

Cite this: *Energy Environ. Sci.*, 2011, **4**, 3640

www.rsc.org/ees

PAPER

Combination of large nanostructures and complex band structure for high performance thermoelectric lead telluride

Yanzhong Pei, Nicholas A. Heinz, Aaron LaLonde and G. Jeffrey Snyder*

Received 9th June 2011, Accepted 11th July 2011

DOI: 10.1039/c1ee01928g

The complexity of the valence band structure in p-type PbTe has been shown to enable a significant enhancement of the average thermoelectric figure of merit (zT) when heavily doped with Na. It has also been shown that when PbTe is nanostructured with large nanometer sized Ag₂Te precipitates there is an enhancement of zT due to phonon scattering at the interfaces. The enhancement in zT resulting from these two mechanisms is of similar magnitude but, in principle, decoupled from one another. This work experimentally demonstrates a successful combination of the complexity in the valence band structure with the addition of nanostructuring to create a high performance thermoelectric material. These effects lead to a high zT over a wide temperature range with peak $zT > 1.5$ at $T > 650$ K in Na-doped PbTe/Ag₂Te. This high average zT produces 30% higher efficiency (300–750 K) than pure Na-doped PbTe because of the nanostructures, while the complex valence band structure leads to twice the efficiency as the related n -type La-doped PbTe/Ag₂Te without such band structure complexity.

Thermoelectric (TE) applications have attracted increasing interest in the last decade as a means to combat the ever growing rate of energy consumption throughout the world. The two main applications for thermoelectric materials are power generation, which utilizes the Seebeck effect, and solid state cooling, which has its roots in the Peltier effect. In recent years, thermoelectric power generation has been a prime interest to the automotive industry¹ as a sustainable and emission free route to vehicular waste heat recovery. The effectiveness of this process is restricted by the overall efficiency of the thermoelectric materials.

On a materials level, the efficiency of the thermal to electrical energy conversion is determined by the thermoelectric figure of

merit, $zT = S^2\sigma T/(\kappa_E + \kappa_L)$, where S , σ , κ_E and κ_L are the Seebeck coefficient, electrical conductivity, and the electronic and lattice components of the thermal conductivity. Since S , σ and κ_E have an intimate relationship with the carrier density,^{1,2} the grand challenge in designing thermoelectric materials is the decoupling of electronic and thermal properties. Many methods to achieve a large power factor ($S^2\sigma$) or to reduce the thermal conductivity were proven to be successful, but combination of these two effects is difficult. As a result, the current best commercially available thermoelectric materials have a peak zT near unity.

A number of mechanisms have been proposed to explain high zT in p-type PbTe. The Tl-doping in PbTe:Tl produces resonant electronic states that enhances the Seebeck coefficient, but reduces hole mobility³ and has $zT \sim 1.4$ at 500 °C. Na-doped PbTe:Na,⁴ without resonant states,^{5,6} has similarly high zT that

Materials Science, California Institute of Technology, Pasadena, CA, 91125, USA. E-mail: jsnyder@caltech.edu

Broader context

Thermoelectric generators, which directly convert heat into electricity are being considered for industrial or automotive waste heat recovery. However, the relatively low efficiency of thermoelectric materials has limited their use to niche applications. Effective strategies for high performance thermoelectric materials during the past half-century have depended largely on the reduction of heat conduction by the crystal lattice and recently on the proposed band structure engineering for improvements in the electronic properties. The present study shows the effort of combining these two independent research directions, namely reducing the lattice thermal conductivity by scattering phonons through nanostructuring and improving the electronic properties by exploiting complex band structure, as a route to improve the thermoelectric performance. The addition of Ag₂Te in PbTe creates nanosized scattering centers for heat carrying phonons while further Na-doping enables multiple bands for enhanced electronic properties. In this way, an enhancement of 30%~100% on thermoelectric efficiency is achieved as compared with the same material having either multiple bands or nanosized scattering centers alone, enabling a synergistic route in the search of high performance thermoelectric materials and devices.

can be rationalized to be from the multi valley valence band structure^{4,7} of PbTe (a second valence band exists slightly below the principal valence band). Because of the complicated band structure, heavy and precise doping of the valence band is required to reach the high zT doping level.⁴

Because this influence on the valence band structure is purely an electronic effect, there is good potential to further enhance zT for PbTe:Na by reducing the lattice thermal conductivity.

It has been demonstrated in numerous reports that nanostructures in PbTe effectively scatter the long wavelength heat carrying phonons to lower the lattice thermal conductivity.^{8–12} This approach led to a peak zT of 1.4~1.7 in both p-^{10,13,14} and n-type^{8,15,16} PbTe materials.

Attempts to nanostructure p-type PbTe:TI with TlSbTe₂ have not led to higher zT because of reduced mobility despite lower thermal conductivity.¹⁷

Among these reports, PbTe/Ag₂Te nanocomposites prepared from an equilibrium precipitation differ because of the relatively large size of the nanoscale Ag₂Te precipitates and in addition, the resultant composite is an intrinsic semiconductor due to the self-compensated substitutional and interstitial Ag⁺ pairs.^{16,18,19} The lack of extrinsic dopants enables a direct observation of κ_L -reduction due to the embedded nanostructures¹⁶ because there is very little contribution to κ due to the electrons. Moreover, the large Ag₂Te nanostructures were found to reduce the κ_L of PbTe at high temperatures with an equal or greater effect than those of smaller nanosized particles,¹⁶ while the small nanoparticles reduce the thermal conductivity better than the large particles at low temperatures.

The PbTe matrix in the PbTe/Ag₂Te nanocomposites can be doped n-type with La¹⁶ or excess Ag¹⁹ which allows as an independent control of carrier density while maintaining the composite microstructure.

Additionally, introducing embedded Ag₂Te nanoparticles into the matrix may mechanically strengthen these brittle materials²⁰ for practical applications.

The present work shows our effort toward Na-doping PbTe/Ag₂Te nanocomposites and the resulting thermoelectric properties. The low lattice thermal conductivity due to the Ag₂Te nanoinclusions and the excellent electronic properties due to the complex valence band structure is combined to result in a $zT > 1.5$ at high temperatures. Moreover, there are significant improvements in the average zT and thermoelectric efficiency throughout the entire temperature range studied, as compared with analogous materials without nanostructures or band structure complexity, or with small nanostructures.

P-type nanocomposites were obtained by Na-doping (PbTe:Na/Ag₂Te). PbTe/Ag₂Te nanocomposites with a composition of (PbTe)_{0.945}(Ag₂Te)_{0.055} (Fig.1) were pre-synthesized according the reported method¹⁶ and subsequently used as the starting material for making PbTe:Na/Ag₂Te together with appropriate amounts of Na and Te metals. The nominal concentration of Na [(Na_xPbTe_{1+x})_{0.945}(Ag₂Te)_{0.055}] is $x = 0\sim 3$ at.% and the samples for this study were synthesized with the same method used previously,¹⁶ including sealing, melting (Step 1 in Fig. 1), quenching, annealing (Step 2 followed by Step 3) and hot pressing. Phase components were checked using X-ray diffraction and scanning electron microscopy (SEM) equipped with an energy dispersive spectrometer (EDS).

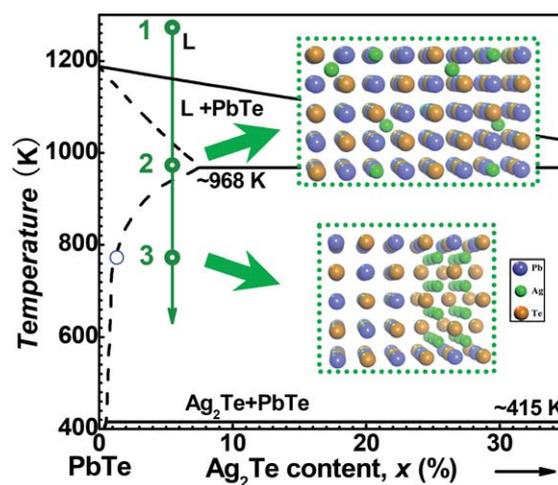


Fig. 1 Phase diagram of PbTe-Ag₂Te system and the procedure used for material preparation. A low-temperature annealing leads to the precipitation of Ag₂Te from the high temperature supersaturated alloy phase.

Hot pressed disk-shape samples with relative densities of 98% or higher were used for the measurements. Details on measuring the transport properties were given elsewhere,^{4,16} where the heat capacity [C_p in k_B per atom = $3.07 + 4.7 \times 10^{-4} \times (T/K-300)$]^{4,7,19,21} was determined by drop calorimetry²⁰ and believed to be accurate for several lead chalcogenides.^{4,22,23} The thermal conductivity for most of the recently reported high zT PbTe materials at 600~800 K, was determined using a heat capacity of,^{3,16} or close to, the Dulong-Petit approximation by $\pm 5\%$.^{10,13,24–26} It should be noted that this equation determines a value $\sim 10\%$ higher than the Dulong-Petit law ($3k_B$ per atom) at $T > 700$ K, and this method is used to recalculate all the literature data used for comparison in this study. The uncertainty for each measurement of S , σ and κ is $\sim 5\%$, resulting in a combined error of $\sim 20\%$ in the determined zT values.

Annealing the high temperature supersaturated solid solution phase (Step 2) of PbTe/Ag₂Te^{27,28} in the low temperature, two-phase region (Step 3) creates homogeneously distributed Ag₂Te nanoinclusions in a PbTe matrix.¹⁶ Fig. 2a shows the typical nanostructure for the obtained as-cast ingots. The plate-like Ag₂Te precipitates have preferred orientation parallel with $\langle 001 \rangle$ directions in the PbTe matrix, consistent with previous reports.^{16,19,29} The Ag₂Te precipitates are nanometer sized in one direction but micron-size in the other directions, which is larger than those generally observed in PbTe/AgSbTe₂ and its analogs.^{9,30} No Ag₂Te precipitates smaller than ~ 50 nm could be found in PbTe/Ag₂Te system with similar synthesis conditions.¹⁶

Similar to PbTe:Na,⁴ Na is found to be an effective p-type dopant in PbTe/Ag₂Te as indicated by the Hall coefficient (R_H) and Seebeck coefficient measurements. However, the doping solubility of Na in PbTe/Ag₂Te is found to be much smaller than in pure PbTe, since the maximum measured room temperature Hall density ($\rho_H = 1/eR_H$, e is the electron charge) in PbTe:Na/Ag₂Te is much smaller than that in PbTe:Na, this is presumably due to the presence of a small amount of soluble Ag₂Te ($\sim 1\%$)¹⁶ in PbTe matrix reducing the solubility of Na. Even though the nominal concentrations of Na are comparable in PbTe:Na/Ag₂Te and PbTe:Na,⁴ the measured ρ_H in PbTe:Na/Ag₂Te does

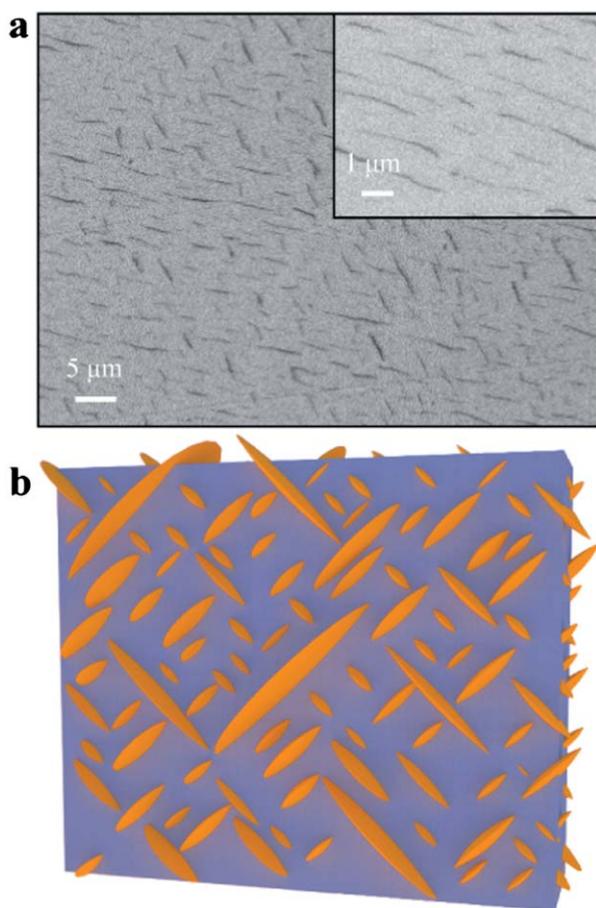


Fig. 2 (a) Typical SEM image of as cast PbTe:Na/Ag₂Te ingot. The short plates with a darker contrast indicate the Ag₂Te phase with different orientations in different grains. The inset is a close up image of the nanostructures indicating the size to be on the order of a few hundred nanometers. (b) A 3D representation of the observed Ag₂Te (orange) nanostructures embedded in the PbTe matrix (blue).

not exceed $4 \times 10^{19} \text{ cm}^{-3}$, while p_H in PbTe:Na can be as high as $14 \times 10^{19} \text{ cm}^{-3}$. Since it was found that the optimal p_H was close to 10^{20} cm^{-3} , due to the dominance of the heavy band behavior in p-PbTe,⁴ the most heavily doped PbTe:Na/Ag₂Te samples with nominal Na content of 1%, 2% and 3% having p_H of 2.5, 3.1 and $3.7 \times 10^{19} \text{ cm}^{-3}$ are used for the following discussions and marked as 2.5e19, 3.1e19 and 3.7e19, respectively.

As has been well established in p-PbTe, the room temperature Seebeck coefficient *versus* Hall density (Pisarenko plot, solid line in Fig. 3a) provides a powerful characterization of the band structure for PbTe.³¹ The majority of published measurements on bulk p-type PbTe fall on this Pisarenko line, as do the present results, within the measurement uncertainty (Fig. 3a). Similar to the findings in La-doped PbTe/Ag₂Te (PbTe:La/Ag₂Te),¹⁶ neither doping with Na nor introducing Ag₂Te alters the band structure of the PbTe matrix.

The flattening of S at high doping levels ($p_H > \sim 3 \times 10^{19} \text{ cm}^{-3}$) can be well understood by the complex valence band structure shown in the inset of Fig. 3a.^{4,31} Furthermore, the energy of the light valence band reduces with increasing temperature and moves below the heavy valence band at $\sim 450 \text{ K}$,^{31–38} making the

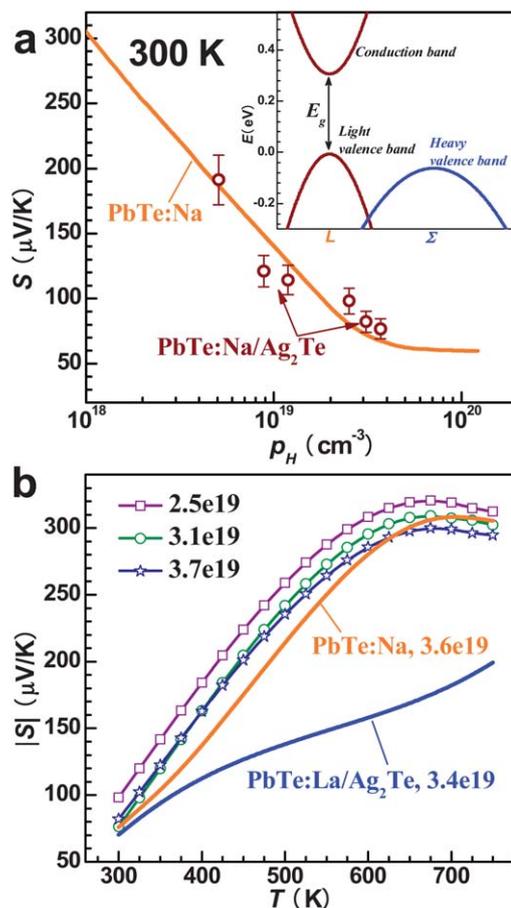


Fig. 3 (a) Room temperature Seebeck coefficient *versus* Hall density for PbTe:Na/Ag₂Te with an overlaid Pisarenko plot for PbTe:Na. The inset shows the schematic band structure at 300 K. (b) Temperature dependent Seebeck coefficient along with PbTe:Na and n-type PbTe:La/Ag₂Te. Complexity of valence band structure significantly increases the Seebeck coefficient.

heavy valence band the dominant band for hole transport at high temperatures where the zT peaks. The electronic transport, optical spectroscopy and other properties of p-PbTe can be well described by this two-valence-band mode.^{31–38}

Rather than the Seebeck coefficient being proportional to absolute temperature (e.g. n-type PbTe:La/Ag₂Te in Fig. 3b), a significant increase in S (Fig. 3b) is observed, particularly at high temperatures, due to the increasing contribution of the heavy band carriers resulting from the temperature dependent band structure and Fermi broadening. The analogous n-type material (PbTe:La/Ag₂Te¹⁶) has a much lower Seebeck coefficient than p-type PbTe:Na/Ag₂Te, due to the lack of conduction band complexity.^{4,31} These PbTe:Na/Ag₂Te materials show a roughly unchanged Seebeck coefficient with respect to PbTe:Na (without Ag₂Te) at similar doping levels, which further supports the above discussion (S vs. p_H) that Ag₂Te inclusions do not affect the band structure of PbTe.

While the Seebeck coefficient is similar, the PbTe:Na/Ag₂Te composites have lower charge carrier mobility compared with PbTe:Na, due to both the increased scattering of carriers at the phase boundaries and the increased point defect scattering in the

matrix phase (dilute solid solution of PbTe with $\sim 1\%$ $\text{Ag}_2\text{Te}^{16,27}$). Shown in the inset of Fig. 4 is the Hall mobility ($\mu_H = R_H\sigma$) of PbTe:Na/ Ag_2Te with $p_H > \sim 3 \times 10^{19} \text{ cm}^{-3}$ along with a comparison to PbTe:Na with similar p_H . The PbTe:Na/ Ag_2Te composites have similar Hall mobility over the entire temperature range, and converge with the PbTe:Na sample at high temperatures. This behavior is understood by the scattering of charge carriers in these materials being dominated by acoustic phonons at high temperatures,³¹ while interface and alloy scattering contribute at lower temperatures. As a result, the electrical resistivity of PbTe:Na/ Ag_2Te nanocomposites is higher than that of PbTe:Na (Fig. 4) at low temperatures, when the doping levels are nearly the same ($p_H = 3.6 \times 10^{19} \text{ cm}^{-3}$ for PbTe:Na vs. $p_H = 3.7 \times 10^{19} \text{ cm}^{-3}$ for PbTe:Na/ Ag_2Te). Similar to PbTe:Na⁴ (and n-PbTe^{16,31}), the electrical resistivity (ρ) rises faster than normally expected for a system dominated by acoustic scattering ($\rho \propto T^{-1.5}$). This can be explained by the increasing effective mass of carriers because of the temperature dependent mass of the light bands as well as the transition of holes from the light to the heavy band that has lower mobility.^{19,31–38} Increasing Hall density results in lower electrical resistivity in PbTe:Na/ Ag_2Te nanocomposites, as shown in Fig. 4.

The thermal conductivity (κ) is reduced by $\sim 50\%$ in roughly the entire measured temperature range due to the Ag_2Te nano-inclusions (Fig. 5). It is important to note that the Hall density of all the samples was similar and found to be $\sim 3.5 \times 10^{19} \text{ cm}^{-3}$. To avoid the ambiguity of bipolar thermal conductivity and for the sake of clarity, only heavily doped samples are shown in Fig. 5.

The observed κ -reduction is only partially attributed to the presence of the nano-inclusions. The reduced electrical conductivity and thus a reduced electronic component to the thermal conductivity (κ_E) also contributes to lower κ . It is difficult to accurately estimate the κ_E via the Wiedemann-Franz law ($\kappa_E = L\sigma T$) because of the difficulty in determining the Lorenz number (L) in p-PbTe due to the complex valence band structure and the band non-parabolicity.^{31,39,40} For simplicity, an estimation of L is made using a single parabolic band (SPB) model assuming an

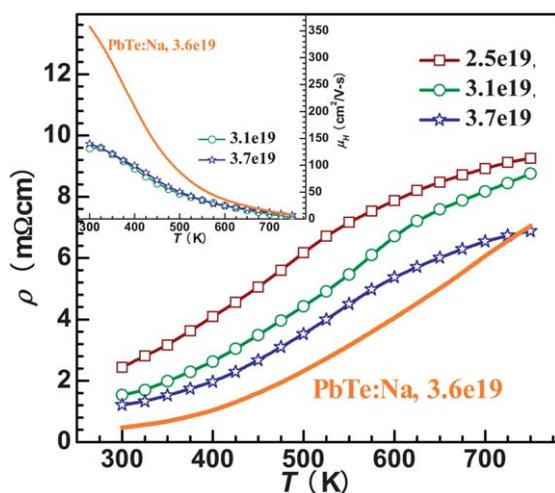


Fig. 4 Temperature dependent electrical resistivity for PbTe:Na/ Ag_2Te and PbTe:Na. Ag_2Te nano-inclusions simultaneously scatter the carriers thereby reducing the Hall mobility, particularly at low temperatures as shown in the inset.

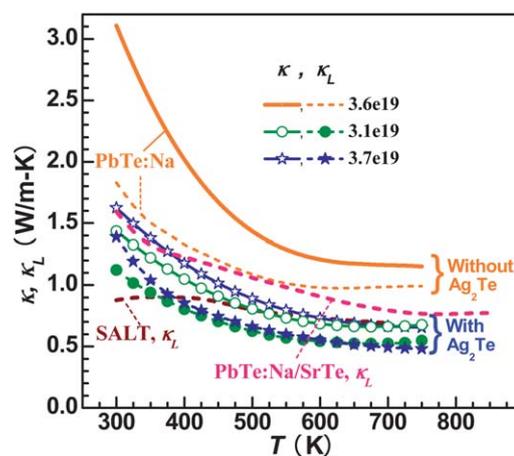


Fig. 5 Temperature dependent thermal conductivity and its lattice component for PbTe:Na/ Ag_2Te compared to PbTe:Na.⁴ Relatively large Ag_2Te nano-inclusions effectively scatter phonons thus reducing the lattice thermal conductivity (κ_L) to $\sim 0.5 \text{ W/m-K}$ at $T > \sim 600 \text{ K}$, well below the value found in high zT PbTe:Na/SrTe¹⁰ and PbTe/NaSbTe₂ (SALT)¹³ containing smaller nanostructures.

acoustic phonon scattering mechanism,⁴¹ which results in an L with a deviation of less than 10% when compared to a more rigorous single non-parabolic band and multiple band model calculation.^{7,16,39,42} And the same method of L -estimation is used to recalculate the literature data for the following discussion.

By subtracting the electronic component from the total thermal conductivity, the lattice component (κ_L) is calculated and shown in Fig. 5. It can be expected that the introduction of nano-inclusions effectively reduces the lattice thermal conductivity due to the enhanced scattering of phonons at boundaries.^{8,9,16,19} The extremely low κ_L of $\sim 0.5 \text{ W/m-K}$ at $T > 600 \text{ K}$ is approaching the theoretical minimum value of 0.36 W/m-K .^{16,43} It should be noted that the total thermal conductivity of PbTe:Na/ Ag_2Te composites is significantly lower than the lattice thermal conductivity of PbTe:Na. As compared with other PbTe nanocomposites having smaller structure features, such as PbTe:Na/SrTe¹⁰ and PbTe/NaSbTe₂¹³ (SALT) and their analogs, where the κ_L is now recalculated with the same estimation of both C_p and L , PbTe with relatively large Ag_2Te precipitates has even lower lattice thermal conductivity at high temperatures, similar to that observed in n-type materials.¹⁶ This is presumably due to at least partial precipitate dissolution of small nanoparticles observed at high temperatures,⁴⁴ while the large Ag_2Te precipitates in our current and previous study^{16,19} are obtained by annealing at high temperature (Fig. 1) and thus are likely to be more stable.

When compared with n-type La-doping¹⁶ (using the same estimation of C_p), p-type Na-doping in PbTe/ Ag_2Te nanocomposites has significantly higher zT (Fig. 6a) at low temperatures because of the favorable electronic effect of the complex band structure.⁴⁷ Alternatively, compared with PbTe:Na having similar p_H of $3.6 \times 10^{19} \text{ cm}^{-3}$,⁴ introducing Ag_2Te nano-inclusions significantly reduces the thermal conductivity (Fig. 5) and thus increases zT (Fig. 6a). This is further evidenced by the excellent agreement between the measured zT of the PbTe:Na/ Ag_2Te sample with $p_H = 3.7 \times 10^{19} \text{ cm}^{-3}$ and the projected zT curve for

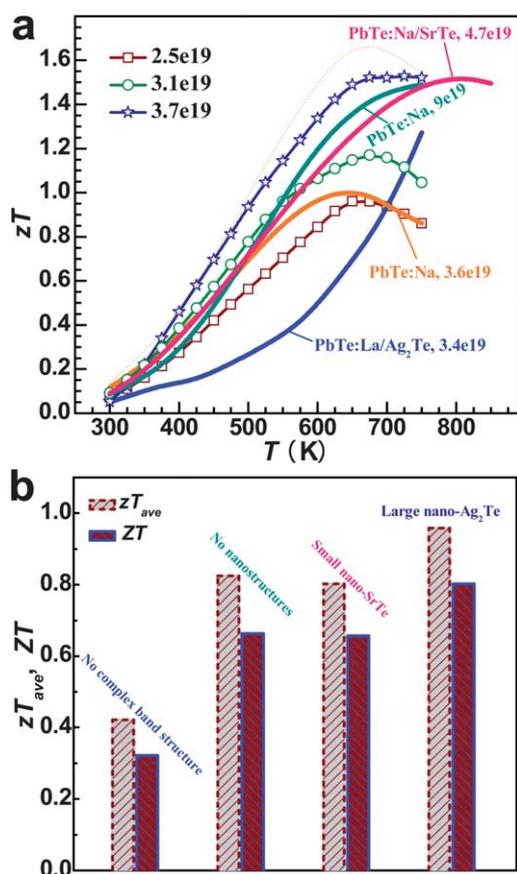


Fig. 6 (a) Temperature dependent thermoelectric figure of merit for PbTe:Na/Ag₂Te, PbTe:Na,⁴ PbTe:La/Ag₂Te¹⁶ and PbTe:Na/SrTe.¹⁰ The improvement in zT is attributed to the reduction of the lattice thermal conductivity in PbTe:Na (calculated dashed curve). (b) Comparison of the device ZT and the average zT (zT_{ave}) from 300–750 K for PbTe:Na/Ag₂Te, PbTe:Na,⁴ PbTe:La/Ag₂Te¹⁶ and PbTe:Na/SrTe.¹⁰ There is roughly a 40–100% enhancement in both zT_{ave} and ZT due to the band structure complexity and nanostructuring, when having similar room temperature Hall density of $\sim 3.5 \times 10^{19} \text{ cm}^{-3}$. A $\sim 20\%$ enhancement on zT_{ave} and ZT can be also seen when compared with similar materials having small structure features such as PbTe:Na/SrTe.¹⁰

PbTe:Na with $p_H = 3.6 \times 10^{19} \text{ cm}^{-3}$ (dashed orange), which is generated by simply assuming the κ_L value to be that of PbTe:Na/Ag₂Te with $p_H = 3.7 \times 10^{19} \text{ cm}^{-3}$. However, it is unlikely that these materials are fully optimized as the optimal carrier density, as observed in PbTe:Na with $p_H = 9 \times 10^{19} \text{ cm}^{-3}$ (cyan curve),⁴ was not obtained in this study due to the small dopant solubility.

The current effort of combining both complex band structure and nanostructures enables a zT higher than 1.5 at $T > \sim 650 \text{ K}$. More importantly, as shown in Fig. 6b, both the average zT (zT_{ave}) and the device ZT that is calculated from the theoretically available power generation efficiency (η)^{45–48} (between 300 K and 750 K), are found to be enhanced by 40–100% in PbTe:Na/Ag₂Te, when compared to PbTe:La/Ag₂Te¹⁶ and PbTe:Na⁴ having similar p_H . Here, the estimations of device ZT also take into account the thermoelectric compatibility effect.^{46–48} It should also be noted that the zT_{ave} enables a fairly accurate evaluation of the relative change on the efficiency⁴⁵ even though real device ZT is lower than zT_{ave} due to the variable

compatibility factor.^{46–48} Even though PbTe:Na/Ag₂Te composite is not fully optimized for carrier concentration, 15–20% higher ZT and zT_{ave} have already been achieved as compared with the best PbTe:Na ($p_H = 9 \times 10^{19} \text{ cm}^{-3}$)⁴ and PbTe:Na/SrTe¹⁰ having smaller and endotaxial nanostructures, indicating the effectiveness of nanostructures for high performance thermoelectrics.

In summary, band structure complexity and nanostructure effects are simultaneously considered as an effective approach for improving thermoelectric performance. These combined approaches have been demonstrated in PbTe:Na/Ag₂Te resulting in a peak zT higher than 1.5, and more importantly, this results in significant enhancements of the average zT and thermoelectric efficiency. Further study of the combination of carrier density and nanostructure control should result in even higher thermoelectric performance in similar PbTe materials.

The authors gratefully acknowledge NASA-JPL and DARPA for funding support of this work.

References

- 1 A. F. Ioffe, *Semiconductor thermoelements, and Thermoelectric cooling*, Infosearch, London, 1957.
- 2 G. J. Snyder and E. S. Toberer, *Nat. Mater.*, 2008, **7**, 105–114.
- 3 J. Heremans, V. Jovovic, E. Toberer, A. Saramat, K. Kurosaki, A. Charoenphakdee and G. J. Snyder, *Science*, 2008, **321**, 554–557.
- 4 Y. Pei, A. LaLonde, S. Iwanaga and G. J. Snyder, *Energy Environ. Sci.*, 2011, **4**, 2085–2089.
- 5 K. Hoang, S. D. Mahanti and M. G. Kanatzidis, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2010, **81**, 115106.
- 6 S. Ahmad, S. D. Mahanti, K. Hoang and M. G. Kanatzidis, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2006, **74**, 155205.
- 7 Y. Pei, X. Shi, A. LaLonde, H. Wang, L. Chen and G. J. Snyder, *Nature*, 2011, **473**, 66–69.
- 8 K. F. Hsu, S. Loo, F. Guo, W. Chen, J. S. Dyck, C. Uher, T. Hogan, E. K. Polychroniadis and M. G. Kanatzidis, *Science*, 2004, **303**, 818–821.
- 9 M. G. Kanatzidis, *Chem. Mater.*, 2010, **22**, 648–659.
- 10 K. Biswas, J. He, Q. Zhang, G. Wang, C. Uher, V. P. Dravid and M. G. Kanatzidis, *Nat. Chem.*, 2011, **3**, 160–166.
- 11 A. Minnich, M. Dresselhaus, Z. Ren and G. Chen, *Energy Environ. Sci.*, 2009, **2**, 466–479.
- 12 J. Baxter, Z. Bian, G. Chen, D. Danielson, M. Dresselhaus and A. Fedorov, *Energy Environ. Sci.*, 2009, **2**, 559–588.
- 13 P. F. P. Poudeu, J. D'Angelo, A. Downey, J. Short, T. Hogan and M. M. Kanatzidis, *Angew. Chem., Int. Ed.*, 2006, **45**, 3835–3839.
- 14 J. Androulakis, K. F. Hsu, R. Pcionek, H. Kong, C. Uher, J. J. D'Angelo, A. Downey, T. Hogan and M. G. Kanatzidis, *Adv. Mater.*, 2006, **18**, 1170.
- 15 J. Androulakis, C. H. Lin, H. J. Kong, C. Uher, C. I. Wu, T. Hogan, B. A. Cook, T. Caillat, K. M. Paraskevopoulos and M. G. Kanatzidis, *J. Am. Chem. Soc.*, 2007, **129**, 9780–9788.
- 16 Y. Pei, J. Lensch-Falk, E. S. Toberer, D. L. Medlin and G. J. Snyder, *Adv. Funct. Mater.*, 2011, **21**, 241–249.
- 17 H. Wang, A. Charoenphakdee, K. Kurosaki, S. Yamanaka and G. J. Snyder, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2011, **83**, 024303.
- 18 A. J. Strauss, *J. Electron. Mater.*, 1973, **2**, 553–569.
- 19 Y. Pei, A. F. May and G. J. Snyder, *Adv. Energy Mater.*, 2011, **1**, 291–296.
- 20 E. A. Skrabek and D. S. Trimmer, in *CRC handbook of thermoelectrics*, ed. D. M. Rowe, CRC Press, Boca Raton, Fla., 1995, p. 272.
- 21 H. Wang, Y. Pei, A. D. LaLonde and G. J. Snyder, *Adv. Mater.*, 2011, **23**, 1366–1370.
- 22 R. Blachnik and R. Igel, *Z. Naturforsch., B*, 1974, **29**, 625–629.
- 23 M. Zhou, J. F. Li and T. Kita, *J. Am. Chem. Soc.*, 2008, **130**, 4527–4532.

- 24 P. F. P. Poudeu, A. Gueguen, C. I. Wu, T. Hogan and M. G. Kanatzidis, *Chem. Mater.*, 2010, **22**, 1046–1053.
- 25 K. Ahn, M. K. Han, J. Q. He, J. Androulakis, S. Ballikaya, C. Uher, V. P. Dravid and M. G. Kanatzidis, *J. Am. Chem. Soc.*, 2010, **132**, 5227–5235.
- 26 A. Gueguen, P. F. P. Poudeu, C. P. Li, S. Moses, C. Uher, J. Q. He, V. Dravid, K. A. Paraskevopoulos and M. G. Kanatzidis, *Chem. Mater.*, 2009, **21**, 1683–1694.
- 27 F. Wald, *J. Less Common Met.*, 1967, **13**, 579–590.
- 28 R. Blachnik and B. Gather, *J. Less Common Met.*, 1978, **60**, 25–32.
- 29 J. Lensch-Falk, J. Sugar, M. Hekmaty and D. Medlin, *J. Alloys Compd.*, 2010, **504**, 37–44.
- 30 C. Vineis, A. Shakouri, A. Majumdar and M. Kanatzidis, *Adv. Mater.*, 2010, **22**, 3970–3980.
- 31 Y. I. Ravich, B. A. Efimova and I. A. Smirnov, *Semiconducting Lead Chalcogenides*, Plenum Press, New York, 1970.
- 32 G. Nimtz and B. Schlicht, *Springer Tracts in Modern Physics*, 1983, **98**, 1–117.
- 33 Y. I. Ravich, in *Lead Chalcogenides: Physics and Applications*, ed. D. Khokhlov, Taylor & Francis Group, New York, 2003, pp. 1–34.
- 34 R. S. Allgaier, *J. Appl. Phys.*, 1961, **32**, 2185–2189.
- 35 A. J. Crocker and L. M. Rogers, *J. Phys. Colloq.*, 1968, **29**, 129–132.
- 36 V. I. Kaidanov, R. B. Melnik, I. A. Chernik and A. A. Kosulina, *Sov. Phys. Semiconductors*, 1969, **2**, 1474.
- 37 R. N. Tauber, A. A. Machonis and I. B. Cadoff, *J. Appl. Phys.*, 1966, **37**, 4855–4860.
- 38 J. Androulakis, I. Todorov, D. Y. Chung, S. Ballikaya, G. Y. Wang, C. Uher and M. Kanatzidis, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2010, **82**, 115209.
- 39 I. A. Smirnov, M. N. Vinogradova, N. V. Kolomoets and L. M. Syssoeva, *Sov. Phys. Solid State*, 1968, **9**, 2074–2079.
- 40 I. A. Smirnov and Y. I. Ravich, *Sov. Phys. Semiconductors*, 1967, **1**, 739–741.
- 41 C. M. Bhandari and D. M. Rowe, in *CRC handbook of thermoelectrics*, ed. D. M. Rowe, CRC Press, Boca Raton, Fla., 1995, ch. 5, pp. 43–53.
- 42 S. Ahmad and S. D. Mahanti, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2010, **81**, 165203.
- 43 Y. K. Koh, C. J. Vineis, S. D. Calawa, M. P. Walsh and D. G. Cahill, *Appl. Phys. Lett.*, 2009, **94**, 153101.
- 44 J. He, J. R. Sootsman, L. Xu, S. Girard, J.-C. Zheng, M. G. Kanatzidis and V. P. Dravid, *J. Am. Chem. Soc.*, 2011, **133**, 8786.
- 45 R. W. Ure and R. R. Heikes, in *Thermoelectricity: science and engineering*, ed. R. R. Heikes and R. W. Ure, Interscience Publishers, New York, 1961, ch. 15, pp. 458–517.
- 46 G. J. Snyder, in *Thermoelectrics handbook: macro to nano*, ed. D. M. Rowe, CRC/Taylor & Francis, Boca Raton, 2006, ch. 9, pp. 1–26.
- 47 G. J. Snyder, *Appl. Phys. Lett.*, 2004, **84**, 2436–2438.
- 48 G. J. Snyder and T. Ursell, *Phys. Rev. Lett.*, 2003, **91**, 148301.