

1 **The effects of orbital precession on tropical precipitation: Mechanisms**
2 **controlling precipitation changes over land and ocean**

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24

Abstract

25 The tropical precipitation response to precessional forcing is investigated using
26 idealized precession experiments from the Geophysical Fluid Dynamics Laboratory
27 Coupled Model version 2.1 and mid-Holocene experiments from ten general circulation
28 models participating in the Paleoclimate Modeling Intercomparison Project Phase III.
29 Both sets of experiments show a seasonal land-ocean asymmetry in the tropical
30 precipitation response: precipitation increases over land and decreases over ocean in the
31 season with increased insolation and the opposite is true in the season with decreased
32 insolation. This response is examined using a framework that describes how changes in
33 net top-of-atmosphere radiation affect the atmosphere and surface energy balances. Over
34 land, surface energy storage is small and changes in precipitation are balanced by
35 changes in moist static energy flux divergence. Over ocean, surface energy storage is
36 large, moist static energy flux divergence is small, and changes in precipitation are
37 ultimately driven by changes in circulation and atmospheric stability.

38 **1. Introduction**

39 Earth's orbital variations impact climate over glacial-interglacial timescales by
40 altering the amount of insolation reaching the top of the atmosphere. Orbital precession
41 modulates seasonal insolation and impacts climate on cycles of about 20 kyr.
42 Precessional signals have been observed in a variety of tropical precipitation records
43 including speleothems from Asia and South America [*Cruz et al.*, 2005; *X Wang et al.*,
44 2006; *X Wang et al.*, 2007; *Y Wang et al.*, 2008] and sediment cores from Asia, South
45 America, and Africa [*Street and Grove*, 1979; *F. Gasse et al.*, 1991; *deMenocal et al.*,
46 2000; *Françoise Gasse*, 2000; *Bush et al.*, 2002; *Trauth et al.*, 2003; *Herzschuh*, 2006]
47 Paleorecords also provide evidence that intensified monsoon precipitation occurred
48 during the mid-Holocene, a period about 6 ka that was characterized by enhanced
49 insolation due to precessional forcing [*Winkler and Wang*, 1993; *Yu and Harrison*, 1996;
50 *Jolly et al.*, 1998; *Yu et al.*, 1998; *Kohfeld and Harrison*, 2000; *Baker et al.*, 2001; *Haug*
51 *et al.*, 2001; *Marchant et al.*, 2009]. Understanding the relationship between precessional
52 changes in insolation and precipitation can help explain how precipitation changed in past
53 climates on glacial-interglacial timescales and how it might change in the future.

54 A variety of modeling studies show that a land-ocean shift in precipitation exists
55 as a result of mid-Holocene forcing [*Braconnot et al.*, 2007; *Braconnot et al.*, 2008; *Hsu*
56 *et al.*, 2010; *Bosmans et al.*, 2012; *Zhao and Harrison*, 2012] and idealized precessional
57 forcing [*Tuenter et al.*, 2003; *Clement et al.*, 2004]. Here, the term “idealized” refers to
58 experiments that do not correspond to any specific past climates. This land-ocean shift in
59 precipitation was originally described as land and sea breezes caused by differences in
60 the heating of land and ocean [*Kutzbach and Ottobliesner*, 1982; *COHMAP Members*,

61 1988; *Tuenter et al.*, 2003; *Ruddiman*, 2008; *Bosmans et al.*, 2012]. *Bosmans et al.*
62 [2012] summarize this mechanism and show that increased summer insolation
63 strengthens the thermal low over the warming continents, which increases the land-sea
64 pressure gradient and the monsoon winds. However, *Merlis et al.* [2013a] suggest that
65 changes in precipitation cannot simply be explained through differences in the heating of
66 land and ocean, because changes in the monsoonal circulation are constrained by changes
67 in the atmospheric energy balance rather than by the surface temperature gradient. In an
68 aquaplanet configuration, this can lead to the counterintuitive result of a weaker
69 circulation in the summer with perihelion, despite an increase in the precipitation there
70 [*Merlis et al.*, 2013b]. Hence, *Merlis et al.* [2013a] argue for a more fundamental role of
71 the energetic framework in studies of the tropical precipitation response to radiative
72 perturbations. *Merlis et al.* [2012] develop a framework describing precessional changes
73 in monsoonal circulation in terms of changes in the net top-of-atmosphere (TOA)
74 radiation with idealized aquaplanet simulations, and *Merlis et al.* [2013a] apply the
75 framework to simulations with a zonally symmetric land surface. The purpose of this
76 study is to extend the framework of *Merlis et al.* [2012] and *Merlis et al.* [2013a] to a
77 climate model with a realistic distribution of land and a more complete set of physical
78 processes in order to better understand the tropical hydrological response to precessional
79 forcing.

80 *Hsu et al.* [2010] also discusses the land-ocean asymmetry of tropical
81 precipitation changes in the mid-Holocene using an energetic approach similar to the one
82 used by *Merlis et al.* [2013a]. They use a balance of moist static energy (MSE) and
83 radiative fluxes at the TOA and surface to explain changes in precipitation over land and

84 ocean, and find that the previous season insolation forcing is important for the magnitude
85 of the land-ocean precipitation change; the different heat capacity of land and ocean and
86 the associated slow response time of the ocean allows the insolation forcing of one season
87 to cancel out the insolation forcing of the following season, causing changes in
88 precipitation over ocean to lag the changes in precipitation over land. While this
89 previous-season insolation mechanism accounts for part of the magnitude of the
90 precipitation change, the direct-season insolation is still significant for the presence of the
91 land-ocean contrast. Our study focuses on the mechanism controlling the direct-season
92 response of the precipitation changes. Applying a TOA energetic framework to idealized
93 precession experiments can help elucidate how tropical precipitation changes in response
94 to radiative forcing and highlight the differences in the mechanisms controlling
95 precipitation changes over land and ocean. Using the atmospheric energy budget
96 sidesteps some of the large but canceling terms, such as the convective diabatic heating,
97 that appear in the individual moisture and temperature equations [*Neelin, 2007*]. This
98 approach emphasizes the importance of the net energy input into the atmospheric column
99 and its implications for understanding the precipitation response to changes in the TOA
100 energy.

101 We also examine the changes in tropical precipitation over land and ocean in the
102 Paleoclimate Modeling Intercomparison Project Phase III (PMIP3) mid-Holocene
103 experiments. Analyzing the previously documented land-ocean asymmetry in a suite of
104 climate models with a standardized experimental design can help assess the robustness of
105 this precipitation signal and motivate the need for a complete mechanistic understanding
106 of precessional changes in precipitation.

107 Section 2 describes the models and experiments used in this study. In Section 3,
108 we discuss the seasonal precipitation change over land and ocean in the PMIP3 models,
109 while results from the idealized experiments and mechanisms used to constrain
110 precipitation changes over land and ocean are examined in Section 4. A brief discussion
111 follows in Section 5.

112

113 **2. Methodology**

114 We use two sets of experiments for this study. The first is from PMIP3, a project
115 that aims to coordinate paleoclimate modeling and model evaluation using standardized
116 experiments from various models and research institutions across the globe. Here, we use
117 climatological monthly data from the PMIP3 mid-Holocene and pre-industrial control
118 experiments. Table 1 lists the 10 models used in this study, the model years used to
119 compute each climatology, and the reference for each model.

120 We use mid-Holocene experiments (MH) in which the autumnal equinox is set
121 close to perihelion ($\omega-180=0.87$, where ω is the longitude of perihelion), and
122 preindustrial control experiments (PI) in which the Northern Hemisphere (NH) winter
123 solstice occurs near perihelion ($\omega-180=102.04$). Further details of the forcings and
124 boundary conditions can be found on the PMIP3 website (<https://pmip3.lsce.ipsl.fr/>).
125 Although there are slight variations in the greenhouse gas concentrations prescribed for
126 MH and PI, the changes resulting from those forcings are small compared to those
127 resulting from the orbital forcing [*Bosmans et al.*, 2012]. We examine the multi-model
128 mean of the difference in the monthly climatologies between the two experiments (MH-
129 PI).

130 To examine in more detail the mechanisms controlling the precession-induced
131 shift in precipitation between land and ocean, we use idealized simulations from the
132 Geophysical Fluid Dynamics Laboratory (GFDL) Coupled Model version 2.1 (CM2.1).
133 The atmosphere and land components have a horizontal resolution of 2° latitude x 2.5°
134 longitude, and the atmosphere has 24 vertical levels. The ocean model has a horizontal
135 resolution of 1° latitude x 1° longitude – with a latitudinal resolution that gradually
136 decreases equatorward of 30° so that it is $1/3^\circ$ at the equator – and 50 vertical levels
137 [Delworth *et al.*, 2006].

138 We use two experiments in which the precessional signal is stronger than in the
139 mid-Holocene experiments. In the first experiment (hereafter referred to as the summer
140 solstice experiment "SS"), perihelion is set at the NH summer solstice ($\omega-180=270$) and
141 in the second experiment (hereafter referred to as the winter solstice experiment "WS"),
142 perihelion is set at the NH winter solstice ($\omega-180=90$). Because eccentricity modulates
143 the strength of precessional forcing [Jackson and Broccoli, 2003], eccentricity is
144 increased to 0.0493 in order to maximize the model's response to the precessional
145 forcing. All other parameters are set to pre-industrial levels. A complete description of
146 the experimental design can be found in Erb *et al.* [2013]. The simulations were run for
147 600 years of which we examine the last 100 years and compute the monthly mean climate
148 changes between the two experiments (SS-WS).

149

150 **3. Precipitation Changes During the Mid-Holocene**

151 It is of interest to first determine if today's climate models show a robust
152 precipitation signal over land and ocean in response to changes in orbital precession. We

153 do this by examining the multi-model mean seasonal cycle of precipitation change over
 154 tropical land and ocean in the PMIP3 models (Fig. 1). In boreal summer, the precipitation
 155 change is positive over land and negative over ocean, indicating that the season with
 156 enhanced insolation also receives enhanced precipitation over land. The opposite is true
 157 in boreal winter: the earth receives less insolation, and the change in precipitation is
 158 negative over land and positive over ocean. This result shows that ten of the climate
 159 models that participated in PMIP3 respond with a robust land-ocean shift in precipitation,
 160 reinforcing the importance of understanding the mechanisms that drive this distinct
 161 precipitation signal.

162

163 **4. Precipitation Changes in Idealized Experiments**

164 **4.1. TOA Energetic Mechanism**

165 We now adopt the energetic framework used by *Merlis et al.* [2013a] who
 166 constrained the response of the monsoonal Hadley circulation to precessional forcing
 167 with an energy budget at the TOA. Their framework was applied to idealized model
 168 simulations with a zonally symmetric continent. Here, we extend this framework to
 169 GFDL-CM2.1, which has a realistic land configuration.

170 From *Merlis et al.* [2013a], the energy balance of the atmosphere and surface is

$$171 \quad \left\{ \frac{\partial \bar{E}}{\partial t} \right\} + \rho_o c_{po} d \frac{\partial \bar{T}_s}{\partial t} = S_{TOA} - L_{TOA} - \nabla \cdot \{ \overline{\mathbf{v}h} \} - \nabla \cdot \mathbf{F}_o, \quad (1)$$

172

173 where the time-mean is denoted by $\overline{(\cdot)}$ and the vertically mass-weighted integral is
 174 denoted by $\{\cdot\}$. $E = c_v T + gz + Lq$ is the total atmospheric energy and $h = c_p T + gz +$
 175 Lq is the MSE. S_{TOA} is the net TOA shortwave radiation, L_{TOA} is the net TOA longwave

176 radiation, ρ_o is the ocean density, c_{po} is the ocean heat capacity, d is the mixed layer
 177 depth, T_s is the surface temperature, $\nabla \cdot F_o$ is the ocean energy flux divergence, and
 178 $\nabla \cdot \{\overline{\mathbf{v}h}\}$ is the atmospheric MSE flux divergence, where \mathbf{v} is the horizontal winds. If
 179 atmospheric and land energy storage are considered negligible, the balance in (1)
 180 becomes

$$181 \quad S_{TOA} - L_{TOA} \approx A\rho_o c_{po} d \frac{\delta \overline{T_s}}{\delta t} + \nabla \cdot \{\overline{\mathbf{v}h}\} + \nabla \cdot F_o, \quad (2)$$

182 where

$$A = \begin{cases} 1 & \text{over ocean} \\ 0 & \text{over land.} \end{cases}$$

183

184 If we consider insolation changes, the perturbation energy balance becomes

$$185 \quad \delta(S_{TOA} - L_{TOA}) \approx A\rho_o c_{po} d \delta \frac{\delta \overline{T_s}}{\delta t} + \delta \nabla \cdot \{\overline{\mathbf{v}h}\} + \delta \nabla \cdot F_o. \quad (3)$$

186

187 Over land, the surface storage term and the ocean energy flux divergence are
 188 zero, so changes in the net (shortwave minus longwave) TOA radiation are balanced
 189 solely by changes in MSE flux divergence. Over ocean, changes in net TOA radiation can
 190 be balanced by changes in MSE flux divergence, ocean energy flux divergence, and
 191 surface energy storage.

192 While Eq. (3) is useful because it describes the energy balance of the atmosphere
 193 and the surface, it does not explicitly account for changes in precipitation. In order to
 194 understand how changes in net TOA radiation lead to changes in precipitation, we link
 195 the energy and moisture budgets by considering the atmosphere in equilibrium. The
 196 surface warms through radiative heating and convection heats the upper atmosphere

197 through precipitation and latent heat release. Therefore, in monsoon regions where
198 sufficient surface moisture is available, changes in net TOA radiation over land are
199 balanced by changes in precipitation:

$$200 \quad \delta P_L \approx \delta(S_{TOA} - L_{TOA}) \approx \delta \nabla \cdot \{\overline{vh}\}, \quad (4)$$

201

202 where P_L is the precipitation over land.

203

204 **4.2. Results**

205 Similar to the mid-Holocene experiments, a land-ocean asymmetry in
206 precipitation exists as a result of the idealized precessional forcing in the GFDL-CM2.1
207 simulations (Fig. 2). In order to examine the mechanism driving this precipitation signal,
208 we examine the tropical-mean changes in net TOA radiation, net surface fluxes, and MSE
209 flux divergence over land and ocean for both models. As expected from Eqs. (3) and (4),
210 changes in precipitation have the same sign as the changes in net TOA radiation over
211 land surfaces. Energy is redistributed from the surface to the atmosphere through vertical
212 motion, which results in convection and latent heat release. Therefore, over land,
213 increases in net TOA radiation correspond to increases in precipitation and decreases in
214 net TOA radiation correspond to decreases in precipitation.

215 Over ocean, the sign of the change in net TOA radiative forcing is the same as it
216 is over land, but the change in precipitation is opposite. This difference arises as a result
217 of the ocean energy storage: Fig. 2 indicates that over ocean, changes in net TOA
218 radiation are almost entirely balanced by changes in the surface energy fluxes. The
219 surface storage term dominates over the ocean heat flux divergence term (not shown)

220 indicating that the energy flux is warming or cooling the ocean surface. The ocean
221 responds to the increase in net TOA radiation during boreal summer by absorbing energy
222 and warming the surface, and it responds to the decrease in net TOA radiation during
223 boreal winter by releasing energy and cooling the surface. As a result, the change in MSE
224 flux divergence over ocean is very small; the atmosphere over ocean does much less
225 work to adjust to the TOA forcing than the atmosphere over land.

226 While this analysis is useful for understanding the atmospheric energy budget
227 over both land and ocean, it does not explicitly explain why precipitation decreases over
228 ocean when net TOA radiation is increasing, and vice versa. To further understand the
229 oceanic precipitation response, we examine the seasonal cycle of the moisture flux
230 divergence over ocean.

231 As in *Clement et al.* [2004], we decompose the moisture flux convergence into the
232 effects of the circulation changes, $-\nabla \cdot \{\mathbf{v}'\bar{q}\}$, and the effects of the moisture changes,
233 $-\nabla \cdot \{\bar{\mathbf{v}}q'\}$, where q is the specific humidity, $\overline{(\cdot)}$ denotes the mean (WS) field and $(\cdot)'$
234 denotes a deviation from the mean. Fig. 3 shows the change in the seasonal cycle of these
235 terms over ocean. It is evident that the seasonal change in circulation, i.e., the dynamic
236 component, dominates the seasonal change in the specific humidity, i.e., the
237 thermodynamic component (this also true over land, not shown). The effect of the
238 circulation change is such that moisture is diverging over ocean in boreal summer and
239 converging over ocean in boreal winter. This is consistent with the change in
240 precipitation over ocean. We deduce from Fig. 3 that a dynamic mechanism is controlling
241 precipitation changes over ocean (i.e., subsidence and decreased precipitation during
242 boreal summer and ascent and increased precipitation during boreal winter). This differs

243 from the mechanism controlling precipitation changes over land, which is dictated by the
244 atmospheric energy response to a change in net TOA radiation.

245 The thermodynamic structure of the atmosphere may be necessary to explain how
246 precipitation can change over ocean despite little change in the MSE flux divergence.
247 Fig. 4 shows the change in the vertical profiles of the dry and latent energy over land and
248 ocean in boreal winter (DJF) and boreal summer (JJA). The dry energy is taken as $c_p\theta$,
249 where θ is the potential temperature and c_p is the specific heat for dry air, and latent
250 energy is taken as L_vq , where L_v is the latent heat of vaporization. The latent energy is
251 most important in the lower levels of the troposphere, and it decreases during DJF when
252 the surface is cooling and increases in JJA when the surface is warming. The profile of
253 $c_p\theta$ has small seasonal change over land and a larger change over ocean. Here, the
254 atmosphere becomes more unstable ($d\theta/dp > 0$) in DJF and more stable ($d\theta/dp < 0$) in
255 JJA.

256 We suggest that a reorganization of the thermodynamic structure of the
257 atmosphere over ocean, which is driven by circulation changes over land, may ultimately
258 allow for a change in precipitation. For example, in boreal summer when there is a net
259 gain in the TOA radiation, the atmosphere over land responds to this forcing through an
260 export of MSE, which is achieved through anomalous vertical motion and an increase in
261 precipitation. It is this circulation change over land that drives anomalous descent over
262 ocean. As we saw in the previous section, the atmosphere over ocean does not import
263 more energy; it may instead adjust by becoming more stable, consistent with less rainfall
264 over ocean.

265

266 **5. Discussion**

267 We examined the precipitation response to mid-Holocene and idealized
268 precessional forcing using a variety of climate models. All models agree that a land-
269 ocean shift in precipitation exists as a result of precessional forcing. The idealized
270 simulations were used to examine the mechanisms controlling this distinct precipitation
271 response. We found that while the TOA energetic mechanism outlined in *Merlis et al.*
272 [2013a] explains how precipitation over land responds to precessional forcing, it does not
273 fully explain the precipitation response over ocean.

274 The fundamental constraint dictating the energy balance and resulting
275 precipitation change over land and ocean lies in the surface boundary condition. Over
276 land, the surface stores very little energy and the atmosphere is required to do all the
277 work to compensate for the net TOA energy imbalance. We see this in the MSE flux
278 divergence: as the net TOA radiation increases, the atmosphere responds by diverging
279 MSE and latent heat is released through an increase in precipitation, and as net TOA
280 radiation decreases, the atmosphere converges MSE to maintain equilibrium and
281 precipitation decreases. Over ocean, the energy balance is not so simple because surface
282 energy storage plays an important role in balancing the TOA energy fluxes. When the net
283 TOA radiation increases over ocean, the surface warms and the change in MSE flux
284 divergence does not need to be as large as it is over land. When there is a decrease in the
285 net TOA radiation, the ocean surface cools and, again, the change in MSE flux
286 divergence over ocean is small. Therefore, it is the differing surface properties of the land
287 and ocean that allow for very different atmospheric responses to precessional forcing.

288 While these changes in the energy budgets of the atmosphere and surface are
289 useful for understanding the precipitation response to precessional forcing over land, they
290 do not explicitly explain the land-ocean asymmetry. Analyses of the change in moisture
291 flux convergence and the dry and latent vertical energy profiles of the atmosphere
292 suggest that the atmospheric energy response over land may drive changes in circulation
293 that affect the thermodynamic structure over ocean, allowing for changes in precipitation
294 to occur.

295 Further insight into this mechanism for describing precipitation changes may
296 provide more details of the climate response to precession than mechanisms traditionally
297 invoked in monsoon dynamics. Understanding the mechanisms that drive changes in
298 precipitation due to precessional forcing can not only help explain precipitation changes
299 in past and future climates, but it can also help explain how the precipitation signal
300 manifests in regional climate records. Moreover, it is of interest to examine regional
301 precipitation changes because over glacial-interglacial timescales, the regional response
302 to solar forcing may differ significantly from the global response [*Clement et al.*, 2004].
303 This has important implications for the interpretation of climate records and for regional
304 climate sensitivity.

305

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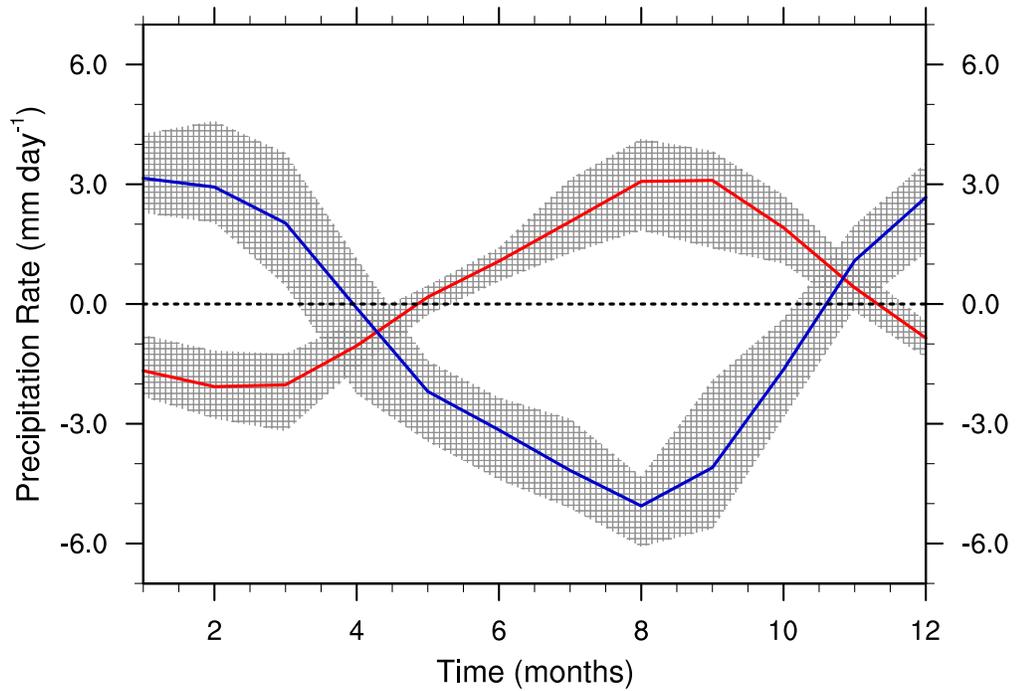
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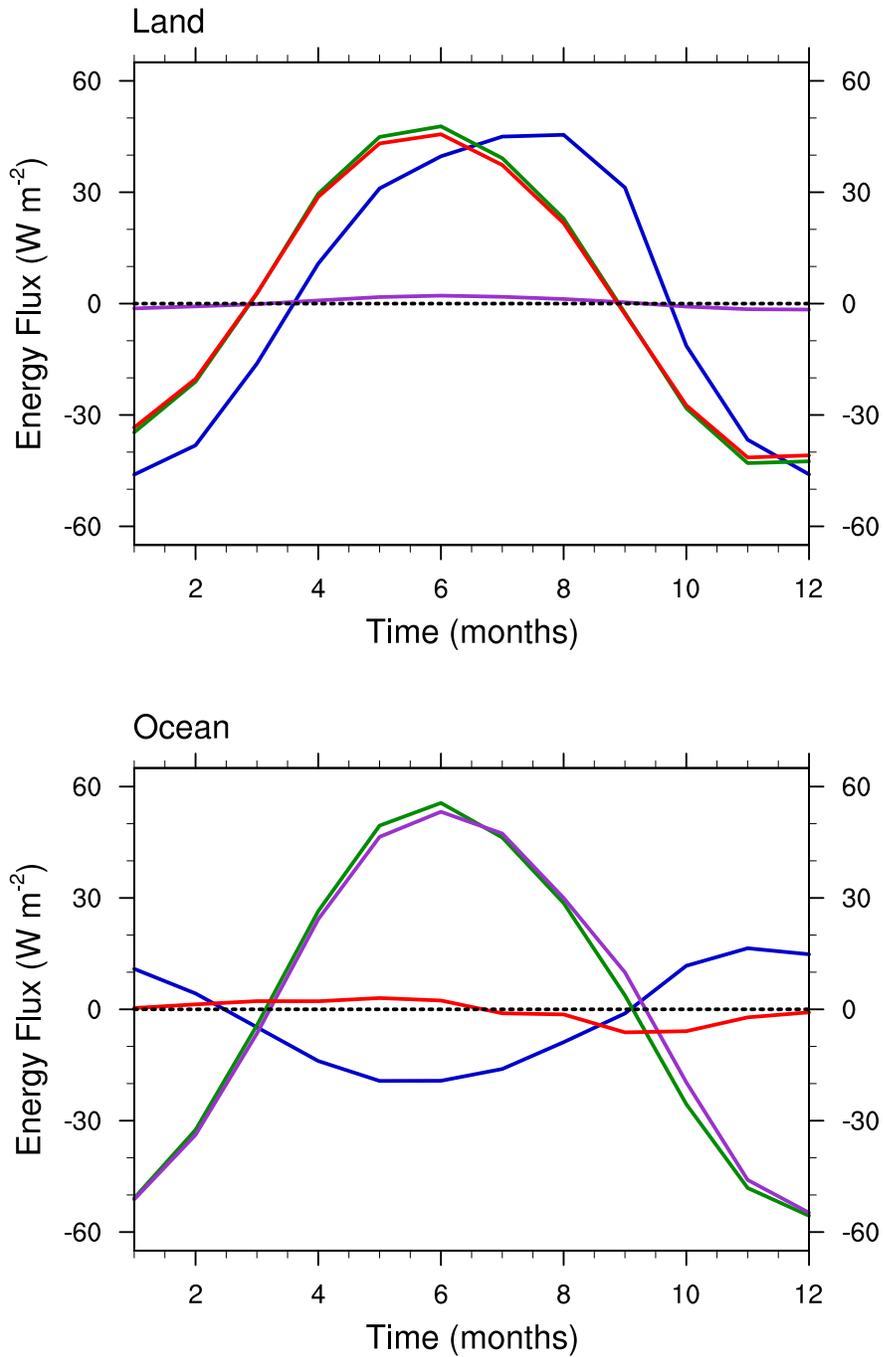
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Model Name	Model Years for Climatology		Reference
	MH	PI	
CCSM4	1000-1300	250-1300	<i>Gent et al.</i> [2011]
CNRM-CM5	1950-2149	1850-2699	<i>Voldoire et al.</i> [2013]
CSIRO-Mk3L-1-2	1-500	1-1000	<i>Phipps et al.</i> [2011]
FGOALS-g2	340-1024	201-900	<i>Li et al.</i> [2013]
FGOALS-s2	1-100	1850-2350	<i>Bao et al.</i> [2013]
GISS-E2-R	2500-2599	3331-4530	<i>Schmidt et al.</i> [2014]
IPSL-CM5A-LR	2301-2800	1800-2799	<i>Dufresne et al.</i> [2013]
MIROC-ESM	2330-2429	1800-2429	<i>Watanabe et al.</i> [2011]
MPI-ESM-P	1850-1949	1850-3005	<i>Giorgetta et al.</i> [2013]
MRI-CGCM3	1951-2050	1851-2350	<i>Yukimoto et al.</i> [2012]

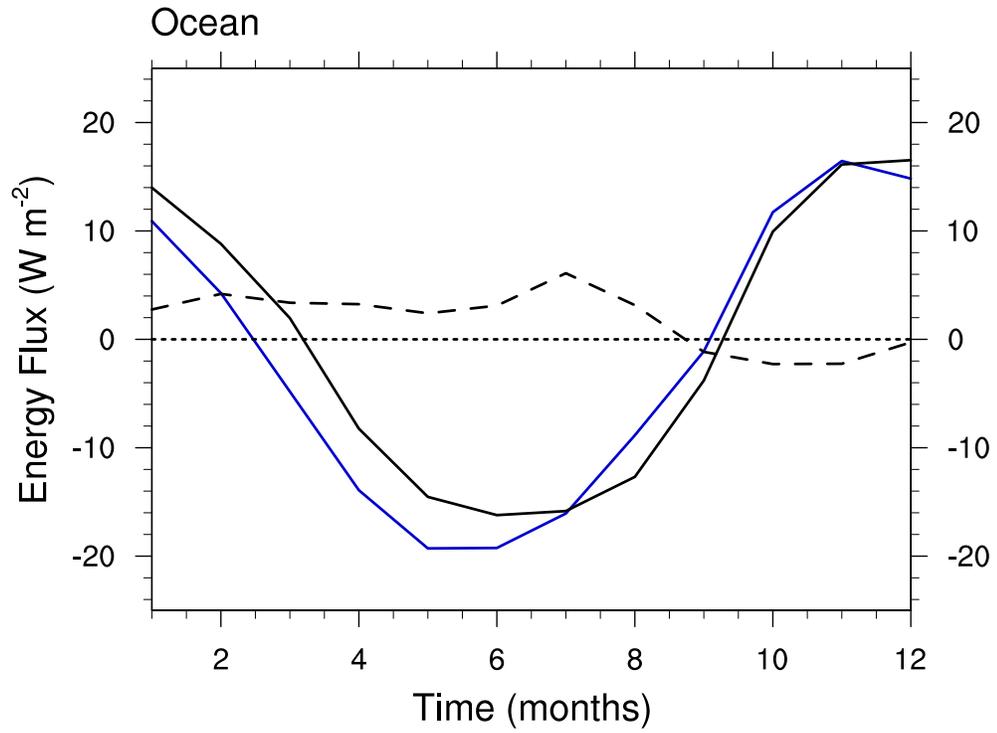
458 Table 1. PMIP3 models used in this analysis, model years averaged over for the mid-
459 Holocene and pre-industrial climatologies, and references.



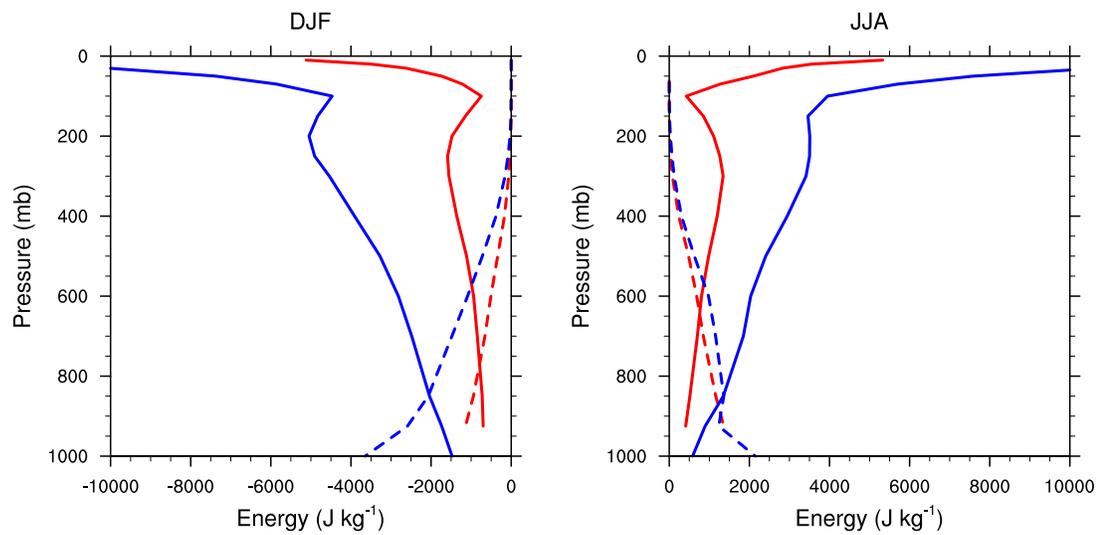
460 Figure 1: Multi-model mean seasonal cycle of the precipitation difference as a function of
 461 time between experiments (MH-PI) for the 10 PMIP3 models (listed in Table 1) over land
 462 (red) and ocean (blue). Shading indicates the model spread. Curves are scaled by the
 463 percentage of land and ocean in the tropics.



464 Figure 2: Seasonal cycle of the difference (SS-WS) in tropical-mean (30°N to 30°S)
 465 precipitation (blue), net TOA radiation (green), net surface energy fluxes (purple), and
 466 MSE flux divergence (red) as a function of time over land and ocean for GFDL-CM2.1.



467 Figure 3: Seasonal cycle of the difference (SS-WS) in tropical-mean (30°N to 30°S)
 468 precipitation (blue) and moisture flux convergence broken down into circulation changes
 469 (solid) and specific humidity changes (dashed) as a function of time over ocean for
 470 GFDL-CM2.1.



471 Figure 4: Vertical profile of the difference (SS-WS) in tropical-mean (30°N to 30°S) dry
 472 (solid) and latent (dashed) energy as a function of pressure over land (red) and ocean
 473 (blue) for GFDL-CM2.1. Curves are scaled by the percentage of land and ocean in the
 474 tropics.