Theorem 9.4. Let p_n be a sequence of points in M converging to p_∞ in Σ . For large values of n there will be balls B_n centered at p_n and embedded in M such that $\|\Delta w_c\|$ and $\|\Delta(\text{curl } w_c)\|$ are less than b_n on B_n with $b_n \rightarrow 0$. Then by Theorem 9.4

$$\|w_c(p_n)\| \leq \frac{1}{\sqrt{\text{vol}(B_n)}} \sqrt{\int_{B_n} \|w_c\|^2} + b_n/2$$

and

$$\|\text{curl } w_c(p_n)\| \leq \frac{1}{\sqrt{\text{vol}(B_n)}} \sqrt{\int_{B_n} \|\text{curl } w_c\|^2} + b_n/2.$$

Since both w_c and $\text{curl } w_c$ are in L^2 , the right-hand side of these inequalities limits to zero at $n \rightarrow \infty$. Therefore $\|w_c(p_n)\|$ and $\|\text{curl } w_c(p_n)\|$ limit to zero on Σ . \square

9.1. Strain fields. Next we prove a mean value inequality for Hodge forms. The real and imaginary parts of a Hodge form are strain fields so this is equivalent to proving a mean value inequality for harmonic strain fields.

We begin by defining a strain field. If v is a vector field on a Riemannian manifold with covariant derivative ∇ , then ∇v is a vector valued 1-form. The traceless symmetric part of ∇v is the *strain* $\text{str } v$ of v which measures the conformal distortion of v . The real and imaginary parts of an E -valued 1-form are vector valued 1-forms. It is shown in [HK1] that if v is a harmonic, divergence free vector field and if ω is the associated Hodge form, then $\text{Re } \omega = \text{str } v$ and $\text{Im } \omega = -\text{str } \text{curl } v$.

Lemma 9.6. *If η is a strain field with $\Delta\eta = 0$, then $-\Delta\|\eta\| + 2\|\eta\| \geq 0$.*

Proof. The proof is a bit more complicated than Lemma 9.1 because the Weitzenböck formula for strain fields is more involved. The Laplacian $\Delta = d\delta + \delta d$ for E -valued forms has the decomposition $\Delta = D^*D + DD^* + T^*T + TT^*$ where $d = D + T$ and $\delta = D^* + T^*$ are the decomposition of d and δ into their real and imaginary parts. We deal with the first two terms $\Delta_D = D^*D + DD^*$ and the last two terms $H = T^*T + TT^*$ separately. For Δ_D we have (see [Wu])

$$\Delta_D\eta = -\sum_i \nabla_{e_i e_i}^2 \eta - \sum_{i,j} \omega^i \wedge (R(e_i, e_j)\eta)(e_j).$$

Let

$$\mathcal{R}\eta = -\sum_{i,j} \omega^i \wedge (R(e_i, e_j)\eta)(e_j)$$

so

$$\Delta_D = \nabla^* \nabla + \mathcal{R}.$$

Since $R(e_i, e_j)$ is a tensor, \mathcal{R} is purely algebraic and therefore easy to calculate.

Any strain field η can be written as a linear combination $\eta = \sum_{i,j} f_k^l e_k \otimes \omega^l$. Then

$\mathcal{R}\eta = \sum_{k,l} f_k^l \mathcal{R}(e_k \otimes \omega^l)$ so we need to calculate

$$\mathcal{R}(e_k \otimes \omega^l) = \sum_{i,j} \omega^i \wedge (\omega^l(e_j)R(e_i, e_j)e_k + (R(e_i, e_j)\omega^l)(e_j)e_k).$$

To calculate these terms, we recall that for hyperbolic space

$$R(e_i, e_j)e_k = \delta_k^i e_j - \delta_k^j e_i$$

where δ_j^i is the Kronecker delta function. For the first term we have

$$\omega^i \wedge \omega^l(e_j)R(e_i, e_j)e_k = \delta_j^l (\delta_k^i e_j - \delta_k^j e_i) \otimes \omega^i.$$

If $k \neq l$, this is only non-zero if $i = j$ and $k = l$ in which case we get $e_l \otimes \omega^k$. If $k = l$, there are two non-zero terms $-e_i \otimes \omega^i$ when $i = j \neq k$. Therefore

$$\sum_{i,j,k,l} f_k^l \omega^i \wedge \omega^l(e_j)R(e_i, e_j)e_k = \eta^T - \text{tr } \eta I.$$

For the second term we calculate

$$(R(e_i, e_j)\omega^l)(e_j) = \omega^l(-R(e_i, e_j)e_j) = (1 - \delta_j^i)\omega^l(e_i) = (1 - \delta_j^i)\delta_i^l$$

so

$$\sum_{i,j} \omega^i \wedge (R(e_i, e_j)\omega^l)(e_j)e_k = 2e_k \otimes \omega^l.$$

Since η is a strain field, it is traceless and symmetric; therefore the two terms combine to give

$$\mathcal{R}\eta = -3\eta.$$

For a harmonic strain field η it is shown that $H\eta = \eta$ in [HK1]. Combining our work so far we have

$$\Delta\eta = \nabla^* \nabla \eta + \mathcal{R}\eta + H\eta = \nabla^* \nabla \eta - 3\eta + \eta = 0$$

and therefore

$$\nabla^* \nabla \eta = 2\eta.$$

The remainder of the proof is exactly like the proof of Lemma 9.1 and we will not repeat it. Note that for a harmonic vector field $\nabla^* \nabla v = -2v$ which accounts for the sign change from the bound we get for vector fields to the bound for strain fields. \square

Now let $Lu = (-\Delta + 2)u$. We then can restate Lemmas 9.2 and 9.3 and Theorem 9.4 for this new definition of L . The proofs are so similar that we leave the details to the reader.

Lemma 9.7. *The function*

$$v = \frac{-\cos(\sqrt{2}r) + \cot(\sqrt{2}R) \sin(\sqrt{2}r)}{4\pi \sinh r}$$

is a fundamental solution for L if $R < \frac{\pi}{\sqrt{2}}$. Furthermore $v(R) = 0$, $v(r) \leq 0$ for $0 < r < R$ and

$$v'(R) = \frac{\sqrt{2}}{4\pi \sinh(R) \sin(\sqrt{2}R)}.$$

Lemma 9.8. *Let u be a smooth function on B_R such that $Lu \geq 0$. Then*

$$u(p) \leq \frac{3\sqrt{2} \operatorname{vol}(B_R)}{4\pi f(R)} \sqrt{\int_{B_R} u^2 dV}$$

where $f(R) = \cosh(R) \sin(\sqrt{2}R) - \sqrt{2} \sinh(R) \cos(\sqrt{2}R)$ for $R < \frac{\pi}{\sqrt{2}}$.

Using the fact that $\|\omega(p)\|^2 = \|\operatorname{Re} \omega(p)\|^2 + \|\operatorname{Im} \omega(p)\|^2$, we then have

Theorem 9.9. *Let ω be a Hodge form on a ball B_R of radius R centered at a point p . Then*

$$\|\omega(p)\| \leq \frac{3\sqrt{2} \operatorname{vol}(B_R)}{4\pi f(R)} \sqrt{\int_{B_R} \|\omega\|^2 dV}$$

for $R < \frac{\pi}{\sqrt{2}}$.

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