

Are all maximally entangled states pure?

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We study if all maximally entangled states are pure through several entanglement monotones. In the bipartite case, we find that the same conditions which lead to the uniqueness of the entropy of entanglement as a measure of entanglement, exclude the existence of maximally mixed entangled states. In the multipartite scenario, our conclusions allow us to generalize the idea of monogamy of entanglement: we establish the *polygamy of entanglement*, expressing that if a general state is maximally entangled with respect to some kind of multipartite entanglement, then it is necessarily factorized of any other system.

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One of the most striking differences between classical and quantum correlations is the restricted capability of quantum states to share entanglement. This so called *monogamy of entanglement* has been increasingly studied in the last years [1], and is related to the security of quantum cryptographic protocols based on entanglement (it limits the amount of correlations which an eavesdropper can have with the honest parties).

The discussion about the monogamy of entanglement usually begins with the apparent straightforward fact that maximally entangled states are pure. This means that when two systems are as much entangled with each other as it is possible, they cannot be entangled and even classically correlated with any one else. In this paper we analyze the trustiness of this “common sense” under the view of several entanglement measures. On one hand, it is shown that it is not true for general entanglement monotones, failing at least in the best separable approximation measure [2] and in the indicator measure [3]. On the other hand, we prove that for the majority of entanglement quantifiers maximally entangled states are indeed pure. In particular, we consider the quantifiers related to entanglement witnesses and, in special, the generalized robustness of entanglement. With the help of the *witnessed entanglement* [4], we introduce the idea of polygamy of entanglement, which states that *if a multipartite state is maximally entangled with respect to a given kind of multipartite entanglement, then it must be pure* [23].

In order to avoid future confusion it is important to stress that the idea of *mixed maximally entangled states (MMES)* discussed here is, although related, different from the idea of *maximally entangled mixed states (MEMS)* presented in the Refs. [6] (this is the reason for the exchange of words). In their articles the authors address the following question: what is the highest value of entanglement that states with a given purity (mixing) can present? In our work we study if, given the maximum value of entanglement, there is some mixed state

which reaches it.

Before proceeding to show the main result of this Letter, we present two simple results, albeit important, valid for every convex quantifier.

Theorem 1. *According to all convex entanglement quantifiers there is at least one maximally entangled pure state.*

Proof. Given a density operator $\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$ and a convex quantity E , it holds $E(\rho) \leq \sum_i p_i E(|\psi_i\rangle\langle\psi_i|)$, for every ensemble decomposition $\{p_i, \psi_i\}$ of ρ . Thus, we see that there must be a $|\psi_i\rangle$ such that $E(|\psi_i\rangle\langle\psi_i|) \geq E(\rho)$. \square

Moreover, from the convex condition it is easily seen that for a mixed state to be maximally entangled (with respect to E), there must be an ensemble description with all $|\psi_i\rangle$ maximally entangled.

Theorem 2. *If ρ is a mixed maximally entangled state with respect to the convex measure E , then all states in the subspace spanned by the eigenvectors of ρ are maximally entangled.*

Proof. According to the unitary freedom in the ensemble for density matrices theorem [7], the sets $\{p_i, |\psi_i\rangle\}$ and $\{q_j, |\phi_j\rangle\}$ generate ρ , i.e.

$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i| = \sum_j q_j |\phi_j\rangle\langle\phi_j|, \quad (1)$$

if and only if $\sqrt{p_i}|\psi_i\rangle = \sum_j u_{ij} \sqrt{q_j}|\phi_j\rangle$, with $\{|\psi_i\rangle\}$ and $\{|\phi_j\rangle\}$ being normalized vectors, u_{ij} a complex unitary matrix and one can ‘pad’ whichever set of vectors $\sqrt{p_i}|\psi_i\rangle$ or $\sqrt{q_j}|\phi_j\rangle$ is smaller with additional null vectors so that the two sets have the same number of elements.

Since each pure state term in any convex decomposition of ρ must be maximally entangled, we find that the state $|\psi_i\rangle = \sum_j c_{ij} |\phi_j\rangle$, with coefficients $c_{ij} = u_{ij} \sqrt{\frac{q_j}{p_i}}$, must be maximally entangled as well. The result follows

letting $\{|\phi_j\rangle\}$ be the eigenvectors of ρ and noting that for a fixed i , the vector c_{ij} can have arbitrary elements belonging to the hypersphere $\sum_j c_{ij}^2 = 1$. \square

It is possible to extend entanglement measures defined for pure states to the whole state space with the convex-roof construction. Given the quantity E , its convex-roof is

$$E(\rho) = \min_{\{p_i, \psi_i\}} \sum p_i E(\psi_i). \quad (2)$$

It can be shown that $E(\rho)$ is an entanglement monotone whenever $E(\psi)$ is. From Eq. (2) we see that, for convex-roof based measures, Theorem (2) gives necessary and sufficient conditions for the existence of mixed maximally entangled states.

We believe the existence of a n -dimensional subspace, with $n \geq 2$, formed only by maximally entangled states is a very demanding condition, so that for general convex entanglement measures the maximally entangled states are pure. One might then conjecture that this is true for all entanglement monotones [8]. However, for the indicator measure [3] defined as 1 for entangled states and 0 for separable states, which is obviously an entanglement monotone, every entangled state is maximally entangled. Furthermore, using the result of Ref. [9] which for every k -partite Hilbert space H there exists an entangled subspace of dimension $d_1 d_2 \dots d_k - (d_1 + d_2 + \dots + d_k) + k - 1$, we find that also for the convex-roof indicator measure maximally entangled states can be mixed.

Theorem 3. *According to every bipartite entanglement measure such that all its maximally entangled pure states have maximum Schmidt rank possible [24], there do not exist maximally entangled mixed states.*

Proof. By theorem (2), it must exist a subspace of maximally entangled pure states. By hypothesis, they all have maximum Schmidt rank. Take two of them,

$$|\psi\rangle = \sum_{ij} c_{ij} |ij\rangle, \quad |\phi\rangle = \sum_{ij} d_{ij} |ij\rangle. \quad (3)$$

If we look at c_{ij} and d_{ij} as coefficients of square matrices C and D , maximum Schmidt rank is equivalent to invertibility of the matrix. However, it always exists $\alpha, \beta \in \mathcal{C}$ such that $\alpha C + \beta D$ is not invertible (take $\frac{-\alpha}{\beta}$ as an eigenvalue of $C^{-1}D$), and $\alpha|\psi\rangle + \beta|\phi\rangle$ does not have maximum Schmidt rank. \square

This theorem applies to a number of important and well-studied entanglement measures. Consider first the entanglement of formation [10] and the relative entropy of entanglement [3]. They are both convex and equal to the entropy of entanglement (E_E) in pure states. As all maximally entangled states of E_E are singlets (which have maximum Schmidt rank), it follows from theorem 3

that neither of them allows MMES. The same argumentation is valid for the negativity [11] and the concurrence [12]. We can go even further and establish the following result:

Theorem 4. *For all asymptotic continuous and partially additive entanglement monotones, all maximally entangled states are pure.*

Proof. From the uniqueness theorem for entanglement measures [13], we have that every entanglement measure E fulfilling the conditions of the theorem obey $E_D \leq E \leq E_F$, where E_D and E_F are the distillable entanglement and the entanglement of formation respectively. Hence the result follows straightforwardly from the fact that E_F does not have MMES and $E_D = E_F$ to pure states. \square

The situation is much more subtle when we are dealing with multi-partite entanglement. In this case we have to specify which kind of entanglement we are talking about [14]. This is because we could be interested in studying the entanglement among different partitions of the whole system. Furthermore in the multi-partite context, other relevant questions arise in order to classify entangled states, as there are different classes of inequivalent states under SLOCC [15]. Thus, answering if a state is more entangled than other will depend on what criterion one is adopting.

Consider an m -partite state with Hilbert space $H = \bigotimes_{i=1}^m H_i$. We call $P_k^m = \{A_j\}_{j=1}^N$ a k -partition of $\{1, 2, \dots, m\}$ if: **1.** $A_j \subset \{1, 2, \dots, m\}$; **2.** $A_i \cap A_j = \emptyset, \forall i \neq j$; **3.** $\bigcup_i A_i = \{1, 2, \dots, m\}$; **4.** $\#A_j \leq k$. The number k is called the diameter of the partition P_k^m . The set of all k -partitions of $\{1, 2, \dots, m\}$ will be denoted by \mathcal{P}_k^m . With this concept, one can define factorizability and separability subjected to a partition, and also subjected only to the diameter of the partitions.

Definition 1. *We say that a state ρ is P_k^m -factorizable if, for a fixed P_k^m , it can be written as $\rho = \rho_{A_1} \otimes \dots \otimes \rho_{A_n}$, where ρ_{A_j} is a density operator on $H_{A_j} = \bigotimes_{i \in A_j} H_i$. A state is P_k^m -separable if it can be written as a convex combination of P_k^m -factorizable states. Finally, we call k -separable a state ρ which can be written as a convex combination of P_k^m -factorizable states, where P_k^m may vary for each pure state.*

Let us denote $S_k(H)$ the set of k -separable states on H . Clearly they form a chain $S_1(H) \subset S_2(H) \subset \dots \subset S_{m-1}(H) \subset S_m(H) = D(H)$, where $D(H)$ denotes the set of density operators on H . As each of these sets is closed and convex, there exists an Hermitian operator W such that $\text{tr}(W\rho) < 0$, and $\text{tr}(W\sigma) \geq 0 \forall \sigma \in S_k(H)$ [16]. One call such W a k -entanglement witness.

Although several entanglement monotones applicable to multipartite states are known [1, 3, 4, 17, 18], only two

approaches, up to now, can be applied to the quantification of the different kinds of multipartite entanglement discussed above [25]: the relative entropy of entanglement [3] and its related measures [3, 17] and the witnessed entanglement [4]. The first, of great importance in the bipartite scenario, is based on the minimization of some distance between the state under question and the sets $S_k(H)$. The second, recently introduced in Ref. [4], includes several well studied bipartite and multipartite entanglement measures and quantifies entanglement based on the concept of optimal entanglement witnesses. In this paper, due to some particular properties, such as the linearity of the objective function, we will consider the witnessed entanglement:

Definition 2. For a m -partite state $\rho \in D(H)$, its witnessed k -partite entanglement is given by

$$E_W^k(\rho) = \max\{0, -\min_{W \in \mathcal{M}} \text{Tr}(W\rho)\}, \quad (4)$$

where $\mathcal{M} = \mathcal{W}_k \cap \mathcal{C}$, \mathcal{W}_k is the set of k -entanglement witnesses and \mathcal{C} is some set such that \mathcal{M} is compact.

Having this definition in mind we can see what are the restrictions imposed by the existence of a MME-state ρ on its optimal entanglement witness W .

Let $\rho = \sum_j q_j |\phi_j\rangle\langle\phi_j|$ be the spectral decomposition of ρ , and $\{p_i, |\psi_i\rangle\}$ another ensemble describing it. Then, $|\psi_i\rangle = \sum_j c_{ij} |\phi_j\rangle$, with coefficients $c_{ij} = u_{ij} \sqrt{\frac{q_i}{p_i}}$. In the case where ρ is maximally entangled with entanglement E , W must be optimal for every $|\psi_i\rangle$ and $|\phi_j\rangle$, which allows us to write for one specific element $|\psi_k\rangle$,

$$\begin{aligned} -E &= \text{Tr}(W|\psi_k\rangle\langle\psi_k|) \\ &= \sum_i |c_{ki}|^2 \langle\phi_i|W|\phi_i\rangle + \sum_{i \neq j} c_{ki}^* c_{kj} \langle\phi_i|W|\phi_j\rangle \\ &= -E + \sum_{i \neq j} c_{ki}^* c_{kj} \langle\phi_i|W|\phi_j\rangle, \end{aligned} \quad (5)$$

which implies

$$\sum_{i \neq j} c_{ki}^* c_{kj} \langle\phi_i|W|\phi_j\rangle = 0.$$

As this equality must be true for every ensemble describing ρ , $\langle\phi_i|W|\phi_j\rangle = 0$ and W is proportional to the identity matrix in the support of ρ , with $-E$ as the proportionality constant. Being E the highest value of entanglement and, therefore, the modulus of the lowest eigenvalue possible among all entanglement witnesses, each eigenvector $|\phi_j\rangle$ of ρ is an eigenvector of W too. Thus W can be written as

$$W = \underbrace{(-E)I}_{\text{Supp}(\rho)} \oplus \underbrace{D}_{\text{Supp}^\perp(\rho)}, \quad (6)$$

where D is some matrix such that the constraints imposed by \mathcal{C} are satisfied. Here again this demanding

condition is not sufficient to rule out the existence of mixed maximally entangled states. As a counterexample, consider the best separable approximation measure [2], $BSA^k(\rho) = 1 - \lambda$, where λ is the optimal value of the following optimization problem:

$$\max \lambda, \quad \text{s.t.} \quad \rho = \lambda\sigma + (1 - \lambda)\pi, \quad (7)$$

with $\sigma \in S_k(H)$, $\pi \in D(H)$ and $\lambda \in [0, 1]$. It can be written alternatively as Eq. (4) with $\mathcal{C} = \{W \mid W \geq -I\}$ [4]. For the following family of mixed states

$$\rho_q = q|W\rangle\langle W| + (1 - q)|GHZ\rangle\langle GHZ|, \quad (8)$$

BSA was calculated in Ref. [4], using the numerical method presented in Ref. [19], and shown to be composed only of maximally entangled states, with respect either to 1 and 2-entanglement. Note, nonetheless, that despite $BSA(\rho)$ being an entanglement monotone [2], it is a quite odd quantity, as every entangled pure state is maximally entangled.

An important measure of multipartite entanglement is the generalized robustness of entanglement [20],

$$R^k(\rho) = \min s, \quad \text{s.t.} \quad \frac{1}{1+s}\rho + \frac{s}{1+s}\pi = \sigma, \quad (9)$$

where $\sigma \in S_k(H)$ and $\pi \in D(H)$. It gives good bounds for the maximum fidelity of teleportation, the distillable entanglement, and the entanglement of formation [4], and has important applications in the study of threshold of errors in quantum gates [21]. In Ref. [4], it was shown that R can be written as Eq. (4), with $\mathcal{C} = \{W \mid W \leq I\}$.

Lemma 1. For every state $\rho \in D(H)$,

$$\max_{\sigma \in S_k(H)} \text{tr}(\rho\sigma) \geq \frac{\text{tr}(\rho^2)}{1 + R^k(\rho)}. \quad (10)$$

Proof. From the theory of convex optimization and Lagrange duality [22], the optimal value of the L.H.S. of Eq. (10) is given by the solution of the following convex problem

$$\min \lambda \quad \text{s.t.} \quad \lambda I - \rho \in \mathcal{W}_k. \quad (11)$$

Let $W = \lambda_{opt}I - \rho$ be an optimal solution of (11). Since, $W/\lambda_{opt} \leq I$, we find that $R^k(\rho) \geq \text{tr}(\rho^2)/\lambda - 1$, from which the result follows. \square

Theorem 5. There do not exist, for any k , mixed maximally k -entangled states according to the generalized robustness of entanglement.

Proof. We will prove by contradiction that it does not exist an optimal entanglement witness of the form (6). Assume that ρ' is a mixed maximally entangled state with spectral decomposition $\rho' = \sum_{i=1}^m \lambda_i |i\rangle\langle i|$. Then, by the dual definition of R , it is easily seen that $\rho =$

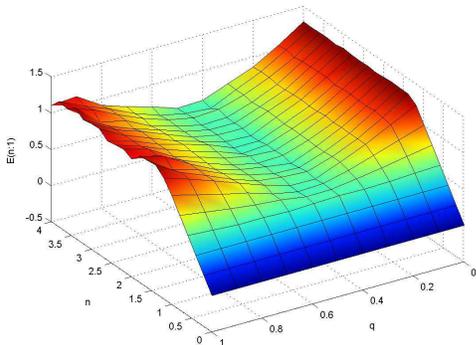


FIG. 1: (Color online) $E_{n:1}^1(\rho_q)$ for $0 \leq n \leq 4$ and $0 \leq q \leq 1$. When $n \ll 1$, $E_{n:1}^1(\rho) = nBSA(\rho)$ and MME-states are possible. In the other limit $E_{n:1}^1(\rho) = R^1(\rho)$ [4], and all maximally entangled states are pure. The irregularities in the figure are due to the approximative nature of the numerical method used [19].

$\lambda|1\rangle\langle 1| + (1-\lambda)|2\rangle\langle 2|$ is also maximal entangled for every $\lambda \in [0, 1]$. Thus we consider, without loss of generality, states of rank two. For the generalized robustness, the matrix D of Eq. (6) must satisfy $D \leq I$. Hence, since for every entanglement witness $W = (-E)I \oplus D$, with $D \leq I$, $W' = (-E)I \oplus I$ is another witness as optimal as W for ρ , we may assume throughout this canonical form. Letting P represents the projector onto the support of ρ , $W = I - (1+E)P$. By assumption, $R(P) = 2E$. From Lemma 1, we find that for some $\sigma \in S_k(H)$:

$$\text{tr}(P\sigma) \geq \frac{\text{tr}(P^2)}{1 + R(P)} = \frac{2}{1 + 2E}. \quad (12)$$

Therefore,

$$\text{Tr}(W\sigma) = 1 - (1+E)\text{tr}(\sigma P) \leq 1 - \frac{2(1+E)}{1+2E} < 0, \quad (13)$$

which contradicts the fact that W is a k -entanglement witness. \square

In Ref. [4], a family $E_{m:n}$ of infinite entanglement monotones which interpolates between the best separable approximation measure and the generalized robustness was proposed. For fixed m and n , $E_{m:n}$ is given by Eq. (4), with $\mathcal{C} = \{W \mid -mI \leq W \leq nI\}$. They provide a tool to the observation of the (smoothly) transition between the regime where there exist mixed maximally entangled states and the regime where all MMES are pure. Fig. 1 shows the transition for the family of states given by Eq. (8).

In conclusion, we have analyzed the existence of maximally mixed entangled states in the bipartite and multipartite scenarios. In the first, we showed that, although monotonicity under LOCC do not exclude MMES, partial additivity and asymptotic continuity together with

monotonicity do. We then extended this result to multipartite systems by showing that maximally entangled (multiparticle-)states are pure. It is now time to ask: what are the physical consequences of this result? One can easily note that every pure state must be completely uncorrelated with any other system (if not it should be written as a non-trivial convex combination, what characterizes mixed states). This notion gives us a solid background to propose the *faithful polygamy of entanglement*, which states that all maximally entangled states are (classically and quantumly) uncorrelated with any other system. One can even propose this condition as another requisite for a good multipartite entanglement quantifier. Furthermore, it is also important to stress that this polygamy holds for all kinds of entanglement, it is, every time the system reaches a maximum amount of entanglement according to any partition, it becomes “free” of its environment.

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- [23] Note that we use the term *polygamy* in the sense of a marriage among multiple partners. This word has appeared before, in Ref. [5], in the context of symmetric multipartite Gaussian states, as opposing to monogamy, in the sense that states which maximizes a certain pair-entanglement quantifier can also maximize the multipartite version of it.
- [24] The Schmidt rank of a bipartite pure state is the number of non-null Schmidt coefficients in its Schmidt decomposition.
- [25] The other measures are either based on bipartite entanglement concepts, such as the localizable entanglement, or can only distinguish entangled from fully-separable states.