NARROWBAND LYMAN-CONTINUUM IMAGING OF GALAXIES AT $z \sim 2.85$

R. E. MOSTARDI¹, A. E. SHAPLEY¹, D. B. NESTOR¹, C. C. STEIDEL², N. A. REDDY³, AND R. F. TRAINOR²

¹ Department of Physics & Astronomy, University of California, Los Angeles, 430 Portola Plaza, Los Angeles, CA 90095, USA

² Cahill Center for Astrophysics, California Institute of Technology, MS 249-17, Pasadena, CA 91125, USA

³ Department of Physics and Astronomy, University of California, Riverside, 900 University Avenue, Riverside, CA 92521, USA

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ABSTRACT

We present results from a survey for $z \sim 2.85$ Lyman-continuum (LyC) emission in the HS1549+1933 field and place constraints on the amount of ionizing radiation escaping from star-forming galaxies. Using a custom narrowband filter (NB3420) tuned to wavelengths just below the Lyman limit at $z \gtrsim 2.82$, we probe the LyC spectral region of 49 Lyman break galaxies (LBGs) and 91 Ly α emitters (LAEs) spectroscopically confirmed at $z \gtrsim 2.82$. Four LBGs and seven LAEs are detected in NB3420. Using *V*-band data probing the rest-frame nonionizing UV, we observe that many NB3420-detected galaxies exhibit spatial offsets between their LyC and nonionizing UV emission and are characterized by extremely blue NB3420–*V* colors, corresponding to low ratios of nonionizing to ionizing radiation ($F_{\text{UV}}/F_{\text{LyC}}$) that are in tension with current stellar population synthesis models. We measure average values of $(F_{\text{UV}}/F_{\text{LyC}})$ for our LBG and LAE samples, correcting for foreground galaxy contamination and H I absorption in the intergalactic medium. We find $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{LBG}} = 82 \pm 45$ and $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{LAE}} = 7.4 \pm 3.6$. These flux density ratios correspond, respectively, to relative LyC escape fractions of $f_{\text{esc, rel}}^{\text{LBG}} = 5\%–8\%$ and $f_{\text{esc, rel}}^{\text{LAE}} = 18\%–49\%$, absolute LyC escape fractions of $f_{\text{esc}}^{\text{LBG}} = 1\%–2\%$ and $f_{\text{esc}}^{\text{LAE}} = 5\%–15\%$, and a comoving LyC emissivity from star-forming galaxies of $8.8–15.0 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$. In order to study the differential properties of galaxies with and without LyC detections, we analyze narrowband Ly α imaging and rest-frame near-infrared imaging, finding that while LAEs with LyC detections have lower Ly α equivalent widths on average, there is no substantial difference in the rest-frame near-infrared colors of LBGs or LAEs with and without LyC detections. These preliminary results are consistent with an orientation-dependent model where LyC emission escapes through cleared paths in a patchy interstellar medium.

Key words: cosmology: observations – diffuse radiation – galaxies: high-redshift – intergalactic medium

Online-only material: color figures

1. INTRODUCTION

One of the foremost goals in the study of cosmological reionization is determining the sources of the ionizing photons. Quasars (QSOs), while able to maintain an ionized universe from $z \sim 0–2$ (Cowie et al. 2009), rapidly fall in number density at redshifts greater than $z \sim 2$ (Hopkins et al. 2007). Recent studies of the faint-end slope of the QSO luminosity function indicate that while the QSO contribution to the overall ionizing budget at $z \gtrsim 4$ may not be negligible, QSOs alone are still unable to sufficiently account for all of the ionizing radiation during the epoch of reionization (Fontanot et al. 2012; Glikman et al. 2011; Siana et al. 2008). It is generally assumed that star-forming galaxies fill in the remainder of the gap in the ionizing budget. Therefore, studying the ionizing Lyman continuum (LyC) properties of high-redshift star-forming galaxies can provide vital information about the evolution of the intergalactic medium (IGM). Because the IGM at $z \gtrsim 6$ is opaque to LyC photons, we cannot directly observe the ionizing radiation at the redshifts corresponding to reionization. In order to directly measure the ionizing radiation escaping galaxies and study the galactic properties that give rise to significant LyC escape fractions, we must locate and study lower-redshift analogs to the star-forming galaxies that reionized the universe.

There have been several recent studies of LyC emission at both low and intermediate redshifts. To probe $0 < z < 2$, space-based far-UV observing facilities such as *Hubble Space Telescope* (HST)/STIS, GALEX, and FUSE have searched for

LyC radiation from star-forming galaxies (Leitherer et al. 1995; Malkan et al. 2003; Siana et al. 2007, 2010; Grimes et al. 2007, 2009; Cowie et al. 2009; Bridge et al. 2010) only to obtain null results (but see, e.g., Leitner et al. 2013). At $z \sim 3$, ground-based optical studies have shown that roughly 10% of star-forming galaxies have a moderately high escape fraction of ionizing radiation ($f_{\text{esc}} > 0.2$; Shapley et al. 2006; Iwata et al. 2009; Nestor et al. 2011, 2013). At even higher redshift ($z \sim 4$), Vanzella et al. (2012) find only 1 LyC emitter out of 102 Lyman break galaxies (LBGs), although this small number of LyC detections might reasonably be attributed to the rapidly increasing IGM opacity rather than to processes internal to galaxies. Siana et al. (2010) investigate the conspicuous lack of LyC-emitting galaxies at low redshift, where LyC transmission through the IGM is high. Given the similar stellar populations for UV-luminous galaxies at $z \sim 1.3$ and $z \sim 3$, Siana et al. (2010) infer that LyC production does not change with redshift, but the mechanism governing LyC escape must vary.

One key observational method used to measure LyC flux is deep imaging through a narrowband filter tuned to wavelengths just blueward of the Lyman limit. Narrowband imaging provides a very effective way to simultaneously probe the LyC of large samples of galaxies at the same redshift. We have designed a narrowband imaging program to study the LyC properties of galaxies in the HS1549+1919 field. This field was observed as part of a larger survey of UV-selected star-forming galaxies at $z \sim 2–3$ (Steidel et al. 2003, 2004, 2011; Reddy et al. 2008) and contains a galaxy protocluster with a redshift-space overdensity

Figure 1. Redshift histogram of galaxies in the HS1549 field. The yellow curve represents the overall redshift selection function of the LBG survey (Steidel et al. 2003, 2004), normalized to the observed number of galaxies in HS1549. The overdensity of objects at $z = 2.85$ is indicated by the dashed line, representing the redshift of the HS1549 protocluster.

(A color version of this figure is available in the online journal.)

of $\delta_{\text{gal}} \sim 5$ at $z = 2.85 \pm 0.03$ (Figure 1). The “spike” redshift coincides with that of a hyperluminous QSO (Trainor & Steidel 2012). More than 350 UV-selected galaxies have been identified in the HS1549 field, ~ 160 of which have been spectroscopically confirmed at $1.5 \leq z \leq 3.5$. Additionally, narrowband imaging with a 4670 Å filter tuned to the wavelength of Ly α at the redshift spike has revealed ~ 300 potential Ly α emitters (LAEs) and several Ly α “blobs” (Steidel et al. 2000, 2011). Such a large sample of star-forming galaxies at approximately the same redshift greatly facilitates the systematic narrowband imaging search for leaking LyC emission.

Our work in the HS1549 field parallels that of Nestor et al. (2011, 2013) and Iwata et al. (2009), who investigated another high-redshift protocluster (SSA22a; $z = 3.09$). Several questions emerged from these initial narrowband LyC studies, including the nature of galaxies with large offsets (occasionally reaching several kiloparsecs) between the centroids of their nonionizing UV-continuum and LyC emission. While significant offsets have been predicted in some simulations modeling the escape of ionizing photons (e.g., Gnedin et al. 2008), in practice it is difficult to determine whether the observed offsets provide information about the interstellar medium (ISM) of a LyC-leaking galaxy or are simply the result of a foreground contaminant (Vanzella et al. 2010b, 2012). Another key result presented in the SSA22a studies is the high apparent ratio of escaping ionizing to nonionizing UV radiation measured for several LAEs. For some objects, this ratio exceeded unity (Nestor et al. 2011; Inoue et al. 2011). If such measurements are free of foreground contamination, they are at odds with standard models for the intrinsic spectral energy distribution (SED) of star-forming galaxies (Bruzual & Charlot 2003). Consequently, as discussed by Vanzella et al. (2012), the most critical goal for LyC studies is to minimize the possibility that candidate LyC-leaking galaxies are contaminated by low-redshift interlopers. While previous $z \sim 3$ LyC studies have been plagued by small samples and lack of spectroscopic redshifts, in this work we present a sample of 131 spectroscopically confirmed galaxies

(49 LBGs and 91 LAEs, 9 of which are constituents of both samples). While our seeing-limited, ground-based imaging makes it difficult to distinguish individual cases of foreground contamination, we have performed simulations to model the amount of expected contamination in our samples as a whole (as in Nestor et al. 2011, 2013).

In this work, we build upon previous high-redshift LyC studies by considering a large spectroscopic sample of galaxies in an independent field from the SSA22a observations. We present imaging for 12 new $z \sim 2.85$ galaxies with putative LyC detections and correct for foreground contamination and IGM absorption. For our LBG and LAE samples, we calculate the escape fraction of ionizing photons both in an absolute sense and relative to the escape fraction of nonionizing UV photons. We also explore the differential multiwavelength properties of objects with and without leaking LyC radiation with regard to their Ly α equivalent widths (EWs), rest-frame near-infrared photometry, and stellar populations. The article is organized as follows. In Section 2, we describe our observations, data reduction, and galaxy sample. In Section 3, we present the LyC and broadband photometric measurements and error analysis. Section 4 contains a discussion of the complexities involved in identifying foreground contaminants and our methods for correcting measured LyC magnitudes for both foreground contamination and IGM absorption. In Section 5, we discuss the colors of individual LyC-detected objects, the average properties of the LBG and LAE samples, and the implied LyC escape fraction and comoving emissivity. We present the multiwavelength properties of our targets in Section 6 and summarize our results in Section 7. Throughout the article we employ the AB magnitude system and assume a cosmology with $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$, and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$. At $z = 2.85$, 1'' corresponds to 7.8 proper kpc.

2. OBSERVATIONS AND DATA REDUCTION

2.1. Photometric Observations

For observations of LyC emission, we used a custom narrowband filter manufactured by Barr Associates with central wavelength 3420 Å and FWHM 105 Å. This filter (hereafter referred to as NB3420) is designed to probe the LyC region blueward of 912 Å for galaxies at $z \geq 2.82$ such that no contaminating flux ($\leq 0.4\%$) from the galaxy’s nonionizing UV continuum is transmitted. The filter lies well within one LyC mean free path for galaxies at $z = 2.85$, minimizing the effect of intervening Lyman Limit systems and allowing for a more direct probe of the LyC absorption properties internal to galaxies. At a redshift of 2.85, current estimates place the LyC mean free path at ~ 100 proper Mpc (Rudie et al. 2013; Faucher-Giguère et al. 2008; Songaila & Cowie 2010), which corresponds to a rest-frame wavelength interval of $\sim 830\text{--}912$ Å. At $z = 2.85$, the range of rest-frame wavelengths probed by the NB3420 filter is 872–904 Å. Figure 2 shows the wavelengths probed by the NB3420 filter with respect to a typical LBG spectrum at $z = 2.85$, along with the locations of the NB4670 and V filters relevant to the identification of LAEs and the photometry of the nonionizing UV spectral region.

Our imaging was centered on the HS1549 field, at $(\alpha, \delta) = (15:51:53.7, +19:10:42.3)$. Observations were taken through the NB3420 filter using the blue side of the Low Resolution Imaging Spectrometer (LRIS; Oke et al. 1995; Steidel et al. 2004) on the Keck I telescope. During observing runs on 2009 June 23–24, 2010 July 11–12, and 2010 August 9, we acquired a total of

Figure 2. Narrowband filter transmission curves overlaid on a composite LBG spectrum (Shapley et al. 2006) redshifted to $z = 2.85$. The NB3420 filter (shown in blue) is located just blueward of the Lyman limit at $z = 2.85$ and probes LyC emission for galaxies at $z \geq 2.82$. The V filter (shown in green) probes the nonionizing UV continuum (rest-frame ~ 1400 Å) of galaxies at $z \sim 2.85$. We note that the V image was taken with the d500 dichroic, which blocks light blueward of 5000 Å, slightly truncating the transmission curve shown in this figure. The NB4670 filter (shown in red) is centered on the Ly α emission line for objects at $z \sim 2.85$ and was used along with the V band to select the sample of LAEs. Dashed lines indicate the locations of the Lyman limit and Ly α feature at $z = 2.85$.

(A color version of this figure is available in the online journal.)

19.2 hr of data comprising thirty-seven 1800 s exposures, one 1500 s exposure, and one 900 s exposure. We dithered the telescope between exposures to minimize the effects of bad pixels and to cover the gap between the two LRIS-B CCDs. Because these two $2\text{ K} \times 4\text{ K}$ detectors have slightly different quantum efficiencies below 4000 Å, we obtained half of the exposures at a sky position angle of 0° and half at 180° in order to even out systematics between the two chips. Conditions were photometric during all observing runs, and the effective seeing FWHM in the final stacked NB3420 image is $0''.7$ with a 1σ surface brightness limit of $29.4\text{ mag arcsec}^{-2.4}$.

Additional data in the HS1549 field includes broadband ground-based optical (U , G , V , $\mathcal{R}$) and near-IR (J , K) imaging, *Spitzer* IRAC and MIPS ($24\text{ }\mu\text{m}$) photometry, along with *HST*/WFC3 $F160W$ and $F475W$ imaging and morphologies for a small portion of the field. At $z = 2.85$, the ground-based V and $\mathcal{R}$ images both probe the nonionizing UV continuum in the vicinity of 1500 Å and do not suffer from contamination by the Ly α emission line or Ly α forest line blanketing. Given the significant redshift spike in the HS1549 field, narrowband 4670 Å imaging (NB4670; $\lambda_{\text{eff}} = 4667$ Å, FWHM = 88 Å) probing Ly α at $z = 2.85$ was also obtained. We used the combination of NB4670 and broadband V imaging to select LAEs near the redshift of the HS1549 protocluster (Steidel et al. 2011) and the NB4670–V continuum subtracted image to examine the morphology of the Ly α emission. Replacing the V-band image with G also yields information about LAEs,⁵ so we experimented with selecting LAEs using NB4670– G colors and examined the Ly α morphology in the NB4670– G image as well. Table 1 gives a summary of the imaging in the HS1549 field most relevant to the LyC observations. Detailed photometry and error analysis are performed on the NB3420, V, and $\mathcal{R}$ images, as described in Section 3.

⁴ The surface brightness limit quoted is a rough estimation made by evaluating sky counts in $1''.0$ apertures placed randomly on blank areas of the image. We conduct a more detailed study of the NB3420 image properties and our photometric accuracy using Monte Carlo simulations, as described in Section 3.2.

⁵ While the G -band image is contaminated by emission from the Ly α feature, the V-band image is offset from Ly α in wavelength and thus may include a color term if the continuum is not flat in f_ν .

Table 1
Keck/LRIS Imaging Observations

Filter ^a	λ_{eff} (Å)	Seeing FWHM ($''$)	Exposure (s)
NB3420	3420	0.7	69000
NB4670	4670	0.7	18000
V	5506	1.0	10800
$\mathcal{R}$	6830	1.3	4800

Note. ^a The NB3420 and NB4670 imaging were taken on LRIS-B, while the V and $\mathcal{R}$ imaging were taken on LRIS-R. Additionally, the V-band transmission is slightly affected by the use of the d500 dichroic.

2.2. NB3420 Imaging Reduction

We used standard IRAF image reduction procedures to reduce individual narrowband exposures and create the final NB3420 stacked image. For each exposure, we divided the image by a flatfield constructed from images of the twilight sky, subtracted the background, and masked out cosmic rays and satellite trails. In order to stack the 39 NB3420 frames, we registered >500 objects in the NB3420 frames with their counterparts in the astrometrically corrected $\mathcal{R}$ -band image and resampled the pixels to the $\mathcal{R}$ -band plate scale of $0''.2119\text{ pixel}^{-1}$. Accurate image registration is necessary for creating a spatial map of the relative positions and morphology of escaping LyC and nonionizing UV continuum radiation and identifying likely low-redshift contaminants.

In the process of combining the individual exposures, which ranged in airmass from 1.00 to 1.71, we scaled the flux in each frame such that each exposure was effectively observed at the minimum airmass. Narrowband images were calibrated onto the AB magnitude system using observations of spectrophotometric standard stars from the list of Massey et al. (1988). These AB magnitudes were also corrected for a Galactic extinction of $E(B - V) = 0.045$, based on *IRAS* $100\text{ }\mu\text{m}$ cirrus emission maps from Schlegel et al. (1998) and the extinction law from O'Donnell (1994). The final stacked image has an area of 35.5 arcmin^2 .

2.3. Spectroscopic Observations, Data Reduction, and Analysis

Previous spectroscopy has already been performed in the HS1549 field, resulting in spectroscopic redshifts for a significant number of LBGs and LAEs (Section 2.4). In order to augment the existing spectroscopic sample and confirm the redshifts of potential LyC-leaking galaxies, we obtained additional spectra, favoring objects with NB3420 detections. Another of our original intentions was to acquire deep spectroscopy in the LyC region of galaxies with NB3420 detections, but we were limited by poor weather conditions.

We performed multi-object spectroscopy in 2011 May on the Keck 1 telescope, using the blue side of LRIS. We observed four slitmasks with exposure times of 16,560, 9000, 8400, and 8100 s, respectively. For all masks, we used the 400 line mm^{-1} grism blazed at 3400 Å, achieving a spectral resolution of $R = 800$ for $1''.2$ slits. The “d500” dichroic beam splitter was used for the first mask (originally designed for deep LyC spectroscopy) and the “d560” dichroic was used for the three additional masks (designed to acquire redshifts). The conditions during the observing run were suboptimal, with intermittent clouds and a seeing FWHM of $0''.7\text{--}1''.0$ during clear spells.

When designing the slitmasks, we targeted both LBGs and LAEs with NB3420 detections. Slits were centered on the

coordinates of the *V* (NB4670) centroid for LBGs (LAEs). While most LAEs were selected using the *V*-band image as the continuum band, a small fraction (20%) were selected using the *G*-band (henceforth referred to as GNBs). Overall, we observed 46 objects on the four slitmasks, 29 of which had repeat observations.

Standard IRAF tasks were used to cut up the multi-object slitmask images into individual slitlets, flatfield the spectra using twilight sky flats, mask out cosmic rays, subtract the sky background, and average individual exposures to make stacked 2D spectra. These spectra were then extracted to one dimension, wavelength calibrated, and shifted into the vacuum frame. Details of these spectroscopic reduction procedures are discussed in Steidel et al. (2003). The centroid of the Ly α emission feature ($\lambda = 1215.67$) was used to estimate redshifts for LAEs. Both Ly α emission and interstellar absorption lines (when detected) were used for LBGs, yielding separate emission and absorption redshifts when both types of features were detected. For objects with spectra taken on multiple masks, we averaged these spectra in order to determine redshifts in cases when doing so increased the signal-to-noise ratio (S/N).

We successfully measured redshifts for 9 of the 12 LBG candidates, confirming three objects (MD5, M16, and D24) to be at $z > 2.82$ and identifying one as a star. We also measured redshifts for 11 out of 27 LAEs, thus providing us with five new LAE candidate LyC emitters. Of the remaining six LAE candidates for which we measured redshifts, two already had redshifts determined from previous spectroscopic studies of Steidel et al. (*lae4680*, *lae7832*; these objects were on the mask that was designed to directly detect LyC emission spectroscopically), one object had an NB3420 detection that was flagged as contamination after the spectroscopy was completed (*lae3208*), one object (*lae6856*) was at $z = 2.807$ (slightly too low redshift for NB3420 filter to probe uncontaminated LyC emission), and two objects (*lae4152*, $z = 2.447$; *lae5165*, $z = 1.873$) were at too low a redshift to be members of the HS1549 protocluster. Out of the seven GNBs, we acquired three redshifts (GNB2861, GNB4769, GNB5270), all of which placed the objects at $z \sim 2.85$, i.e., in the correct redshift range for our study. We note that in cases where no redshift could be measured, we did not draw conclusions about the quality of the object; the poor weather conditions, combined with the faintness of our targets, made it impossible to remove objects from our sample on the basis of a nondetection.

In summary, the analysis of the spectra allowed us to confirm $z \geq 2.82$ redshifts for three LBGs, five LAEs, and three GNBs with NB3420 detections. With these new redshifts, we were able to include the three additional LBGs in our LyC analysis. However, all five of the LAEs for which we confirmed redshifts have $m_{4670} > 26.0$. Because we do not have a complete and unbiased spectroscopic sample of LAEs with $m_{4670} > 26.0$ or GNBs (see Section 2.4), we did not include these objects in the LyC analysis. In the Appendix A.1, we present their postage stamp images and uncorrected NB3420 magnitudes in Figure 13 and Table 8, respectively.

2.4. Sample

Our initial sample consisted of 363 UV-selected galaxies and 289 narrowband-color selected LAEs. The UV-selected galaxies were identified using the *UGR* color selection criteria discussed in Steidel et al. (2003, 2004). Spectroscopic redshifts of 145 of these UV-selected galaxies were previously measured in this field (Reddy et al. 2008; Steidel et al. 2011) and our

follow-up spectroscopy (described in Section 2.3) yielded an additional eight redshifts. For the purposes of studying LyC emission, we kept only galaxies that were spectroscopically confirmed to be at $z \geq 2.82$ with no active galactic nucleus (AGN) signatures in their spectra. This sample consists of 49 LBGs with $z_{\text{spec}} \geq 2.82$.⁶ We note that while the spectroscopic studies of Reddy et al. (2008) and Steidel et al. (2011) were conducted without reference to the LyC properties of the LBGs, our follow-up spectroscopy was aimed at confirming redshifts of LBGs with NB3420 detections. While the addition of these objects may introduce a slight bias in the average nonionizing to ionizing UV flux density ratio for the full LBG sample, it allows us to study the individual flux density ratios of a larger number of LyC-emitting LBGs and perform a more useful differential analysis of the stellar populations of LBGs with and without LyC emission.

LAEs were selected via a broadband filter (*V*) and a narrowband filter (NB4670) designed to probe the Ly α emission line for galaxies in the range of $2.803 < z < 2.876$. In previous work (e.g., Steidel et al. 2000; Nestor et al. 2011), a broadband minus narrowband color excess of 0.7 mag was used to identify LAEs, corresponding to an EW threshold of 80 Å (corresponding to rest-frame 20 Å). Our HS1549 photometric LAE sample, however, is comprised of objects with $V - \text{NB4670} > 0.6$, a slightly lower threshold designed to increase the sample size by including LAEs with observed Ly α EWs slightly less than 80 Å.⁷ Spectroscopic follow-up of 116 of these narrowband excess objects confirmed 99 to be LAEs at the redshift of the protocluster (R. F. Trainor et al., in preparation). Of the 17 objects not confirmed to be at $z \sim 2.85$, 2 were found to be low-redshift galaxies (at $z = 1.983$ and $z = 2.7773$) and the other 15 did not yield spectroscopic redshifts. With the exclusion of Ly α blobs, galaxies with evidence for AGN emission in their spectra, and two LAEs that lie just below the redshift limit of the NB3420 filter ($z < 2.82$), the spectroscopically confirmed LAE sample consists of 91 galaxies. The spectroscopic follow-up of LAEs was conducted independently of LyC observations; thus, the sample of LAEs with spectroscopic redshifts is unbiased with respect to LyC properties.

We also created an additional photometric LAE sample that includes 33 photometric candidates whose NB4670 magnitudes are in the same range as the LAEs with spectroscopy ($m_{4670} \leq 26$), and include photometry and postage-stamp images for these LAE photometric candidates in Appendix A.2. The lack of spectroscopic confirmation of these photometric LAE candidates creates a small likelihood of contamination due to lower-redshift galaxies. The first source of contamination arises from the fact that the NB4670 filter selects for Ly α emission at a range that extends down to $z \sim 2.80$, while the NB3420 filter only measures uncontaminated LyC emission for galaxies at $z \geq 2.82$. However, only 2 out of 99 LAEs confirmed by spectroscopy has $z_{\text{spec}} < 2.82$; these objects (at $z = 2.811$ and $z = 2.801$) lie in the tail of the redshift overdensity

⁶ We note that one of the 49 spectroscopically confirmed LBGs (D27) has an absorption redshift of 2.814 and a Ly α emission redshift of 2.816. Although it is at a lower redshift than our conservative cutoff of $z = 2.82$, less than 1% of the flux redward of the Lyman limit is transmitted through the NB3420 filter at $z = 2.814$. Accordingly, we include it in the sample of spectroscopically confirmed LBGs.

⁷ Three of the LAE candidates in the HS1549 field identified on the basis of an early generation of NB4670 and *V*-band photometry by Steidel et al. have been shown by subsequent deeper NB4670 data to have colors not quite red enough to satisfy $V - \text{NB4670} > 0.6$. However, as these objects have already been confirmed spectroscopically to be at $z \sim 2.85$, we include them in the LAE sample regardless of $V - \text{NB4670}$ color.

centered at $z = 2.85$. With the assumption that the redshift distribution of the LAEs without spectroscopy matches the distribution of those with spectroscopy, contamination from additional galaxies in the low-redshift tail should be negligible. The second potential source of contamination arises from [O II] emitters at $z \sim 0.24\text{--}0.26$ whose emission lines fall within the NB4670 filter bandpass, but this type of contamination is also unlikely. Not only is the volume probed at $z \sim 2.80\text{--}2.88$ 40 times larger than that probed at $z \sim 0.24\text{--}0.26$, but the photometrically measured EWs of the LAE candidates, while typical for LAEs, would be considered exceptionally large if the objects were in fact low-redshift [O II] emitters (see, e.g., Hogg et al. 1998). Additionally, as described in Section 3.4, the relative spatial positions of the NB3420, V, and NB4670-V emission point toward the LAE candidates being high-redshift objects. However, because of the increased likelihood of contaminated NB3420 detections within the photometric LAE sample, we do not include these objects in the LAE analysis.

In summary, the samples of LBGs and LAEs with $z_{\text{spec}} \geq 2.82$ consist of 49 and 91 galaxies, respectively. There is some overlap in the final LBG and LAE samples: nine LBGs are also LAEs. We exclude one of these overlap objects (MD12/lae3540) from our analysis of global LyC properties because of its complex morphology, but discuss its photometric properties in detail in Section 4.1. With the exclusion of MD12/lae3540, the final LBG and LAE samples consist of 48 and 90 galaxies, respectively.

3. DETECTING LyC EMISSION

3.1. Photometric Measurements

Source identification and photometry were performed using SExtractor (Bertin & Arnouts 1996). Objects were detected independently in all images to allow for different spatial distributions of light at different wavelengths. For $\mathcal{R}$ -band and V-band images, we ran SExtractor in single-image mode using a Gaussian smoothing kernel with a FWHM of 2 pixels and a detection threshold of 1.0 standard deviation above the local smoothed background. For the NB3420 image, we ran SExtractor in dual-image mode with separate detection and measurement images, following the methodology of Nestor et al. (2011). For detection, we smoothed the NB3420 image by a Gaussian kernel with $\text{FWHM} = 2.35$ pixels and used a detection threshold of 2.25 standard deviations. For measurement, we used the unsmoothed NB3420 image. In the case of the NB3420 photometry, SExtractor parameters were chosen to reflect the higher S/N of the NB3420 image, to produce object number counts similar to those in Vanzella et al. (2010b), and to be complete in the faint magnitudes of interest for studying LyC emission. In all images, magnitudes were computed using “Kron-like” elliptical apertures (i.e., MAG_AUTO in SExtractor).

Because there has been some contention in this field concerning the most appropriate way to photometrically measure the flux density ratio between ionizing (NB3420) and nonionizing (V) UV emission (Vanzella et al. 2012; Nestor et al. 2011, 2013), we explain our methods here. Galaxies at $z \sim 3$ are known to exhibit clumpy morphologies (see, e.g., Law et al. 2007). As LyC photons may not escape isotropically from all portions of the galaxy, it is possible that LyC emission will be observed emanating from only one clump of a given galaxy. With the seeing-limited resolution of our images, it is often impossible to distinguish individual star-forming clumps; thus LyC emission may appear to be offset from the centroid of the V-band

emission. The idea of offset LyC emission is also supported by results from galaxy formation simulations. Simulations by Gnedin et al. (2008) claim that increased star-formation efficiencies in the more luminous LBGs result in higher scale heights of young stars relative to the H I in the disk. Therefore, minor interactions can disrupt the H I sufficiently to reveal the young stars and allow significant emission of ionizing photons to escape the galaxy. With the proper orientation, this configuration can be recognized as LyC emission offset from the primary UV-continuum emission. Alternatively, Ricotti (2002) suggests that the formation of globular clusters may have reionized the universe, since luminous OB associations in the outer haloes may have escape fractions approaching unity. Globular clusters are still forming at $z \sim 3$ (Stetson et al. 1996), and would be observable in our data as LyC emission offset spatially from the primary galaxy. Because these models support the possibility that the spatial distribution of LyC emission may be offset from V-band emission or have different morphology (e.g., compact versus diffuse), we must choose photometric methods that can account for such scenarios.

While Vanzella et al. (2012) suggest that flux density ratios should be measured using both LyC and nonionizing UV emission *only* from regions of escaping LyC emission, we argue here that the entire nonionizing and ionizing UV flux should be measured. Measuring the flux density ratio within an individual LyC-emitting region (defined by the isophote of the LyC emission) may provide useful information about the stellar populations of that region, but it does not provide information about the average LyC escape fraction among star-forming galaxies or the global LyC emissivity produced by these galaxies. The average escape fraction must be computed from the nonionizing to ionizing flux density ratio of an ensemble of $z \sim 3$ galaxies, including both galaxies without observed LyC emission and the non-LyC-emitting regions from galaxies with LyC emission. The calculation of the ionizing emissivity, moreover, relies entirely on the $z \sim 3$ UV luminosity function constructed by measuring nonionizing UV light from *entire* galaxies, not from their constituent clumps. Alternatively, we find that a method in which LyC is measured solely within an isophote defined by the galaxy’s nonionizing UV also does not accomplish our photometric goals. As this “isophotal” method does not allow for LyC emission to have a different spatial distribution from that of the nonionizing UV, it will result in the systematic loss of LyC flux. A differential analysis of LAE magnitudes using this isophotal photometric method versus the “Kron-like” apertures to measure V-band and NB3420 magnitudes separately (i.e., the method presented here) corroborates the idea that LyC emission is systematically missed in the isophotal method; the average LAE nonionizing to ionizing flux density ratio using the isophotal method is a factor of ~ 1.3 larger than the methods we use. While a factor of 1.3 is well within the 1σ uncertainty,⁸ it is an unnecessary loss of NB3420 flux. Therefore, we maintain that using the large “Kron-like” apertures to measure the full V-band and NB3420 magnitudes is the correct photometric method for the measurement of the global LyC escape fraction and emissivity.

3.2. Characterizing Photometric Uncertainties

We characterized both the statistical and systematic photometric uncertainties by running Monte Carlo simulations

⁸ The two photometric methods agree so well because the NB3420 detections in the LAE sample have very small offsets from the corresponding V-band detections.

Table 2
Uncertainties in Simulated Photometry for $\mathcal{R}$, V , and NB3420

Magnitude Bin ^a	$\Delta\mathcal{R}^b$	$\sigma_{\mathcal{R}}^{+c}$	$\sigma_{\mathcal{R}}^{-c}$	ΔV^b	σ_V^{+c}	σ_V^{-c}	ΔNB^b	σ_{NB}^{+c}	σ_{NB}^{-c}
22.5–23.0	0.00	0.03	0.03	0.00	0.02	0.02	0.00	0.02	0.02
23.0–23.5	0.00	0.05	0.05	0.00	0.03	0.03	0.00	0.03	0.03
23.5–24.0	0.00	0.08	0.07	0.00	0.05	0.05	0.00	0.05	0.04
24.0–24.5	0.00	0.12	0.11	0.00	0.07	0.06	0.00	0.07	0.07
24.5–25.0	0.00	0.17	0.15	0.00	0.10	0.09	0.00	0.11	0.10
25.0–25.5	0.00	0.25	0.21	0.00	0.14	0.13	0.00	0.17	0.14
25.5–26.0	0.01	0.34	0.26	0.00	0.22	0.18	0.01	0.25	0.20
26.0–26.5	0.06	0.45	0.31	0.01	0.32	0.25	0.03	0.34	0.26
26.5–27.0	0.17 ^d	0.54	0.36	0.04	0.38	0.28	0.09	0.44	0.31
27.0–27.5	0.37 ^d	0.71	0.43	0.13	0.51	0.34	0.24 ^d	0.60	0.38
27.5–28.0	0.70 ^d	1.22	0.56	0.25 ^d	0.60	0.38	0.40 ^d	0.79	0.45

Notes.

^a Recovered magnitude range of simulated galaxies. The statistics in this table are calculated for an object whose flux lies in the center of the bin; for example, this flux corresponds to a magnitude of 22.72 for the brightest bin.

^b Average value of the difference between recovered and input magnitudes. In the fainter bins, significant departures from zero imply systematic biases in the photometry.

^c The standard deviation in the difference between recovered and input magnitudes, used to calculate photometric uncertainties. Note that uncertainties in magnitudes are asymmetric because the standard deviation is calculated from the simulated flux distribution (Section 3.2).

^d Magnitudes of objects with systematic biases greater than one third of the object’s uncertainty are adjusted to reflect the systematic bias, as discussed in Section 3.2.

designed to reproduce our photometric measurement procedures. In the simulations, we added to each image fake galaxies of known magnitude and a range of radial profiles representative of the observed galaxies. In each iteration, 100 fake galaxies were placed in random, empty positions chosen to avoid image edges and photometric blending with existing objects or previously placed fake galaxies. We ran SExtractor on each filter ($\mathcal{R}$, V , NB3420) using the same parameters as for the actual photometry. The process of adding fake galaxies was repeated until we recovered 50,000 in each simulation, enough to populate all magnitude bins of interest with a statistically significant number of objects. The bins of recovered magnitude are 0.5 mag wide and span a range of $22.5 \leq m \leq 28$ in each band to encompass the magnitude range of the observed LBGs and LAEs.

Table 2 lists the systematic bias and photometric uncertainty associated with each magnitude bin, as derived from the results of the simulations. In each bin, the systematic bias is defined to be the average difference between input and recovered fluxes of fake galaxies, while the photometric uncertainty is defined to be the standard deviation of the distribution of differences between input and recovered fluxes. The uncertainty associated with an object of a given flux is a local quadratic interpolation between the uncertainties of neighboring bins. Because the errors in flux are Gaussian, all errors quoted in magnitudes are double-sided. Using the uncertainties estimated by our simulations, we define a 2σ photometric limit that corresponds to a magnitude limit of 27.33 in $\mathcal{R}$, 27.58 in V , and 27.30 in NB3420. In all filters, the simulations show that SExtractor misses a higher percentage of flux from fainter objects, resulting in a systematic offset in magnitude that is larger for fainter objects. In order to correct for this bias, we added the systematic offset determined from the simulation analysis to the flux of an object if the offset was greater than one third of the object’s associated uncertainty. This threshold was chosen in order to avoid the unnecessary addition of noise into our measurements by adding systematic offsets

much smaller than the 1σ error. In practice, this correction only affected a few of the fainter bins (see Table 2).

3.3. Object Matching

As we have two broadband filters (V and $\mathcal{R}$) that probe the nonionizing rest-frame UV ($\sim 1500 \text{ \AA}$) for galaxies at $z = 2.85$, it is necessary to choose one to represent the nonionizing UV flux. The spatial distributions of the V and $\mathcal{R}$ detections are well matched and their centroids agree within $0''.2$. The deeper V -band image has $1''.0$ seeing, while the $\mathcal{R}$ image has $1''.3$ seeing, larger photometric errors (see Table 2), and an elongated point-spread function (PSF). Although the $\mathcal{R}$ -band image was originally used to identify LBGs, we adopt the V -band image for our nonionizing UV photometry based on its superior quality and the fact that more LAEs are detected in V than in $\mathcal{R}$. We note that all but one of the LBGs with LyC detections is detected in V ; M4 does not have a V detection because it falls off the edge of the V -band image. For this object, we calculate a V magnitude by linearly interpolating between its measured $\mathcal{R}$ magnitude and a G magnitude computed from its $G - \mathcal{R}$ color obtained from the parent LBG survey (Reddy et al. 2008; Steidel et al. 2004).

The coordinates of LBGs are defined by V -band centroids, while the coordinates of LAEs are defined by NB4670 centroids. In order to find nonionizing rest-frame UV counterparts for LAEs, we compiled the positions and photometric measurements for any SExtractor detection in the V -band image within a 5 pixel ($1''.1 = 8.6 \text{ kpc}$ at $z = 2.85$) offset from the NB4670 coordinates of the target and successfully found counterparts for 77 out of 91 LAEs. Given that the largest galaxies reach sizes of $\sim 10 \text{ kpc}$ at $z \sim 2-3$ (Förster Schreiber et al. 2009; Law et al. 2012a) and that we may be looking for LyC emission in the outskirts of these galaxies, we used a similar matching technique with a larger radius (9 pixel = $1''.9 = 14.8 \text{ kpc}$ at $z = 2.85$) to find SExtractor detections in the NB3420 image near our LBGs and LAEs. This generous matching radius guarantees the

inclusion of all NB3420 detections potentially associated with the galaxies in the sample. Our candidate with the largest offset between the galaxy centroid and the associated NB3420 detection (M16) has $\Delta_{\text{Ly}\alpha, \text{LyC}} = 1''.26 = 9.8$ kpc, but the majority of our NB3420 detections (especially around LAEs) are at much smaller offset (< 3 kpc).

In order to remove any obvious false matches corresponding to neighboring sources, we visually inspected each NB3420 source matched with a known LBG or LAE in the $\mathcal{R}$, V , and G images (probing rest-frame nonionizing UV), and the $F160W$ image (probing rest-frame optical), if available. We also inspected LAEs with NB3420 matches in the continuum-subtracted NB4670– V image, which probes the spatial distribution of $\text{Ly}\alpha$ emission. NB3420 matches were removed if they corresponded spatially to a visible counterpart in another image. We also removed matches where the LyC emission was further than $1''.9$ away from the nonionizing UV-continuum emission, which occasionally occurred because the LAE matching radius was centered around the $\text{Ly}\alpha$ emission, not the nonionizing UV. Altogether, we retained matches to 4 LBGs and 7 LAEs, labeling the NB3420 detections corresponding to 4 LBGs and 16 LAEs as invalid.

3.4. Targets with NB3420 Detections

We report NB3420 detections in 4 out of 48 LBGs and 7 out of 90 LAEs. MD12/*lae3540* also has an NB3420 detection and is discussed in Section 4.1. Although they are not the focus of our study, we note that two of the four AGNs in the field at $z \geq 2.82$ have NB3420 detections. Tables 3 and 4 display photometric information for the NB3420-detected LBGs and LAEs. Postage stamp images of targets with NB3420 detections are shown in Figures 3 and 4, indicating the NB3420 (LyC), NB4670– V ($\text{Ly}\alpha$), and V (nonionizing UV continuum) morphologies.

As indicated by studies of the SSA22a field (Nestor et al. 2013, 2011; Inoue et al. 2011; Iwata et al. 2009) and implied by the simulations of Gnedin et al. (2008) and Ricotti (2002) (see Section 3.1), the centroids of the UV-continuum and LyC emission may not always coincide. The NB3420 versus V -band offsets for the four LBGs with NB3420 detections span a large range: $0''.34$, $0''.85$, $1''.06$, and $1''.26$. The offsets for the seven LAEs with NB3420 detections, however, are $\leq 0''.36 = 2.8$ kpc for all but one LAE (*lae6662*), which has an offset of $1''.09$. The smaller offsets of the NB3420 detections around LAEs strengthen the argument that these detections are truly associated with the LAEs in question, and also may indicate that LAEs are more compact galaxies than LBGs. For LAEs, we may also examine the offset between $\text{Ly}\alpha$ emission and LyC emission ($\Delta_{\text{Ly}\alpha, \text{LyC}}$), which is on average larger than both the offset between the UV-continuum and LyC ($\Delta_{\text{UV}, \text{LyC}}$) and the offset between UV-continuum and $\text{Ly}\alpha$ ($\Delta_{\text{UV}, \text{Ly}\alpha}$). While some spatial discrepancy between $\text{Ly}\alpha$ and LyC emission is to be expected due to the resonant scattering of $\text{Ly}\alpha$ photons, large offsets point to an increased probability of contamination and weaken the argument that the NB3420 detection is LyC emission associated with the high-redshift galaxy (Nestor et al. 2011, 2013; Vanzella et al. 2010b, 2012). However, while it is a good rule of thumb to assume that objects with $\Delta_{\text{UV}, \text{LyC}} \gtrsim 1''$ are likely contaminants, Nestor et al. (2013) demonstrated with high-resolution spectroscopy in the SSA22a field individual cases where putative LyC detections at large offset (for example, $\Delta_{\text{UV}, \text{LyC}} = 1''.0$) were found to be associated with the high-redshift galaxy, and detections at small offset (for example, $\Delta_{\text{UV}, \text{LyC}} = 0''.3$) were found to be contaminants.

Figure 3. $10'' \times 10''$ postage stamp images of the 4 LBGs with NB3420 detections. Each object is displayed in two bands: NB3420 (indicating the LyC) and V (indicating the nonionizing UV continuum). All postage stamps are centered on the V -band centroid, and blue circles ($1''$ radius) indicate the centroid of the NB3420 emission. All postage stamps follow the conventional orientation, with north up and east to the left.

(A color version of this figure is available in the online journal.)

The fact that $\Delta_{\text{Ly}\alpha, \text{LyC}}$ is often greater than $\Delta_{\text{UV}, \text{LyC}}$ provides additional evidence against significant contamination from low-redshift [O II] emitters in the LAE sample, corroborating the conclusion drawn in Section 2.4. Figure 5 compares the offsets $\Delta_{\text{Ly}\alpha, \text{LyC}}$ and $\Delta_{\text{UV}, \text{LyC}}$ for LAEs, demonstrating the tendency for NB3420 detections to be more closely associated with V rather than NB4670 detections. If the LAEs were low-redshift [O II] emitters rather than high-redshift LAEs, the NB3420 image would probe the rest-frame UV, the NB4670– V image would probe [O II] emission, and the V image would probe the rest-frame optical. In this case, both the NB3420 and NB4670– V images would probe active star formation, so a smaller offset would be expected between the centroids of the detections in these two images. However, larger values of $\Delta_{\text{Ly}\alpha, \text{LyC}}$ support the opposite interpretation; namely, the LAE candidates are truly high-redshift objects whose LyC and $\text{Ly}\alpha$ emission correlate more strongly with the nonionizing UV-continuum than with each other. In Tables 3 and 4, we list the offsets $\Delta_{\text{UV}, \text{LyC}}$ and $\Delta_{\text{Ly}\alpha, \text{LyC}}$ for galaxies with NB3420 detections.

In addition to exploring the properties of individual objects leaking LyC radiation, we created stacked NB3420 and V images of subsamples of our targets to examine the average LyC emission properties and attempt a LyC measurement in the stacks of objects without NB3420 detections. These subsamples include LBGs, LAEs, and LAEs in bins of V magnitude. For each subsample, we also created two additional stacks comprising objects with and without NB3420 detections,

Table 3
LBG Photometry

ID	R.A. (J2000)	Decl. (J2000)	z_{em}^a	z_{abs}^b	V	NB3420	$\Delta_{\text{UV, LyC}}^c$	$\frac{F_{\text{UV}}}{F_{\text{LyC}}}_{\text{obs}}^d$
BX84	15:51:53.696	19:12:24.64	...	2.823	24.41	>27.30	...	>14.3
MD5	15:51:45.211	19:11:05.13	3.146	3.139	24.96	26.89	0''3	5.9 ± 2.0
MD7	15:51:47.521	19:10:13.56	2.857	2.851	25.11	>27.30	...	>7.5
MD9	15:51:50.711	19:09:38.20	2.852	2.843	24.13	>27.30	...	>18.6
MD12 ^e	15:51:51.882	19:10:41.16	2.856	2.849	24.56	26.74	1''0	7.5 ± 2.5
MD34	15:52:06.315	19:12:48.60	...	2.849	24.23	26.56	0''8	8.6 ± 2.7
C1	15:51:39.653	19:10:40.67	2.845	...	24.77	>27.30	...	>10.3
C2	15:51:40.960	19:13:12.77	3.100	...	23.76	>27.30	...	>26.1
C4	15:51:44.413	19:11:24.75	2.863	2.857	25.01	>27.30	...	>8.2
C5	15:51:44.627	19:10:59.73	3.173	...	25.14	>27.30	...	>7.3
C6	15:51:44.760	19:10:32.84	...	2.828	24.36	>27.30	...	>15.0
C7	15:51:45.247	19:12:13.21	...	2.841	24.53	>27.30	...	>12.8
C8	15:51:45.386	19:08:49.84	...	2.935	25.03	>27.30	...	>8.1
C9	15:51:46.707	19:11:52.39	2.925	2.919	23.67	>27.30	...	>28.2
C10	15:51:48.424	19:09:24.93	3.193	3.183	24.56	>27.30	...	>12.4
C12	15:51:49.360	19:09:52.55	...	2.835	24.93	>27.30	...	>8.8
C13	15:51:49.685	19:10:58.10	2.843	...	24.92	>27.30	...	>9.0
C14	15:51:50.616	19:09:18.46	...	2.841	24.87	>27.30	...	>9.4
C15	15:51:51.352	19:10:19.50	2.849	2.849	25.39	>27.30	...	>5.8
C17	15:51:55.283	19:12:19.48	2.941	2.934	24.92	>27.30	...	>9.0
C19	15:52:00.196	19:10:08.73	3.166	3.158	24.41	>27.30	...	>14.3
C20	15:52:00.402	19:08:40.75	3.115	...	25.09	>27.30	...	>7.6
C22	15:52:03.833	19:09:43.21	...	2.960	24.65	>27.30	...	>11.5
C24	15:52:05.618	19:13:11.73	2.834	2.826	25.24	>27.30	...	>6.7
C25	15:52:06.069	19:11:28.37	3.159	...	24.60	>27.30	...	>12.0
C27	15:52:07.041	19:12:19.29	2.931	2.922	24.32	>27.30	...	>15.5
D3	15:51:43.712	19:09:12.37	2.942	2.934	24.06	>27.30	...	>19.7
D4	15:51:43.976	19:11:39.67	2.863	2.856	24.43	>27.30	...	>14.1
D6	15:51:45.191	19:09:05.31	2.849	2.840	24.64	>27.30	...	>11.5
D7	15:51:46.246	19:09:50.11	2.943	2.932	24.51	>27.30	...	>13.0
D11	15:51:49.764	19:09:02.94	...	2.837	23.72	>27.30	...	>27.1
D13	15:51:51.724	19:10:15.89	2.852	2.842	24.39	>27.30	...	>14.5
D14	15:51:53.262	19:11:01.00	2.851	2.851	24.87	>27.30	...	>9.3
D16	15:51:54.848	19:11:31.18	3.139	3.130	24.00	>27.30	...	>20.8
D17	15:51:57.435	19:11:02.56	2.841	2.825	25.02	>27.30	...	>8.1
D18	15:51:59.695	19:09:39.25	2.850	...	24.60	>27.30	...	>12.1
D19	15:52:00.270	19:09:40.75	2.847	2.844	25.28	>27.30	...	>6.4
D20	15:52:00.484	19:10:27.55	...	2.825	24.11	>27.30	...	>18.8
D23	15:52:03.743	19:09:24.47	2.902	2.893	24.49	>27.30	...	>13.3
D24	15:52:05.278	19:09:45.17	2.951	2.942	24.24	27.01	1''1	12.8 ± 4.0
D25	15:52:07.999	19:08:55.80	...	2.825	24.79	>27.30	...	>10.1
D27	15:52:08.314	19:09:48.82	2.816	2.814	23.95	>27.30	...	>21.8
M2	15:51:41.355	19:10:06.11	2.875	2.867	25.32	>27.30	...	>6.2
M5	15:51:41.970	19:08:22.21	2.936	2.931	25.00	>27.30	...	>8.3
M6	15:51:43.678	19:09:44.58	2.891	...	23.77	>27.30	...	>25.9
M16	15:51:53.636	19:09:29.49	2.955	2.953	25.28	26.56	1''3	3.2 ± 1.4
M21	15:52:01.356	19:13:00.78	2.834	...	25.14	>27.30	...	>7.3
M22	15:52:02.705	19:09:40.06	3.159	3.149	24.77	>27.30	...	>10.3
M23	15:52:05.748	19:12:08.72	...	3.409	24.85	>27.30	...	>9.6

Notes.^a Emission redshift of Ly α .^b Interstellar absorption redshift.^c Spatial offset between the centroids of V-band and NB3420 flux-densities.^d Observed ratio and uncertainty in the ratio of nonionizing UV to LyC emission, inferred from the NB3420–V color. This value has not been corrected for either contamination by foreground sources or IGM absorption.^e MD12 is not included in the LBG sample.

respectively. Stacked images were made by averaging postage stamps of individual galaxies centered on the coordinates of either their V (LBGs) or NB4670 (LAEs) detections. Pixels contaminated by nearby objects were excluded from the stacks by creating a mask from the object detection isophotes in the SExtractor segmentation image. Stacked image photometry was

performed using the IRAF PHOT routine using a 1''9 aperture, corresponding to the matching radius used to find detections. Although the postage stamps for individual galaxies were previously background subtracted as part of the data reduction process, we found that a second pass of background subtraction on the stacked image was necessary to remove remaining sky

Table 4
LAE Photometry

ID	R.A. ^a (J2000)	Decl. ^a (J2000)	z	NB4670	V	NB3420	$\Delta_{UV, LyC}^b$	$\Delta_{Ly\alpha, LyC}^c$	$\frac{F_{UV}}{F_{LyC}^{obs}}^d$
lae32	15:51:38.692	19:10:04.89	2.846	25.32	26.78	>27.30	...	...	>1.6
lae274	15:51:39.406	19:09:42.66	2.851	24.41	25.47	>27.30	...	...	>5.4
lae367	15:51:39.993	19:09:00.13	2.882	25.45	25.96	>27.30	...	...	>3.4
lae413	15:51:39.673	19:13:15.43	2.846	24.47	27.21	>27.30	...	...	>1.1
lae447	15:51:41.908	19:10:16.77	2.844	25.93	27.24	>27.30	...	...	>1.1
lae576	15:51:40.601	19:10:58.09	2.846	25.17	26.24	>27.30	...	...	>2.6
lae599	15:51:40.758	19:11:00.26	2.834	25.07	>27.58	>27.30	...	...	...
lae618	15:51:39.812	19:10:53.58	2.847	24.07	27.58	>27.30	...	...	>0.8
lae633	15:51:39.624	19:11:04.99	2.852	23.91	26.78	>27.30	...	...	>1.6
lae661	15:51:41.061	19:11:54.91	2.836	25.43	26.31	>27.30	...	...	>2.5
lae1012	15:51:42.495	19:10:34.26	2.839	24.21	25.69	>27.30	...	...	>4.4
lae1206	15:51:43.025	19:10:39.98	2.863	24.78	26.61	>27.30	...	...	>1.9
lae1261	15:51:43.596	19:11:57.22	2.851	25.70	26.99	>27.30	...	...	>1.3
lae1359	15:51:43.780	19:11:40.20	2.861	25.02	26.49	>27.30	...	...	>2.1
lae1375	15:51:43.969	19:11:02.19	2.864	25.83	27.20	>27.30	...	...	>1.1
lae1500 ^e	15:51:43.955	19:11:39.63	2.863	23.93	24.43	>27.30	...	...	>14.1
lae1528	15:51:44.682	19:09:24.59	2.846	25.66	>27.58	>27.30	...	...	...
lae1540	15:51:44.490	19:11:47.22	2.845	24.90	26.88	>27.30	...	...	>1.5
lae1552	15:51:44.709	19:08:44.92	2.838	25.09	>27.58	>27.30	...	...	...
lae1599 ^f	15:51:44.424	19:11:24.64	2.863	23.96	25.01	>27.30	...	...	>8.2
lae1610	15:51:44.873	19:11:02.12	2.840	24.77	26.54	>27.30	...	...	>2.0
lae1679	15:51:45.164	19:12:34.40	2.842	25.40	26.34	>27.30	...	...	>2.4
lae1751	15:51:45.449	19:11:11.71	2.840	25.50	27.12	>27.30	...	...	>1.2
lae1765	15:51:45.323	19:09:29.49	2.847	24.04	26.20	>27.30	...	...	>2.7
lae1774	15:51:53.129	19:10:57.03	2.852	23.20	25.04	>27.30	...	...	>8.0
lae1787	15:51:45.646	19:11:28.60	2.836	25.98	26.63	>27.30	...	...	>1.8
lae1803	15:51:45.697	19:11:53.03	2.852	25.55	>27.58	>27.30	...	...	...
lae1835	15:51:45.882	19:09:45.11	2.836	25.45	27.50	>27.30	...	...	>0.8
lae1843	15:51:45.676	19:09:58.21	2.847	25.81	27.18	>27.30	...	...	>1.1
lae2015	15:51:46.504	19:12:38.62	2.845	25.81	>27.58	>27.30	...	...	...
lae2063	15:51:45.966	19:08:22.13	2.849	24.49	25.91	>27.30	...	...	>3.6
lae2158	15:51:47.005	19:11:02.98	2.847	25.34	26.11	>27.30	...	...	>3.0
lae2174	15:51:47.001	19:08:22.06	2.885	25.89	26.69	>27.30	...	...	>1.8
lae2183	15:51:47.235	19:09:53.38	2.869	25.76	27.41	>27.30	...	...	>0.9
lae2292	15:51:47.635	19:10:00.50	2.851	25.86	27.19	27.11	0''3	0''6	0.9 ± 0.5
lae2306 ^g	15:51:47.509	19:10:13.34	2.857	24.60	25.11	>27.30	...	...	>7.5
lae2358	15:51:47.945	19:09:03.21	2.841	25.31	>27.58	>27.30	...	...	...
lae2436	15:52:03.231	19:12:52.43	2.832	23.44	24.11	25.21	0''3	0''3	2.7 ± 0.4
lae2489	15:51:48.187	19:08:30.15	2.845	25.24	27.05	>27.30	...	...	>1.3
lae2551	15:51:48.539	19:12:03.11	2.849	25.59	27.08	>27.30	...	...	>1.2
lae2561	15:51:48.346	19:13:13.88	2.828	24.71	26.13	>27.30	...	...	>2.9
lae2668	15:51:49.100	19:11:22.54	2.841	24.73	27.19	>27.30	...	...	>1.1
lae2747	15:51:49.311	19:08:44.33	2.842	24.77	26.84	>27.30	...	...	>1.5
lae2796	15:51:49.579	19:10:41.28	2.844	25.66	26.69	>27.30	...	...	>1.8
lae2854	15:51:49.793	19:12:47.89	2.836	25.84	>27.58	>27.30	...	...	...
lae2856	15:51:49.718	19:10:49.05	2.843	25.15	25.85	>27.30	...	...	>3.8
lae2949	15:51:49.870	19:10:53.62	2.840	25.14	27.45	>27.30	...	...	>0.9
lae2966	15:51:49.995	19:10:41.59	2.841	24.61	26.87	26.82	0''1	0''7	1.0 ± 0.4
lae2984 ^h	15:51:49.695	19:10:57.98	2.843	23.83	24.92	>27.30	...	...	>9.0
lae3167	15:51:50.781	19:10:34.29	2.846	25.62	26.89	>27.30	...	...	>1.5
lae3208	15:51:50.908	19:11:16.18	2.844	25.28	27.32	>27.30	...	...	>1.0
lae3339 ⁱ	15:51:51.354	19:10:19.71	2.849	24.69	25.39	>27.30	...	...	>5.8
lae3354	15:51:51.525	19:10:47.02	2.843	24.70	27.29	>27.30	...	...	>1.0
lae3540 ^j	15:51:51.878	19:10:41.09	2.856	23.93	24.56	26.74	1''0	0''9	7.5 ± 2.5
lae3763	15:51:51.530	19:10:58.22	2.842	23.27	26.60	>27.30	...	...	>1.9
lae3798	15:51:53.127	19:10:34.46	2.852	25.47	27.01	>27.30	...	...	>1.3
lae3808	15:51:52.858	19:11:41.05	2.839	24.03	25.75	>27.30	...	...	>4.2
lae3866	15:51:52.742	19:11:39.09	2.839	23.32	25.42	>27.30	...	...	>5.7
lae3922	15:51:53.457	19:11:43.89	2.856	24.63	25.97	>27.30	...	...	>3.4
lae4147	15:51:54.158	19:11:05.14	2.841	24.70	27.12	>27.30	...	...	>1.2
lae4366	15:51:54.775	19:11:06.42	2.844	24.44	25.75	>27.30	...	...	>4.2
lae4680	15:51:56.206	19:09:56.72	2.848	25.24	27.07	27.05	0''1	0''3	1.0 ± 0.4
lae4684	15:51:56.167	19:11:56.65	2.849	25.50	26.15	>27.30	...	...	>2.9
lae4730	15:51:56.413	19:10:42.03	2.843	25.46	26.12	>27.30	...	...	>3.0
lae4796	15:51:56.367	19:11:06.15	2.868	24.68	26.30	>27.30	...	...	>2.5

Table 4
(Continued)

ID	R.A. ^a (J2000)	Decl. ^a (J2000)	z	NB4670	V	NB3420	$\Delta_{UV, LyC}^b$	$\Delta_{Ly\alpha, LyC}^c$	$\frac{F_{UV}}{F_{LyC}}^d$ obs
lae4804	15:51:56.621	19:12:26.86	2.860	25.49	26.99	>27.30	...	...	>1.3
lae4882	15:51:56.873	19:13:06.72	2.835	25.23	27.35	>27.30	...	...	>1.0
lae4947	15:51:57.132	19:08:48.24	2.857	24.95	25.49	>27.30	...	...	>5.3
lae5132	15:52:01.057	19:09:47.89	2.849	24.76	27.18	>27.30	...	...	>1.1
lae5193	15:52:00.760	19:10:53.36	2.853	25.26	26.21	>27.30	...	...	>2.7
lae5322	15:52:00.102	19:10:17.14	2.832	23.53	25.44	>27.30	...	...	>5.5
lae5371 ^k	15:52:00.265	19:09:40.41	2.848	24.54	25.28	>27.30	...	...	>6.4
lae5458 ^l	15:51:59.696	19:09:39.30	2.849	24.04	24.60	>27.30	...	...	>12.1
lae5470	15:51:59.582	19:11:40.17	2.843	25.85	>27.58	>27.30	...	...	...
lae5713	15:51:58.999	19:09:20.60	2.857	25.57	>27.58	>27.30	...	...	...
lae5720	15:51:58.755	19:13:00.43	2.833	25.56	>27.58	>27.30	...	...	...
lae5740	15:51:59.315	19:10:35.10	2.831	24.96	25.90	>27.30	...	...	>3.6
lae5900	15:51:58.057	19:11:21.84	2.845	24.24	25.42	>27.30	...	...	>5.6
lae5995 ^m	15:51:57.456	19:11:02.63	2.833	24.54	25.02	>27.30	...	...	>8.1
lae6193	15:52:07.992	19:11:23.26	2.836	25.79	26.10	>27.30	...	...	>3.0
lae6312	15:52:07.610	19:08:47.48	2.827	25.72	26.43	>27.30	...	...	>2.2
lae6662	15:52:06.357	19:10:42.74	2.833	25.16	25.94	27.23	1''1	1''2	3.3 $\pm$ 1.3
lae6774	15:52:06.090	19:11:36.81	2.845	25.80	>27.58	>27.30	...	...	...
lae6979	15:52:04.940	19:09:53.42	2.863	25.63	26.25	>27.30	...	...	>2.6
lae7180	15:52:04.662	19:11:42.19	2.930	25.86	26.38	25.85	0''3	0''1	0.6 $\pm$ 0.2
lae7542	15:52:03.575	19:12:50.73	2.841	24.28	25.48	>27.30	...	...	>5.4
lae7577	15:52:03.189	19:09:09.08	2.828	25.35	>27.58	>27.30	...	...	...
lae7803	15:52:02.304	19:09:03.58	2.865	25.97	26.89	>27.30	...	...	>1.4
lae7830	15:52:02.146	19:09:55.06	2.826	23.77	>27.58	>27.30	...	...	...
lae7832	15:52:02.201	19:10:48.59	2.829	24.23	24.83	25.11	0''4	0''7	1.3 $\pm$ 0.2
lae7893	15:52:01.861	19:12:49.92	2.858	25.32	>27.58	>27.30	...	...	...

Notes.^a Coordinates of LAEs are based on NB4670 centroids.^b Spatial offset between the centroids of V and NB3420 emission.^c Spatial offset between the centroids of NB4670 and NB3420 emission.^d Observed ratio and uncertainty in the ratio of nonionizing UV to LyC flux-densities, inferred from the NB3420– V color. This value has not been corrected for either contamination by foreground sources or IGM absorption.^e D4^f C4^g D7^h C13ⁱ C15^j MD12. This object is not included in either the LAE or LBG samples.^k D19^l D18^m D17

subtraction systematics. Bootstrap resampling was employed for each stack in order to include both sample variance and photometric error in the calculation of stack uncertainties. The results of the stacking analysis are presented in Table 5. None of the NB3420 stacks of individual LBGs and LAEs without LyC detections exhibited any significant flux; the NB3420 stack of LBG nondetections reached a 3σ limiting magnitude of 28.62 and the NB3420 stack of LAE nondetections reached a 3σ limiting magnitude of 28.97. These limits reflect only photometric errors. We note that for the NB3420 stacks of all LBGs and all LAEs, the addition of noise from a large number of galaxies undetected in NB3420 overpowered the signal from the few detected galaxies. As the resulting signal in the stack was detected at less than 3σ , we quote a lower limit in magnitude in Table 5.

4. ACCOUNTING FOR CONTAMINATION

In narrowband imaging studies of $z \sim 3$ LyC emission, it is critical to determine whether the detection in the narrowband

filter is actually high-redshift LyC emission or contaminating radiation from a lower-redshift object. In this section, we analyze the morphology of several candidate LyC emitters for which we have high-resolution *HST* imaging and discuss the complexities of identifying contaminants. Because we do not have high-resolution, multi-band imaging for all of our candidate LyC emitters, we also discuss statistical corrections applied to our samples in order to account for foreground contamination by low-redshift galaxies (which artificially boosts the LyC signal) and LyC absorption by the IGM (which decreases the observed LyC emission).

4.1. Morphology of NB3420 Detections

In order to cull potential LyC detections from obvious neighbors, we have visually inspected each NB3420 detection in all available bands (as discussed in Section 3.3). This process of visual inspection, however, is limited by the depth and resolution of the data. Within our dataset, the *HST*/WFC3 imaging has the best resolution, followed by the ground-based optical data from

Figure 4. $10'' \times 10''$ postage stamp images of the 7 LAEs with NB3420 detections. Each object is displayed in three bands: NB3420 (indicating the LyC), NB4670–V (indicating Ly α emission and labeled LyA), and V (indicating the nonionizing UV continuum). The V–NB4670 color of each LAE is indicated below the object name. All postage stamps are centered on the V-band centroid, and blue circles ($1''$ radius) indicate the centroid of the NB3420 emission. All postage stamps follow the conventional orientation, with north up and east to the left. We note that LAEs with more diffuse Ly α emission may be difficult to distinguish in the NB4670–V image, even though their V–NB4670 colors identify them as LAEs and their redshifts have been confirmed by spectroscopy; for such objects, we have increased the stretch of the NB4670–V image to make the diffuse emission more easily visible.

(A color version of this figure is available in the online journal.)

Keck. The WFC3/UVIS $F475W$ and WFC3/IR $F160W$ images have PSF FWHMs of $0''.08$ and $0''.18$, respectively (Law et al. 2012b), while the seeing FWHMs of the Keck data are larger ($0''.7$ – $1''.0$). Unfortunately, only one (two) of the targets in our samples with NB3420 matches fall within the small $2''.3 \times 2''.1$ ($2''.9 \times 2''.7$) field of view of the WFC3 IR (UVIS) imaging. Three additional objects with NB3420 detections (two in the photometric LAE sample and one LAE/LBG that we removed from both samples for its complicated morphology) also fall in the field of view of the WFC3 IR and UVIS images. Examining the morphologies of these objects across the ground-based and *HST* images highlights the difficulties associated with object matching.

Figure 6 shows four LAEs with *HST* imaging and NB3420 detections (the LAE photometric candidates *lae4070* and *lae2158*,

Figure 5. Comparison of the offsets between UV-continuum and LyC emission ($\Delta_{UV, LyC}$) and between Ly α and LyC emission ($\Delta_{Ly\alpha, LyC}$) for LAEs with NB3420 detections. On average, $\Delta_{UV, LyC} < \Delta_{Ly\alpha, LyC}$, reflecting the fact that Ly α photons are resonantly scattered; therefore, Ly α emission is unlikely to coincide exactly with the location of LyC emission. This observed trend supports the claim that the LAEs are not low-redshift [O II] emitters (see Section 3.4).

along with spectroscopically confirmed LAEs *lae2966* and *lae4680*) displayed in five filters: NB3420 (LyC), NB4670–V (Ly α), V (UV continuum), $F475W$ (rest-frame $1240 \text{ \AA}$), and $F160W$ (rest-frame $4000 \text{ \AA}$). *Lae4070* is an example of an object with a simple morphology. In all images, a single detection is visible, and the offsets between the centroids of the detection in each image are very small. Similarly, *lae4680* can be identified with a single source in all images except $F475W$, where the object breaks into two clumps located $0''.3$ apart. As this offset is very small, these clumps are likely part of the same system, although foreground contamination is still a possibility. For *lae4070* and *lae4680*, the only possibilities of contamination arise from either the small probability of [O II] emission being misidentified as Ly α emission or a lower-redshift foreground galaxy coincident with the LAE along the line of sight. We discussed the first possibility in Section 2.4 and concluded that it is small enough to ignore. The second possibility is quantified in Section 4.2 in the discussion of the contamination simulations. The analysis of the remaining two objects is more complex. In the case of *lae2966*, multiple clumps are visible in the V-band, $F475W$, and $F160W$ imaging, and the NB3420 emission is associated with only one clump. The spectroscopically confirmed Ly α emission is extended and not distinctly associated with any single clump. Therefore, it is unclear whether the NB3420 detection is LyC emission from one clump of a $z \sim 2.85$ galaxy or a low-redshift interloper along the line of sight. In the case of *lae2158*, the *HST* imaging can be used to identify the NB3420 detection as contamination. While it is not possible to determine in the V-band image whether the UV-continuum light associated with the LAE belongs to multiple objects or whether it is merely extended, there are two distinct galaxies visible in the $F475W$ image. Since the northeastern galaxy is associated with the NB3420 detection while the southwestern galaxy is associated with the Ly α emission, we classify this system as a case of contamination.

We also highlight one object whose complex morphology led us to remove it from our analysis completely. This object (shown in Figure 7, with multiple clumps indicated) was originally identified in our catalogs as both an LBG (MD12)

Figure 6. $10'' \times 10''$ postage stamp images of four LAEs with *HST* imaging and NB3420 detections (the LAE photometric candidates *lae4070* and *lae2158*, along with spectroscopically confirmed LAEs *lae2966* and *lae4680*). The postage stamps highlight the difficulties associated with interpreting the morphologies of these objects with ground-based resolution. Objects are displayed in up to five filters: NB3420 (indicating the LyC), NB4670–V (indicating Ly α emission and labeled LyA), V (indicating the nonionizing UV continuum), F475W (rest-frame ~ 1200 Å), and F160W (rest-frame ~ 4000 Å). All postage stamps are centered on the V-band centroid and blue circles ($1''$ radius) indicate the centroid of the NB3420 emission. All postage stamps follow the conventional orientation, with north up and east to the left.

(A color version of this figure is available in the online journal.)

Table 5
Photometry in Stacked Images

Sample	N_{gal}	V^a	NB3420 ^a	NB3420–V	$\langle F_{\text{UV}}/F_{\text{LyC}} \rangle_{\text{obs}}^b$
LBG, all	48	$24.59^{+0.09}_{-0.08}$	>28.67	>4.08	>42.7
LBG, detect	4	$24.44^{+0.73}_{-0.43}$	$27.05^{+0.63}_{-0.39}$	$2.60^{+0.55}_{-1.16}$	11.0 ± 7.2
LBG, nondetect	44	$24.60^{+0.09}_{-0.08}$	>28.62	>4.03	>40.8
LAE, all	90	$26.03^{+0.14}_{-0.13}$	>29.01	>2.98	>15.6
LAE, detect	7	$25.49^{+0.40}_{-0.29}$	$25.97^{+0.35}_{-0.26}$	$0.48^{+0.38}_{-0.58}$	1.6 ± 0.6
LAE, nondetect	83	$26.09^{+0.15}_{-0.13}$	>28.97	>2.88	>14.2
LAE, $24 < V < 25$	5	$24.67^{+0.09}_{-0.08}$	$26.51^{+2.01}_{-0.66}$	$1.84^{+0.67}_{-2.04}$	5.5 ± 4.6
LAE, $25 < V < 26$	21	$25.53^{+0.08}_{-0.07}$	>28.22	>2.69	>11.9
LAE, $26 < V < 27$	30	$26.31^{+0.07}_{-0.06}$	>28.41	>2.10	>6.9
LAE, $V > 27$	34	$27.31^{+0.18}_{-0.16}$	>28.48	>1.17	>2.9

Notes.

^a Uncertainties listed are 1σ and include both photometric error and sample variance. All photometric lower limits are 3σ .

^b Observed nonionizing UV to LyC flux density ratios and uncertainties, inferred from the NB3420–V color of each stacked subsample. These values have not been corrected for either contamination by foreground sources or IGM absorption.

and an LAE (*lae3540*), with both the LBG and LAE centroids coincident with the bright V-band detection (indicated by region A). Analysis of the Ly α morphology in both the NB4670–V and NB4670–G images reveals extended emission in both Ly α images. Although both Ly α images show nearly identical morphology for all other LAEs with NB3420 detections, the extended Ly α emission around MD12 appears clumpy in the NB4670–V image (associated with regions B, C, and D) and more diffuse in the NB4670–G image. Also, while the Ly α emission appears to extend down to region C where the NB3420 detection is located ($1''$ 0 to the south of the LBG centroid),

this location is also coincident with detections in the F475W and F160W images that are not visibly connected to the LBG centroid in region A. Unfortunately, due to the slit position of 74° , the LRIS spectrum of this object only provides us with information about region A (where we observe double-peaked Ly α emission in the spectrum) and misses the region associated with the NB3420 detection. We note that region E likely indicates a lower-redshift interloper along the line of sight, as its position does not coincide with Ly α emission and it is located at the large offset of $2''$ 1 from region A (the centroid of MD12).

Figure 7. $3''.2 \times 3''.2$ postage stamp images of the MD12/lae3540, an additional object with a potential LyC detection. Originally identified as both an LBG and LAE, MD12/lae3540 was removed from both samples due to its complex morphology (see Section 4.1). The object is displayed in six filters: NB3420 (indicating the LyC), NB4670–V and NB4670–G (both indicating Ly α emission), V (indicating the nonionizing UV continuum), F475W (rest-frame ~ 1200 Å), and F160W (rest-frame ~ 4000 Å). The clumpy morphology for MD12/lae3540 is indicated by five red circles labeled A–E. Region A corresponds to the LBG and LAE centroids for MD12/lae3540, while regions B, C, and D correspond to areas of Ly α emission indicated by the NB4670–V image. Region C also corresponds to the location of the NB3420 emission. Region E likely indicates a lower redshift galaxy along the line of sight. All postage stamps follow the conventional orientation, with north up and east to the left.

(A color version of this figure is available in the online journal.)

The interpretation of the suite of imaging for MD12 is not straightforward, and it is impossible to fully understand the nature of these clumps with the current data. The clumps may comprise an extended, perhaps interacting, system, with all clumps at the same redshift. Alternatively, as region A has Ly α emission in its spectrum and regions B, C, and D are all associated with fairly compact Ly α emission in the NB4670–V image, these four clumps may simply be several protocluster members along the line of sight located at slightly different redshifts; in this case, they should be treated as separate LAEs. A final possibility (supported by the fact that the Ly α emission appears diffuse in the NB4670–G image) is that the extended Ly α emission actually originates in the central part of the MD12 system and is being projected over a large area; in this case, the clump with the NB3420 emission in region C may be an interloper along the line of sight. Because of the ambiguity regarding the nature of MD12 as a either single, complex system or the superposition of multiple galaxies, we exclude both MD12 and lae3540 from our LBG and LAE analysis. Nevertheless, this object represents another possible LyC-emitter.

Determining conclusively whether the NB3420 detections are contamination or true LyC emission requires imaging of each candidate galaxy with high enough resolution to discern the individual clumps and spectroscopy at high enough resolution that distinct spectra are obtained (and redshifts calculated) for each emitting region. This method has been recently implemented in the SSA22a field using *HST*/WFC3/F336W to acquire high-resolution imaging below the Lyman limit for three $z \sim 3.1$ LBGs. Additionally, near-infrared spectroscopy of rest-frame optical [O III] nebular emission was obtained using NIRSPEC on Keck II (Nestor et al. 2013; B. Siana et al., in preparation). The seeing of $\sim 0''.5$ for these near-infrared observations enabled the spatial separation of clumps. One of three LBGs was confirmed to have escaping LyC radiation, and two others showed evidence of low-redshift contamination. Unfortunately, a similar study for the HS1549 field is not possible using ground-based or current space-based instrumentation. There is no appropriate imaging filter on *HST* that probes the region just below the Lyman limit at $z = 2.85$ without some contamination redward of the limit. Furthermore, spectroscopy from the ground is unfeasible because the rest-frame optical features do not fall within windows of atmospheric transmission at $z = 2.85$. The next best option is to estimate photometric redshifts of each clump by acquiring high-resolution multi-band imaging. Such a technique has been successfully employed by Vanzella et al. (2012) for a sample of 19 LBGs with potential LyC emission at $3.4 \leq z \leq 4.5$ using multi-band *HST* imaging from the GOODS and CANDELS surveys.

4.2. Contamination from Foreground Galaxies

Although we currently do not have the high-resolution imaging or spectroscopic data required to prove the validity of each individual LyC detection, we can characterize the probability of foreground contamination statistically. With the assumption that all of our targets have been correctly identified as high-redshift galaxies, the NB3420 flux we measured can either be associated with LyC emission from the high-redshift galaxy itself or with contaminating radiation from a spatially coincident lower-redshift interloper. In order to statistically characterize the fraction of contaminated NB3420 detections in a given sample, we performed a simulation to calculate the expected number of uncontaminated LyC detections and the contamination-corrected average NB3420 magnitude for that sample. We summarize the simulation below, and further details are described in Section 5.1 of Nestor et al. (2013).

The contamination simulation was run separately for the LBGs and LAEs. Within the simulation, we considered as possible LyC detections all NB3420 SExtractor detections within $3''.5$ of our targets that fall within 1σ of the magnitude range of our reported LyC detections, corresponding to $26.19 \leq m_{3420} \leq 27.59$ for LBGs and $24.98 \leq m_{3420} \leq 27.70$ for LAEs. These SExtractor detections included both the detections considered to be candidates for LyC emission and detections previously identified by eye as belonging to interlopers visible in other wavelength bands. We computed the probability of each NB3420 detection being an interloper, which depends on both the global surface density of objects in the NB3420 image and the local surface density at the offset of the NB3420 detection from each LBG and LAE. In each iteration of the simulation, we used these probabilities to randomly determine whether or not each NB3420 detection was flagged as an interloper. Detections at large radial offset—where the local surface density approaches the global surface density—are more likely to be flagged as interlopers, as demonstrated in Figure 8. We removed the NB3420 magnitudes of objects flagged as interlopers before the average NB3420 magnitude was calculated, resulting in the contamination-corrected average NB3420 magnitude. The simulation was repeated 1000 times, and we recorded the average number of uncontaminated NB3420 detections and the average contamination-corrected NB3420 magnitude for both samples. Uncertainties in these average values were computed from the distribution of simulation results.

In Figure 8, we present histograms of the radial surface density of NB3420 detections around galaxies in the LBG and LAE samples, compared to the global surface density. The surface density of NB3420 detections is plotted using

Figure 8. Radial surface density of NB3420 detections within $3''.5$ of galaxies in our LBG and LAE samples (solid red histograms). The radial bin size for the LBG histogram is $0''.5$, and the bin size for the LAE histogram is $0''.4$. Red histogram error bars represent 1σ Poisson uncertainties in the number of NB3420 detections in each bin. The blue solid line represents the average surface density of NB3420 detections within the magnitude range ($\pm 1\sigma$) of the LAEs/LBGs, corresponding to the expected amount of contamination. Blue dotted lines indicate the 1σ uncertainties in the average surface density in each bin. Black hashed regions on the red histogram represent NB3420 detections that have been rejected by eye as contaminants and are not considered candidates for LyC emission. As described in Section 4.2, the surface density of NB3420 detections around LAEs is plotted using the NB3420 offsets calculated from LAE position defined by either a combination of NB4670 and V-band centroids (panel (b)) or NB4670 centroids only (panel (c)). In both cases, we observe a significant excess in surface density of NB3420 detections at small offset from LAEs, which indicates that many of the LAE NB3420 detections are likely uncontaminated LyC emission. We also note that panels (b) and (c) confirm that LyC emission is on average more closely associated with V-band emission than with NB4670, as the average offset in panel (b) is less than that in panel (c). As there are only four LBGs with NB3420 detections, the LBG excess over the average surface density at low spatial offset is much smaller.

(A color version of this figure is available in the online journal.)

NB3420 offsets computed relative to V-band (nonionizing UV continuum) for LBGs. For LAEs, offsets were computed relative to both V-band and NB4670 (a combination of Ly α and continuum emission in the vicinity of the Ly α wavelength). As discussed in Section 3.4, LyC emission is spatially more closely associated with the UV-continuum than with Ly α emission. Therefore, we adopt V-band coordinates to represent the LAE centroids, except in cases where the LAE is undetected in V and we must use NB4670 coordinates. Panels (b) and (c) of Figure 8, displaying V-band and NB4670 offsets, respectively, confirm that LyC emission is on average more closely associated with V-band emission than with NB4670, as the average offset in panel (b) is less than that in panel (c). All panels show an excess surface density at small offsets (although this excess is only statistically significant for the LAEs), suggesting that a large number of the NB3420 detections are physically associated with the LBG and LAE targets; the contamination simulation statistically quantifies this number. For the four LBGs detected in NB3420, the simulations predict on average 1.5 ± 1.0 uncontaminated detections. For the seven NB3420-detected LAEs, the simulations predict 4.3 ± 1.3 uncontaminated detections. For both samples, the simulation

yields the contamination-corrected average magnitude of the NB3420 detections. These values are used to compute the sample-averaged NB3420–V colors and flux density ratios presented in Table 6 and discussed in Section 5. We note that the removal of contaminants must increase the nonionizing UV to LyC flux density ratio for the ensembles containing the full sample of LBGs (or LAEs), as the total amount of LyC flux is decreased when a contaminant is removed while the total amount of nonionizing UV flux stays the same. For ensembles that contain only objects with NB3420 detections, however, the flux density ratio may not necessarily increase with the removal of the contaminants. In these cases, when an NB3420 detection is identified as a contaminant, both its LyC and nonionizing UV emission are omitted from the sample when calculating the flux density ratio as the object is no longer considered to have a true LyC detection.

4.3. Correction for IGM Absorption

In addition to correcting for contamination by low-redshift galaxies, we must also correct the NB3420 photometry for the absorption of LyC photons by neutral hydrogen in the IGM. In order to account for such attenuation, we ran a second set of simulations to determine the mean IGM correction factor and its associated uncertainty. We note that, as discussed in Nestor et al. (2011), our simulations do not take into account the possible environmental effects due to the presence of the $z \sim 2.85$ protocluster and the proximity to the hyperluminous QSO (Q1549), and it is not clear whether such effects would culminate in an increase or decrease of neutral hydrogen absorbers. We summarize the methods employed below, and further details are described in Nestor et al. (2011, 2013), the latter of which describes the current version of the methodology.

First, we constructed 500 model sightlines simulating the distribution of H I absorbers in the IGM. For each model sightline, absorbers were drawn randomly from their redshift and column density distributions (Rudie et al. 2013), spanning redshifts from $z = 1.7$ to the redshift of each object of interest. The model sightlines for the LAEs were created at the mean redshift of the LAEs ($z = 2.85$) since the narrow width of the NB4670 filter implies that the LAE targets lie within a small redshift range. Given that the LBGs span a wider range in redshift ($2.815 \leq z \leq 3.414$), we created a set of 500 model sightlines at the redshift of each LBG. We then calculated the mean transmission of each model sightline in the LyC region. For the purposes of these simulations, the LyC region consists of the fixed observed-frame bandpass of the NB3420 filter, taking into account the filter transmission profile. Figure 9 shows the probability distributions of the LyC transmission factor for model sightlines at $z = 2.85$ (representing typical LAE redshifts) and $z = 3.41$ (corresponding to our highest-redshift LBG with the most extreme case of attenuation).

The sample average transmission through the IGM ($\bar{t}_{\text{sample}}$) is equal to the mean transmission of the 500 $z = 2.85$ sightlines for the LAEs and the mean transmission of all model sightlines for the LBGs at different redshifts. The uncertainty in $\bar{t}_{\text{sample}}$ is estimated by first assuming an exponential distribution of unattenuated LyC flux (the parameters of this exponential function are fit to our data via a maximum likelihood method; see Nestor et al. 2013) and then creating 1000 realizations of our sample by randomly choosing for each galaxy an unattenuated LyC flux from our exponential distribution and an attenuation factor from one of the simulated model sightlines at the redshift of the galaxy. We set the uncertainty in $\bar{t}_{\text{sample}}$ equal to the

Table 6
Average UV to LyC Flux Density Ratios

Correction	LBGs		LAEs	
	$\langle \text{NB3420} \rangle - \langle V \rangle^a$	$\langle F_{\text{UV}}/F_{\text{LyC}} \rangle^b$	$\langle \text{NB3420} \rangle - \langle V \rangle^a$	$\langle F_{\text{UV}}/F_{\text{LyC}} \rangle^b$
Full Ensembles				
None	$4.91^{+0.47}_{-0.84}$	92 ± 49	$2.62^{+0.43}_{-0.72}$	11.2 ± 5.4
Contamination ^c	$5.93^{+0.47}_{-0.84}$	230 ± 130	$3.06^{+0.43}_{-0.72}$	16.8 ± 8.1
IGM + contamination ^d	$4.78^{+0.48}_{-0.88}$	82 ± 45	$2.17^{+0.43}_{-0.73}$	7.4 ± 3.6
Sources with NB3420 Detections Only				
None	$2.15^{+0.29}_{-0.39}$	7.2 ± 2.2	$0.54^{+0.46}_{-0.80}$	1.6 ± 0.9
Contamination ^c	$2.18^{+0.49}_{-0.91}$	7.5 ± 4.2	$0.50^{+0.56}_{-1.22}$	1.6 ± 1.1
IGM + contamination ^d	$0.86^{+0.60}_{-1.44}$	2.2 ± 1.6	$-0.39^{+0.59}_{-1.36}$	0.7 ± 0.5

Notes.

^a Color determined from average NB3420 and V-band fluxes. Uncertainties include photometric and sample uncertainties.

^b Ratio and uncertainty in nonionizing UV and LyC flux densities inferred from $\langle \text{NB3420} \rangle - \langle V \rangle$ color.

^c Color and flux density ratio after statistically correcting sample for foreground contamination of NB3420 fluxes. For the ensembles containing all LBGs (or LAEs), the removal of contaminants must increase the nonionizing UV to LyC flux density ratio as the total amount of LyC flux is decreased. For ensembles that contain only objects with NB3420 detections, however, the flux density ratio may not necessarily increase with the removal of the contaminants.

^d Color and flux density ratio after correcting sample for both foreground contamination and IGM absorption of NB3420 fluxes.

standard deviation of the 1000 simulated $\bar{t}_{\text{sample}}$ values. We find that $\bar{t}_{\text{sample,LAE}} = 0.44 \pm 0.03$ and $\bar{t}_{\text{sample,LBG}} = 0.35 \pm 0.04$. For both the LBG and LAE samples, we multiply the contamination-corrected nonionizing to ionizing flux density ratios by the transmission factor to obtain the IGM-and-contamination-corrected values presented in Table 6. Unlike the contamination correction discussed in Section 4.2, which decreases the average NB3420 flux, the IGM attenuation correction acts to increase it.

5. RESULTS

In order to study the amount of ionizing radiation escaping star-forming galaxies at high redshift, we have imaged a large sample of $z \sim 2.85$ galaxies in a narrowband filter designed to probe LyC emission. As described in Section 3.4, we have detected 4 out of 48 LBGs and 7 out of 90 LAEs in our NB3420 filter. After application of the contamination and IGM corrections discussed in Sections 4.2 and 4.3, the average NB3420– V properties of the galaxies in our sample provide information about the ratio of nonionizing to ionizing UV flux density and, with some assumptions, the LyC escape fraction.

5.1. NB3420 minus V Colors and Flux Density Ratios

In Figure 10, we plot NB3420– V color versus V magnitude for LBGs and LAEs with NB3420 detections. The colors of the LBG and LAE stacks are also plotted. All of the data presented in Figure 10 represent observed values, uncorrected for foreground galaxy contamination and IGM absorption. Qualitatively, the color–magnitude diagram (CMD) for objects in the HS1549 field agrees with the CMD for objects in the SSA22a field (Nestor et al. 2011); both samples of galaxies exhibit narrowband minus continuum colors across a wide range of continuum magnitudes that are extremely blue with respect to expectations from standard stellar population models. For example, Bruzual & Charlot (2003) models with solar metallicity and a constant star formation rate (SFR) predict the intrinsic ratio of nonionizing to ionizing radiation for LAEs to

be roughly 3–6, which corresponds to an NB3420– V color of 1.2–2. Out of the seven LAEs with NB3420 detections, only one LAE (*lae6662*) has a color redder than 1.2. Even more surprising, especially given the small offsets ($\Delta_{\text{UV,LyC}} \leq 0''.36 = 2.8$ kpc) of all LAEs except *lae6662*, four of the six remaining LAEs have NB3420– $V < 0$, implying a complete absence of a Lyman break (i.e., a flat spectrum). For typical LBG stellar populations (e.g., Shapley et al. 2001; Kornei et al. 2010), the intrinsic ratio is predicted to be >6 (corresponding to an NB3420– V color >2); however, two out of the four NB3420-detected LBGs have NB3420– $V < 2$. Thus, the HS1549 field provides the first confirmation that the blue nonionizing to ionizing UV colors observed in the SSA22a field are common among $z \sim 3$ galaxies and not simply an unusual property of the SSA22a field itself. Tables 3 and 4 display the raw nonionizing to ionizing flux density ratios ($F_{\text{UV}}/F_{\text{LyC}}$) for individual LBGs and LAEs with NB3420 detections, and Table 5 shows the same values for the stacks.

The average corrected and uncorrected NB3420– V colors and flux density ratios for the samples of LBGs and LAEs are presented in Table 6. Because photometry of the stacks of LBGs and LAEs without NB3420 detections does not yield any NB3420 signal, we determine average NB3420 and V magnitudes by calculating the mean flux of all the galaxies in each sample and assuming zero flux for undetected objects. These average magnitudes agree with stacked photometry within the errors. As our V and NB3420 filters probe rest-frame wavelengths of 1430 Å and 888 Å, respectively, we can convert our observed NB3420– V colors directly into nonionizing to ionizing flux density ratios (i.e., $F_{\text{UV}}/F_{\text{LyC}}$). We find uncorrected flux density ratios of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{obs}} = 92 \pm 49$ for LBGs and $(F_{\text{UV}}/F_{\text{LyC}})_{\text{obs}} = 11.2 \pm 5.4$ for LAEs. Considering only objects with NB3420 detections in each sample, we find $(F_{\text{UV}}/F_{\text{LyC}})_{\text{obs}} = 7.2 \pm 2.2$ for LBGs and $(F_{\text{UV}}/F_{\text{LyC}})_{\text{obs}} = 1.6 \pm 0.9$ for LAEs. In order to correct these colors and flux-density ratios for foreground contamination, we use adjusted NB3420 and V magnitudes derived from the contamination

Figure 9. The probability distribution of the LyC transmission factor through the NB3420 filter for an object at $z = 2.85$ (top) and an object at $z = 3.41$ (bottom). The mean transmission at each redshift is indicated by the dashed lines. The top panel represents the vast majority of LBGs and LAEs that lie near the redshift spike at $z = 2.85$. The wide range of possible transmission values at $z = 2.85$ reflects the variation in IGM absorption at $z = 2.85$, and the slight peak around zero transmission corresponds to sightlines that encounter a Lyman limit system. The bottom panel corresponds to the redshift of M23, an LBG with a redshift far removed from the spike, and represents the most extreme case of attenuation. This large amount of attenuation is due to a combination of two effects. Not only is the sightline through the IGM longer for the photons coming from $z = 3.41$ (thus providing each photon with more time to encounter an absorber), but the $z = 3.41$ IGM has a higher fraction of neutral hydrogen—and thus more absorbers—than the $z = 2.85$ IGM.

(A color version of this figure is available in the online journal.)

simulation described in Section 4.2. Correcting for IGM attenuation, however, only affects NB3420 magnitudes because rest-frame $\sim 1500 \text{ \AA}$ flux is not affected by H I absorption in the IGM. Thus, for the IGM correction we simply multiply the contamination-corrected flux density ratio by the IGM transmission factor. With the application of these two corrections to the colors and flux density ratios, we find $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 82 \pm 45$ for LBGs and $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 7.4 \pm 3.6$ for LAEs. For objects with NB3420 detections, we find corrected values of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 2.2 \pm 1.6$ for LBGs and $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 0.7 \pm 0.5$ for LAEs.

It has so far proven difficult to reconcile the observed low non-ionizing to ionizing flux density ratios with intrinsic luminosity density ratios predicted by current stellar-population synthesis models. Theoretical values of the intrinsic luminosity density ratio featured in previous works range from $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}} = 3\text{--}6$ (Steidel et al. 2001; Shapley et al. 2006; Siana et al. 2007). Many of the values in Table 6—and especially those that consider only the NB3420-detected galaxies—do not fall within the theoretical range. We also note that if we were to consider the flux density ratio in the region defined by the NB3420 isophote (thus only including V-band light in the vicinity of the NB3420 emission,

Figure 10. Color-magnitude diagram of observed NB3420- V vs. V magnitude, with the equivalent values for $F_{\text{UV}}/F_{\text{LyC}}$ indicated on the right-hand axis. LBGs are indicated by red squares and LAEs by blue circles. LBGs (LAEs) without NB3420 detections are plotted as 2σ lower limits, indicated by the red (blue) arrows. The typical V -band uncertainty associated with these lower limits is of similar size to the V -band uncertainty of data points at a given V magnitude. The red (blue) diamond indicates the 3σ limit in NB3420- V color for the stack of all LBGs (LAEs). All data represent observed values, uncorrected for foreground galaxy contamination and IGM absorption. We note that an NB3420- V color of zero ($F_{\text{UV}}/F_{\text{LyC}} = 1$) corresponds to a flat spectrum with no Lyman break. Such a blue spectrum cannot be reasonably explained by current stellar population synthesis models, and many of the LAEs on this plot lie uncomfortably close to or below this limit.

(A color version of this figure is available in the online journal.)

which more closely emulates a single stellar population), there would be even more tension between the measurements and the models. Nestor et al. (2013) examined a wider range of theoretical $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ values using stellar population synthesis models from Bruzual & Charlot (2003) and the BPASS models from Eldridge & Stanway (2009), which include a more detailed treatment of massive stellar binaries (specifically, Wolf-Rayet stars) and nebular emission. The models were used to describe galaxies with an array of ages and metallicities, assuming a constant SFR. The largest source of variation in $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ is due to the age of the galaxy, as $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ increases quickly as the stellar population evolves. The choice of model is also important, as the BPASS models predict ratios of $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ a factor of ~ 1.5 lower than the Bruzual & Charlot (2003) models for a given age and metallicity. Changes in stellar metallicity cause smaller variations in $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ on the order of a few percent, but the magnitude of these changes increases with the age of the stellar population. It is important to note that, if the models are correct, the predicted value for $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ should serve as a lower limit to the observed nonionizing to ionizing flux density ratio of the galaxy. The observed flux density ratio will likely be higher than this limit, as ionizing radiation may be absorbed by neutral hydrogen in both the ISM of the galaxy and the IGM. Thus, for any given galaxy, the observed value of $F_{\text{UV}}/F_{\text{LyC}}$ should be greater than or equal to $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$. We note that the two sets of stellar population synthesis models used in this analysis, as with all such models, are limited by the absence of direct observations in the LyC region of the O- and B-type stars that produce most of the ionizing radiation (Zastrow et al. 2013). Without direct observations to verify that the models are accurate, the fact that our LyC observations disagree with the models may indicate that there is something systematically incorrect with the model spectra at ultraviolet wavelengths. Alternatively, the tension between the models and

our observations may be due to the uncertainties involved in removing contaminants. Any low-redshift contaminants left in our sample (e.g., an [O II]-emitter or foreground galaxy) would also produce colors different from those predicted for models of $z \sim 2.85$ galaxies.

Here, we consider our inferred values of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ with respect to the model predictions of $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$. We first discuss the interpretation of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for samples of LBGs and LAEs with NB3420 detections because these are the samples for which LyC emission is actually measured. We consider the ensembles of NB3420-detected galaxies rather than individual galaxies because our statistical corrections for foreground contamination and IGM absorption do not apply to individual galaxies.

For the sample of LBGs with NB3420 detections, $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 2.2 \pm 1.6$. This value of 2.2 implies the unphysical ages of $\lesssim 10$ Myr using the Bruzual & Charlot (2003) and BPASS models. Adding the 1σ error to the LBG flux density ratio implies $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 3.8$, which is consistent with a 100 Myr galaxy using the BPASS model. Typical LBGs are characterized by ages > 100 Myr (Kornei et al. 2010; Shapley et al. 2005, 2001; although in Section 6.1 we show that this average age may be somewhat over-estimated, as it is based only on LBG samples with near- or mid-infrared detections) and not younger than ~ 50 Myr (given the typical LBG dynamical timescale; Reddy et al. 2012). It is worth noting that using an average IGM correction for the NB3420-detections-only samples probably constitutes an overcorrection, as NB3420-detected galaxies are likely to have clearer sightlines through the IGM than the sample as a whole. Thus, the true value of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for LBGs is likely somewhere between 2.2 (corrected for contamination and IGM absorption) and 7.5 (corrected for contamination only). Values $\gtrsim 4.5$ are consistent with galaxy ages > 100 Myr using BPASS models, and values $\gtrsim 6.1$ are consistent with the same limit in age using Bruzual & Charlot (2003) models.

As for the LAE sample, the value of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 0.7 \pm 0.5$ is inconsistent with all models presented in Nestor et al. (2013). For models with metallicity $Z = 0.2 Z_{\odot}$ and ages of 1, 10, and 100 Myr, respectively, $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}} = 1.33, 2.10$, and 3.16 for BPASS models and $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}} = 1.98, 3.59$, and 6.17 for Bruzual & Charlot (2003) models. Only by considering values of the LAE flux density ratio at $> 1\sigma$ from our measured value can we reconcile our flux density ratio with those predicted by young BPASS models with ages of several Myr. Because the dynamical timescale argument (based on galaxy sizes and velocity dispersions) that limits LBG ages to > 50 Myr may not apply to LAEs, it is possible that the extremely low nonionizing to ionizing flux density ratios we measure imply that LAEs are on average very young galaxies. Young ages for LAEs have been previously determined by SED fits of stacked photometry (e.g., Gawiser et al. 2007). We also note that, once again, the correction for IGM absorption may constitute an overcorrection, which would place the true value of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for LAEs between 0.7 (corrected for contamination and IGM absorption) and 1.6 (corrected for contamination only).

Until now, we have only discussed the interpretation of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for samples with NB3420 detections. We now consider the full samples of LBGs and LAEs, which include both objects with and without NB3420 detections. For the full sample of LBGs, $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 82 \pm 45$. This number is significantly higher than the lower limit required by stellar population synthesis models for typical LBGs with ages > 100 Myr, namely $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}} > 4.5$ for BPASS models and $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}} >$

6.1 for Bruzual & Charlot (2003) models. This consistency with the models, however, arises not because all 48 LBGs have values of $F_{\text{UV}}/F_{\text{LyC}} \sim 82$, which would not conflict with models, but rather because 4 LBGs have low values of $F_{\text{UV}}/F_{\text{LyC}}$ and the rest are undetected in LyC. For the full sample of LAEs, $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}} = 7.4 \pm 3.6$. As with the LBGs, these values would not significantly conflict with models if they represented the typical LAE in the sample. However, the LAE sample is also comprised of some objects with very strong LyC emission and some objects with no observed emission. For both LBGs and LAEs, we must consider how values of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for individual galaxies compare with stellar population synthesis models, as the ensemble-averaged value does not take into account variation within the sample.

5.2. The LyC Escape Fraction

Using the values of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for the full LBG and LAE samples presented in Table 6, we can determine the relative and absolute escape fractions for each sample. The relative escape fraction, a measure of how the observed flux density ratio $F_{\text{UV}}/F_{\text{LyC}}$ compares to the theoretical ratio, is defined to be

$$f_{\text{esc, rel}}^{\text{LyC}} = \frac{(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}}{(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}}, \quad (1)$$

where $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ is the intrinsic ratio of UV to LyC luminosity densities produced in star-forming regions. The absolute escape fraction includes an additional term for the escape fraction of nonionizing UV photons ($f_{\text{esc}}^{\text{UV}}$) due to dust extinction:

$$f_{\text{esc}}^{\text{LyC}} = f_{\text{esc, rel}}^{\text{LyC}} \times f_{\text{esc}}^{\text{UV}}. \quad (2)$$

Thus,

$$f_{\text{esc}}^{\text{LyC}} = \left(\frac{F_{\text{UV}}}{F_{\text{LyC}}} \right)_{\text{corr}}^{-1} \left(\frac{L_{\text{UV}}}{L_{\text{LyC}}} \right)_{\text{intr}} (f_{\text{esc}}^{\text{UV}}). \quad (3)$$

Both unknowns $f_{\text{esc}}^{\text{UV}}$ and $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$ are uncertain and likely to vary from object to object. Following Nestor et al. (2013), we estimate these values based on the observed ages and $E(B - V)$ values of $z \sim 3$ LBGs and LAEs. For the escape fraction of nonionizing UV photons, we adopt $f_{\text{esc, LBG}}^{\text{UV}} = 0.2$ and $f_{\text{esc, LAE}}^{\text{UV}} = 0.3$. For the intrinsic ratio of UV to LyC luminosity densities, we quote a range of values bracketed by the BPASS and Bruzual & Charlot (2003) models. For LBGs, we use 10^8 Myr, $Z = Z_{\odot}$ models to obtain $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr, LBG}} = 4.43\text{--}6.38$. For LAEs, which have been shown to be younger (Gawiser et al. 2007) and more metal-poor, we use $10^6\text{--}10^7$ Myr, $Z = 0.2 Z_{\odot}$ models to obtain $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr, LAE}} = 1.33\text{--}3.59$. With these assumptions, we derive relative LyC escape fractions of $f_{\text{esc, rel}}^{\text{LyC}} = 5\%\text{--}8\%$ and $f_{\text{esc, rel}}^{\text{LyC}} = 18\%\text{--}49\%$, and absolute escape fractions of $f_{\text{esc}}^{\text{LyC}} = 1\%\text{--}2\%$ and $f_{\text{esc}}^{\text{LyC}} = 5\%\text{--}15\%$. As our measured values of $F_{\text{UV}}/F_{\text{LyC}}$ have uncertainties of roughly fifty percent, the uncertainty in f_{esc} is at minimum fifty percent and likely higher due to uncertainties in our assumed values of $f_{\text{esc}}^{\text{UV}}$ and $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$.

Our values of f_{esc} are consistent with, though slightly lower than, those calculated by Nestor et al. (2013) for the SSA22a field using very similar methods: $f_{\text{esc, LBG}}^{\text{LyC}} = 5\%\text{--}7\%$, $f_{\text{esc, LAE}}^{\text{LyC}} =$

10%–30%. Using IGM-corrected values of $F_{\text{UV}}/F_{\text{LyC}}$ from the literature and our assumptions for $f_{\text{esc}}^{\text{UV}}$ and $(L_{\text{UV}}/L_{\text{LyC}})_{\text{intr}}$, we find additional values of f_{esc} in the literature that range from 19%–27% (Steidel et al. 2001, considering 29 averaged $z = 3.4$ LBG spectra), 24%–35% (Iwata et al. 2009, for $z \sim 3$ LBGs in the SSA22a field), and 4%–6% (Shapley et al. 2006), although the two LBGs in Shapley et al. (2006) with putative LyC emission were later shown to be a spurious detection (Nestor et al. 2011; Iwata et al. 2009) and a foreground contaminant (Nestor et al. 2013) such that f_{esc} calculated from this work would be consistent with zero. Using a slightly different method to make a direct measurement of f_{esc} for a sample of $3.4 < z < 4.5$ LBGs, Vanzella et al. (2010a) calculated $f_{\text{esc}} < 5\%$ –20%. Finally, by measuring H I opacity along 32 gamma-ray burst sightlines in the redshift range $2.0 < z < 5.5$, Fynbo et al. (2011) determined $\langle f_{\text{esc}} \rangle = 0.02 \pm 0.02$ with a 95% confidence level upper limit of $\langle f_{\text{esc}} \rangle \leq 0.07$. This wide range of values highlights how difficult it has proven to determine an accurate value for f_{esc} . These difficulties stem from many factors, including determination of redshifts, spatial resolution, foreground contamination, and the fact that only $\sim 10\%$ of galaxies in the sample will be detected in the LyC. In our study, only 12 out of 131 spectroscopically confirmed galaxies were observed to have NB3420 detections, and some of these are probably contaminated by foreground galaxies. Such a small sample size of NB3420 detections, while an improvement upon many previous studies of spectroscopically confirmed galaxies, limits the precision of our f_{esc} measurements. Until there exist uncontaminated samples of LyC-emitting galaxies an order of magnitude larger, small sample sizes will be a constant source of uncertainty.

5.3. The LyC Emissivity of Star-forming Galaxies

We estimate the comoving specific emissivity of ionizing photons as

$$\epsilon_{\text{LyC}} = \left(\frac{F_{\text{UV}}}{F_{\text{LyC}}} \right)_{\text{corr}}^{-1} \int_{L_{\text{min}}}^{L_{\text{max}}} L \Phi dL, \quad (4)$$

following the assumptions of Nestor et al. (2013). In estimating ϵ_{LyC} , we assume that the UV luminosity functions of LBGs and LAEs do not change significantly from $z = 3.09$ to $z = 2.85$ and the only change in the value of $\int L \Phi dL$ is due to the integration bounds (i.e., L_{min} and L_{max}). Comparing our observed values for $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ to those of Nestor et al. (2013), we obtain $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{HS1549}} = 4.6(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{SSA22a}}$ for LBGs and $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{HS1549}} = 2.0(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{SSA22a}}$ for LAEs with spectroscopic redshifts. Nestor et al. (2013) suggest two different models for combining the LBG and LAE luminosity functions (see Equations (4) and (5) in Nestor et al. 2013) to calculate the global ionizing emissivity. In the luminosity-dependent model, LAEs are assumed to represent galaxies with faint UV continuum magnitudes ($0.06 L^* < L < 0.34 L^*$, corresponding to $25.5 < V < 27.5$) and LBGs represent brighter galaxies ($L > 0.34 L^*$). However, the fact that our data show no change in average LAE NB3420– V color across a range in V magnitude (Figure 10) supports the idea that LAEs are a population of galaxies with properties distinct from those of LBGs, and not simply faint LBG analogs. This scenario is described by the LAE-dependent model, in which LAEs are assumed to comprise 23% of the LBG population (see Nestor et al. 2013), galaxies identified both as LBGs and LAEs are

Table 7
Contributions to the Ionizing Background

	LF ^a	$F_{\text{UV}}/F_{\text{LyC}}^b$	Magnitude Range ^c	ϵ_{LyC}^d
(i)	LBG	82 ± 45	$M_{\text{AB}} \leq -19.7$	1.4 ± 0.8
(ii)	LAE	7.4 ± 3.6	$-19.7 < M_{\text{AB}} \leq -17.7$	3.2 ± 1.6
(iii)	LBG	7.4 ± 3.6	$-19.7 < M_{\text{AB}} \leq -17.7$	13.6 ± 6.7
(iv)	LBG	82 ± 45	$M_{\text{AB}} \leq -17.7$	2.6 ± 1.5
(v)	LAE	7.4 ± 3.6	$M_{\text{AB}} \leq -17.7$	6.8 ± 3.3
	Total (lum.-dep.) ^e	...	$M_{\text{AB}} \leq -17.7$	15.0 ± 6.7
	Total (LAE-dep.) ^f	...	$M_{\text{AB}} \leq -17.7$	8.8 ± 3.5

Notes.

^a Luminosity function parameters are the same as those described in Nestor et al. (2013).

^b Sample average flux density ratio corrected for foreground contamination and IGM absorption, from Table 6.

^c Magnitude range over which the first moment of the luminosity function is determined. $M_{\text{AB}} = -19.7$ and -17.7 correspond to $0.34 L_*$ and $0.06 L_*$, respectively.

^d Comoving specific emissivity of ionizing radiation in units of $10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$.

^e Totals for the luminosity-dependent model, determined by summing rows (i) and (iii).

^f Total for the LAE-dependent model, determined by summing $0.77 \times$ row (iv) and row (v).

treated as LAEs, and the luminosity function is integrated over the full luminosity range ($0.06 L^* < L < \infty$) for both LBGs and LAEs. Based on the values of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ derived for the HS1549 field, $\epsilon_{\text{LyC}} = 15.0 \pm 6.7 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$ for the luminosity-dependent model and $\epsilon_{\text{LyC}} = 8.8 \pm 3.5 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$ for the LAE-dependent model. The uncertainties in ϵ_{LyC} reflect only uncertainties in $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$, which dominate over uncertainties in the luminosity function. Our values of ϵ_{LyC} are roughly half of those calculated by Nestor et al. (2013) for star-forming galaxies in the SSA22a field: $\epsilon_{\text{LyC}}^{\text{SSA22a}} = 32.2^{+12.0}_{-11.4} \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$ for the luminosity-dependent model and $\epsilon_{\text{LyC}}^{\text{SSA22a}} = 16.8^{+6.9}_{-6.5} \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$ for the LAE-dependent model. In Table 7, we summarize the contributions to ϵ_{LyC} as determined from galaxies in the HS1549 field.

To place these values of ϵ_{LyC} for star-forming galaxies in context, we can compare to values of the total ionizing emissivity ($\epsilon_{\text{LyC}}^{\text{tot}}$) derived from Ly α -forest studies, which represent an upper bound on the ionizing emissivity from star-forming galaxies. Using the formulation described in Nestor et al. (2011), we define $\epsilon_{\text{LyC}}^{\text{tot}}$ to be

$$\epsilon_{\text{LyC}}^{\text{tot}} = \frac{\Gamma_{\text{H I}} h (3 - \alpha_s)}{\sigma_{\text{H I}} \lambda_{\text{mfp}} (1 + z)^3}, \quad (5)$$

where $\Gamma_{\text{H I}}$ is the total hydrogen photoionization rate in the IGM at $z = 2.85$ (measured by Ly α -forest studies), h is Planck's constant, α_s is the power-law index of the UV spectral slope in the LyC region ($f_\nu \propto \nu^{\alpha_s}$; we adopt $\alpha_s = -3$), $\sigma_{\text{H I}} = 6.3 \times 10^{-18} \text{ cm}^2$ is the atomic hydrogen photoionization cross section, and $\lambda_{\text{mfp}} = 100 \text{ Mpc}$ is the ionizing photon mean free path through the IGM at $z = 2.85$ (Rudie et al. 2013; Faucher-Giguère et al. 2008; Songaila & Cowie 2010). Using $\Gamma_{\text{H I}} = 0.92 \times 10^{-12} \text{ s}^{-1}$ inferred from Bolton & Haehnelt (2007), we derive the total ionizing photon emissivity at $z = 2.85$ to be $\epsilon_{\text{LyC}}^{\text{tot}} = 9.8 \pm 4.1 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$. Using $\Gamma_{\text{H I}} = 0.53 \times 10^{-12} \text{ s}^{-1}$ inferred from Faucher-Giguère et al. (2008), we

derive $\epsilon_{\text{LyC}}^{\text{tot}} = 5.6 \pm 1.6 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$. Estimates of the contribution of QSOs to the ionizing background at $z = 2.85$ are lower, and range from $\epsilon_{\text{LyC}}^{\text{QSO}} \sim 1.5 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}$ (Cowie et al. 2009) to $\epsilon_{\text{LyC}}^{\text{QSO}} \sim 5.5 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}$ (Hopkins et al. 2007). Within their errors, the values of ϵ_{LyC} that we report for star-forming galaxies in the HS1549 field are consistent with the inferred total ionizing emissivity from the Ly α -forest studies. This agreement between the total ionizing emissivity and the ionizing emissivity from star-forming galaxies leaves little room for the contribution to the emissivity from QSOs or from fainter star-forming galaxies not probed by our observations. Although the LyC emissivity we measure from star-forming galaxies may still be an overestimation, it is in better agreement with $\epsilon_{\text{LyC}}^{\text{tot}}$ determined from Ly α -forest studies than the higher values determined in past work (e.g., Nestor et al. 2013, 2011; Steidel et al. 2001). Even with the large uncertainties in $\epsilon_{\text{LyC}}^{\text{tot}}$, $\epsilon_{\text{LyC}}^{\text{QSO}}$, and $\epsilon_{\text{LyC}}^{\text{gal}}$, our results may point to star-forming galaxies providing the dominant contribution of ionizing photons at $z = 2.85$.

6. PROPERTIES OF LyC-EMITTING GALAXIES

The measurements of the ionizing to nonionizing flux density ratios of the galaxies in our samples point to a large spread in the amount of LyC emission escaping from galaxy to galaxy. While the majority of galaxies in our sample ($\sim 90\%$) appear to have no leaking LyC radiation, the remaining 10% exhibit very blue NB3420–V colors, indicating a high escape fraction. Two possible scenarios may describe our data. One possibility is that galaxies with and without observed LyC emission have an intrinsically different physical property governing whether LyC radiation escapes from or is absorbed by the ISM. A second possibility is that LyC emission escapes from all star-forming galaxies, but only over a small solid angle where neutral hydrogen has been cleared away (e.g., by stellar winds or supernovae). In this section, we explore the differential properties of galaxies with and without observed LyC radiation, with the goal of distinguishing between these two scenarios. Specifically, we examine the rest-frame near-infrared properties of both LBGs and LAEs and the Ly α EWs of the LAEs. The ultimate goal is to be able to identify galaxies associated with strong LyC emission by some other galactic property and search for analogs to such galaxies at higher redshifts, in regimes where the IGM is opaque to LyC photons.

6.1. Rest-frame Near-infrared Properties of LyC-emitting Galaxies

In addition to the optical and near-UV data in the HS1549 field, there exists imaging in several infrared bands. A small fraction of our objects lie in the footprint of the *HST*/WFC3 *F160W* image (rest-frame 4150 Å) and the *Spitzer*/IRAC Channels 1 and 3 images (*Ch1*, 3.6 μm ; *Ch3*, 5.8 μm ; *Spitzer* program G03, PI: Steidel) corresponding, respectively, to rest-frame 0.9 μm and 1.5 μm . The entire field has been imaged by Palomar/WIRC in *J* and *K* (rest-frame 3250 Å and 5700 Å, respectively) and *Spitzer*/IRAC Channels 2 and 4 (*Ch2*, 4.5 μm ; *Ch4*, 8 μm) corresponding respectively to rest-frame 1.2 μm and 2.1 μm . As the *Ch2* image is the deepest of these images, we focus our subsequent analysis on *Ch2*. In the deepest part of the mosaics, the *Ch2* image has an exposure time of 18500 s and a 3σ limiting magnitude of 24.0 (AB). The typical *Ch2* IRAC

PSF FWHM ($\sim 2''.5$) is significantly larger than the seeing in the LRIS optical imaging. All of the LBG and LAE targets lie within the footprint of the *Ch2* image, and the majority lie within the deepest, central regions. We performed PSF-fitting photometry of the LBGs and LAEs in the *Ch2* image using procedures described in Reddy et al. (2006). Results of the photometry indicate that 24 out of the 48 LBGs are detected in *Ch2* at the 3σ level, while there are *Ch2* detections for only 7 out of the 90 LAEs (2 of the 7 LAEs are also LBGs). The median limits in *Ch2* magnitude and $V - \text{Ch2}$ color for the nondetections are $\text{Ch2} > 24.0$ and $V - \text{Ch2} < 0.7$ for LBGs, and $\text{Ch2} > 24.0$ and $V - \text{Ch2} < 2.1$ for LAEs.

In order to understand the typical properties of LBG stellar populations in the HS1549 field, we obtained SED fits for each LBG with at least one infrared datapoint. While all 48 of the LBGs in our sample have *U*, *G*, and *R* photometry, only 33 of these also have detections in one or more infrared bands (*J*, *K*, *F160W*, and/or IRAC *Ch1* – *Ch4*). For these 33 objects, we derived SED fits based on the Bruzual & Charlot (2003) models using a Chabrier initial mass function (IMF). Following Reddy et al. (2012), we adopted constant SFR models, the Calzetti et al. (2000) dust attenuation curve, and 50 Myr lower limits in age (representing the typical LBG dynamical timescale). The ages of the galaxies range from 50 to 2200 Myr, with a median age of 202 Myr. Roughly one-third of the sample was assigned the lowest allowed age value (50 Myr). The SFRs range from 5 to 214 $M_{\odot} \text{ yr}^{-1}$, with a median SFR of 43 $M_{\odot} \text{ yr}^{-1}$. The stellar masses range from $\log(M/M_{\odot}) = 9$ to $\log(M/M_{\odot}) = 11$ with a median of $\log(M/M_{\odot}) = 9.8$. Finally, the values of $E(B - V)$ range from 0 to 0.31, with a median value of 0.2. These values agree well with typical values of LBG properties quoted in the literature (e.g., Kornei et al. 2010). The 15 LBGs without IRAC photometry ($\sim 30\%$ of our sample) for which SED fits could not be calculated are most likely undetected in IRAC bands because of their lower stellar mass and dust content. Thus, the inclusion of these objects would likely change the average distributions of SED fit parameters. A similar bias exists in the average LBG properties reported by Kornei et al. (2010); of the 321 $z \sim 3$ LBGs in their sample, roughly 25% of these LBGs did not have photometric detections redward of the Balmer break, and were therefore not modeled with SED fits. Out of the 33 LBGs in our sample that were modeled with SED fits, only 1 (MD34) has an NB3420 detection. The SED fit to MD34 produced the following values: age = 50 Myr (the minimum allowed), SFR = 137 $M_{\odot} \text{ yr}^{-1}$, $\log(M/M_{\odot}) = 9.8$, $E(B - V) = 0.28$. While it is not possible to draw conclusions about the global population of LBGs emitting LyC based on only one object, we note that MD34 has been assigned the lowest possible age allowed by our modeling and the second largest SFR of the 33 modeled galaxies.

Because only a fraction of the LBGs and LAEs have sufficient infrared photometry to obtain SED fits, we use $V - \text{Ch2}$ colors and limits to constrain the range of possible stellar populations of our full sample of galaxies. For galaxies at $z \sim 2.85$, the *V* and *Ch2* filters lie on either side of the Balmer Break and thus are a direct probe of the stellar mass-to-light ratio of the galaxy. Thus, we can convert $V - \text{Ch2}$ colors to stellar masses by assuming a range of values for $E(B - V)$ and using *V*-band imaging to estimate galaxy luminosity. We present a $V - \text{Ch2}$ versus *V* CMD in Figure 11. Photometry for individual LBGs and LAEs is plotted for objects with and without NB3420 detections along with curves of constant stellar mass for various assumptions of $E(B - V)$. In the sample of LBGs, there is only

Figure 11. Color-magnitude diagram plotting $V-Ch2$ color against V magnitude for LBGs and LAEs with detections in V . Individual LBGs are plotted using filled blue squares for objects with NB3420 detections and open green squares for objects without NB3420 detections. LAEs without NB3420 detections are plotted with open black circles, and the two LAE detections brightest in V are also LBGs. Objects undetected in $Ch2$ are plotted as 3σ upper limits in $V-Ch2$ color following the color scheme for the detected objects. All LAEs with NB3420 detections are undetected in $Ch2$; the upper limits on these objects are indicated by red arrows. The open purple diamond represents the color of the LBG stack of NB3420 nondetections. The filled (open) purple inverted triangle represents the upper limit in $V-Ch2$ color for the stack of LAEs with NB3420 detections (nondetections). Gray curves indicate lines of constant stellar mass, with $M_{\text{stellar}} = 10^8 M_{\odot}$, $10^9 M_{\odot}$, $10^{10} M_{\odot}$, and $10^{11} M_{\odot}$ shown in black, dark gray, medium gray, and light gray, respectively. Values of $E(B-V)$ are indicated by solid, dashed, and dot-dashed lines, corresponding to $E(B-V) = 0.0$, 0.15 , and 0.3 , respectively. The curves of constant mass are produced by a Bruzual & Charlot (2003) SED fit to a $z = 2.85$ galaxy using a Chabrier IMF. The length of each curve corresponds to ages ranging from 50 Myr to 2300 Myr (the age of the universe at $z = 2.85$).

(A color version of this figure is available in the online journal.)

one galaxy (MD34) with both a $Ch2$ detection and an NB3420 detection; MD34 does not differ significantly from the other $Ch2$ -detected LBGs in its $V-Ch2$ versus V properties. In the sample of LAEs, there only five objects with both $Ch2$ and V detections, none of which have NB3420 detections. While we cannot make any strong statements comparing the $V-Ch2$ properties of LAEs with and without NB3420 detections, we can note that the limits of LAEs with NB3420 detections do not distinguish them from LAEs without NB3420 detections. Deeper $Ch2$ data is necessary to make any stronger inferences. The preliminary results from this analysis seem to show no strong differences between the stellar populations of galaxies with and without NB3420 detections, implying that the detection of LyC emission from a small portion of galaxies stems from the effect of varying observer perspective with respect to the geometry of the ISM of each galaxy, rather than intrinsic physical differences between galaxies. According to such a scenario, galaxies are described by a constant escape fraction that appears to vary between objects based on the perspective of the observer. Another important feature of Figure 11 is that the curves of constant stellar mass indicate that the $V-Ch2$ limits of nearly all the LAEs are consistent with $M_{\text{stellar}} < 10^{10} M_{\odot}$ galaxies with $E(B-V) \lesssim 0.3$. A notable outlier is *lae1843*, the LAE with $V-Ch2 > 4$. *lae1843* is not located near an obvious contaminant in any of our images, but there is still the possibility of a very red foreground galaxy not visible in the shorter-wavelength images or a foreground galaxy in those images that cannot be distinguished from *lae1843* at the current spatial resolution.

Because the $Ch2$ data are not deep enough to characterize the individual rest-frame near-infrared properties of all of our target galaxies, we also performed photometry on stacked images of galaxies with and without LyC detections for the LBG and LAE samples. The stacks were created using methods similar to those described in Appendix D of Reddy et al. (2012), with the requirement that objects included in the stack not be blended with a nearby neighbor in the $Ch2$ image. $V-Ch2$ colors of the LBG and LAE stacks are plotted in Figure 11. We use stellar population synthesis models to estimate masses from the stacked photometry; while this method is not necessarily equivalent to reporting the average properties of individual SED fits, it provides rough insights into the relative properties of different subsamples in the absence of SED fits for every object in our sample.

Unfortunately, it is not possible to draw definitive conclusions from the LBG stacks. While the stack of LBGs without NB3420 detections yields a $Ch2$ magnitude of 24.01 ± 0.10 and $V-Ch2 = 0.58$, we were unable to make a useful stack for the LBGs with NB3420 detections because two of the four objects are blended with bright neighbors in the $Ch2$ image. Of the two LBGs with NB3420 detections and unblended $Ch2$ photometry, MD34 has a magnitude of 22.34 ± 0.12 and MD5 has a 3σ lower limit of >24.22 . While the $Ch2$ photometry of MD34 and MD5 is consistent with that of the LBGs undetected in NB3420, the small sample size precludes us from making any comparisons between objects with and without LyC detections. We can also compare the LBG properties determined by $V-Ch2$ color to those determined by the SED fits. The $V-Ch2$ color of the stack of LBGs without NB3420 detections implies $\log(M/M_{\odot}) = 9.60$ (9.10), assuming $E(B-V) = 0.0$ (0.1). Higher values of $E(B-V)$ imply younger galactic ages, and assumptions of $E(B-V) > 0.1$ do not result in meaningful fits to the stacked color because they imply ages < 50 Myr. The assumption of $E(B-V) = 0.1$ requires an age of ~ 60 Myr, significantly younger than the median age derived for the sample of LBGs with SED fits (202 Myr). However, LBGs have been shown to have typical reddening values of $E(B-V) \sim 0.17$ (Kornei et al. 2010), and assuming such a value further decreases the average age implied by the stacked photometry. This apparent discrepancy (i.e., the fact that the median age implied by the stacks is lower than that implied by the SED fits) arises because the stacks include all LBGs regardless of whether or not they are detected in $Ch2$, while the SED fits only include the older or more massive LBGs with IRAC photometry redward of the Balmer break. We can further demonstrate this discrepancy by contrasting the $V-Ch2$ color of the stack of all LBGs without LyC detections ($V-Ch2 = 0.58$) with the average $V-Ch2$ color of the objects in that stack that have $Ch2$ detections ($V-Ch2 = 1.33$). Because the stack includes all LBGs regardless of whether or not they have infrared data, the $V-Ch2$ color of the stack is much bluer than the $V-Ch2$ color of the subset of objects within the stack that have $Ch2$ detections.

In the case of the stacks of LAEs with and without NB3420 detections, neither stack was detected in $Ch2$; the stacks reached 3σ limiting magnitudes of >24.30 and >25.64 , respectively, corresponding to $V-Ch2 < 1.19$ and $V-Ch2 < 0.45$. While these limits on LAE $V-Ch2$ color do not give any constraints on the differential properties between LAEs with and without NB3420 detections, they do indicate that LAEs are preferentially low-mass galaxies with small values of $E(B-V)$. Converting the limits on LAE $V-Ch2$ color into limits on stel-

lar mass results in $\log(M/M_\odot) < 9.72$ (9.38) for LAEs with NB3420 detections and $\log(M/M_\odot) < 8.87$ (8.35) for LAEs without NB3420 detections, assuming $E(B - V) = 0.0$ (0.1). The $E(B - V) = 0.1$ fit to the deeper LAE stack (LAEs without NB3420 detections) corresponds to an age slightly less than 50 Myr. While such a young age is unphysical for LBGs based on dynamical timescale arguments, LAEs may be more compact systems where such young ages are feasible; such young ages have been determined for LAEs in previous work (Gawiser et al. 2007).

In summary, the analysis of the infrared data does not yield firm results differentiating between the two proposed models of LyC escape presented at the beginning of this section: that objects with and without LyC detections are either intrinsically different or intrinsically similar with anisotropic emission of LyC radiation. Although the individual galaxies and stacks plotted in Figure 11 do not indicate strong differences in the $V - Ch2$ properties of galaxies with and without LyC emission, the fact that only one galaxy (MD34) with an NB3420 detection is actually detected in $Ch2$ severely limits our interpretation of the $Ch2$ data. At the same time, we find that MD34 (the only object with a LyC detection modeled by an SED fit) has been assigned the lowest possible age allowed by our modeling and the second largest SFR of the 33 modeled LBGs. A larger sample of LBGs with both LyC detections and IRAC photometry is needed to test whether or not LBGs with LyC detections are preferentially young. In terms of the overall sample properties, the $V - Ch2$ color of our LBG stacks (which include all LBGs in the sample) indicate an age for LBGs that is younger than the median age of LBGs modeled with SED fits, suggesting a possible bias in previous LBG stellar population studies (e.g., Shapley et al. 2001; Kornei et al. 2010) limited to objects with near- or mid-infrared detections. Finally, the $V - Ch2$ colors and limits we observe for the LAEs in our sample indicate that they are preferentially low-mass galaxies with small values of $E(B - V)$, consistent with previous results in $z \sim 3$ LAE stellar population studies. The two obstacles hindering our understanding of the rest-frame near-infrared properties of the galaxies in our sample are the low $Ch2$ detection rate of both LBGs and LAEs and the lack of spatial resolution, resulting in the blending of several galaxies with bright nearby neighbors. Deeper and higher resolution data enabling robust SED fits for all galaxies in our sample would constitute a much more powerful tool to distinguish between the proposed models of LyC escape.

6.2. The Relationship between LyC Emission and Ly α Emission

In addition to exploring the differential rest-frame UV–near-infrared colors of galaxies with and without LyC emission, we studied their relative Ly α properties. Figure 12 shows the $V - NB4670$ versus $NB4670$ CMD for LAEs with and without NB3420 detections (shown in red and black, respectively), indicating the rest-frame Ly α EW on the right-hand axis. Although the LAEs with NB3420 detections have lower measured values of Ly α EW on average than LAEs without NB3420 detections, they span a range in $NB4670$ magnitude equivalent to that of the full LAE sample. In order to account for the fact that for several LAEs the measured EWs are only lower limits, we used the package ASURV (“Astronomy SURVival Analysis”) Rev 1.2 (Isobe & Feigelson 1990; Lavalley et al. 1992) to calculate the Ly α EW mean and standard error of LAEs with and without NB3420 detections using survival analysis. This analysis shows that the mean Ly α EW for objects with NB3420 detections is

Figure 12. Color-magnitude diagram plotting $V - NB4670$ vs. $NB4670$ for LAEs. The rest-frame Ly α equivalent width corresponding to a given $V - NB4670$ color is indicated on the right-hand axis. Galaxies without NB3420 detections are shown in black, and galaxies with NB3420 detections are shown as red with a circle around each point. Lower limits correspond to objects undetected in V and follow the same color scheme as the detections. The dashed diagonal line indicates the observational limit on $V - NB4670$ color, which is determined by our V -band magnitude limit of 27.58. The vertical dashed line indicates the cutoff of $m_{4670} < 26.0$ applied in order to ensure the robustness of our sample (see Section 2.4). The horizontal dashed line corresponds to the color cut of $V - NB4670 > 0.6$ used for LAE selection. A few LAEs fall below this line because they were selected as LAEs on the basis of earlier, shallower photometry, and they were already spectroscopically confirmed to lie at $z = 2.85$.

(A color version of this figure is available in the online journal.)

$34.0 \pm 5.9 \text{ \AA}$ and the mean Ly α EW for objects without NB3420 detections is $80.1 \pm 8.3 \text{ \AA}$. These results confirm that galaxies with NB3420 detections tend to have lower Ly α EWs, on average, than galaxies without NB3420 detections, a correlation that has been previously reported in the SSA22a field by Nestor et al. (2011).

In testing the connection between Ly α and LyC emission, we must also consider changes in LyC emission when Ly α EW is the independent variable. To determine if $(F_{UV}/F_{LyC})_{\text{corr}}$ is positively correlated with Ly α EW, we calculated $(F_{UV}/F_{LyC})_{\text{corr}}$, the nonionizing to ionizing flux density ratio corrected for contamination and IGM absorption, for LAEs in two bins of Ly α EW separated by the mean LAE Ly α EW ($46 \text{ \AA}$). For LAEs with smaller and larger EWs, respectively, we find $(F_{UV}/F_{LyC})_{\text{corr}}^{\text{EW} \leq 46 \text{ \AA}} = 4.9 \pm 2.6$ and $(F_{UV}/F_{LyC})_{\text{corr}}^{\text{EW} > 46 \text{ \AA}} = 33 \pm 23$; LAEs with lower EWs have smaller values of $(F_{UV}/F_{LyC})_{\text{corr}}$ (indicating stronger LyC emission), and LAEs with larger EWs have larger values of $(F_{UV}/F_{LyC})_{\text{corr}}$. These two values are consistent with the results of the survival analysis: LAEs with LyC detections are characterized by a smaller average Ly α EW than those without LyC detections.

One explanation of the observed anticorrelation between Ly α EW and LyC emission strength consists of the fact that Ly α photons are simply reprocessed LyC photons. As a LyC photon propagates through the ISM of a galaxy, one possibility is that the photon escapes without being absorbed by a neutral hydrogen atom, and LyC emission is observed from the galaxy. A second possibility is that the photon encounters a neutral hydrogen atom and is absorbed, ionizing the atom. When the free electron recombines with the hydrogen ion, the emission cascade spectrum will often culminate in the release of a Ly α photon. Thus, for a fixed production rate of LyC photons, galaxies with larger Ly α EWs should tend to exhibit less LyC emission.

The fact that the stellar populations of galaxies with and without LyC detections do not appear to differ significantly

(see Section 6.1) suggests that detecting LyC in a given galaxy depends on the orientation of the galaxy with respect to the observer. If the hypothesis that LyC detectability depends on observer orientation is correct, then the anticorrelation between observed LyC and Ly α emission implies that Ly α EW must also depend on orientation, and in such a way to support the anticorrelation. Results from several recent galaxy simulations (of both isolated galaxies and those set in cosmological context) suggest that the observed Ly α EW is strongly dependent on the orientation of the observer (Verhamme et al. 2012; Barnes et al. 2011; Yajima et al. 2012). Specifically, these simulations predict orientation effects for disk galaxies: a disk galaxy, when viewed face-on, will have a much greater observed Ly α EW than the same galaxy viewed edge-on. This orientation effect stems from the fact that Ly α photons are resonantly scattered; for a Ly α photon, the direction of final escape from the galaxy is nearly independent of its original trajectory. As trajectories perpendicular to the plane of the galaxy present the paths of least opacity, more Ly α photons will escape perpendicular to the galactic plane. It is still unclear how the spatial redistribution of Ly α photons relates to that of LyC photons, and how these distributions are affected by the clumpy morphologies of high-redshift galaxies. If, for example, Ly α and LyC photons are both more likely to escape along lines of sight with lower hydrogen column densities, a positive correlation between observed Ly α and LyC emission might be expected—contrary to the anticorrelation observed in the LAE populations of the HS1549 and SSA22a fields. In any case, it would be extremely useful for simulations to examine the effect that galaxy orientation produces on the relationship between observed Ly α and LyC emission. It would also be useful for such simulations to consider morphologies that more closely represent those of high-redshift galaxies (i.e., clumpy and irregular morphologies; Law et al. 2007). If orientation effects *cannot* explain the observed trend that smaller observed Ly α EWs in LAEs coincide with more LyC emission, then the difference in Ly α EWs must indicate some intrinsic difference between galaxies exhibiting strong or weak LyC emission that can be traced by Ly α emission. Along these lines, we note that the mean rest-frame Ly α EW for LAEs without NB3420 detections (80 Å) is close to that expected by the “Case B” assumption for 100% absorption of LyC photons in a dust-free environment (~ 100 Å; Steidel et al. 2011). In the presence of dust, the resonant scattering of Ly α photons would tend to decrease the expected “Case B” Ly α EW. If Ly α orientation effects do not correlate with LyC in the manner described above, then the larger Ly α EWs in our sample of LAEs without NB3420 detections indicates a higher rate of LyC absorption within the ISM of these objects.

Also relevant for understanding the connection between Ly α and LyC emission, we find a lower average nonionizing to ionizing UV flux density ratio among LAEs (which have a median Ly α EW of 46 Å) than among LBGs (which have a much lower median Ly α EW of 5 Å). Considering LAEs in the same magnitude range as LBGs ($V < 25.4$), we derive $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{LAE}, V < 25.4} = 4.1 \pm 2.7$, which is significantly lower than the value for LBGs, $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}^{\text{LBG}} = 82 \pm 45$. The relative values of $(F_{\text{UV}}/F_{\text{LyC}})_{\text{corr}}$ for LAEs and LBGs apparently suggest a *positive* correlation between Ly α EW and LyC emission, in the opposite sense of the trend traced by the LyC versus Ly α emission properties of LAEs. However, LBGs have been shown to differ from LAEs in several galactic properties, including dust content and gas/dust covering fraction. As dust destroys UV

photons, and Ly α photons travel particularly long path lengths before they escape their galaxy due to resonant scattering, Ly α photons suffer greater dust attenuation than other UV photons. Thus, while the amount of dust in a galaxy will not affect the creation of Ly α photons from LyC photons, it will affect the observed Ly α EW. Observational studies have confirmed that LBGs with larger Ly α EWs have smaller values of $E(B - V)$ (e.g., Atek et al. 2009; Shapley et al. 2003) and simulations have reproduced the same result (e.g., Yajima et al. 2012; Dayal et al. 2009). Also, LBGs with larger Ly α EWs have weaker interstellar absorption lines, indicating lower gas covering fraction (Shapley et al. 2003). Thus, it is possible that the low median Ly α EW observed in LBGs compared to LAEs is due to increased dust attenuation (LBGs have median values of $E(B - V)$ that are typically higher than those of LAEs; Gronwall et al. 2007; Blanc et al. 2011) and larger gas covering fraction. In summary, the observed strengths of Ly α and LyC emission may be anticorrelated (as seen in the LAE sample) except in the cases of higher interstellar extinction where Ly α photons are systematically destroyed (as in the case of the LBGs).

7. SUMMARY

We have presented an analysis of the $z \sim 3$ LyC properties of the HS1549 field. Multiple LyC studies of star-forming galaxies have been conducted in the SSA22a field (e.g., Inoue et al. 2005; Shapley et al. 2006; Iwata et al. 2009; Nestor et al. 2011, 2013) and HS1549 therefore provides an important independent set of $z \sim 3$ LyC measurements. The HS1549 field contains a galaxy protocluster at $z = 2.85$, and our narrowband (NB3420) imaging has targeted the LyC spectral region of 131 spectroscopically confirmed galaxies at $z \geq 2.82$ (49 LBGs and 91 LAEs, 9 of which are constituents of both samples). We have detected 12 galaxies in NB3420 and implemented simulations to correct the average NB3420 magnitudes of subsamples of these galaxies for both foreground galaxy contamination and absorption from neutral hydrogen in the IGM. In addition to more than 19 hr of imaging in the NB3420 filter, we have analyzed ancillary narrowband Ly α imaging and broadband imaging in the rest-frame nonionizing UV, optical, and near-infrared in order to determine nonionizing to ionizing UV flux density ratios and place limits on other galactic properties. Our main conclusions are as follows.

1. We find four LBGs and seven LAEs with NB3420 detections, along with one additional object (MD12/lae3540) that we have removed from both samples. The NB3420 detections of the LBGs range in magnitude from $26.56 \leq m_{3420} \leq 27.01$, and those of the LAEs range from $25.11 \leq m_{3420} \leq 27.23$. Our contamination simulations predict that 1.5 ± 1.0 LBG NB3420 detections and 4.3 ± 1.3 LAE NB3420 detections are uncontaminated by foreground galaxies.
2. The ratio of nonionizing to ionizing UV radiation of the LBGs and LAEs is traced by the observed NB3420– V color. The colors of our targets vary widely, with NB3420-detected galaxies at a given V magnitude having NB3420 magnitudes up to a factor of ten brighter than the NB3420 limits of undetected galaxies; i.e., a small percentage of galaxies are detected strongly in NB3420 and the rest remain undetected. These results confirm those presented in Nestor et al. (2011) for galaxies in the SSA22a field, and lend further support to the hypothesis that LyC radiation escapes galaxies through cleared holes in the ISM and can

Table 8
Photometry for GNBs and LAEs with $m_{4670} > 26.0$

ID	R.A. ^a (J2000)	Decl. ^a (J2000)	z	NB4670	V	NB3420	$\Delta_{UV, LyC}^b$	$\Delta_{Ly\alpha, LyC}^c$	$\frac{F_{UV}}{F_{LyC}}^d_{obs}$
GNB2861	15:51:53.603	19:11:36.53	2.844	26.80	26.08	26.86	0.28	0.33	2.1 ± 1.0
GNB4769	15:51:59.873	19:08:42.73	2.849	25.30	26.54	27.33	0.34	0.52	2.1 ± 0.9
GNB5270	15:51:57.364	19:09:53.38	2.847	24.84	25.14	26.43	0.40	0.30	3.3 ± 1.0
lae1670	15:51:45.121	19:10:15.34	2.846	26.18	27.27	26.45	0.25	0.65	0.5 ± 0.1
lae3506	15:51:52.242	19:11:41.01	2.841	26.71	27.58	27.09	0.11	0.08	0.6 ± 0.3
lae3828	15:51:53.228	19:13:08.50	2.892	26.64	26.85	26.27	0.28	0.33	0.6 ± 0.3
lae5404	15:52:00.008	19:08:54.42	2.816	26.60	27.16	26.53	0.21	0.15	0.6 ± 0.3
lae7890	15:52:01.943	19:12:42.47	2.850	26.02	26.13	25.91	0.23	0.17	0.8 ± 0.3

Notes.^a GNB and LAE coordinates are based on NB4670 centroids.^b Spatial offset between the centroids of V and NB3420 emission.^c Spatial offset between the centroids of NB4670 and NB3420 emission.^d Observed ratio and uncertainty in the ratio of nonionizing UV to LyC flux densities, inferred from the NB3420– V color. This value has not been corrected for either contamination by foreground sources or IGM absorption.

Table 9
Photometry for LAE Photometric Candidates

ID	R.A. ^a (J2000)	Decl. ^a (J2000)	NB4670	V	NB3420	$\Delta_{UV, LyC}^b$	$\Delta_{Ly\alpha, LyC}^c$	$\frac{F_{UV}}{F_{LyC}}^d_{obs}$
lae759	15:51:41.039	19:09:57.61	25.54	26.79	>27.30	...	...	>1.6
lae803	15:51:41.598	19:10:19.54	25.68	27.02	>27.30	...	...	>1.3
lae1043	15:51:42.623	19:13:00.70	26.00	27.07	>27.30	...	...	>1.2
lae1058	15:51:42.791	19:11:07.95	26.00	27.14	>27.30	...	...	>1.2
lae1080	15:51:42.918	19:12:33.53	25.92	26.37	26.03	0''3	0''4	0.7 ± 0.2
lae1569	15:51:44.416	19:08:43.92	25.16	25.84	25.51	0''3	0''4	0.7 ± 0.2
lae1840	15:51:45.875	19:11:55.76	25.98	27.19	>27.30	...	...	>1.1
lae1883	15:51:45.893	19:12:23.30	25.49	26.28	>27.30	...	...	>2.6
lae2278	15:51:47.324	19:12:01.54	25.77	>27.58	>27.30	...	...	...
lae2369	15:51:47.863	19:11:24.63	25.93	>27.58	>27.30	...	...	...
lae2431	15:51:48.048	19:09:04.50	25.97	>27.58	>27.30	...	...	...
lae2482	15:51:48.129	19:09:48.31	25.29	>27.58	>27.30	...	...	...
lae2664	15:51:49.090	19:11:00.51	25.52	>27.58	>27.30	...	...	...
lae2666	15:51:49.161	19:13:17.59	25.80	26.89	>27.30	...	...	>1.4
lae3038	15:51:50.091	19:09:02.19	25.74	26.25	26.44	0''7	0''2	1.2 ± 0.4
lae3348	15:51:51.561	19:11:20.54	25.90	26.98	>27.30	...	...	>1.3
lae3365	15:51:51.691	19:11:20.50	25.90	26.78	>27.30	...	...	>1.6
lae3616	15:51:52.218	19:08:21.79	25.10	26.96	27.17	1''9	1''4	1.2 ± 0.6
lae4070	15:51:54.054	19:10:26.76	25.93	27.47	25.88	0''2	0''1	0.2 ± 0.1
lae4079	15:51:54.023	19:10:05.98	25.19	25.71	>27.30	...	...	>4.3
lae4468	15:51:55.367	19:12:57.55	25.82	26.46	25.97	0''6	0''4	0.6 ± 0.2
lae5157	15:52:01.016	19:08:50.34	26.00	>27.58	>27.30	...	...	...
lae5200	15:52:01.083	19:11:25.97	25.83	27.19	26.70	0''2	0''1	0.6 ± 0.3
lae5252	15:52:00.519	19:08:31.67	26.00	26.38	25.72	0''0	0''1	0.5 ± 0.2
lae5446	15:51:58.774	19:12:46.40	25.43	26.01	>27.30	...	...	>3.3
lae5661	15:51:59.092	19:12:39.74	25.89	26.31	>27.30	...	...	>2.5
lae5665	15:51:57.738	19:12:27.59	23.96	25.31	24.78	0''0	0''7	0.6 ± 0.1
lae6041	15:51:57.864	19:11:14.43	25.56	25.58	>27.30	...	...	>4.9
lae6436	15:52:07.285	19:11:54.07	25.50	26.30	>27.30	...	...	>2.5
lae6510	15:52:06.976	19:12:03.51	25.80	25.94	25.41	0''1	0''1	0.6 ± 0.2
lae7110	15:52:02.523	19:12:07.71	25.98	26.96	>27.30	...	...	>1.4
lae7247	15:52:04.452	19:10:48.30	25.96	26.83	>27.30	...	...	>1.5
lae7642	15:52:03.478	19:12:58.32	25.26	26.20	>27.30	...	...	>2.7

Notes.^a Coordinates for LAEs are based on NB4670 centroids.^b Spatial offset between the centroids of V and NB3420 emission.^c Spatial offset between the centroids of NB4670 and NB3420 emission.^d Observed ratio and uncertainty in the ratio of nonionizing UV to LyC flux densities, inferred from the NB3420– V color. This value has not been corrected for either contamination by foreground sources or IGM absorption.

Figure 13. $10'' \times 10''$ postage stamp images of the three GNBS and five faint LAEs with NB3420 detections. Objects are displayed in the following bands: NB3420 (indicating the LyC), V (indicating the nonionizing UV continuum), NB4670–V (indicating Ly α emission), and NB4670–G (also indicating Ly α emission; GNBS only). All postage stamps are centered on the NB4670 centroid, and blue circles ($1''$ radius) indicate the centroid of the NB3420 emission. All postage stamps follow the conventional orientation, with north up and east to the left.

(A color version of this figure is available in the online journal.)

only be viewed by the observer with fortuitous orientation. The large variation in individual galaxy NB3420 properties makes it necessary to consider properties of the full ensemble of LBGs and LAEs in order to determine the average LyC properties of $z \sim 3$ galaxies.

3. One of the principal difficulties in performing accurate LyC photometry is the presence of low-redshift foreground galaxies that introduce contaminating light into the NB3420 filter. At the spatial resolution of *HST*, galaxies at $z \sim 3$ tend to exhibit clumpy morphology. Separating a $z \sim 3$ clump from a foreground contaminant is impossible unless the individual redshifts of each clump are known. While simulations of the amount of foreground contamination can help

correct the observed nonionizing to ionizing UV flux density ratios, it is very difficult to study the properties of LyC-emitting galaxies when the simulations predict that roughly half of the NB3420-detected sample may be contaminated. To obtain spectroscopic or photometric redshifts at the high spatial resolution necessary to distinguish each clump requires imaging from a high-resolution, space-based observing facility such as *HST* coupled with deep high-resolution spectroscopy from the ground.

4. Reconciling the low observed nonionizing to ionizing UV flux density ratios (F_{UV}/F_{LyC}) with intrinsic values predicted by stellar population synthesis models has been a challenge in all $z \sim 3$ LyC studies. As in the SSA22a field, the observed flux density ratios of many of the LBGs and LAEs with putative LyC detections are too low to be reasonably explained by current models. There are possible ways of solving the problem; for example, one could make the assumption of no IGM absorption along the line of sight for galaxies with observed LyC detections. However, this condition is unlikely to hold for all galaxies with putative LyC detections. As of now, the discrepancy between the LyC observations and models remains unresolved.
5. Using estimates of the UV escape fraction from observational studies of $z \sim 3$ galaxies and a range of estimated values for the theoretical ratio of intrinsic UV to LyC luminosities given by the Bruzual & Charlot (2003) and BPASS (Eldridge & Stanway 2009) models, we derive relative LyC escape fractions of $f_{esc, rel}^{LBG} = 5\% - 8\%$ and $f_{esc, rel}^{LAE} = 18\% - 49\%$, and absolute escape fractions of $f_{esc}^{LBG} = 1\% - 2\%$ and $f_{esc}^{LAE} = 5\% - 15\%$. The uncertainties on the escape fractions that we calculate are on the order of 50%. We also determine a comoving specific ionizing emissivity of LyC photons (ϵ_{LyC}) in the range of $8.8 - 15.0 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$. This range of emissivity values was determined by using two different models to combine the LBG and LAE samples (as discussed in Section 5.3); the luminosity-dependent model yields $\epsilon_{LyC} = 15.0 \pm 6.7 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$ and the LAE-dependent model yields $\epsilon_{LyC} = 8.8 \pm 3.5 \times 10^{24} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ Mpc}^{-3}$. These values of ϵ_{LyC} for star-forming galaxies are consistent with the total ionizing emissivity inferred from the Ly α -forest studies of Bolton & Haehnelt (2007) and Faucher-Giguère et al. (2008). The rough agreement between the total ionizing emissivity and the ionizing emissivity from star-forming galaxies leaves little room for the contribution to the emissivity from QSOs or from fainter star-forming galaxies not probed by our observations. While the LyC emissivity we measure from star-forming galaxies may still be an over-estimation, it is in better agreement with the total LyC emissivity determined from Ly α -forest studies than with the higher values determined in past work.
6. Examining the rest-frame near-infrared properties of galaxies in our sample based on IRAC *Ch2* imaging, we find no significant difference in the *Ch2* or *V–Ch2* properties of objects with and without NB3420 detections. This agreement indicates that objects with and without LyC detections are drawn from a populations of star-forming galaxies with similar distributions of stellar mass, age, $E(B - V)$, and SFR. We caution, however, that our interpretation is limited by the fact that only one galaxy (MD34) with an NB3420 detection is detected in *Ch2*. While the *V–Ch2* versus *V* properties of this object do not distinguish it from

Figure 14. $10'' \times 10''$ postage stamp images of the 10 LAE photometric candidates with NB3420 detections. Objects are displayed in the following bands: NB3420 (indicating the LyC), NB4670–V (indicating Ly α emission), and V (indicating the nonionizing UV continuum). The V–NB4670 color of each LAE is indicated below the object name. All postage stamps are centered on the V centroid, and blue circles ($1''$ radius) indicate the centroid of the NB3420 emission. All postage stamps follow the conventional orientation, with north up and east to the left. We note that LAEs with more diffuse Ly α emission may be difficult to distinguish in the NB4670–V image, even though their V–NB4670 colors identify them as LAEs; for such objects, we have increased the stretch of the NB4670–V image to make the diffuse emission more easily visible.

(A color version of this figure is available in the online journal.)

LBGs without NB3420 detections, in the SED modeling it has been assigned the lowest possible age allowed and the second largest SFR of the 33 modeled LBGs. A much larger sample of LBGs with both LyC detections and IRAC photometry is needed to test whether LBGs with LyC detections are preferentially young, or whether LBGs with and without LyC detections come from similar populations.

7. While we do not find significant differences in the stellar populations of objects with and without LyC detection, we do find that LAEs with LyC detections have smaller Ly α EWs on average than those without. A comparison of $(F_{UV}/F_{LyC})_{corr}$ for LAEs in two bins of EW yields $(F_{UV}/F_{LyC})_{corr}^{EW \leq 46 \text{ \AA}} = 4.9 \pm 2.6$ and $(F_{UV}/F_{LyC})_{corr}^{EW > 46 \text{ \AA}} = 33 \pm 23$. These results imply an inverse relationship between the amount of Ly α and LyC emission observed from LAEs. One possibility is that this relationship stems from the fact that Ly α photons are reprocessed LyC photons, so an overall increase in Ly α photons must correlate with a decrease in LyC photons. The details of LyC and Ly α radiative transfer through the ISM of clumpy, high-redshift galaxies, however, are still unclear, as are the effects of varying the observer's orientation while observing Ly α and LyC emission. We also note that despite this apparent anticorrelation between the strength of Ly α and LyC emission, LBGs (which have lower median Ly α EWs than LAEs) display much weaker LyC emission than LAEs in the same magnitude range. This trend may be explained by

the higher dust content in LBGs systematically destroying both LyC photons and resonantly scattered Ly α photons.

Future progress in this field is contingent upon amassing a large and well-defined sample of LyC-emitting star-forming galaxies and eradicating the possibility that LyC measurements are contaminated by foreground interlopers. A detailed, multi-wavelength analysis of such a sample would reveal definitively whether LyC-emitting objects possess specific properties that facilitate LyC escape or whether LyC emission escapes from all galaxies through randomly oriented paths cleared through the ISM. With the goals of removing contamination and obtaining multiwavelength photometry, we are currently pursuing follow-up observations of our NB3420-detected galaxies in the HS1549 field with *HST*/WFC3.

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APPENDIX

A.1. GNBs and Faint LAEs

Here we present photometry (Table 8) and postage stamp images (Figure 13) for eight additional Ly α -emitting galaxies with NB3420 detections for which we have obtained spectroscopic redshifts (see Section 2.3). This sample comprises three GNBs and five faint LAEs ($m_{4670} > 26$). As we do not have a complete and unbiased sample of these objects, we do not include them in the main LAE sample.

A.2. Photometric LAE Sample

Here we present photometry of 33 LAE photometric candidates without spectroscopically confirmed redshifts (see Section 2.4). Ten of these LAEs have NB3420 detections. Table 9 contains photometric information for all 33 LAE photometric candidates, and Figure 14 displays postage stamp images for the 10 candidates with NB3420 detections. While some of these NB3420 detections may be true LyC detections, the high rate of NB3420 detections in this sample may be due to the increased probability of contamination for objects without spectroscopic redshifts. For this reason, we do not include these objects in the LAE analysis.

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