

Supporting Information

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SI Materials and Methods

Composition Dependence of Glass-Forming Ability. Series I–IV of the alloy family $\text{Ni}_{100-w-x-y-z}\text{Cr}_w\text{Nb}_x\text{P}_y\text{B}_z$ are as follows: $\text{Ni}_{77.5-w}\text{Cr}_w\text{Nb}_3\text{P}_{16.5}\text{B}_3$ (series I), $\text{Ni}_{69}\text{Cr}_{11.5-x}\text{Nb}_x\text{P}_{16.5}\text{B}_3$ (series II), $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{19.5-z}\text{B}_z$ (series III), and $(\text{Ni}_{0.8541}\text{Cr}_{0.1085}\text{Nb}_{0.0374})_{100-(y+z)}(\text{P}_{0.8376}\text{B}_{0.1624})_{(y+z)}$ (series IV). The eutectic composition is $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{16.5}\text{B}_3$, and hence the transition from hypo- to hypereutectic occurs at $w = 8.5$, $x = 3$, $z = 3$, and $(y + z) = 19.5$, for series I–IV, respectively.

The d_{cr} vs. composition data presented in Fig. 4 for series I–IV are fitted using an exponential fitting function, as follows:

$$d_{cr} = a \exp(bc) + d. \quad [\text{S1}]$$

The fitting coefficients a , b , and d for series I–IV are presented in Table S1 for the corresponding hypo- and hypereutectic regions. The composition variable c is in units of w , x , z , and $(y + z)$, for series I–IV, i.e., in atomic percentages of Cr, Nb, B, and total metalloid (i.e., combined P and B). d_{cr} is presented in millimeters.

The d_{cr}^2 vs. B concentration data presented in Fig. 5A for $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{19.5-z}\text{B}_z$ (series III) are fitted using the exponential fitting function.

$$d_{cr}^2 = a \exp(bc). \quad [\text{S2}]$$

The composition variable c is in units of z , i.e., in atomic percentages of B, and d_{cr}^2 in millimeters squared. The fitting coefficients a and b are 800 and -0.693 , respectively.

Calorimetry. Differential scanning calorimetry was performed to evaluate the glass transition temperature T_g , crystallization temperature T_x , solidus temperature T_S , and liquidus temperature T_L for $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{19.5-z}\text{B}_z$ (series III) metallic glass alloys for z between 1.5 and 6. T_g and T_x were evaluated as the onset of the glass transition at a scan rate of 20 K/min. T_S and T_L were evaluated at a scan rate of 5 K/min to reduce scan rate effects and instrumental broadening of the melting transition. Scans around the glass transition are presented in Fig. S1, and around the melting transition are presented in Fig. S2. Data for T_g , T_x , T_S , T_L , and t_{rg} are presented in Table S2.

Rheometry. The equilibrium (Newtonian) viscosity for $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{19.5-z}\text{B}_z$ (series III) metallic glass alloys for $z = 1.5, 3$, and 5 was measured by performing three-point beam bending of amorphous rods 2 mm in diameter and 10 mm in length using a thermomechanical analyzer (Perkin–Elmer; TMA 7), as described by Hagy (1). Specifically, the isothermal viscosity at a given temperature is determined by the following equation:

$$\eta = -\frac{g L^3}{144 I_c \nu} \left(M + \frac{5 \rho A L}{8} \right), \quad [\text{S2}]$$

where η is the apparent viscosity (Pa·s), g the gravitational constant (meters per second squared), L the support span length (meters), I_c the cross-sectional moment of inertia (meters to the fourth power), ν the midpoint deflection velocity (meters per second), M the applied load (kilograms), ρ the density (kilograms per meter cubed), and A is the cross-sectional area (meters squared). Loads ranging from 20 to 1,000 mN were applied. Measured data for $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{18}\text{B}_{1.5}$, $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{16.5}\text{B}_3$, and $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{14.5}\text{B}_5$ are presented in Fig. S3.

To describe the temperature dependence of viscosity $\eta(T)$, the cooperative shear model is used (2, 3):

$$\frac{\eta(T)}{\eta_\infty} = \exp \left\{ \frac{W_g}{kT} \exp \left[2n \left(1 - \frac{T}{T_{go}} \right) \right] \right\}, \quad [\text{S3}]$$

where η_∞ is the high-temperature limit of viscosity, which is assumed to be 10^{-3} Pa·s for all alloys, W_g is the activation energy barrier at the glass transition, approximated by $W_g \approx kT_{go} \log(\eta_g/\eta_\infty)$, where $\eta_g \equiv 1 \times 10^{12}$ Pa·s, n is the effective fragility parameter, k the Boltzmann constant, T temperature, and T_{go} the glass transition temperature associated with a viscosity value of 10^{12} Pa·s. The fragility parameter m is related to n via

$$m = (1 + 2n) \log(\eta_g/\eta_\infty). \quad [\text{S4}]$$

The fitting curves for $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{18}\text{B}_{1.5}$, $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{16.5}\text{B}_3$, and $\text{Ni}_{69}\text{Cr}_{8.5}\text{Nb}_3\text{P}_{14.5}\text{B}_5$ are presented in Fig. S3, whereas the fit parameters T_{go} and n and the calculated fragility m are given in Table S3.

Mathematical Analysis. The actual dependence of $\ln \tau_\alpha^*(c)$ on W and ΔG_α depends on the theoretical model used to describe these barriers. In this analysis, we assume only that $\ln \tau_\alpha^*(c)$ is a well-behaved function of composition that depends on the dimensionless parameters t_{rg} , m , and any other potentially relevant material parameters (e.g., the interfacial free energy, σ_{La} , of the liquid–crystal interface). In other words we take $\ln \tau_\alpha^*(c) = \ln[f(c, t_{rg}, m, \sigma_{La}, \dots)]$, where f is an unknown function of the parameters $t_{rg}(c)$, $m(c)$, etc. that are themselves functions of composition.

For a 1D composition variation z as in our alloy series, the compositional derivative of $\ln f$ [where $f = \tau_\alpha^*(c)$] for α -crystallization becomes

$$\frac{d \ln f}{dz} = \frac{1}{f} \frac{df}{dz} + \frac{1}{f} \left[\frac{df}{dt_{rg}} \frac{dt_{rg,\alpha}}{dz} + \frac{df}{dm} \frac{dm}{dz} + \frac{df}{d\sigma_{La}} \frac{d\sigma_{La}}{dz} + \dots \right]. \quad [\text{S5}]$$

The first term on the right includes any explicit dependence of f on composition not described by our parameters. For a model of crystal nucleation, this term can be assumed to vanish because the independent parameters taken together with their dependence on z and T are assumed to give a full description of the nucleation rate. The series expansion of Eq. S5 along the z direction then becomes to first order:

$$\ln(f(z)_\alpha) = \ln(f(z_0)_\alpha) + \left[\frac{df}{dt_{rg}} \frac{dt_{rg,\alpha}}{dz} + \frac{df}{dm} \frac{dm}{dz} + \frac{df}{d\sigma_{La}} \frac{d\sigma_{La}}{dz} + \dots \right]_{z_0} (z - z_0), \quad [\text{S6}]$$

Referring to Eq. S6, one sees that the term in the square bracket can be identified as the z component of the λ_α vector, where $z = \lambda_\alpha \cdot (c - c_0)$. Further, one observes that the λ_α for the alloy series naturally separates into contributions arising from each independent variable t_{rg} , m , etc. We write:

$$\lambda_\alpha = \lambda_{t_{rg,\alpha}} + \lambda_{m,\alpha} + \dots, \quad [\text{S7}]$$

where $\lambda_{t_{rg,\alpha}} = (df/dt_{rg})(dt_{rg,\alpha}/dz)$, $\lambda_{m,\alpha} = (df/dm)(dm/dz)$, etc. for each independent parameter. Eq. S7 is essentially Eq. 3a for analyzing the glass-forming ability (GFA) cusp.

