

Measurement of J/ψ and $\psi(2S)$ Prompt Double-Differential Cross Sections in pp Collisions at $\sqrt{s} = 7$ TeV

V. Khachatryan *et al.**

(CMS Collaboration)

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The double-differential cross sections of promptly produced J/ψ and $\psi(2S)$ mesons are measured in pp collisions at $\sqrt{s} = 7$ TeV, as a function of transverse momentum p_T and absolute rapidity $|y|$. The analysis uses J/ψ and $\psi(2S)$ dimuon samples collected by the CMS experiment, corresponding to integrated luminosities of 4.55 and 4.90 fb⁻¹, respectively. The results are based on a two-dimensional analysis of the dimuon invariant mass and decay length, and extend to $p_T = 120$ and 100 GeV for the J/ψ and $\psi(2S)$, respectively, when integrated over the interval $|y| < 1.2$. The ratio of the $\psi(2S)$ to J/ψ cross sections is also reported for $|y| < 1.2$, over the range $10 < p_T < 100$ GeV. These are the highest p_T values for which the cross sections and ratio have been measured.

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Studies of heavy-quarkonium production are of central importance for an improved understanding of non-perturbative quantum chromodynamics (QCD) [1]. The nonrelativistic QCD (NRQCD) effective-field-theory framework [2], arguably the best formalism at this time, factorizes high- p_T quarkonium production in short-distance and long-distance scales. First, a heavy quark-antiquark pair, $Q\bar{Q}$, is produced in a Fock state $^{2S+1}L_J^{[a]}$, with spin S , orbital angular momentum L , and total angular momentum J that are either identical to (color singlet, $a = 1$) or different from (color octet, $a = 8$) those of the corresponding quarkonium state. The $Q\bar{Q}$ cross sections are determined by short-distance coefficients (SDCs), kinematic-dependent functions calculable perturbatively as expansions in the strong-coupling constant α_s . Then this “preresonant” $Q\bar{Q}$ pair binds into the physically observable quarkonium through a nonperturbative evolution that may change L and S , with bound-state formation probabilities proportional to long-distance matrix elements (LDMEs). The LDMEs are conjectured to be constant (i.e., independent of the $Q\bar{Q}$ momentum) and universal (i.e., process independent). The color-octet terms are expected to scale with powers of the heavy-quark velocity in the $Q\bar{Q}$ rest frame. In the nonrelativistic limit, an S -wave vector quarkonium state should be formed from a $Q\bar{Q}$ pair produced as a color singlet ($^3S_1^{[1]}$) or as one of three color octets ($^1S_0^{[8]}$, $^3S_1^{[8]}$, and $^3P_J^{[8]}$).

Three “global fits” to measured quarkonium data [3–5] obtained incompatible octet LDMEs, despite the use of essentially identical theory inputs: next-to-leading-order (NLO) QCD calculations of the singlet and octet SDCs. The disagreement stems from the fact that different sets of measurements were considered. In particular, the results crucially depend on the minimum p_T of the fitted measurements [6], because the octet SDCs have different p_T dependences. Fits including low- p_T cross sections lead to the conclusion that, at high p_T , quarkonium production should be dominated by transversely polarized octet terms. This prediction is in stark contradiction with the unpolarized production seen by the CDF [7,8] and CMS [9,10] experiments, an observation known as the “quarkonium polarization puzzle.” As shown in Ref. [6], the puzzle is seemingly solved by restricting the NRQCD global fits to high- p_T quarkonia, indicating that the presently available fixed-order calculations provide SDCs that are unable to reproduce reality at lower p_T values or that NRQCD factorization only holds for p_T values much larger than the quarkonium mass. The polarization measurements add a crucial dimension to the global fits because the various channels have remarkably distinct polarization properties: in the helicity frame, $^3S_1^{[1]}$ is longitudinally polarized, $^1S_0^{[8]}$ is unpolarized, $^3S_1^{[8]}$ is transversely polarized, and $^3P_J^{[8]}$ has a polarization that changes significantly with p_T . Bottomonium and prompt charmonium polarizations reaching or exceeding $p_T = 50$ GeV were measured by CMS [9,10], using a very robust analysis framework [11,12], on the basis of event samples collected in 2011. Instead, the differential charmonium cross sections published by CMS [13] are based on data collected in 2010 and have a much lower p_T reach. Measurements of prompt charmonium cross sections extending well beyond $p_T = 50$ GeV will trigger improved NRQCD global fits, restricted to a kinematic domain where the factorization formalism is

*Full author list given at the end of the article.

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unquestioned, and will provide more accurate and reliable LDMEs.

This Letter presents measurements of the double-differential cross sections of J/ψ and $\psi(2S)$ mesons promptly produced in pp collisions at a center-of-mass energy of 7 TeV, based on dimuon event samples collected by CMS in 2011. They complement other prompt charmonium cross sections measured at the LHC, by ATLAS [14,15], LHCb [16,17], and ALICE [18]. The analysis is made in four bins of absolute rapidity ($|y| < 0.3$, $0.3 < |y| < 0.6$, $0.6 < |y| < 0.9$, and $0.9 < |y| < 1.2$) and in the p_T ranges 10–95 GeV for the J/ψ and 10–75 GeV for the $\psi(2S)$. A rapidity-integrated result in the range $|y| < 1.2$ is also provided, extending the p_T reach to 120 GeV for the J/ψ and 100 GeV for the $\psi(2S)$. The corresponding $\psi(2S)$ over J/ψ cross section ratios are also reported. The dimuon invariant mass distribution is used to separate the J/ψ and $\psi(2S)$ signals from other processes, mostly pairs of uncorrelated muons, while the dimuon decay length is used to separate the non-prompt charmonia, coming from decays of b hadrons, from the prompt component. Feed-down from decays of heavier charmonium states, approximately 33% of the prompt J/ψ cross section [19], is not distinguished from the directly produced charmonia.

The CMS apparatus is based on a superconducting solenoid of 6 m internal diameter, providing a 3.8 T field. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter. Muons are measured with three kinds of gas-ionization detectors: drift tubes, cathode strip chambers, and resistive-plate chambers. The main subdetectors used in this analysis are the silicon tracker and the muon system, which enable the measurement of muon momenta over the pseudorapidity range $|\eta| < 2.4$. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [20].

The events were collected using a two-level trigger system. The first level, made of custom hardware processors, uses data from the muon system to select events with two muon candidates. The high-level trigger, adding information from the silicon tracker, reduces the rate of stored events by requiring an opposite-sign muon pair of invariant mass $2.8 < M < 3.35$ GeV, $p_T > 9.9$ GeV, and $|y| < 1.25$ for the J/ψ trigger, and $3.35 < M < 4.05$ GeV and $p_T > 6.9$ GeV for the $\psi(2S)$ trigger. No p_T requirement is imposed on the single muons at trigger level. Both triggers require a dimuon vertex fit χ^2 probability greater than 0.5% and a distance of closest approach between the two muons less than 5 mm. Events where the muons bend towards each other in the magnetic field are rejected to lower the trigger rate while retaining the highest-quality dimuons. The J/ψ and $\psi(2S)$ analyses are conducted independently, using event samples separated at the trigger

level. The $\psi(2S)$ sample corresponds to an integrated luminosity of 4.90 fb^{-1} , while the J/ψ sample has a reduced value, 4.55 fb^{-1} , because the p_T threshold of the J/ψ trigger was raised to 12.9 GeV in a fraction of the data-taking period; the integrated luminosities have an uncertainty of 2.2% [21].

The muon tracks are required to have hits in at least eleven tracker layers, with at least two in the silicon pixel detector, and to be matched with at least one segment in the muon system. They must have a good track fit quality (χ^2 per degree of freedom smaller than 1.8) and point to the interaction region. The selected muons must also match in pseudorapidity and azimuthal angle with the muon objects responsible for triggering the event. The analysis is restricted to muons produced within a fiducial phase-space window where the muon detection efficiencies are accurately measured: $p_T > 4.5$, 3.5, and 3.0 GeV for the regions $|\eta| < 1.2$, $1.2 < |\eta| < 1.4$, and $1.4 < |\eta| < 1.6$, respectively. The combinatorial dimuon background is reduced by requiring a dimuon vertex fit χ^2 probability larger than 1%. After applying the event selection criteria, the combined yields of prompt and nonprompt charmonia in the range $|y| < 1.2$ are 5.45 M for the J/ψ and 266 k for the $\psi(2S)$. The prompt charmonia are separated from those resulting from decays of b hadrons through the use of the dimuon pseudo-proper-decay-length [22], $\ell = L_{xy}M/p_T$, where L_{xy} is the transverse decay length in the laboratory frame, measured after removing the two muon tracks from the calculation of the primary vertex position. For events with multiple collision vertices, L_{xy} is calculated with respect to the vertex closest to the direction of the dimuon momentum, extrapolated towards the beam line.

For each $(|y|, p_T)$ bin, the prompt charmonium yields are evaluated through an extended unbinned maximum-likelihood fit to the two-dimensional (M, ℓ) event distribution. In the mass dimension, the shape of each signal peak is represented by a Crystal Ball (CB) function [23], with free mean (μ_{CB}) and width (σ_{CB}) parameters. Given the strong correlation between the two CB tail parameters, α_{CB} and n_{CB} , they are fixed to values evaluated from fits to event samples integrated in broader p_T ranges. A single CB function provides a good description of the signal mass peaks, given that the dimuon mass distributions are studied in narrow $(|y|, p_T)$ bins, within which the dimuon invariant mass resolution has a negligible variation. The mass distribution of the underlying continuum background is described by an exponential function. Concerning the pseudo-proper-decay-length variable, the prompt signal component is modeled by a resolution function, which exploits the per-event uncertainty information provided by the vertex reconstruction algorithm, while the nonprompt charmonium term is modeled by an exponential function convolved with the resolution function. The continuum background component is represented by a sum of prompt

and nonprompt empirical forms. The distributions are well described with a relatively small number of free parameters.

Figure 1 shows the J/ψ and $\psi(2S)$ dimuon invariant mass and pseudo-proper-decay-length projections for two representative $(|y|, p_T)$ bins. The decay length projections are shown for events with dimuon invariant mass within $\pm 3\sigma_{CB}$ of the pole mass. In the highest p_T bins, where the number of dimuons is relatively small, stable results are obtained by fixing μ_{CB} and the slope of the exponential-like function describing the nonprompt combinatorial background to values extrapolated from the trend found from the lower- p_T bins. The systematic uncertainties in the signal yields are evaluated by repeating the fit with different functional forms, varying the values of the fixed parameters, and allowing for more free parameters in the fit. The fit results are robust with respect to changes in the procedure; the corresponding systematic uncertainties are negligible at low p_T and increase to $\approx 2\%$ for the J/ψ and $\approx 6\%$ for the $\psi(2S)$ in the highest p_T bins.

The single-muon detection efficiencies ϵ_μ are measured with a “tag-and-probe” (T&P) technique [24], using event samples collected with triggers specifically designed for

this purpose, including a sample enriched in dimuons from J/ψ decays where a muon is combined with another track and the pair is required to have an invariant mass within the range 2.8–3.4 GeV. The procedure was validated in the phase-space window of the analysis with detailed Monte Carlo (MC) simulation studies. The measured efficiencies are parametrized as a function of muon p_T , in eight bins of muon $|\eta|$. Their uncertainties, reflecting the statistical precision of the T&P samples and possible imperfections of the parametrization, are $\approx 2\%$ – 3% . The efficiency of the dimuon vertex fit χ^2 probability requirement is also measured with the T&P approach, using a sample of events collected with a dedicated (prescaled) trigger. It is around 95%–97%, improving with increasing p_T , with a 2% systematic uncertainty. At high p_T , when the two muons might be emitted relatively close to each other, the efficiency of the dimuon trigger $\epsilon_{\mu\mu}$ is smaller than the product of the two single-muon efficiencies [13], $\epsilon_{\mu\mu} = \epsilon_{\mu_1}\epsilon_{\mu_2}\rho$. The correction factor ρ is evaluated with MC simulations, validated from data collected with single-muon triggers. For $p_T < 35$ GeV, ρ is consistent with being unity, within a systematic uncertainty estimated as

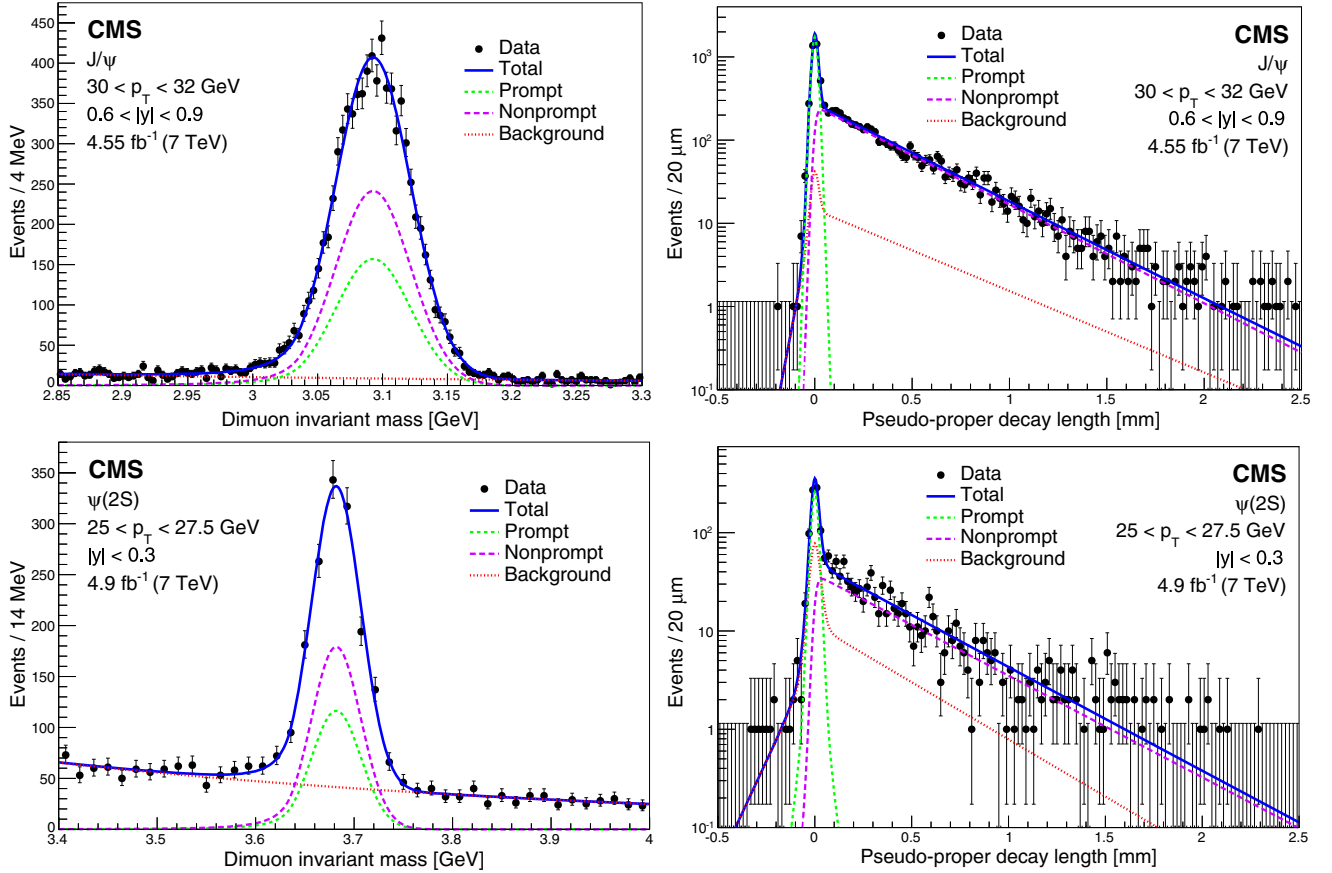


FIG. 1 (color online). Projections on the dimuon invariant mass (left) and pseudo-proper-decay-length (right) axes, for the J/ψ (top) and $\psi(2S)$ (bottom) events in the kinematic bins given in the plots. The right panels show dimuons of invariant mass within $\pm 3\sigma_{CB}$ of the pole masses. The curves, identified in the legends, represent the result of the fits described in the text. The vertical bars on the data points show the statistical uncertainties.

2%, except in the $0.9 < |y| < 1.2$ bin, where the uncertainty increases to 4.3% for the J/ψ if $p_T < 12$ GeV, and to 2.7% for the $\psi(2S)$ if $p_T < 11$ GeV. For $p_T > 35$ GeV, ρ decreases approximately linearly with p_T , reaching 60%–70% for $p_T \sim 85$ GeV, with systematic uncertainties evaluated by comparing the MC simulation results with estimations made using data collected with single-muon triggers: 5% up to $p_T = 50$ (55) GeV for the J/ψ [$\psi(2S)$] and 10% for higher p_T . The total dimuon detection efficiency increases from $\epsilon_{\mu\mu} \approx 78\%$ at $p_T = 15$ GeV to $\approx 85\%$ at 30 GeV, and then decreases to $\approx 65\%$ at 80 GeV.

To obtain the charmonium cross sections in each $(|y|, p_T)$ bin without any restrictions on the kinematic variables of the two muons, we correct for the corresponding dimuon acceptance, defined as the fraction of dimuon decays having both muons emitted within the single-muon fiducial phase space. These acceptances are calculated using a detailed MC simulation of the CMS experiment. Charmonia are generated using a flat rapidity distribution and p_T distributions based on previous measurements [13]; using flat p_T distributions leads to negligible changes. The particles are decayed by EVTGEN [25] interfaced to PYTHIA 6.4 [26], while PHOTOS [27] is used to simulate final-state radiation. The fractions of J/ψ and $\psi(2S)$ dimuon events in a given $(|y|, p_T)$ bin with both muons surviving the fiducial selections depend on the decay kinematics and, in particular, on the polarization of the mother particle. Acceptances are calculated using polarization scenarios corresponding to different values of the polar anisotropy parameter in the helicity frame, λ_θ^{HX} : 0 (unpolarized), +1 (transverse), and −1 (longitudinal). A fourth scenario, corresponding to $\lambda_\theta^{HX} = +0.10$ for the J/ψ and +0.03 for the $\psi(2S)$, reflects the results published by CMS [10]. The two other parameters characterizing the dimuon angular distributions [28], λ_ϕ and $\lambda_{\phi\psi}$, have been measured to be essentially zero [10] and have a negligible influence on the acceptance. The acceptances are essentially identical for the two charmonia and are almost rapidity independent for $|y| < 1.2$. The two-dimensional acceptance maps are calculated with large MC simulation samples, so that statistical fluctuations are small, and in narrow $|y|$ bins, so that variations within the bins can be neglected. Since the efficiencies and acceptances are evaluated for events where the two muons bend away from each other, a factor of 2 is applied to obtain the final cross sections.

The double-differential cross sections of promptly produced J/ψ and $\psi(2S)$ in the dimuon channel, $\mathcal{B} d^2\sigma/dp_T dy$, where \mathcal{B} is the J/ψ or $\psi(2S)$ dimuon branching fraction, are obtained by dividing the fitted prompt-signal yields, already corrected on an event-by-event basis for efficiencies and acceptance, by the integrated luminosity and the widths of the p_T and $|y|$ bins. The numerical values, including the relative statistical and systematic uncertainties, are reported for both charmonia, five rapidity intervals, and four polarization scenarios in Tables 1–4 of the Supplemental Material [29]. Figure 2

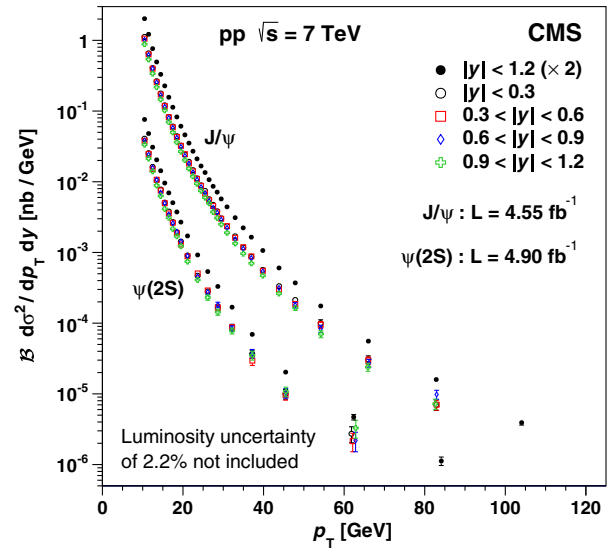


FIG. 2 (color online). The J/ψ and $\psi(2S)$ differential p_T cross sections times the dimuon branching fractions for four rapidity bins and integrated over the range $|y| < 1.2$ (scaled up by a factor of 2 for presentation purposes), assuming the unpolarized scenario. The vertical bars show the statistical and systematic uncertainties added in quadrature.

shows the results obtained in the unpolarized scenario. With respect to the $|y| < 0.3$ bin, the cross sections drop by $\approx 5\%$ for $0.6 < |y| < 0.9$ and $\approx 15\%$ for $0.9 < |y| < 1.2$. Measuring the charmonium production cross sections in the broader rapidity range $|y| < 1.2$ has the advantage that the increased statistical accuracy allows the measurement to be extended to higher- p_T values, where comparisons with theoretical calculations are particularly informative. Figure 3 compares the rapidity-integrated (unpolarized) cross sections, after rescaling with the branching fraction \mathcal{B} of the dimuon decay channels [30], with results reported by ATLAS [14,15]. The curve represents a fit of the J/ψ cross section measured in this analysis to a power-law function [31]. The band labeled FKLSW represents the result of a global fit [6] comparing SDCs calculated at NLO [3] with $\psi(2S)$ cross sections and polarizations previously reported by CMS [10,13] and LHCb [17]. According to that fit, $\psi(2S)$ mesons are produced predominantly unpolarized. At high p_T , the values reported in this Letter tend to be higher than the band, which is essentially determined from results for $p_T < 30$ GeV.

The ratio of the $\psi(2S)$ to J/ψ differential cross sections is also measured in the $|y| < 1.2$ range, recomputing the J/ψ values in the p_T bins of the $\psi(2S)$ analysis. The measured values are reported in Table 5 of the Supplemental Material [29]. The corrections owing to the integrated luminosity, acceptances, and efficiencies cancel to a large extent in the measurement of the ratio. The total systematic uncertainty, dominated by the ρ correction for $p_T > 30$ GeV and by the acceptance and

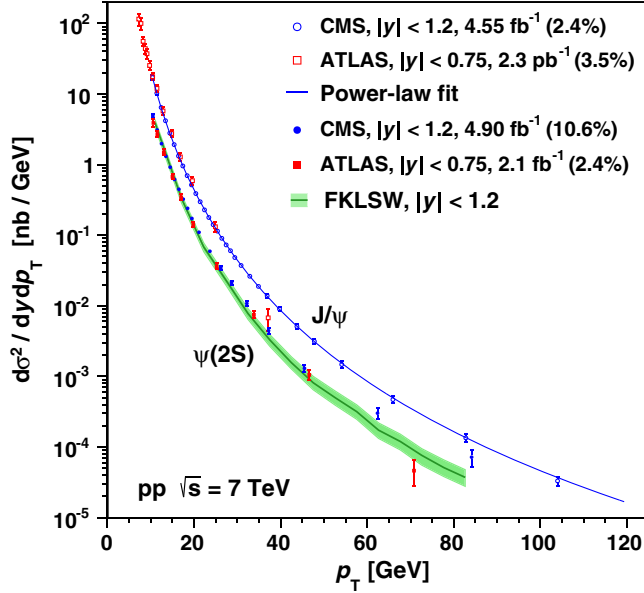


FIG. 3 (color online). The J/ψ (open symbols) and $\psi(2S)$ (closed symbols) differential (unpolarized) cross sections from this analysis (circles) and from ATLAS (squares) [14,15]. The vertical bars show the statistical and systematic uncertainties added in quadrature, not including the uncertainties from integrated luminosities and branching fractions, which are indicated by the percentages given in the legend. The curve shows a fit of the J/ψ cross section measured in this analysis to a power-law function, while the band labeled FKLSW represents a calculation of the $\psi(2S)$ cross section using LDMEs determined with lower- p_T LHC data [6].

efficiency corrections for $p_T < 20$ GeV, does not exceed 3%, except for $p_T > 75$ GeV, where it reaches 5%. Larger event samples are needed to clarify the trend of the ratio for p_T above ≈ 35 GeV.

In summary, the double-differential cross sections of the J/ψ and $\psi(2S)$ mesons promptly produced in pp collisions at $\sqrt{s} = 7$ TeV have been measured as a function of p_T in four $|y|$ bins, as well as integrated over the $|y| < 1.2$ range, extending up to or beyond $p_T = 100$ GeV. New global fits of cross sections and polarizations, including these high- p_T measurements, will probe the theoretical calculations in a kinematical region where NRQCD factorization is believed to be most reliable. The new data should also provide input to stringent tests of recent theory developments, such as those described in Refs. [32–34].

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V. Khachatryan,¹ A. M. Sirunyan,¹ A. Tumasyan,¹ W. Adam,² T. Bergauer,² M. Dragicevic,² J. Erö,² M. Friedl,² R. Frühwirth,^{2,b} V. M. Ghete,² C. Hartl,² N. Hörmann,² J. Hrubec,² M. Jeitler,^{2,b} W. Kiesenhofer,² V. Knünz,² M. Krammer,^{2,b} I. Krätschmer,² D. Liko,² I. Mikulec,² D. Rabady,^{2,c} B. Rahbaran,² H. Rohringer,² R. Schöffbeck,² J. Strauss,² W. Treberer-Treberspurg,² W. Waltenberger,² C.-E. Wulz,^{2,b} V. Mossolov,³ N. Shumeiko,³ J. Suarez Gonzalez,³ S. Alderweireldt,⁴ S. Bansal,⁴ T. Cornelis,⁴ E. A. De Wolf,⁴ X. Janssen,⁴ A. Knutsson,⁴ J. Lauwers,⁴ S. Luyckx,⁴ S. Ochesanu,⁴ R. Rougny,⁴ M. Van De Klundert,⁴ H. Van Haevermaet,⁴ P. Van Mechelen,⁴ N. Van Remortel,⁴ A. Van Spillbeeck,⁴ F. Blekman,⁵ S. Blyweert,⁵ J. D'Hondt,⁵ N. Daci,⁵ N. Heracleous,⁵ J. Keaveney,⁵ S. Lowette,⁵ M. Maes,⁵ A. Olbrechts,⁵ Q. Python,⁵ D. Strom,⁵ S. Tavernier,⁵ W. Van Doninck,⁵ P. Van Mulders,⁵ G. P. Van Onsem,⁵ I. Villella,⁵ C. Caillol,⁶ B. Clerbaux,⁶ G. De Lentdecker,⁶ D. Dobur,⁶ L. Favart,⁶ A. P. R. Gay,⁶ A. Grebenyuk,⁶ A. Léonard,⁶ A. Mohammadi,⁶ L. Perniè,^{6,c} A. Randle-conde,⁶ T. Reis,⁶ T. Seva,⁶ L. Thomas,⁶ C. Vander Velde,⁶ P. Vanlaer,⁶ J. Wang,⁶ F. Zenoni,⁶ V. Adler,⁷ K. Beernaert,⁷ L. Benucci,⁷ A. Cimmino,⁷ S. Costantini,⁷ S. Crucy,⁷ A. Fagot,⁷ G. Garcia,⁷ J. McCartin,⁷ A. A. Ocampo Rios,⁷ D. Poyraz,⁷ D. Ryckbosch,⁷ S. Salva Diblen,⁷ M. Sigamani,⁷ N. Strobbe,⁷ F. Thyssen,⁷ M. Tytgat,⁷ E. Yazgan,⁷ N. Zaganidis,⁷ S. Basegmez,⁸ C. Beluffi,^{8,d} G. Bruno,⁸ R. Castello,⁸ A. Caudron,⁸ L. Ceard,⁸ G. G. Da Silveira,⁸ C. Delaere,⁸ T. du Pree,⁸ D. Favart,⁸ L. Forthomme,⁸ A. Giammanco,^{8,e} J. Hollar,⁸ A. Jafari,⁸ P. Jez,⁸ M. Komm,⁸ V. Lemaître,⁸ C. Nuttens,⁸ D. Pagano,⁸ L. Perrini,⁸ A. Pin,⁸ K. Piotrkowski,⁸ A. Popov,^{8,f} L. Quertenmont,⁸ M. Selvaggi,⁸ M. Vidal Marono,⁸ J. M. Vizán García,⁸ N. Belyi,⁹ T. Caebergs,⁹ E. Daubie,⁹ G. H. Hammad,⁹

W. L. Aldá Júnior,¹⁰ G. A. Alves,¹⁰ L. Brito,¹⁰ M. Correa Martins Junior,¹⁰ T. Dos Reis Martins,¹⁰ J. Molina,¹⁰ C. Mora Herrera,¹⁰ M. E. Pol,¹⁰ P. Rebello Teles,¹⁰ W. Carvalho,¹¹ J. Chinellato,^{11,g} A. Custódio,¹¹ E. M. Da Costa,¹¹ D. De Jesus Damiao,¹¹ C. De Oliveira Martins,¹¹ S. Fonseca De Souza,¹¹ H. Malbouisson,¹¹ D. Matos Figueiredo,¹¹ L. Mundim,¹¹ H. Nogima,¹¹ W. L. Prado Da Silva,¹¹ J. Santaolalla,¹¹ A. Santoro,¹¹ A. Sznajder,¹¹ E. J. Tonelli Manganote,^{11,g} A. Vilela Pereira,¹¹ C. A. Bernardes,^{12b} S. Dogra,^{12a} T. R. Fernandez Perez Tomei,^{12a} E. M. Gregores,^{12b} P. G. Mercadante,^{12b} S. F. Novaes,^{12a} Sandra S. Padula,^{12a} A. Aleksandrov,¹³ V. Genchev,^{13,c} R. Hadjiiska,¹³ P. Iaydjiev,¹³ A. Marinov,¹³ S. Piperov,¹³ M. Rodozov,¹³ S. Stoykova,¹³ G. Sultanov,¹³ M. Vutova,¹³ A. Dimitrov,¹⁴ I. Glushkov,¹⁴ L. Litov,¹⁴ B. Pavlov,¹⁴ P. Petkov,¹⁴ J. G. Bian,¹⁵ G. M. Chen,¹⁵ H. S. Chen,¹⁵ M. Chen,¹⁵ T. Cheng,¹⁵ R. Du,¹⁵ C. H. Jiang,¹⁵ R. Plestina,^{15,h} F. Romeo,¹⁵ J. Tao,¹⁵ Z. Wang,¹⁵ C. Asawatangtrakuldee,¹⁶ Y. Ban,¹⁶ W. Guo,¹⁶ S. Liu,¹⁶ Y. Mao,¹⁶ S. J. Qian,¹⁶ D. Wang,¹⁶ Z. Xu,¹⁶ F. Zhang,^{16,i} L. Zhang,¹⁶ W. Zou,¹⁶ C. Avila,¹⁷ A. Cabrera,¹⁷ L. F. Chaparro Sierra,¹⁷ C. Florez,¹⁷ J. P. Gomez,¹⁷ B. Gomez Moreno,¹⁷ J. C. Sanabria,¹⁷ N. Godinovic,¹⁸ D. Lelas,¹⁸ D. Polic,¹⁸ I. Puljak,¹⁸ Z. Antunovic,¹⁹ M. Kovac,¹⁹ V. Brigljevic,²⁰ K. Kadija,²⁰ J. Luetic,²⁰ D. Mekterovic,²⁰ L. Sudic,²⁰ A. Attikis,²¹ G. Mavromanolakis,²¹ J. Mousa,²¹ C. Nicolaou,²¹ F. Ptochos,²¹ P. A. Razis,²¹ H. Rykaczewski,²¹ M. Bodlak,²² M. Finger,²² M. Finger Jr.,^{22,j} Y. Assran,^{23,k} A. Ellithi Kamel,^{23,l} M. A. Mahmoud,^{23,m} A. Radi,^{23,n,o} M. Kadastik,²⁴ M. Murumaa,²⁴ M. Raidal,²⁴ A. Tiko,²⁴ P. Eerola,²⁵ M. Voutilainen,²⁵ J. Härkönen,²⁶ V. Karimäki,²⁶ R. Kinnunen,²⁶ M. J. Kortelainen,²⁶ T. Lampén,²⁶ K. Lassila-Perini,²⁶ S. Lehti,²⁶ T. Lindén,²⁶ P. Luukka,²⁶ T. Mäenpää,²⁶ T. Peltola,²⁶ E. Tuominen,²⁶ J. Tuominiemi,²⁶ E. Tuovinen,²⁶ L. Wendland,²⁶ J. 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Carrillo Montoya,³² J. Chasserat,³² R. Chierici,³² D. Contardo,^{32,c} B. Courbon,³² P. Depasse,³² H. El Mamouni,³² J. Fan,³² J. Fay,³² S. Gascon,³² M. Gouzevitch,³² B. Ille,³² T. Kurca,³² M. Lethuillier,³² L. Mirabito,³² A. L. Pequegnot,³² S. Perries,³² J. D. Ruiz Alvarez,³² D. Sabes,³² L. Sgandurra,³² V. Sordini,³² M. Vander Donckt,³² P. Verdier,³² S. Viret,³² H. Xiao,³² Z. Tsamalaidze,^{33,j} C. Autermann,³⁴ S. Beranek,³⁴ M. Bontenackels,³⁴ M. Edelhoff,³⁴ L. Feld,³⁴ A. Heister,³⁴ K. Klein,³⁴ M. Lipinski,³⁴ A. Ostapchuk,³⁴ M. Preuten,³⁴ F. Raupach,³⁴ J. Sammet,³⁴ S. Schael,³⁴ J. F. Schulte,³⁴ H. Weber,³⁴ B. Wittmer,³⁴ V. Zhukov,^{34,f} M. Ata,³⁵ M. Brodski,³⁵ E. Dietz-Laursonn,³⁵ D. Duchardt,³⁵ M. Erdmann,³⁵ R. Fischer,³⁵ A. Güth,³⁵ T. Hebbeker,³⁵ C. Heidemann,³⁵ K. Hoepfner,³⁵ D. Klingebiel,³⁵ S. Knutzen,³⁵ P. Kreuzer,³⁵ M. Merschmeyer,³⁵ A. Meyer,³⁵ P. Millet,³⁵ M. Olschewski,³⁵ K. Padeken,³⁵ P. Papacz,³⁵ H. Reithler,³⁵ S. A. Schmitz,³⁵ L. Sonnenschein,³⁵ D. Teyssier,³⁵ S. Thüer,³⁵ V. Cherepanov,³⁶ Y. Erdogan,³⁶ G. Flügge,³⁶ H. Geenen,³⁶ M. Geisler,³⁶ W. Haj Ahmad,³⁶ F. Hoehle,³⁶ B. Kargoll,³⁶ T. Kress,³⁶ Y. Kuessel,³⁶ A. Künsken,³⁶ J. Lingemann,^{36,c} A. Nowack,³⁶ I. M. Nugent,³⁶ C. Pistone,³⁶ O. Pooth,³⁶ A. Stahl,³⁶ M. Aldaya Martin,³⁷ I. Asin,³⁷ N. Bartosik,³⁷ J. Behr,³⁷ U. Behrens,³⁷ A. J. Bell,³⁷ A. Bethani,³⁷ K. Borras,³⁷ A. Burgmeier,³⁷ A. Cakir,³⁷ L. Calligaris,³⁷ A. Campbell,³⁷ S. Choudhury,³⁷ F. Costanza,³⁷ C. Diez Pardos,³⁷ G. Dolinska,³⁷ S. Dooling,³⁷ T. Dorland,³⁷ G. Eckerlin,³⁷ D. Eckstein,³⁷ T. Eichhorn,³⁷ G. Flucke,³⁷ J. Garay Garcia,³⁷ A. Geiser,³⁷ A. Gikhko,³⁷ P. Gunnellini,³⁷ J. Hauk,³⁷ M. Hempel,^{37,q} H. Jung,³⁷ A. Kalogeropoulos,³⁷ O. Karacheban,^{37,q} M. Kasemann,³⁷ P. Katsas,³⁷ J. Kieseler,³⁷ C. Kleinwort,³⁷ I. Korol,³⁷ D. Krücker,³⁷ W. Lange,³⁷ J. Leonard,³⁷ K. Lipka,³⁷ A. Lobanov,³⁷ W. Lohmann,^{37,q} B. Lutz,³⁷ R. Mankel,³⁷ I. Marfin,^{37,q} I.-A. Melzer-Pellmann,³⁷ A. B. Meyer,³⁷ G. Mittag,³⁷ J. Mnich,³⁷ A. Mussgiller,³⁷ S. Naumann-Emme,³⁷ A. Nayak,³⁷ E. Ntomari,³⁷ H. Perrey,³⁷ D. Pitzl,³⁷ R. Placakyte,³⁷ A. Raspereza,³⁷ P. M. Ribeiro Cipriano,³⁷ B. Roland,³⁷ E. Ron,³⁷ M. Ö. Sahin,³⁷ J. Salfeld-Nebgen,³⁷ P. Saxena,³⁷ T. Schoerner-Sadenius,³⁷ M. Schröder,³⁷ C. Seitz,³⁷ S. Spannagel,³⁷ A. D. R. Vargas Trevino,³⁷ R. Walsh,³⁷ C. Wissing,³⁷ V. Blobel,³⁸ M. Centis Vignali,³⁸ A. R. Draeger,³⁸ J. Erfle,³⁸ E. Garutti,³⁸ K. Goebel,³⁸ M. Görner,³⁸ J. Haller,³⁸ M. Hoffmann,³⁸ R. S. Höing,³⁸ A. Junkes,³⁸ H. Kirschenmann,³⁸ R. Klanner,³⁸ R. Kogler,³⁸ T. Lapsien,³⁸ T. Lenz,³⁸ I. Marchesini,³⁸ D. Marconi,³⁸ J. Ott,³⁸ T. Peiffer,³⁸ A. Perieanu,³⁸ N. Pietsch,³⁸ J. Poehlsen,³⁸ T. Poehlsen,³⁸ D. Rathjens,³⁸ C. Sander,³⁸ H. Schettler,³⁸ P. Schleper,³⁸ E. Schlieckau,³⁸ A. Schmidt,³⁸ M. Seidel,³⁸ V. Sola,³⁸ H. Stadie,³⁸ G. Steinbrück,³⁸ D. Troendle,³⁸ E. Usai,³⁸ L. Vanelderen,³⁸ A. Vanhoefer,³⁸ C. Barth,³⁹ C. Baus,³⁹ J. Berger,³⁹ C. Böser,³⁹ E. Butz,³⁹

T. Chwalek,³⁹ W. De Boer,³⁹ A. Descroix,³⁹ A. Dierlamm,³⁹ M. Feindt,³⁹ F. Frensch,³⁹ M. Giffels,³⁹ A. Gilbert,³⁹ F. Hartmann,^{39,c} T. Hauth,³⁹ U. Husemann,³⁹ I. Katkov,^{39,f} A. Kornmayer,^{39,c} P. Lobelle Pardo,³⁹ M. U. Mozer,³⁹ T. Müller,³⁹ Th. Müller,³⁹ A. Nürnberg,³⁹ G. Quast,³⁹ K. Rabbertz,³⁹ S. Röcker,³⁹ H. J. Simonis,³⁹ F. M. Stober,³⁹ R. Ulrich,³⁹ J. Wagner-Kuhr,³⁹ S. Wayand,³⁹ T. Weiler,³⁹ R. Wolf,³⁹ G. Anagnostou,⁴⁰ G. Daskalakis,⁴⁰ T. Geralis,⁴⁰ V. A. Giakoumopoulou,⁴⁰ A. Kyriakis,⁴⁰ D. Loukas,⁴⁰ A. Markou,⁴⁰ C. Markou,⁴⁰ A. Psallidas,⁴⁰ I. Topsis-Giotis,⁴⁰ A. Agapitos,⁴¹ S. Kesisoglou,⁴¹ A. Panagiotou,⁴¹ N. Saoulidou,⁴¹ E. Stiliaris,⁴¹ E. Tziaferi,⁴¹ X. Aslanoglou,⁴² I. Evangelou,⁴² G. Flouris,⁴² C. Foudas,⁴² P. Kokkas,⁴² N. Manthos,⁴² I. Papadopoulos,⁴² E. Paradas,⁴² J. Strologas,⁴² G. Bencze,⁴³ C. Hajdu,⁴³ P. Hidas,⁴³ D. Horvath,^{43,r} F. Sikler,⁴³ V. Veszpremi,⁴³ G. Vesztergombi,^{43,s} A. J. Zsigmond,⁴³ N. Beni,⁴⁴ S. Czellar,⁴⁴ J. Karancsi,^{44,t} J. Molnar,⁴⁴ J. Palinkas,⁴⁴ Z. Szillasi,⁴⁴ A. Makovec,⁴⁵ P. Raics,⁴⁵ Z. L. Trocsanyi,⁴⁵ B. Ujvari,⁴⁵ S. K. Swain,⁴⁶ S. B. Beri,⁴⁷ V. Bhatnagar,⁴⁷ R. Gupta,⁴⁷ U. Bhawandeep,⁴⁷ A. K. Kalsi,⁴⁷ M. Kaur,⁴⁷ R. Kumar,⁴⁷ M. Mittal,⁴⁷ N. Nishu,⁴⁷ J. B. Singh,⁴⁷ Ashok Kumar,⁴⁸ Arun Kumar,⁴⁸ S. Ahuja,⁴⁸ A. Bhardwaj,⁴⁸ B. C. Choudhary,⁴⁸ A. Kumar,⁴⁸ S. Malhotra,⁴⁸ M. Naimuddin,⁴⁸ K. Ranjan,⁴⁸ V. Sharma,⁴⁸ S. Banerjee,⁴⁹ S. Bhattacharya,⁴⁹ K. Chatterjee,⁴⁹ S. Dutta,⁴⁹ B. Gomber,⁴⁹ Sa. Jain,⁴⁹ Sh. Jain,⁴⁹ R. Khurana,⁴⁹ A. Modak,⁴⁹ S. Mukherjee,⁴⁹ D. Roy,⁴⁹ S. Sarkar,⁴⁹ M. Sharan,⁴⁹ A. Abdulsalam,⁵⁰ D. Dutta,⁵⁰ V. Kumar,⁵⁰ A. K. Mohanty,^{50,c} L. M. Pant,⁵⁰ P. Shukla,⁵⁰ A. Topkar,⁵⁰ T. Aziz,⁵¹ S. Banerjee,⁵¹ S. Bhowmik,^{51,u} R. M. Chatterjee,⁵¹ R. K. Dewanee,⁵¹ S. Dugad,⁵¹ S. Ganguly,⁵¹ S. Ghosh,⁵¹ M. Guchait,⁵¹ A. Gurtu,^{51,v} G. Kole,⁵¹ S. Kumar,⁵¹ M. Maity,^{51,u} G. Majumder,⁵¹ K. Mazumdar,⁵¹ G. B. Mohanty,⁵¹ B. Parida,⁵¹ K. Sudhakar,⁵¹ N. Wickramage,^{51,w} S. Sharma,⁵² H. Bakhshiansohi,⁵³ H. Behnamian,⁵³ S. M. Etesami,^{53,x} A. Fahim,^{53,y} R. Goldouzian,⁵³ M. Khakzad,⁵³ M. Mohammadi Najafabadi,⁵³ M. Naseri,⁵³ S. Paktinat Mehdiabadi,⁵³ F. Rezaei Hosseinabadi,⁵³ B. Safarzadeh,^{53,z} M. Zeinali,⁵³ M. Felcini,⁵⁴ M. Grunewald,⁵⁴ M. Abbrescia,^{55a,55b} C. Calabria,^{55a,55b} S. S. Chhibra,^{55a,55b} A. Colaleo,^{55a} D. Creanza,^{55a,55c} L. Cristella,^{55a,55b} N. De Filippis,^{55a,55c} M. De Palma,^{55a,55b} L. Fiore,^{55a} G. Iaselli,^{55a,55c} G. Maggi,^{55a,55c} M. Maggi,^{55a} S. My,^{55a,55c} S. Nuzzo,^{55a,55b} A. Pompili,^{55a,55b} G. Pugliese,^{55a,55c} R. Radogna,^{55a,55b,c} G. Selvaggi,^{55a,55b} A. Sharma,^{55a} L. Silvestris,^{55a,c} R. Venditti,^{55a,55b} P. Verwilligen,^{55a} G. Abbiendi,^{56a} A. C. Benvenuti,^{56a} D. Bonacorsi,^{56a,56b} S. Braibant-Giacomelli,^{56a,56b} L. Brigliadori,^{56a,56b} R. Campanini,^{56a,56b} P. Capiluppi,^{56a,56b} A. Castro,^{56a,56b} F. R. Cavallo,^{56a} G. Codispoti,^{56a,56b} M. Cuffiani,^{56a,56b} G. M. Dallavalle,^{56a} F. Fabbri,^{56a} A. Fanfani,^{56a,56b} D. Fasanella,^{56a,56b} P. Giacomelli,^{56a} C. Grandi,^{56a} L. Guiducci,^{56a,56b} S. Marcellini,^{56a} G. Masetti,^{56a} A. Montanari,^{56a} F. L. Navarria,^{56a,56b} A. Perrotta,^{56a} A. M. Rossi,^{56a,56b} T. Rovelli,^{56a,56b} G. P. Siroli,^{56a,56b} N. Tosi,^{56a,56b} R. Travaglini,^{56a,56b} S. Albergo,^{57a,57b} G. Cappello,^{57a} M. Chiorboli,^{57a,57b} S. Costa,^{57a,57b} F. Giordano,^{57a,c} R. Potenza,^{57a,57b} A. Tricomi,^{57a,57b} C. Tuve,^{57a,57b} G. Barbagli,^{58a} V. Ciulli,^{58a,58b} C. Civinini,^{58a} R. D'Alessandro,^{58a,58b} E. Focardi,^{58a,58b} E. Gallo,^{58a} S. Gonzi,^{58a,58b} V. Gori,^{58a,58b} P. Lenzi,^{58a,58b} M. Meschini,^{58a} S. Paoletti,^{58a} G. Sguazzoni,^{58a} A. Tropiano,^{58a,58b} L. Benussi,⁵⁹ S. Bianco,⁵⁹ F. Fabbri,⁵⁹ D. Piccolo,⁵⁹ R. Ferretti,^{60a,60b} F. Ferro,^{60a} M. Lo Vetere,^{60a,60b} E. Robutti,^{60a} S. Tosi,^{60a,60b} M. E. Dinardo,^{61a,61b} S. Fiorendi,^{61a,61b} S. Gennai,^{61a,c} R. Gerosa,^{61a,61b,c} A. Ghezzi,^{61a,61b} P. Govoni,^{61a,61b} M. T. Lucchini,^{61a,61b,c} S. Malvezzi,^{61a} R. A. Manzoni,^{61a,61b} A. Martelli,^{61a,61b} B. Marzocchi,^{61a,61b,c} D. Menasce,^{61a} L. Moroni,^{61a} M. Paganoni,^{61a,61b} D. Pedrini,^{61a} S. Ragazzi,^{61a,61b} N. Redaelli,^{61a} T. Tabarelli de Fatis,^{61a,61b} S. Buontempo,^{62a} N. Cavallo,^{62a,62c} S. Di Guida,^{62a,62d,c} F. Fabozzi,^{62a,62c} A. O. M. Iorio,^{62a,62b} L. Lista,^{62a} S. Meola,^{62a,62d,c} M. Merola,^{62a} P. Paolucci,^{62a,c} P. Azzi,^{63a} N. Bacchetta,^{63a} D. Bisello,^{63a,63b} R. Carlin,^{63a,63b} P. Checchia,^{63a} M. Dall'Osso,^{63a,63b} T. Dorigo,^{63a} U. Dosselli,^{63a} F. Gasparini,^{63a,63b} U. Gasparini,^{63a,63b} A. Gozzelino,^{63a} S. Lacaprara,^{63a} M. Margoni,^{63a,63b} A. T. Meneguzzo,^{63a,63b} F. Montecassiano,^{63a} M. Passaseo,^{63a} J. Pazzini,^{63a,63b} N. Pozzobon,^{63a,63b} P. Ronchese,^{63a,63b} F. Simonetto,^{63a,63b} E. Torassa,^{63a} M. Tosi,^{63a,63b} P. Zotto,^{63a,63b} A. Zucchetta,^{63a,63b} G. Zumerle,^{63a,63b} M. Gabusi,^{64a,64b} S. P. Ratti,^{64a,64b} V. Re,^{64a} C. Riccardi,^{64a,64b} P. Salvini,^{64a} P. Vitulo,^{64a,64b} M. Biasini,^{65a,65b} G. M. Bilei,^{65a} D. Ciangottini,^{65a,65b,c} L. Fanò,^{65a,65b} P. Lariccia,^{65a,65b} G. Mantovani,^{65a,65b} M. Menichelli,^{65a} A. Saha,^{65a} A. Santocchia,^{65a,65b} A. Spiezia,^{65a,65b,c} K. Androsov,^{66a,aa} P. Azzurri,^{66a} G. Bagliesi,^{66a} J. Bernardini,^{66a} T. Boccali,^{66a} G. Broccolo,^{66a,66c} R. Castaldi,^{66a} M. A. Ciocci,^{66a,aa} R. Dell'Orso,^{66a} S. Donato,^{66a,66c,c} G. Fedi,^{66a} F. Fiori,^{66a,66c} L. Foà,^{66a,66c} A. Giassi,^{66a} M. T. Grippo,^{66a,aa} F. Ligabue,^{66a,66c} T. Lomtadze,^{66a} L. Martini,^{66a,66b} A. Messineo,^{66a,66b} C. S. Moon,^{66a,bb} F. Palla,^{66a,c} A. Rizzi,^{66a,66b} A. Savoy-Navarro,^{66a,cc} A. T. Serban,^{66a} P. Spagnolo,^{66a} P. Squillacioti,^{66a,aa} R. Tenchini,^{66a} G. 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A. Degano,^{68a,68b} N. Demaria,^{68a} L. Finco,^{68a,68b,c} C. Mariotti,^{68a} S. Maselli,^{68a} E. Migliore,^{68a,68b} V. Monaco,^{68a,68b} M. Musich,^{68a} M. M. Obertino,^{68a,68c} L. Pacher,^{68a,68b} N. Pastrone,^{68a} M. Pelliccioni,^{68a} G. L. Pinna Angioni,^{68a,68b} A. Potenza,^{68a,68b} A. Romero,^{68a,68b} M. Ruspà,^{68a,68c} R. Sacchi,^{68a,68b} A. Solano,^{68a,68b} A. Staiano,^{68a} U. Tamponi,^{68a} S. Belforte,^{69a} V. Candelise,^{69a,69b,c} M. Casarsa,^{69a} F. Cossutti,^{69a} G. Della Ricca,^{69a,69b} B. Gobbo,^{69a} C. La Licata,^{69a,69b} M. Marone,^{69a,69b} A. Schizzi,^{69a,69b} T. Umer,^{69a,69b} A. Zanetti,^{69a} S. Chang,⁷⁰ A. Kropivnitskaya,⁷⁰ S. K. Nam,⁷⁰ D. H. Kim,⁷¹ G. N. Kim,⁷¹ M. S. Kim,⁷¹ D. J. Kong,⁷¹ S. Lee,⁷¹ Y. D. Oh,⁷¹ H. Park,⁷¹ A. Sakharov,⁷¹ D. C. Son,⁷¹ T. J. Kim,⁷² M. S. Ryu,⁷² J. Y. Kim,⁷³ D. H. Moon,⁷³ S. Song,⁷³ S. Choi,⁷⁴ D. Gyun,⁷⁴ B. Hong,⁷⁴ M. Jo,⁷⁴ H. Kim,⁷⁴ Y. Kim,⁷⁴ B. Lee,⁷⁴ K. S. Lee,⁷⁴ S. K. Park,⁷⁴ Y. Roh,⁷⁴ H. D. Yoo,⁷⁵ M. Choi,⁷⁶ J. H. Kim,⁷⁶ I. C. Park,⁷⁶ G. Ryu,⁷⁶ Y. Choi,⁷⁷ Y. K. Choi,⁷⁷ J. Goh,⁷⁷ D. Kim,⁷⁷ E. Kwon,⁷⁷ J. Lee,⁷⁷ I. Yu,⁷⁷ A. Juodagalvis,⁷⁸ J. R. Komaragiri,⁷⁹ M. A. B. Md Ali,^{79,dd} W. A. T. Wan Abdullah,⁷⁹ E. Casimiro Linares,⁸⁰ H. Castilla-Valdez,⁸⁰ E. De La Cruz-Burelo,⁸⁰ I. Heredia-de La Cruz,⁸⁰ A. Hernandez-Almada,⁸⁰ R. Lopez-Fernandez,⁸⁰ A. Sanchez-Hernandez,⁸⁰ S. Carrillo Moreno,⁸¹ F. Vazquez Valencia,⁸¹ I. Pedraza,⁸² H. A. Salazar Ibarguen,⁸² A. Morelos Pineda,⁸³ D. Krofcheck,⁸⁴ P. H. Butler,⁸⁵ S. Reucroft,⁸⁵ A. Ahmad,⁸⁶ M. Ahmad,⁸⁶ Q. Hassan,⁸⁶ H. R. Hoorani,⁸⁶ W. A. Khan,⁸⁶ T. Khurshid,⁸⁶ M. Shoaib,⁸⁶ H. Bialkowska,⁸⁷ M. Bluj,⁸⁷ B. Boimska,⁸⁷ T. Frueboes,⁸⁷ M. Górski,⁸⁷ M. Kazana,⁸⁷ K. Nawrocki,⁸⁷ K. Romanowska-Rybinska,⁸⁷ M. Szleper,⁸⁷ P. Zalewski,⁸⁷ G. Brona,⁸⁸ K. Bunkowski,⁸⁸ M. Cwiok,⁸⁸ W. Dominik,⁸⁸ K. Doroba,⁸⁸ A. Kalinowski,⁸⁸ M. Konecki,⁸⁸ J. Krolikowski,⁸⁸ M. Misiura,⁸⁸ M. Olszewski,⁸⁸ P. Bargassa,⁸⁹ C. Beirão Da Cruz E Silva,⁸⁹ P. Faccioli,⁸⁹ P. G. Ferreira Parracho,⁸⁹ M. Gallinaro,⁸⁹ L. Lloret Iglesias,⁸⁹ F. Nguyen,⁸⁹ J. Rodrigues Antunes,⁸⁹ J. Seixas,⁸⁹ D. Vadrucio,⁸⁹ J. Varela,⁸⁹ P. Vischia,⁹⁰ S. Afanasiev,⁹⁰ I. Golutvin,⁹⁰ V. Karjavin,⁹⁰ V. Konoplyanikov,⁹⁰ V. Korenkov,⁹⁰ G. Kozlov,⁹⁰ A. Lanev,⁹⁰ A. Malakhov,⁹⁰ V. Matveev,^{90,ee} V. V. Mitsyn,⁹⁰ P. Moiseenz,⁹⁰ V. Palichik,⁹⁰ V. Perelygin,⁹⁰ S. Shmatov,⁹⁰ N. Skatchkov,⁹⁰ V. Smirnov,⁹⁰ E. Tikhonenko,⁹⁰ A. Zarubin,⁹⁰ V. Golovtsov,⁹¹ Y. Ivanov,⁹¹ V. Kim,^{91,ff} E. Kuznetsova,⁹¹ P. Levchenko,⁹¹ V. Murzin,⁹¹ V. Oreshkin,⁹¹ I. Smirnov,⁹¹ V. Sulimov,⁹¹ L. Uvarov,⁹¹ S. Vavilov,⁹¹ A. Vorobyev,⁹¹ An. Vorobyev,⁹¹ Yu. Andreev,⁹² A. Dermenev,⁹² S. Gninenko,⁹² N. Golubev,⁹² M. Kirsanov,⁹² N. Krasnikov,⁹² A. Pashenkov,⁹² D. Tlisov,⁹² A. Toropin,⁹² V. Epshteyn,⁹³ V. Gavrilov,⁹³ N. Lychkovskaya,⁹³ V. Popov,⁹³ I. Pozdnyakov,⁹³ G. Safronov,⁹³ S. Semenov,⁹³ A. Spiridonov,⁹³ V. Stolin,⁹³ E. Vlasov,⁹³ A. Zhokin,⁹³ V. Andreev,⁹⁴ M. Azarkin,⁹⁴ I. Dremin,⁹⁴ M. Kirakosyan,⁹⁴ A. Leonidov,⁹⁴ G. Mesyats,⁹⁴ S. V. Rusakov,⁹⁴ A. Vinogradov,⁹⁴ A. Belyaev,⁹⁵ E. Boos,⁹⁵ M. Dubinin,^{95,gg} L. Dudko,⁹⁵ A. Ershov,⁹⁵ A. Gribushin,⁹⁵ V. Klyukhin,⁹⁵ O. Kodolova,⁹⁵ I. Lokhtin,⁹⁵ S. Obraztsov,⁹⁵ S. Petrushanko,⁹⁵ V. Savrin,⁹⁵ A. Snigirev,⁹⁵ I. Azhgirey,⁹⁶ I. Bayshev,⁹⁶ S. Bitioukov,⁹⁶ V. Kachanov,⁹⁶ A. Kalinin,⁹⁶ D. Konstantinov,⁹⁶ V. Krychkin,⁹⁶ V. Petrov,⁹⁶ R. Ryutin,⁹⁶ A. Sobol,⁹⁶ L. Tourtchanovitch,⁹⁶ S. Troshin,⁹⁶ N. Tyurin,⁹⁶ A. Uzunian,⁹⁶ A. Volkov,⁹⁶ P. Adzic,^{97,hh} M. Ekmedzic,⁹⁷ J. Milosevic,⁹⁷ V. Rekovic,⁹⁷ J. Alcaraz Maestre,⁹⁸ C. Battilana,⁹⁸ E. Calvo,⁹⁸ M. Cerrada,⁹⁸ M. Chamizo Llatas,⁹⁸ N. Colino,⁹⁸ B. De La Cruz,⁹⁸ A. Delgado Peris,⁹⁸ D. Domínguez Vázquez,⁹⁸ A. Escalante Del Valle,⁹⁸ C. Fernandez Bedoya,⁹⁸ J. P. Fernández Ramos,⁹⁸ J. Flix,⁹⁸ M. C. Fouz,⁹⁸ P. Garcia-Abia,⁹⁸ O. Gonzalez Lopez,⁹⁸ S. Goy Lopez,⁹⁸ J. M. Hernandez,⁹⁸ M. I. Josa,⁹⁸ E. Navarro De Martino,⁹⁸ A. Pérez-Calero Yzquierdo,⁹⁸ J. Puerta Pelayo,⁹⁸ A. Quintario Olmeda,⁹⁸ I. Redondo,⁹⁸ L. Romero,⁹⁸ M. S. Soares,⁹⁸ C. Albajar,⁹⁹ J. F. de Trocóniz,⁹⁹ M. Missiroli,⁹⁹ D. Moran,⁹⁹ H. Brun,¹⁰⁰ J. Cuevas,¹⁰⁰ J. Fernandez Menendez,¹⁰⁰ S. Folgueras,¹⁰⁰ I. Gonzalez Caballero,¹⁰⁰ J. A. Brochero Cifuentes,¹⁰¹ I. J. Cabrillo,¹⁰¹ A. Calderon,¹⁰¹ J. Duarte Campderros,¹⁰¹ M. Fernandez,¹⁰¹ G. Gomez,¹⁰¹ A. Graziano,¹⁰¹ A. Lopez Virto,¹⁰¹ J. Marco,¹⁰¹ R. Marco,¹⁰¹ C. Martinez Rivero,¹⁰¹ F. Matorras,¹⁰¹ F. J. Munoz Sanchez,¹⁰¹ J. Piedra Gomez,¹⁰¹ T. Rodrigo,¹⁰¹ A. Y. Rodríguez-Marrero,¹⁰¹ A. Ruiz-Jimeno,¹⁰¹ L. Scodellaro,¹⁰¹ I. Vila,¹⁰¹ R. Vilar Cortabitarte,¹⁰¹ D. Abbaneo,¹⁰² E. Auffray,¹⁰² G. Auzinger,¹⁰² M. Bachtis,¹⁰² P. Baillon,¹⁰² A. H. Ball,¹⁰² D. Barney,¹⁰² A. Benaglia,¹⁰² J. Bendavid,¹⁰² L. Benhabib,¹⁰² J. F. Benitez,¹⁰² P. Bloch,¹⁰² A. Bocci,¹⁰² A. Bonato,¹⁰² O. Bondu,¹⁰² C. Botta,¹⁰² H. Breuker,¹⁰² T. Camporesi,¹⁰² G. Cerminara,¹⁰² S. Colafranceschi,^{102,ii} M. D'Alfonso,¹⁰² D. d'Enterria,¹⁰² A. Dabrowski,¹⁰² A. David,¹⁰² F. De Guio,¹⁰² A. De Roeck,¹⁰² S. De Visscher,¹⁰² E. Di Marco,¹⁰² M. Dobson,¹⁰² M. Dordevic,¹⁰² B. Dorney,¹⁰² N. Dupont-Sagorin,¹⁰² A. Elliott-Peisert,¹⁰² G. Franzoni,¹⁰² W. Funk,¹⁰² D. Gigi,¹⁰² K. Gill,¹⁰² D. Giordano,¹⁰² M. Girone,¹⁰² F. Glege,¹⁰² R. Guida,¹⁰² S. Gundacker,¹⁰² M. Guthoff,¹⁰² J. Hammer,¹⁰² M. Hansen,¹⁰² P. Harris,¹⁰² J. Hegeman,¹⁰² V. Innocente,¹⁰² P. Janot,¹⁰² K. Kousouris,¹⁰² K. Krajczar,¹⁰² P. Lecoq,¹⁰² C. Lourenço,¹⁰² N. Magini,¹⁰² L. Malgeri,¹⁰² M. Mannelli,¹⁰² J. Marrouche,¹⁰² L. Masetti,¹⁰² F. Meijers,¹⁰² S. Mersi,¹⁰² E. Meschi,¹⁰² F. Moortgat,¹⁰² S. Morovic,¹⁰² M. Mulders,¹⁰² S. Orfanelli,¹⁰² L. Orsini,¹⁰² L. Pape,¹⁰² E. Perez,¹⁰² A. Petrilli,¹⁰² G. Petrucciani,¹⁰² A. Pfeiffer,¹⁰² M. Pimiä,¹⁰² D. Piparo,¹⁰² M. Plagge,¹⁰² A. Racz,¹⁰² G. Rolandi,^{102,jj} M. Rovere,¹⁰² H. Sakulin,¹⁰² C. Schäfer,¹⁰² C. Schwick,¹⁰² A. Sharma,¹⁰² P. Siegrist,¹⁰² P. Silva,¹⁰² M. Simon,¹⁰² P. Sphicas,^{102,kk} D. Spiga,¹⁰² J. Steggemann,¹⁰²

B. Stieger,¹⁰² M. Stoye,¹⁰² Y. Takahashi,¹⁰² D. Treille,¹⁰² A. Tsirou,¹⁰² G. I. Veres,^{102,s} N. Wardle,¹⁰² H. K. Wöhri,¹⁰² H. Wollny,¹⁰² W. D. Zeuner,¹⁰² W. Bertl,¹⁰³ K. Deiters,¹⁰³ W. Erdmann,¹⁰³ R. Horisberger,¹⁰³ Q. Ingram,¹⁰³ H. C. Kaestli,¹⁰³ D. Kotlinski,¹⁰³ U. Langenegger,¹⁰³ D. Renker,¹⁰³ T. Rohe,¹⁰³ F. Bachmair,¹⁰⁴ L. Bäni,¹⁰⁴ L. Bianchini,¹⁰⁴ M. A. Buchmann,¹⁰⁴ B. Casal,¹⁰⁴ N. Chanon,¹⁰⁴ G. Dissertori,¹⁰⁴ M. Dittmar,¹⁰⁴ M. Donegà,¹⁰⁴ M. Dünser,¹⁰⁴ P. Eller,¹⁰⁴ C. Grab,¹⁰⁴ D. Hits,¹⁰⁴ J. Hoss,¹⁰⁴ G. Kasieczka,¹⁰⁴ W. Lustermann,¹⁰⁴ B. Mangano,¹⁰⁴ A. C. Marini,¹⁰⁴ M. Marionneau,¹⁰⁴ P. Martinez Ruiz del Arbol,¹⁰⁴ M. Masciovecchio,¹⁰⁴ D. Meister,¹⁰⁴ N. Mohr,¹⁰⁴ P. Musella,¹⁰⁴ C. Nägeli,^{104,l} F. Nessi-Tedaldi,¹⁰⁴ F. Pandolfi,¹⁰⁴ F. Pauss,¹⁰⁴ L. Perrozzi,¹⁰⁴ M. Peruzzi,¹⁰⁴ M. Quittnat,¹⁰⁴ L. Rebane,¹⁰⁴ M. Rossini,¹⁰⁴ A. Starodumov,^{104,mm} M. Takahashi,¹⁰⁴ K. Theofilatos,¹⁰⁴ R. Wallny,¹⁰⁴ H. A. Weber,¹⁰⁴ C. Amsler,^{105,nn} M. F. Canelli,¹⁰⁵ V. Chiochia,¹⁰⁵ A. De Cosa,¹⁰⁵ A. Hinzmann,¹⁰⁵ T. Hreus,¹⁰⁵ B. Kilminster,¹⁰⁵ C. Lange,¹⁰⁵ J. Ngadiuba,¹⁰⁵ D. Pinna,¹⁰⁵ P. Robmann,¹⁰⁵ F. J. Ronga,¹⁰⁵ S. Taroni,¹⁰⁵ Y. Yang,¹⁰⁵ M. Cardaci,¹⁰⁶ K. H. Chen,¹⁰⁶ C. Ferro,¹⁰⁶ C. M. Kuo,¹⁰⁶ W. Lin,¹⁰⁶ Y. J. Lu,¹⁰⁶ R. Volpe,¹⁰⁶ S. S. Yu,¹⁰⁶ P. Chang,¹⁰⁷ Y. H. Chang,¹⁰⁷ Y. Chao,¹⁰⁷ K. F. Chen,¹⁰⁷ P. H. Chen,¹⁰⁷ C. Dietz,¹⁰⁷ U. Grundler,¹⁰⁷ W.-S. Hou,¹⁰⁷ Y. F. Liu,¹⁰⁷ R.-S. Lu,¹⁰⁷ M. Miñano Moya,¹⁰⁷ E. Petrakou,¹⁰⁷ J. F. Tsai,¹⁰⁷ Y. M. Tzeng,¹⁰⁷ R. Wilken,¹⁰⁷ B. Asavapibhop,¹⁰⁸ G. Singh,¹⁰⁸ N. Srimanobhas,¹⁰⁸ N. Suwonjandee,¹⁰⁸ A. Adiguzel,¹⁰⁹ M. N. Bakirci,^{109,oo} S. Cerci,^{109,pp} C. Dozen,¹⁰⁹ I. Dumanoglu,¹⁰⁹ E. Eskut,¹⁰⁹ S. Girgis,¹⁰⁹ G. Gokbulut,¹⁰⁹ Y. Guler,¹⁰⁹ E. Gurpinar,¹⁰⁹ I. Hos,¹⁰⁹ E. E. Kangal,^{109,qq} A. Kayis Topaksu,¹⁰⁹ G. Onengut,^{109,rr} K. Ozdemir,^{109,ss} S. Ozturk,^{109,oo} A. Polatoz,¹⁰⁹ D. Sunar Cerci,^{109,pp} B. Tali,^{109,pp} H. Topakli,^{109,oo} M. Vergili,¹⁰⁹ C. Zorbilmez,¹⁰⁹ I. V. Akin,¹¹⁰ B. Bilin,¹¹⁰ S. Bilmis,¹¹⁰ H. Gamsizkan,^{110,tt} B. Isildak,^{110,uu} G. Karapinar,^{110,vv} K. Ocalan,^{110,ww} S. Sekmen,¹¹⁰ U. E. Surat,¹¹⁰ M. Yalvac,¹¹⁰ M. Zeyrek,¹¹⁰ E. A. Albayrak,^{111,xx} E. Gülmez,¹¹¹ M. Kaya,^{111,yy} O. Kaya,^{111,zz} T. Yetkin,^{111,aaa} K. Cankocak,¹¹² F. I. Vardarli,¹¹² L. Levchuk,¹¹³ P. Sorokin,¹¹³ J. J. Brooke,¹¹⁴ E. Clement,¹¹⁴ D. Cussans,¹¹⁴ H. Flacher,¹¹⁴ J. Goldstein,¹¹⁴ M. Grimes,¹¹⁴ G. P. Heath,¹¹⁴ H. F. Heath,¹¹⁴ J. Jacob,¹¹⁴ L. Kreczko,¹¹⁴ C. Lucas,¹¹⁴ Z. Meng,¹¹⁴ D. M. Newbold,^{114,bbb} S. Paramesvaran,¹¹⁴ A. Poll,¹¹⁴ T. Sakuma,¹¹⁴ S. Seif El Nasr-storey,¹¹⁴ S. Senkin,¹¹⁴ V. J. Smith,¹¹⁴ K. W. Bell,¹¹⁵ A. Belyaev,^{115,ccc} C. Brew,¹¹⁵ R. M. Brown,¹¹⁵ D. J. A. Cockerill,¹¹⁵ J. A. Coughlan,¹¹⁵ K. Harder,¹¹⁵ S. Harper,¹¹⁵ E. Olaiya,¹¹⁵ D. Petyt,¹¹⁵ C. H. Shepherd-Themistocleous,¹¹⁵ A. Thea,¹¹⁵ I. R. Tomalin,¹¹⁵ T. Williams,¹¹⁵ W. J. Womersley,¹¹⁵ S. D. Worm,¹¹⁵ M. Baber,¹¹⁶ R. Bainbridge,¹¹⁶ O. Buchmuller,¹¹⁶ D. Burton,¹¹⁶ D. Colling,¹¹⁶ N. Cripps,¹¹⁶ P. Dauncey,¹¹⁶ G. Davies,¹¹⁶ M. Della Negra,¹¹⁶ P. Dunne,¹¹⁶ A. Elwood,¹¹⁶ W. Ferguson,¹¹⁶ J. Fulcher,¹¹⁶ D. Futyan,¹¹⁶ G. Hall,¹¹⁶ G. Iles,¹¹⁶ M. Jarvis,¹¹⁶ G. Karapostoli,¹¹⁶ M. Kenzie,¹¹⁶ R. Lane,¹¹⁶ R. Lucas,^{116,bbb} L. Lyons,¹¹⁶ A.-M. Magnan,¹¹⁶ S. Malik,¹¹⁶ B. Mathias,¹¹⁶ J. Nash,¹¹⁶ A. Nikitenko,^{116,mm} J. Pela,¹¹⁶ M. Pesaresi,¹¹⁶ K. Petridis,¹¹⁶ D. M. Raymond,¹¹⁶ S. Rogerson,¹¹⁶ A. Rose,¹¹⁶ C. Seez,¹¹⁶ P. Sharp,^{116,a} A. Tapper,¹¹⁶ M. Vazquez Acosta,¹¹⁶ T. Virdee,¹¹⁶ S. C. Zenz,¹¹⁶ J. E. Cole,¹¹⁷ P. R. Hobson,¹¹⁷ A. Khan,¹¹⁷ P. Kyberd,¹¹⁷ D. Leggat,¹¹⁷ D. Leslie,¹¹⁷ I. D. Reid,¹¹⁷ P. Symonds,¹¹⁷ L. Teodorescu,¹¹⁷ M. Turner,¹¹⁷ J. Dittmann,¹¹⁸ K. Hatakeyama,¹¹⁸ A. Kasmi,¹¹⁸ H. Liu,¹¹⁸ N. Pastika,¹¹⁸ T. Scarborough,¹¹⁸ Z. Wu,¹¹⁸ O. Charaf,¹¹⁹ S. I. Cooper,¹¹⁹ C. Henderson,¹¹⁹ P. Rumerio,¹¹⁹ A. Avetisyan,¹²⁰ T. Bose,¹²⁰ C. Fantasia,¹²⁰ P. Lawson,¹²⁰ C. Richardson,¹²⁰ J. Rohlf,¹²⁰ J. St. John,¹²⁰ L. Sulak,¹²⁰ J. Alimena,¹²¹ E. Berry,¹²¹ S. Bhattacharya,¹²¹ G. Christopher,¹²¹ D. Cutts,¹²¹ Z. Demiragli,¹²¹ N. Dhingra,¹²¹ A. Ferapontov,¹²¹ A. Garabedian,¹²¹ U. Heintz,¹²¹ E. Laird,¹²¹ G. Landsberg,¹²¹ Z. Mao,¹²¹ M. Narain,¹²¹ S. Sagir,¹²¹ T. Sinthuprasith,¹²¹ T. Speer,¹²¹ J. Swanson,¹²¹ R. Breedon,¹²² G. Breto,¹²² M. Calderon De La Barca Sanchez,¹²² S. Chauhan,¹²² M. Chertok,¹²² J. Conway,¹²² R. Conway,¹²² P. T. Cox,¹²² R. Erbacher,¹²² M. Gardner,¹²² W. Ko,¹²² R. Lander,¹²² M. Mulhearn,¹²² D. Pellett,¹²² J. Pilot,¹²² F. Ricci-Tam,¹²² S. Shalhout,¹²² J. Smith,¹²² M. Squires,¹²² D. Stolp,¹²² M. Tripathi,¹²² S. Wilbur,¹²² R. Yohay,¹²² R. Cousins,¹²³ P. Everaerts,¹²³ C. Farrell,¹²³ J. Hauser,¹²³ M. Ignatenko,¹²³ G. Rakness,¹²³ E. Takasugi,¹²³ V. Valuev,¹²³ M. Weber,¹²³ K. Burt,¹²⁴ R. Clare,¹²⁴ J. Ellison,¹²⁴ J. W. Gary,¹²⁴ G. Hanson,¹²⁴ J. Heilman,¹²⁴ M. Ivova Rikova,¹²⁴ P. Jandir,¹²⁴ E. Kennedy,¹²⁴ F. Lacroix,¹²⁴ O. R. Long,¹²⁴ A. Luthra,¹²⁴ M. Malberti,¹²⁴ M. Olmedo Negrete,¹²⁴ A. Shrinivas,¹²⁴ S. Sumowidagdo,¹²⁴ S. Wimpenny,¹²⁴ J. G. Branson,¹²⁵ G. B. Cerati,¹²⁵ S. 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W. T. Ford,¹²⁹ A. Gaz,¹²⁹ M. Krohn,¹²⁹ E. Luiggi Lopez,¹²⁹ U. Nauenberg,¹²⁹ J. G. Smith,¹²⁹ K. Stenson,¹²⁹ S. R. Wagner,¹²⁹ J. Alexander,¹³⁰ A. Chatterjee,¹³⁰ J. Chaves,¹³⁰ J. Chu,¹³⁰ S. Dittmer,¹³⁰ N. Eggert,¹³⁰ N. Mirman,¹³⁰ G. Nicolas Kaufman,¹³⁰ J. R. Patterson,¹³⁰ A. Ryd,¹³⁰ E. Salvati,¹³⁰ L. Skinnari,¹³⁰ W. Sun,¹³⁰ W. D. Teo,¹³⁰ J. Thom,¹³⁰ J. Thompson,¹³⁰ J. Tucker,¹³⁰ Y. Weng,¹³⁰ L. Winstrom,¹³⁰ P. Wittich,¹³⁰ D. Winn,¹³¹ S. Abdullin,¹³² M. Albrow,¹³² J. Anderson,¹³² G. Apollinari,¹³² L. A. T. Bauerdick,¹³² A. Beretvas,¹³² J. Berryhill,¹³² P. C. Bhat,¹³² G. Bolla,¹³² K. Burkett,¹³² J. N. Butler,¹³² H. W. K. Cheung,¹³² F. Chlebana,¹³² S. Cihangir,¹³² V. D. Elvira,¹³² I. Fisk,¹³² J. Freeman,¹³² E. Gottschalk,¹³² L. Gray,¹³² D. Green,¹³² S. Grünendahl,¹³² O. Gutsche,¹³² J. Hanlon,¹³² D. Hare,¹³² R. M. Harris,¹³² J. Hirschauer,¹³² B. Hooberman,¹³² S. Jindariani,¹³² M. Johnson,¹³² U. Joshi,¹³² B. Klima,¹³² B. Kreis,¹³² S. Kwan,^{132,a} J. Linacre,¹³² D. Lincoln,¹³² R. Lipton,¹³² T. Liu,¹³² R. Lopes De Sá,¹³² J. Lykken,¹³² K. Maeshima,¹³² J. M. Marraffino,¹³² V. I. Martinez Outschoorn,¹³² S. Maruyama,¹³² D. Mason,¹³² P. McBride,¹³² P. Merkel,¹³² K. Mishra,¹³² S. Mrenna,¹³² S. Nahn,¹³² C. Newman-Holmes,¹³² V. O'Dell,¹³² O. Prokofyev,¹³² E. Sexton-Kennedy,¹³² A. Soha,¹³² W. J. Spalding,¹³² L. Spiegel,¹³² L. Taylor,¹³² S. Tkaczyk,¹³² N. V. Tran,¹³² L. Uplegger,¹³² E. W. Vaandering,¹³² R. Vidal,¹³² A. Whitbeck,¹³² J. Whitmore,¹³² F. Yang,¹³² D. Acosta,¹³³ P. Avery,¹³³ P. Bortignon,¹³³ D. Bourilkov,¹³³ M. Carver,¹³³ D. Curry,¹³³ S. Das,¹³³ M. De Gruttola,¹³³ G. P. Di Giovanni,¹³³ R. D. Field,¹³³ M. Fisher,¹³³ I. K. Furic,¹³³ J. Hugon,¹³³ J. Konigsberg,¹³³ A. Korytov,¹³³ T. Kypreos,¹³³ J. F. Low,¹³³ K. Matchev,¹³³ H. Mei,¹³³ P. Milenovic,^{133,ddd} G. Mitselmakher,¹³³ L. Muniz,¹³³ A. Rinkevicius,¹³³ L. Shchutska,¹³³ M. Snowball,¹³³ D. Sperka,¹³³ J. Yelton,¹³³ M. Zakaria,¹³³ S. Hewamanage,¹³⁴ S. Linn,¹³⁴ P. Markowitz,¹³⁴ G. Martinez,¹³⁴ J. L. Rodriguez,¹³⁴ J. R. Adams,¹³⁵ T. Adams,¹³⁵ A. Askew,¹³⁵ J. Bochenek,¹³⁵ B. Diamond,¹³⁵ J. Haas,¹³⁵ S. Hagopian,¹³⁵ V. Hagopian,¹³⁵ K. F. Johnson,¹³⁵ H. Prosper,¹³⁵ V. Veeraraghavan,¹³⁵ M. Weinberg,¹³⁵ M. M. Baarmand,¹³⁶ M. Hohlmann,¹³⁶ H. Kalakhety,¹³⁶ F. Yumiceva,¹³⁶ M. R. Adams,¹³⁷ L. Apanasevich,¹³⁷ D. Berry,¹³⁷ R. R. Betts,¹³⁷ I. Bucinskaite,¹³⁷ R. Cavanaugh,¹³⁷ O. Evdokimov,¹³⁷ L. Gauthier,¹³⁷ C. E. Gerber,¹³⁷ D. J. Hofman,¹³⁷ P. Kurt,¹³⁷ C. O'Brien,¹³⁷ I. D. Sandoval Gonzalez,¹³⁷ C. Silkworth,¹³⁷ P. Turner,¹³⁷ N. Varelas,¹³⁷ B. Bilki,^{138,eee} W. Clarida,¹³⁸ K. Dilsiz,¹³⁸ M. Haytmyradov,¹³⁸ V. Khristenko,¹³⁸ J.-P. Merlo,¹³⁸ H. Mermerkaya,^{138,fff} A. Mestvirishvili,¹³⁸ A. Moeller,¹³⁸ J. Nachtman,¹³⁸ H. Ogul,¹³⁸ Y. Onel,¹³⁸ F. Ozok,^{138,xx} A. Penzo,¹³⁸ R. Rahmat,¹³⁸ S. Sen,¹³⁸ P. Tan,¹³⁸ E. Tiras,¹³⁸ J. Wetzel,¹³⁸ K. Yi,¹³⁸ I. Anderson,¹³⁹ B. A. Barnett,¹³⁹ B. Blumenfeld,¹³⁹ S. Bolognesi,¹³⁹ D. Fehling,¹³⁹ A. V. Gritsan,¹³⁹ P. Maksimovic,¹³⁹ C. Martin,¹³⁹ M. Swartz,¹³⁹ M. Xiao,¹³⁹ P. Baringer,¹⁴⁰ A. Bean,¹⁴⁰ G. Benelli,¹⁴⁰ C. Bruner,¹⁴⁰ J. Gray,¹⁴⁰ R. P. Kenny III,¹⁴⁰ D. Majumder,¹⁴⁰ M. Malek,¹⁴⁰ M. Murray,¹⁴⁰ D. Noonan,¹⁴⁰ S. Sanders,¹⁴⁰ J. Sekaric,¹⁴⁰ R. Stringer,¹⁴⁰ Q. Wang,¹⁴⁰ J. S. Wood,¹⁴⁰ I. Chakaberia,¹⁴¹ A. Ivanov,¹⁴¹ K. Kaadze,¹⁴¹ S. Khalil,¹⁴¹ M. Makouski,¹⁴¹ Y. Maravin,¹⁴¹ L. K. Saini,¹⁴¹ N. Skhirtladze,¹⁴¹ I. Svintradze,¹⁴¹ J. Gronberg,¹⁴² D. Lange,¹⁴² F. Rebassoo,¹⁴² D. Wright,¹⁴² C. Anelli,¹⁴³ A. Baden,¹⁴³ A. Belloni,¹⁴³ B. Calvert,¹⁴³ S. C. Eno,¹⁴³ J. A. Gomez,¹⁴³ N. J. Hadley,¹⁴³ S. Jabeen,¹⁴³ R. G. Kellogg,¹⁴³ T. Kolberg,¹⁴³ Y. Lu,¹⁴³ A. C. Mignerey,¹⁴³ K. Pedro,¹⁴³ Y. H. Shin,¹⁴³ A. Skuja,¹⁴³ M. B. Tonjes,¹⁴³ S. C. Tonwar,¹⁴³ A. Apyan,¹⁴⁴ R. Barbieri,¹⁴⁴ K. Bierwagen,¹⁴⁴ W. Busza,¹⁴⁴ I. A. Cali,¹⁴⁴ L. Di Matteo,¹⁴⁴ G. Gomez Ceballos,¹⁴⁴ M. Goncharov,¹⁴⁴ D. Gulhan,¹⁴⁴ M. Klute,¹⁴⁴ Y. S. Lai,¹⁴⁴ Y.-J. Lee,¹⁴⁴ A. Levin,¹⁴⁴ P. D. Luckey,¹⁴⁴ C. Paus,¹⁴⁴ D. Ralph,¹⁴⁴ C. Roland,¹⁴⁴ G. Roland,¹⁴⁴ G. S. F. Stephens,¹⁴⁴ K. Sumorok,¹⁴⁴ D. Velicanu,¹⁴⁴ J. Veverka,¹⁴⁴ B. Wyslouch,¹⁴⁴ M. Yang,¹⁴⁴ M. Zanetti,¹⁴⁴ V. Zhukova,¹⁴⁴ B. Dahmes,¹⁴⁵ A. Gude,¹⁴⁵ S. C. Kao,¹⁴⁵ K. Klapoetke,¹⁴⁵ Y. Kubota,¹⁴⁵ J. Mans,¹⁴⁵ S. Nourbakhsh,¹⁴⁵ R. Rusack,¹⁴⁵ A. Singovsky,¹⁴⁵ N. Tambe,¹⁴⁵ J. Turkewitz,¹⁴⁵ J. G. Acosta,¹⁴⁶ S. Oliveros,¹⁴⁶ E. Avdeeva,¹⁴⁷ K. Bloom,¹⁴⁷ S. Bose,¹⁴⁷ D. R. Claes,¹⁴⁷ A. Dominguez,¹⁴⁷ R. Gonzalez Suarez,¹⁴⁷ J. Keller,¹⁴⁷ D. Knowlton,¹⁴⁷ I. Kravchenko,¹⁴⁷ J. Lazo-Flores,¹⁴⁷ F. Meier,¹⁴⁷ F. Ratnikov,¹⁴⁷ G. R. Snow,¹⁴⁷ M. Zvara,¹⁴⁷ J. Dolen,¹⁴⁸ A. Godshalk,¹⁴⁸ I. Iashvili,¹⁴⁸ A. Kharchilava,¹⁴⁸ A. Kumar,¹⁴⁸ S. Rappoccio,¹⁴⁸ G. Alverson,¹⁴⁹ E. Barberis,¹⁴⁹ D. Baumgartel,¹⁴⁹ M. Chasco,¹⁴⁹ A. Massironi,¹⁴⁹ D. M. Morse,¹⁴⁹ D. Nash,¹⁴⁹ T. Orimoto,¹⁴⁹ D. Trocino,¹⁴⁹ R.-J. Wang,¹⁴⁹ D. Wood,¹⁴⁹ J. Zhang,¹⁴⁹ K. A. Hahn,¹⁵⁰ A. Kubik,¹⁵⁰ N. Mucia,¹⁵⁰ N. Odell,¹⁵⁰ B. Pollack,¹⁵⁰ A. Pozdnyakov,¹⁵⁰ M. Schmitt,¹⁵⁰ S. Stoynev,¹⁵⁰ K. Sung,¹⁵⁰ M. Trovato,¹⁵⁰ M. Velasco,¹⁵⁰ S. Won,¹⁵⁰ A. Brinkerhoff,¹⁵¹ K. M. Chan,¹⁵¹ A. Drozdetskiy,¹⁵¹ M. Hildreth,¹⁵¹ C. Jessop,¹⁵¹ D. J. Karmgard,¹⁵¹ N. Kellams,¹⁵¹ K. Lannon,¹⁵¹ S. Lynch,¹⁵¹ N. Marinelli,¹⁵¹ Y. Musienko,^{151,ee} T. Pearson,¹⁵¹ M. Planer,¹⁵¹ R. Ruchti,¹⁵¹ G. Smith,¹⁵¹ N. Valls,¹⁵¹ M. Wayne,¹⁵¹ M. Wolf,¹⁵¹ A. Woodard,¹⁵¹ L. Antonelli,¹⁵² J. Brinson,¹⁵² B. Bylsma,¹⁵² L. S. Durkin,¹⁵² S. Flowers,¹⁵² A. Hart,¹⁵² C. Hill,¹⁵² R. Hughes,¹⁵² K. Kotov,¹⁵² T. Y. Ling,¹⁵² W. Luo,¹⁵² D. Puigh,¹⁵² M. Rodenburg,¹⁵² B. L. Winer,¹⁵² H. Wolfe,¹⁵² H. W. Wulsin,¹⁵² O. Driga,¹⁵³ P. Elmer,¹⁵³ J. Hardenbrook,¹⁵³ P. Hebda,¹⁵³ S. A. Koay,¹⁵³ P. Lujan,¹⁵³ D. Marlow,¹⁵³ T. Medvedeva,¹⁵³ M. Mooney,¹⁵³ J. Olsen,¹⁵³ P. Piroué,¹⁵³ X. Quan,¹⁵³ H. Saka,¹⁵³ D. Stickland,^{153,e} C. Tully,¹⁵³ J. S. Werner,¹⁵³ A. Zuranski,¹⁵³ E. Brownson,¹⁵⁴ S. Malik,¹⁵⁴ H. Mendez,¹⁵⁴ J. E. Ramirez Vargas,¹⁵⁴ V. E. Barnes,¹⁵⁵ D. Benedetti,¹⁵⁵ D. Bortoletto,¹⁵⁵ L. Gutay,¹⁵⁵ Z. Hu,¹⁵⁵ M. K. Jha,¹⁵⁵ M. Jones,¹⁵⁵ K. Jung,¹⁵⁵ M. Kress,¹⁵⁵ N. Leonardo,¹⁵⁵

D. H. Miller,¹⁵⁵ N. Neumeister,¹⁵⁵ F. Primavera,¹⁵⁵ B. C. Radburn-Smith,¹⁵⁵ X. Shi,¹⁵⁵ I. Shipsey,¹⁵⁵ D. Silvers,¹⁵⁵ A. Svyatkovskiy,¹⁵⁵ F. Wang,¹⁵⁵ W. Xie,¹⁵⁵ L. Xu,¹⁵⁵ J. Zablocki,¹⁵⁵ N. Parashar,¹⁵⁶ J. Stupak,¹⁵⁶ A. Adair,¹⁵⁷ B. Akgun,¹⁵⁷ K. M. Ecklund,¹⁵⁷ F. J. M. Geurts,¹⁵⁷ W. Li,¹⁵⁷ B. Michlin,¹⁵⁷ B. P. Padley,¹⁵⁷ R. Redjimi,¹⁵⁷ J. Roberts,¹⁵⁷ J. Zabel,¹⁵⁷ B. Betchart,¹⁵⁸ A. Bodek,¹⁵⁸ P. de Barbaro,¹⁵⁸ R. Demina,¹⁵⁸ Y. Eshaq,¹⁵⁸ T. Ferbel,¹⁵⁸ M. Galanti,¹⁵⁸ A. Garcia-Bellido,¹⁵⁸ P. Goldenzweig,¹⁵⁸ J. Han,¹⁵⁸ A. Harel,¹⁵⁸ O. Hindrichs,¹⁵⁸ A. Khukhunaishvili,¹⁵⁸ S. Korjenevski,¹⁵⁸ G. Petrillo,¹⁵⁸ M. Verzetti,¹⁵⁸ D. Vishnevskiy,¹⁵⁸ R. Ciesielski,¹⁵⁹ L. Demortier,¹⁵⁹ K. Goulianos,¹⁵⁹ C. Mesropian,¹⁵⁹ S. Arora,¹⁶⁰ A. Barker,¹⁶⁰ J. P. Chou,¹⁶⁰ C. Contreras-Campana,¹⁶⁰ E. Contreras-Campana,¹⁶⁰ D. Duggan,¹⁶⁰ D. Ferencek,¹⁶⁰ Y. Gershtein,¹⁶⁰ R. Gray,¹⁶⁰ E. Halkiadakis,¹⁶⁰ D. Hidas,¹⁶⁰ E. Hughes,¹⁶⁰ S. Kaplan,¹⁶⁰ A. Lath,¹⁶⁰ S. Panwalkar,¹⁶⁰ M. Park,¹⁶⁰ S. Salur,¹⁶⁰ S. Schnetzer,¹⁶⁰ D. Sheffield,¹⁶⁰ S. Somalwar,¹⁶⁰ R. Stone,¹⁶⁰ S. Thomas,¹⁶⁰ P. Thomassen,¹⁶⁰ M. Walker,¹⁶⁰ K. Rose,¹⁶¹ S. Spanier,¹⁶¹ A. York,¹⁶¹ O. Bouhali,^{162,ggg} A. Castaneda Hernandez,¹⁶² M. Dalchenko,¹⁶² M. De Mattia,¹⁶² S. Dildick,¹⁶² R. Eusebi,¹⁶² W. Flanagan,¹⁶² J. Gilmore,¹⁶² T. Kamon,^{162,hhh} V. Khotilovich,¹⁶² V. Krutelyov,¹⁶² R. Montalvo,¹⁶² I. Osipenko,¹⁶² Y. Pakhotin,¹⁶² R. Patel,¹⁶² A. Perloff,¹⁶² J. Roe,¹⁶² A. Rose,¹⁶² A. Safonov,¹⁶² I. Suarez,¹⁶² A. Tatarinov,¹⁶² K. A. Ulmer,¹⁶² N. Akchurin,¹⁶³ C. Cowden,¹⁶³ J. Damgov,¹⁶³ C. Dragoiu,¹⁶³ P. R. Duerdo,¹⁶³ J. Faulkner,¹⁶³ K. Kovitanggoon,¹⁶³ S. Kunori,¹⁶³ S. W. Lee,¹⁶³ T. Libeiro,¹⁶³ I. Volobouev,¹⁶³ E. Appelt,¹⁶⁴ A. G. Delannoy,¹⁶⁴ S. Greene,¹⁶⁴ A. Gurrola,¹⁶⁴ W. Johns,¹⁶⁴ C. Maguire,¹⁶⁴ Y. Mao,¹⁶⁴ A. Melo,¹⁶⁴ M. Sharma,¹⁶⁴ P. Sheldon,¹⁶⁴ B. Snook,¹⁶⁴ S. Tuo,¹⁶⁴ J. Velkovska,¹⁶⁴ M. W. Arenton,¹⁶⁵ S. Boutle,¹⁶⁵ B. Cox,¹⁶⁵ B. Francis,¹⁶⁵ J. Goodell,¹⁶⁵ R. Hirosky,¹⁶⁵ A. Ledovskoy,¹⁶⁵ H. Li,¹⁶⁵ C. Lin,¹⁶⁵ C. Neu,¹⁶⁵ E. Wolfe,¹⁶⁵ J. Wood,¹⁶⁵ C. Clarke,¹⁶⁶ R. Harr,¹⁶⁶ P. E. Karchin,¹⁶⁶ C. Kottachchi Kankanamge Don,¹⁶⁶ P. Lamichhane,¹⁶⁶ J. Sturdy,¹⁶⁶ D. A. Belknap,¹⁶⁷ D. Carlsmith,¹⁶⁷ M. Cepeda,¹⁶⁷ S. Dasu,¹⁶⁷ L. Dodd,¹⁶⁷ S. Duric,¹⁶⁷ E. Friis,¹⁶⁷ R. Hall-Wilton,¹⁶⁷ M. Herndon,¹⁶⁷ A. Hervé,¹⁶⁷ P. Klabbbers,¹⁶⁷ A. Lanaro,¹⁶⁷ C. Lazaridis,¹⁶⁷ A. Levine,¹⁶⁷ R. Loveless,¹⁶⁷ A. Mohapatra,¹⁶⁷ I. Ojalvo,¹⁶⁷ T. Perry,¹⁶⁷ G. A. Pierro,¹⁶⁷ G. Polese,¹⁶⁷ I. Ross,¹⁶⁷ T. Sarangi,¹⁶⁷ A. Savin,¹⁶⁷ W. H. Smith,¹⁶⁷ D. Taylor,¹⁶⁷ C. Vuosalo¹⁶⁷ and N. Woods¹⁶⁷

(CMS Collaboration)

¹*Yerevan Physics Institute, Yerevan, Armenia*

²*Institut für Hochenergiephysik der OeAW, Wien, Austria*

³*National Centre for Particle and High Energy Physics, Minsk, Belarus*

⁴*Universiteit Antwerpen, Antwerpen, Belgium*

⁵*Vrije Universiteit Brussel, Brussel, Belgium*

⁶*Université Libre de Bruxelles, Bruxelles, Belgium*

⁷*Ghent University, Ghent, Belgium*

⁸*Université Catholique de Louvain, Louvain-la-Neuve, Belgium*

⁹*Université de Mons, Mons, Belgium*

¹⁰*Centro Brasileiro de Pesquisas Físicas, Rio de Janeiro, Brazil*

¹¹*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*

^{12a}*Universidade Estadual Paulista, São Paulo, Brazil*

^{12b}*Universidade Federal do ABC, São Paulo, Brazil*

¹³*Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria*

¹⁴*University of Sofia, Sofia, Bulgaria*

¹⁵*Institute of High Energy Physics, Beijing, China*

¹⁶*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*

¹⁷*Universidad de Los Andes, Bogota, Colombia*

¹⁸*University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia*

¹⁹*University of Split, Faculty of Science, Split, Croatia*

²⁰*Institute Rudjer Boskovic, Zagreb, Croatia*

²¹*University of Cyprus, Nicosia, Cyprus*

²²*Charles University, Prague, Czech Republic*

²³*Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt*

²⁴*National Institute of Chemical Physics and Biophysics, Tallinn, Estonia*

²⁵*Department of Physics, University of Helsinki, Helsinki, Finland*

²⁶*Helsinki Institute of Physics, Helsinki, Finland*

²⁷*Lappeenranta University of Technology, Lappeenranta, Finland*

²⁸*DSM/IRFU, CEA/Saclay, Gif-sur-Yvette, France*

²⁹*Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France*

³⁰*Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France*

³¹*Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France*

³²*Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France*

³³*Institute of High Energy Physics and Informatization, Tbilisi State University, Tbilisi, Georgia*

³⁴*RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany*

³⁵*RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany*

³⁶*RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany*

³⁷*Deutsches Elektronen-Synchrotron, Hamburg, Germany*

³⁸*University of Hamburg, Hamburg, Germany*

³⁹*Institut für Experimentelle Kernphysik, Karlsruhe, Germany*

⁴⁰*Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece*

⁴¹*University of Athens, Athens, Greece*

⁴²*University of Ioánnina, Ioánnina, Greece*

⁴³*Wigner Research Centre for Physics, Budapest, Hungary*

⁴⁴*Institute of Nuclear Research ATOMKI, Debrecen, Hungary*

⁴⁵*University of Debrecen, Debrecen, Hungary*

⁴⁶*National Institute of Science Education and Research, Bhubaneswar, India*

⁴⁷*Panjab University, Chandigarh, India*

⁴⁸*University of Delhi, Delhi, India*

⁴⁹*Saha Institute of Nuclear Physics, Kolkata, India*

⁵⁰*Bhabha Atomic Research Centre, Mumbai, India*

⁵¹*Tata Institute of Fundamental Research, Mumbai, India*

⁵²*Indian Institute of Science Education and Research (IISER), Pune, India*

⁵³*Institute for Research in Fundamental Sciences (IPM), Tehran, Iran*

⁵⁴*University College Dublin, Dublin, Ireland*

^{55a}*INFN Sezione di Bari, Bari, Italy*

^{55b}*Università di Bari, Bari, Italy*

^{55c}*Politecnico di Bari, Bari, Italy*

^{56a}*INFN Sezione di Bologna, Bologna, Italy*

^{56b}*Università di Bologna, Bologna, Italy*

^{57a}*INFN Sezione di Catania, Catania, Italy*

^{57b}*Università di Catania, Catania, Italy*

^{57c}*CSFNSM, Catania, Italy*

^{58a}*INFN Sezione di Firenze, Firenze, Italy*

^{58b}*Università di Firenze, Firenze, Italy*

⁵⁹*INFN Laboratori Nazionali di Frascati, Frascati, Italy*

^{60a}*INFN Sezione di Genova, Genova, Italy*

^{60b}*Università di Genova, Genova, Italy*

^{61a}*INFN Sezione di Milano-Bicocca, Milano, Italy*

^{61b}*Università di Milano-Bicocca, Milano, Italy*

^{62a}*INFN Sezione di Napoli, Napoli, Italy*

^{62b}*Università di Napoli 'Federico II', Napoli, Italy*

^{62c}*Università della Basilicata (Potenza), Napoli, Italy*

^{62d}*Università G. Marconi (Roma), Napoli, Italy*

^{63a}*INFN Sezione di Padova, Padova, Italy*

^{63b}*Università di Padova, Padova, Italy*

^{63c}*Università di Trento (Trento), Padova, Italy*

^{64a}*INFN Sezione di Pavia, Pavia, Italy*

^{64b}*Università di Pavia, Pavia, Italy*

^{65a}*INFN Sezione di Perugia, Perugia, Italy*

^{65b}*Università di Perugia, Perugia, Italy*

^{66a}*INFN Sezione di Pisa, Pisa, Italy*

^{66b}*Università di Pisa, Pisa, Italy*

^{66c}*Scuola Normale Superiore di Pisa, Pisa, Italy*

^{67a}*INFN Sezione di Roma, Roma, Italy*

^{67b}*Università di Roma, Roma, Italy*

^{68a}*INFN Sezione di Torino, Torino, Italy*

^{68b}*Università di Torino, Torino, Italy*

- ^{68c}*Università del Piemonte Orientale (Novara), Torino, Italy*
^{69a}*INFN Sezione di Trieste, Trieste, Italy*
^{69b}*Università di Trieste, Trieste, Italy*
⁷⁰*Kangwon National University, Chunchon, Korea*
⁷¹*Kyungpook National University, Daegu, Korea*
⁷²*Chonbuk National University, Jeonju, Korea*
⁷³*Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea*
⁷⁴*Korea University, Seoul, Korea*
⁷⁵*Seoul National University, Seoul, Korea*
⁷⁶*University of Seoul, Seoul, Korea*
⁷⁷*Sungkyunkwan University, Suwon, Korea*
⁷⁸*Vilnius University, Vilnius, Lithuania*
⁷⁹*National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia*
⁸⁰*Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico*
⁸¹*Universidad Iberoamericana, Mexico City, Mexico*
⁸²*Benemerita Universidad Autonoma de Puebla, Puebla, Mexico*
⁸³*Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico*
⁸⁴*University of Auckland, Auckland, New Zealand*
⁸⁵*University of Canterbury, Christchurch, New Zealand*
⁸⁶*National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan*
⁸⁷*National Centre for Nuclear Research, Swierk, Poland*
⁸⁸*Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland*
⁸⁹*Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal*
⁹⁰*Joint Institute for Nuclear Research, Dubna, Russia*
⁹¹*Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia*
⁹²*Institute for Nuclear Research, Moscow, Russia*
⁹³*Institute for Theoretical and Experimental Physics, Moscow, Russia*
⁹⁴*P.N. Lebedev Physical Institute, Moscow, Russia*
⁹⁵*Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia*
⁹⁶*State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia*
⁹⁷*University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia*
⁹⁸*Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain*
⁹⁹*Universidad Autónoma de Madrid, Madrid, Spain*
¹⁰⁰*Universidad de Oviedo, Oviedo, Spain*
¹⁰¹*Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain*
¹⁰²*CERN, European Organization for Nuclear Research, Geneva, Switzerland*
¹⁰³*Paul Scherrer Institut, Villigen, Switzerland*
¹⁰⁴*Institute for Particle Physics, ETH Zurich, Zurich, Switzerland*
¹⁰⁵*Universität Zürich, Zurich, Switzerland*
¹⁰⁶*National Central University, Chung-Li, Taiwan*
¹⁰⁷*National Taiwan University (NTU), Taipei, Taiwan*
¹⁰⁸*Chulalongkorn University, Faculty of Science, Department of Physics, Bangkok, Thailand*
¹⁰⁹*Cukurova University, Adana, Turkey*
¹¹⁰*Middle East Technical University, Physics Department, Ankara, Turkey*
¹¹¹*Bogazici University, Istanbul, Turkey*
¹¹²*Istanbul Technical University, Istanbul, Turkey*
¹¹³*National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine*
¹¹⁴*University of Bristol, Bristol, United Kingdom*
¹¹⁵*Rutherford Appleton Laboratory, Didcot, United Kingdom*
¹¹⁶*Imperial College, London, United Kingdom*
¹¹⁷*Brunel University, Uxbridge, United Kingdom*
¹¹⁸*Baylor University, Waco, Texas 76798, USA*
¹¹⁹*The University of Alabama, Tuscaloosa, Alabama 35487, USA*
¹²⁰*Boston University, Boston, Massachusetts 02215, USA*
¹²¹*Brown University, Providence, Rhode Island 02912, USA*
¹²²*University of California, Davis, Davis, California 95616, USA*
¹²³*University of California, Los Angeles, California 90095, USA*
¹²⁴*University of California, Riverside, Riverside, California 92521, USA*
¹²⁵*University of California, San Diego, La Jolla, California 92093, USA*
¹²⁶*University of California, Santa Barbara, Santa Barbara, California 93106, USA*

- ¹²⁷California Institute of Technology, Pasadena, California 91125, USA
- ¹²⁸Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA
- ¹²⁹University of Colorado at Boulder, Boulder, Colorado 80309, USA
- ¹³⁰Cornell University, Ithaca, New York 14853, USA
- ¹³¹Fairfield University, Fairfield, Connecticut 06430, USA
- ¹³²Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA
- ¹³³University of Florida, Gainesville, Florida 32611, USA
- ¹³⁴Florida International University, Miami, Florida 33199, USA
- ¹³⁵Florida State University, Tallahassee, Florida 32306, USA
- ¹³⁶Florida Institute of Technology, Melbourne, Florida 32901, USA
- ¹³⁷University of Illinois at Chicago (UIC), Chicago, Illinois 60607, USA
- ¹³⁸The University of Iowa, Iowa City, Iowa 52242, USA
- ¹³⁹Johns Hopkins University, Baltimore, Maryland 21218, USA
- ¹⁴⁰The University of Kansas, Lawrence, Kansas 66045, USA
- ¹⁴¹Kansas State University, Manhattan, Kansas 66506, USA
- ¹⁴²Lawrence Livermore National Laboratory, Livermore, California 94551, USA
- ¹⁴³University of Maryland, College Park, Maryland 20742, USA
- ¹⁴⁴Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA
- ¹⁴⁵University of Minnesota, Minneapolis, Minnesota 55455, USA
- ¹⁴⁶University of Mississippi, Oxford, Mississippi 38677, USA
- ¹⁴⁷University of Nebraska-Lincoln, Lincoln, Nebraska 68588, USA
- ¹⁴⁸State University of New York at Buffalo, Buffalo, New York 14260, USA
- ¹⁴⁹Northeastern University, Boston, Massachusetts 02115, USA
- ¹⁵⁰Northwestern University, Evanston, Illinois 60208, USA
- ¹⁵¹University of Notre Dame, Notre Dame, Indiana 46556, USA
- ¹⁵²The Ohio State University, Columbus, Ohio 43210, USA
- ¹⁵³Princeton University, Princeton, New Jersey 08542, USA
- ¹⁵⁴University of Puerto Rico, Mayaguez, Puerto Rico 00681, USA
- ¹⁵⁵Purdue University, West Lafayette, Indiana 47907, USA
- ¹⁵⁶Purdue University Calumet, Hammond, Indiana 46323, USA
- ¹⁵⁷Rice University, Houston, Texas 77251, USA
- ¹⁵⁸University of Rochester, Rochester, New York 14627, USA
- ¹⁵⁹The Rockefeller University, New York, New York 10021, USA
- ¹⁶⁰Rutgers, The State University of New Jersey, Piscataway, New Jersey 08854, USA
- ¹⁶¹University of Tennessee, Knoxville, Tennessee 37996, USA
- ¹⁶²Texas A&M University, College Station, Texas 77843, USA
- ¹⁶³Texas Tech University, Lubbock, Texas 79409, USA
- ¹⁶⁴Vanderbilt University, Nashville, Tennessee 37235, USA
- ¹⁶⁵University of Virginia, Charlottesville, Virginia 22904, USA
- ¹⁶⁶Wayne State University, Detroit, Michigan 48202, USA
- ¹⁶⁷University of Wisconsin, Madison, Wisconsin 53706, USA

^aDeceased.

^bAlso at Vienna University of Technology, Vienna, Austria.

^cAlso at CERN, European Organization for Nuclear Research, Geneva, Switzerland.

^dAlso at Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France.

^eAlso at National Institute of Chemical Physics and Biophysics, Tallinn, Estonia.

^fAlso at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia.

^gAlso at Universidade Estadual de Campinas, Campinas, Brazil.

^hAlso at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France.

ⁱAlso at Université Libre de Bruxelles, Bruxelles, Belgium.

^jAlso at Joint Institute for Nuclear Research, Dubna, Russia.

^kAlso at Suez University, Suez, Egypt.

^lAlso at Cairo University, Cairo, Egypt.

^mAlso at Fayoum University, El-Fayoum, Egypt.

ⁿAlso at British University in Egypt, Cairo, Egypt.

^oAlso at Ain Shams University, Cairo, Egypt.

^pAlso at Université de Haute Alsace, Mulhouse, France.

^qAlso at Brandenburg University of Technology, Cottbus, Germany.

- ^r Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.
- ^s Also at Eötvös Loránd University, Budapest, Hungary.
- ^t Also at University of Debrecen, Debrecen, Hungary.
- ^u Also at University of Visva-Bharati, Santiniketan, India.
- ^v Also at King Abdulaziz University, Jeddah, Saudi Arabia.
- ^w Also at University of Ruhuna, Matara, Sri Lanka.
- ^x Also at Isfahan University of Technology, Isfahan, Iran.
- ^y Also at University of Tehran, Department of Engineering Science, Tehran, Iran.
- ^z Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran.
- ^{aa} Also at Università degli Studi di Siena, Siena, Italy.
- ^{bb} Also at Centre National de la Recherche Scientifique (CNRS) - IN2P3, Paris, France.
- ^{cc} Also at Purdue University, West Lafayette, USA.
- ^{dd} Also at International Islamic University of Malaysia, Kuala Lumpur, Malaysia.
- ^{ee} Also at Institute for Nuclear Research, Moscow, Russia.
- ^{ff} Also at St. Petersburg State Polytechnical University, St. Petersburg, Russia.
- ^{gg} Also at California Institute of Technology, Pasadena, USA.
- ^{hh} Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia.
- ⁱⁱ Also at Facoltà Ingegneria, Università di Roma, Roma, Italy.
- ^{jj} Also at Scuola Normale e Sezione dell'INFN, Pisa, Italy.
- ^{kk} Also at University of Athens, Athens, Greece.
- ^{ll} Also at Paul Scherrer Institut, Villigen, Switzerland.
- ^{mm} Also at Institute for Theoretical and Experimental Physics, Moscow, Russia.
- ⁿⁿ Also at Albert Einstein Center for Fundamental Physics, Bern, Switzerland.
- ^{oo} Also at Gaziosmanpasa University, Tokat, Turkey.
- ^{pp} Also at Adiyaman University, Adiyaman, Turkey.
- ^{qq} Also at Mersin University, Mersin, Turkey.
- ^{rr} Also at Cag University, Mersin, Turkey.
- ^{ss} Also at Piri Reis University, Istanbul, Turkey.
- ^{tt} Also at Anadolu University, Eskisehir, Turkey.
- ^{uu} Also at Ozyegin University, Istanbul, Turkey.
- ^{vv} Also at Izmir Institute of Technology, Izmir, Turkey.
- ^{ww} Also at Necmettin Erbakan University, Konya, Turkey.
- ^{xx} Also at Mimar Sinan University, Istanbul, Istanbul, Turkey.
- ^{yy} Also at Marmara University, Istanbul, Turkey.
- ^{zz} Also at Kafkas University, Kars, Turkey.
- ^{aaa} Also at Yildiz Technical University, Istanbul, Turkey.
- ^{bbb} Also at Rutherford Appleton Laboratory, Didcot, United Kingdom.
- ^{ccc} Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
- ^{ddd} Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia.
- ^{eee} Also at Argonne National Laboratory, Argonne, USA.
- ^{fff} Also at Erzincan University, Erzincan, Turkey.
- ^{ggg} Also at Texas A&M University at Qatar, Doha, Qatar.
- ^{hhh} Also at Kyungpook National University, Daegu, Korea.