Supplementary Material for

*Swift and NuSTAR observations of GW170817: Detection of a blue kilonova*


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Supplementary material for Swift and NuSTAR observations of GW170817: detection of a blue kilonova

Materials and Methods

1 Observations

Advanced LIGO and advanced Virgo (operating as the LIGO/Virgo Consortium: hereafter ‘LVC’) registered GW 170817 at 12:41:04.45 UT on 2017 August 17. This was announced to the followup community first as an LVC/Gamma-ray Coordinates Network (GCN) notice at 13:08 UT, and then by (1), (trigger ‘G298048’). Fermi Gamma-ray Burst Monitor (GBM) triggered on a weak short gamma-ray signal at 12:41:06 UT (4), reported as a short GRB (sGRB, (54)). A weak signal was also observed by the INTEGRAL satellite (55). At this time, the Swift Burst Alert Telescope (BAT; (56)) was pointing at RA, Dec (J2000)= 02h24m18.0s, -52°17′13″, i.e. away from the GBM and GW position. Indeed, the GW localization region was entirely occulted by the Earth (Figure S1), with none of the probability region above the Earth limb (defined as 69° from the center of the Earth, which accounts for the radius of the Earth plus 100 km of atmosphere, up to the Kármán line, which commonly represents the top of the atmosphere). The position of EM 170817A (below) was likewise not visible to Swift at the trigger time, while the very edge of the GBM error region overlapped the edge of the BAT field of view (0.1% of the GBM probability was within the BAT field of view). We found no evidence for a signal in the BAT within 100-s of the GW and GBM triggers.

Swift observations comprising more than one pointing can only be uploaded to the spacecraft in the form of a pre-planned science timeline which requires a ground-station pass. In order to obtain some rapid coverage, we uploaded commands using the NASA Tracking and Data Relay
Satellite System (TDRSS) network to utilize the on-board 37-point tiling pattern. Using this pattern the X-ray Telescope (XRT; (57)) covered a region $\sim 1.1^\circ$ in radius, which was centered on the GBM position. The UV/optical Telescope (UVOT; (58)) has a smaller field of view ($17' \times 17'$) than the XRT ($24.6'$ diameter), thus there are gaps between the tiles in UVOT. These observations began at 13:37 UT, 0.04 d after GW 170817.

The initial GW skymap was created using only a single GW detector (LIGO-Hanford; (59)) and thus covered almost the entire sky. The GBM error had a statistical position error of $11.6^\circ$ radius (1-$\sigma$) with a systematic error estimated to be at least $3.7^\circ$ (4); the GBM localization probability has a 90% confidence area of 1626 deg$^2$. We convolved the GW skymap with the GWGC catalog (60), using the 3D information in the GW skymap and the methodology described by (48). This effectively re-weights the GW probability according to our knowledge of the location of the stellar mass in the local Universe, accounting for our estimate of how incomplete that knowledge is. We next multiplied the resultant probability map by a map of the GBM localization probability and renormalised the result, then identified the most probable XRT fields-of-view from this skymap. Swift-XRT and UVOT observations of these fields, with a nominal 120 s per exposure began at 17:16 UT, 0.19 d after GW 170817. The onboard slewing algorithm is not perfectly modelled, thus exposure times vary someone from the nominal value.

At $\sim 17:30$ UT a revised GW skymap was released which used data from all three LVC interferometers and had a 90% credible region of 31 deg$^2$ (11). Since this error region only overlapped the outer contours of the GBM error region, our planned observations did not significantly cover this revised localisation. We therefore convolved this new skymap with GWGC as above, and generated a new observing plan based on this probability map; these observations began at 20:24 UT, 0.31 d after GW 170817, and are listed in Table S3.

At 01:05 UT on 2017 August 18, (14) reported the detection of a candidate EM counterpart, hereafter denoted ‘EM 170817’, apparently associated with the galaxy NGC 4993 (distance $\approx 40$ Mpc
Shortly thereafter, several others groups reported detection of the source (62, 63), sometimes with differing names (e.g. SSS17a, DLT17ck). For simplicity we refer to the object henceforth as EM 170817.

**Swift**-XRT and UVOT observed this field for 2 ks, starting on 2017 Aug 18 at 03:35 UT ($\Delta t = 0.6$ d), temporarily overriding the wide-area search of the GW error region. The optical source was detected in all UV filters, but no X-ray counterpart was seen (64).

The nature of EM 170817 was not clear at this time, so **Swift** resumed the wide-area search of the GW error region, but regularly interrupted it to carry out repeat observations of EM 170817. These targeted observations continued after the large-area tiling was complete (at $\Delta t = 4.3$ d), apart from the interval 22:04 on August 24 to 22:14 on August 26, at which times the source was too close to the Moon for **Swift** observations.

After the discovery of X-ray emission from EM 170817 by **Chandra** (44–46) we increased the frequency of these observations, until the field entered the **Swift** Sun constraint at 05:46 on September 2, from which it will emerge on December 02. Details of the **Swift** observations of EM 170817 are given in Table S2.

The **Nuclear Spectroscopic Telescope ARray** (**NuSTAR**) (16) also observed the location of EM 170817, taking data at four different epochs starting at 4.8 d after the GW trigger. Details of these observations are also presented in Table S2.

## 2 Data analysis

### 2.1 **Swift**-BAT

While the location of EM 170817A was not visible to **Swift** at the time of the trigger, the initial GW localisation (which covered almost the entire sky) and the edge of the **Fermi** localization error were within the BAT field of view. We therefore searched the raw 15–350 keV BAT light curves for any evidence for GRB emission within 100-s of the GW/GBM trigger. No such
emission was detected. As there was no BAT trigger, this search is limited to the raw 15–350 keV light curves (the so-called ‘rate data’ accumulated onboard), which have time bins of 64ms, 1 s and 1.6 s. Assuming a sGRB with a typical spectrum in the BAT energy range (i.e. a simple power-law model with a power-law index of $-1.32$, from the averaged sGRB power-law indices based on (65)), the 5-$\sigma$ upper limit in the 1.6-s binned light curve corresponds to a 15–350 keV flux of $< 2.3 \times 10^{-8} \text{ erg cm}^{-2} \text{ s}^{-1}$ for an on-axis energy response of the BAT, and $< 4.8 \times 10^{-7} \text{ erg cm}^{-2} \text{ s}^{-1}$ for an extremely off-axis (i.e. within the Fermi error region) energy response of the BAT.

2.2 Swift-UVOT

Swift began observations of EM 170817 with the UVOT at 03:37 on 2017 August 18 ($\Delta t = 0.6 \text{ d}$), initially with just the $uvw2$, $uvm2$, $uvw1$, and $u$ filters. Observations with the full 6-filter complement (including the $v$ and $b$ bands) began at $\Delta t = 2.5 \text{ d}$ using a pointing offset (required to avoid a nearby bright star). The source was detected in all UV filters; $b$ and $v$ band yield only marginal detections or upper limits, in part due to the considerable galaxy background. UVOT analysis was carried out using HEASOFT v6.21 (66) and UVOT CALDB version 20170130 (67).

Absolute photometry was obtained by extracting the photon counts using a circular aperture of $3''$ radius centered on EM 170817, followed by an aperture correction to the calibration standard of $5''$. The use of this smaller aperture reduced the contamination from NGC 4993 and increased the signal-to-noise ratio (68). Background counts were extracted from a region located well outside of the host galaxy. Count rates (including coincidence loss correction) were obtained using the HEASOFT tool UVOTSOURCE.

To determine the host galaxy contribution to the source, we measured the background-subtracted count rate in the source region in the interval $\Delta t = 9.5–16 \text{ d}$, by which time the
source had faded below detectability in all filters. A few images were removed due to an unstable spacecraft attitude. The host count rate was then subtracted from the earlier measurements of EM 170817; this was performed for each filter and exposure individually. The host-subtracted count rates were then converted to AB magnitudes using zero points from (69).

Examination of the spectral energy distribution (SED) derived from the photometry showed that the flux drops off steeply going further into the UV. As a result the effective wavelength of the UV bands is shifted to the red. We calculated the effective wavelengths of the UVOT bands by iteratively computing the weighted average of the wavelength multiplied by the effective area multiplied by the approximate flux (as derived from the SED), normalized by the effective area and flux integral. The results are shown in Table S1.

Finally, light curves were corrected for Galactic extinction, \( E(B-V) = 0.106 \) (70), derived using a region of size 2°. The final, background-subtracted light curves are shown in Figure 2A.

### 2.3 Swift-XRT

The data were analysed using HEASOFT v6.21 (66) and CALDB version 20170501 (67). The XRT data were grouped into “analysis blocks” of overlapping fields, to a maximum size of \( \sim 50' \) in diameter. Source detection was carried out on each of these blocks using the iterative method developed by (71). Any sources found were ranked according to their likelihood of being a transient related to the GW trigger. The method and ranking is described in detail by (72). No sources above rank 3 (i.e. uncataloged, but below catalog limits) were found in our search of the GW error region.

We focus only on the data at the location of EM 170817. No source was found in the first 2-ks observation of this source (64), with a 0.3–10 keV 3-\( \sigma \) upper limit of \( 5.5 \times 10^{-3} \) ct s\(^{-1} \). Assuming an energy conversion of \( 4 \times 10^{-11} \) erg cm s\(^{-1} \) ct\(^{-1} \), which is typical for GRB afterglows (73), this corresponds to a flux of \( 2.2 \times 10^{-13} \) erg cm\(^{-2} \) s\(^{-1} \). Given the distance to
NGC 4993 of $\sim$40 Mpc this equates to $L_X < 4.2 \times 10^{40}$ erg s$^{-1}$.

No X-ray emission was detected near the position of EM 170817 in the next four observations, when a standard GRB afterglow would be at its brightest (see, Figure 4 and (48), their figures 12–13). However, we did detect a source $\sim$ 2 days after the trigger, once we had accumulated 17.5 ks of exposure time with XRT (74). The source was poorly localized due to its faintness, but lay 11.3$''$ from the optical counterpart position, and 8.1$''$ from NGC 4993, with a position uncertainty of 6.4$''$ (90% confidence). Summing all of the X-ray data on this source, and correcting the astrometry by aligning the X-ray sources in the field with objects cataloged in 2MASS as described by (71) (and using only XRT sources detected with flags ‘Good’ or ‘Reasonable’, i.e. unlikely to be spurious detections.), we obtain a final position of RA, Dec (J2000)= 13$^h$09$^m$47.65$^s$, -23$^\circ$23'01.6", with an uncertainty of 3.9$''$ (radius, 90% confidence). This is 10.2$''$ from the position of EM 170817, and 0.9$''$ from the position of NGC 4993, suggesting that the emission detected is from the host galaxy, rather than the electromagnetic counterpart. This interpretation was confirmed by Chandra observations (45, 46, 75, 76). Taken with the detection of radio emission from the galaxy (9, 77–79) this may indicate an active galactic nucleus in NGC 4993; however the X-ray emission appears somewhat diffuse (75) and is thus likely made up of multiple X-ray emitters in that galaxy.

Since this emission may contaminate the location of EM 170817, the predicted count-rate from the host galaxy (assuming point-like emission) was taken into account when calculating upper limits. Summing all of the XRT data (172 ks), we calculate a 3-$\sigma$ upper limit at the location of EM 170817 of $2.8 \times 10^{-4}$ ct s$^{-1}$. Using the same assumptions as above, this corresponds to $L_X < 2.1 \times 10^{39}$ erg s$^{-1}$ (0.3–10 keV). The count rates or upper limit for the host NGC 4993 and EM 170817 for each individual observation are given in Table S2.
2.4 **NuSTAR**

The *NuSTAR* data were reduced using NUSTARDAS version 06Dec16.01.7.1 within HEASOFT (66), and CALDB version 20170614 (67). We used all standard settings in NUPPIPELINE, but included filtering modes for the South Atlantic Anomaly (SAACALC=2, SAAMODE=OPTIMIZED, TENTACLE=YES), to reduce the solar low energy background. For extracting count rates, we used an optimized extraction region of radius 20″, which allows for a possible 1-σ positional offset of 8″ in the *NuSTAR* absolute astrometry (16), and ensures that the *NuSTAR* full-width half-maximum point-spread function (PSF), which has a diameter of 20″, is included in the extraction region. The region was centered on the location of EM 170817 (14). Spectra and responses were obtained for each module and epoch using NUPRODUCTS with default settings.

We combined the data from the two *NuSTAR* focal plane modules (FPM): ‘FPMA’ and ‘FPMB’, and for the first epoch we obtained count rates of $7.8 \times 10^{-4}$ ct s$^{-1}$ (3–10 keV), $1.2 \times 10^{-3}$ ct s$^{-1}$ (10–39 keV) and $7.5 \times 10^{-4}$ ct s$^{-1}$ (39–79 keV). We compared these count rates with predicted background rates obtained from NUSKYBKG, which are based on a model fitted to the background of the entire module (80), and found that they agreed at the 1-σ level, indicating that the neither EM 170817 nor NGC 4993 were detected. We calculated the 3-σ upper limits on emission from EM 170817 by assuming a power-law spectrum with a photon index of $-2$ and absorption Hydrogen column density $3 \times 10^{21}$ cm$^{-2}$, as is typical for GRBs (73), and assuming a distance of 40 Mpc. These are given in Table S2. We analyzed the other epochs in the same way; the source was likewise undetected, with the limits given in Table S2. Combining the data from all epochs yields a limiting luminosity of $L = 2.8 \times 10^{39}$ erg s$^{-1}$ (3–10 keV), $1.2 \times 10^{40}$ erg s$^{-1}$ (10–39 keV), and $1.7 \times 10^{41}$ erg s$^{-1}$ (39–79 keV). The non-detection of NGC 4993 in the > 3 keV band pass of *NuSTAR* suggests a relatively soft spectrum, given the detection by *Swift*-XRT. Indeed, we find no evidence for an elevated count-rate at the location of NGC 4993. Thus, while we cannot rule out that our *NuSTAR* limits on the flux from
EM 170817 are slightly underestimated due to contamination by host-galaxy emission, any such contamination is small and will have little effect.

3 Interpretation and discussion

3.1 Gamma-Ray Emission

If the GRB 170817 was associated with GW 170817, then it is the closest sGRB with known distance: regardless of the connection with the optical counterpart, the prompt GW analysis reported a sky-averaged distance estimate of $39 \pm 17$ Mpc. Assuming a distance of 40 Mpc (i.e. to NGC 4993; (61)), the Fermi-GBM detection corresponds to an isotropic-equivalent 1 keV–10 MeV luminosity $L_{\text{iso}} = 1.6(\pm0.6) \times 10^{47}$ erg s$^{-1}$ (23). This is three orders of magnitude below the lowest-energy sGRB in the sample reported by (13), consistent with the predictions of (48). That work compared the observed flux of sGRBs detected by Swift, Fermi and CGRO without known redshift with that which would have been seen from the Swift sGRBs of known redshift had they been in the range of advanced LIGO/Virgo. They predicted that a binary neutron star merger detected by advanced LIGO/Virgo would be several orders of magnitude less luminous than GRBs detected previously. Given this low luminosity, GW 170817 allows us to probe the parameter space not covered by the sGRBs detected hitherto. While the GBM spectral properties of the object (23, 81) are fairly typical of sGRBs, it should be borne in mind that this event is so underluminous compared to the previous sample of GRBs, that models constructed for those objects may not be simply transferrable to such a low-luminosity event as this. Indeed (see Section 3.3.2), the low-luminosity is difficult to explain as arising from an ultra-relativistic jet at all. The implications of the sub-luminous gamma-ray emission are explored in further detail in (10).
3.2 UV/Optical Emission

3.2.1 SED Construction and Fitting

Since our UVOT detections extend out to only $\Delta t = 1.0 \, \text{d}$, we focus on the early-time behavior of the optical counterpart in this work. We constructed the spectral energy distribution (SED) of EM 170817 using data from our work and the GCN Circulars. We limit the dataset considered here to the $izy$ photometry provided by the Pan-STARRS project (53, 82). At the time of the preparation of this manuscript, this was one of the few datasets with host-subtracted photometry (like the UVOT), and the transient field was previously calibrated in these filters. We linearly interpolated the Pan-STARRS 1 observations to match the two epochs of UVOT photometry.

After correcting for Galactic extinction, we fit the resulting SED to a blackbody function. Adding a free parameter for host extinction did not improve the fit quality. At $\Delta t = 0.6 \, \text{d}$, we find $T_{\text{BB}} = 7300 \pm 200 \, \text{K}$, corresponding to $R_{\text{BB}} = 6.0 \times 10^{14} \, \text{cm}$ and $L_{\text{BB}} = 7 \times 10^{41} \, \text{erg s}^{-1}$. The blackbody model yields a fit statistic of $\chi_r^2 = 2.8$ (5 degrees of freedom). At $\Delta t = 1.0 \, \text{d}$, we derive $T_{\text{BB}} = 6400 \pm 200 \, \text{K}$, corresponding to $R_{\text{BB}} = 6.5 \times 10^{14} \, \text{cm}$ and $L_{\text{BB}} = 5 \times 10^{41} \, \text{erg s}^{-1}$. At this epoch, we derive a goodness of fit of $\chi_r^2 = 1.6$ (4 degrees of freedom). The resulting SED fits are plotted in Figure 2B–C.

For comparison, we also fit the SED at both of these epochs to a power-law model (as would be expected from synchrotron afterglow emission) of the form $f_\lambda \propto \lambda^{-\alpha}$. At $\Delta t = 0.6 \, \text{d}$, we find $\alpha = 1.0 \pm 0.4$. The fit quality is extremely poor: $\chi_r^2 = 102.8$ (5 degrees of freedom). Repeating this analysis at $\Delta t = 1.0 \, \text{d}$, we find $\alpha = 1.0 \pm 0.7$ and $\chi_r^2 = 93.5$ (4 degrees of freedom). Clearly the power-law model provides a much worse fit to the early SED.

3.2.2 Model Comparisons to the Early UV/Optical Light Curve

Our kilonova calculations to model the early ($\Delta t \leq 2 \, \text{d}$) UV/optical emission from EM 170817 are based on the approach of Wollaeger et al. (42). We employ the multigroup, multidimensional
radiative Monte Carlo code SUPERNu (83–85) with the set of opacities produced by the Los Alamos suite of atomic physics codes (86–88).

We studied a range of 1-dimensional and 2-dimensional models, systematically exploring different uncertainties affecting the kilonova light curves: ejecta mass, velocity, composition, and morphology, as well as the model for the energy deposition in post-nucleosynthetic radioactive decays. Our 2-dimensional models are based on the dynamical ejecta morphologies from (43), computed by following long-term evolutions of the ejecta from neutron star mergers, which in turn were simulated in (89). The r-process nucleosynthesis and radioactive heating are computed using the nuclear network code WINNET (90–92). Reaction rates for nucleosynthesis are taken from the compilation of (93) for the finite range droplet model (FRDM; 94), including density-dependent weak reaction rates (95) and fission (96, 97). Coordinate- and time-dependent thermalization of nuclear energy is computed using empirical fits developed in (98) and (99).

The abundances of trans-lead elements depend sensitively on the nuclear mass model and, since they dominate the production of α-particles, the production of these elements can dramatically alter the heating at late times (98, 99). On the other hand, the nuclear heating rate on the timescale up to one to two days is dominated by β-decay, which shows consistency across different nuclear mass models (98). Because the modeling here is focused on this early-time data, we opt to utilize the FRDM nuclear mass model and do not study the uncertainties of the nuclear heating here.

Instead, we focus on the properties of the ejecta, which consist of two components: a dynamical ejecta component and a wind ejecta component. The eight most sophisticated models in (42) have labels of the form γA1, γA2, γB1, γB2 … γD1, γD2, where the letter A–D represents the tidal morphology, are taken from (43), see also (100), their figure 1 for an illustration; the subscripted number represents the wind type (below).
The dynamical ejecta are in an axisymmetric, very neutron-rich outflow \((Y_e \approx 0.04)\) composed almost entirely of the main \(r\)-process elements with significant mass fraction of lanthanides and actinides \((\approx 14\%)\), taken from \((43)\). The wind ejecta are in a spherically-symmetric, moderately neutron-rich outflow with two different degrees of neutron richness: i) (subscript 1 in the model descriptor) \(Y_e = 0.37\) with a composition dominated by iron peak and slightly-above iron peak isotopes, and ii) (subscript 2 in the model descriptor) \(Y_e = 0.27\) with a composition dominated by elements at or just above the first \(r\)-process peak. Astrophysical and thermodynamical conditions to generate the wind compositions were sampled from typical outcomes of the high- and low-latitude wind trajectories, simulated in \((101)\) and \((36)\).

The \(\gamma\)-models in \((42)\) take into account the distinct composition of the two outflow components, which in turn determine the opacities, ionization conditions, radioactive nuclear heating and thermalization. The “\(\gamma\)” notation indicates that the gray \(\gamma\)-ray transport is also included to accurately account for the \(\gamma\)-ray energy escape or deposition.

To fit the observed data, we allowed the mass and velocity to vary for both of our two components. We also considered two different compositions for the wind ejecta based on electron fractions of 0.27 and 0.37. The SEDs at times of 0.6 d and 1.0 d after the launch of the explosion are shown in Figure S2. This figure plots the flux as a function of wavelength for a variety of viewing angles. We focused on the dynamical ejecta configuration A, because only in this morphology the polar region is free from the “lanthanide curtain” created by the dynamical ejecta, and the wind is completely unobscured when viewed on-axis.

Figure S2 demonstrates that in order to fit the observed emission, we must raise the wind ejecta mass to at least \(0.03 \, M_\odot\) (see panels C–D). With a wind ejecta mass of \(0.1 \, M_\odot\), a fit to the data can be made with a viewing angle of \(40^\circ\). Although this is at the high end of the predicted wind ejecta masses, it is not impossible to produce such high-mass accretion disk outflow with e.g. an asymmetric merger \((37, 41, 102)\). It is also possible that we are observing more irregular
wind morphologies which will in any case produce brighter transients (by a factor of a few) than our spherical models.

Figure S3 compares the SEDs for wind ejecta mass $m_{\text{wind}} = 0.03 M_\odot$ between the two different wind compositions that we explore. The high electron-fraction wind ejecta model is unable to match the data without increasing the ejecta mass further, favoring of the neutron-rich wind ejecta, which indeed is close to the values predicted for the polar regions (35, 36, 102). The spectrum for the model with high electron-fraction exhibits much faster reddening between the epochs than its medium electron-fraction counterpart.

We emphasize that the “effective” opacity of our wind model 2 ($Y_e = 0.27$; first $r$-process peak elements) is actually somewhat lower than that of wind model 1 ($Y_e = 0.37$; Fe-peak elements with effective $\kappa \approx 1$ cm$^2$ g$^{-1}$). This is a result of the less complex atomic configurations (42). As a result, the models with the Fe-peak composition for the disk wind become redder much more quickly and are therefore difficult to match with the bright UV emission seen at early times.

The neutron rich dynamical ejecta does not directly contribute significantly to the ultraviolet and optical emission in the first two days. In fact, the spectra and light curves in optical and UV are completely insensitive to the dynamical ejecta mass, and only weakly sensitive to its velocity in the red part of the optical spectrum. Figure S4 represents our best models in which we have increased the dynamical ejecta velocity to produce the early red excess in the three points around 7000–8000 Å (note that velocities here refer to the average speed of the ejecta; the fastest-moving material has a velocity of twice this value). Faster dynamical ejecta ignite the infrared transient earlier, and as a result causes moderate enhancement in the red part of the spectra. We can see that velocity $v_{\text{dyn}} = 0.2 - 0.3 \, c$ is compatible with observations. Faster wind, however, causes the blue transient to expire sooner and shift the spectrum redwards earlier, which apparently excludes wind velocities as high as $v_{\text{dyn}} = 0.16 \, c$. 

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Although the direct contribution of the dynamical ejecta is small, the lanthanides and actinides in its composition can obscure any emission and, for our morphology of the tidal ejecta, the emission in the UV and optical is too obscured for viewing angles more than 40° off axis. If the viewing angle is greater than this value, the neutron-rich dynamical ejecta must be more constrained to the orbital plane. In another scenario an asymmetric merger may end up with ‘kidney-shaped’ dynamical ejecta, which only cover a finite sector in azimuthal angle (as in (100), their figure 1). If the ‘kidney’ is facing away from us, it will not obscure the brighter and bluer view of the wind. In either case, it seems that the favored scenario is the one in which the dynamical ejecta do not strongly obscure the wind.

We conclude with the following optimal range of parameters:

- wind composition: first peak elements;
- wind mass: \( m_{\text{wind}} = 0.03 - 0.1 \, M_\odot \);
- wind velocity: \( v_{\text{wind}} = 0.08 \, c \);
- wind kinetic energy: \( E_{\text{wind}} = 2 \times 10^{50} \, \text{erg} \);
- dynamical ejecta mass: poorly constrained, compatible with the range from \( m_{\text{dyn}} = 0.002 - 0.03 \, M_\odot \);
- dynamical ejecta velocity: \( v_{\text{dyn}} = 0.2 - 0.3 \, c \);
- dynamical ejecta kinetic energy: \( E_{\text{dyn}} \sim 6 \times 10^{50} \, \text{erg} \);
- viewing angle: \(< \sim 40^\circ\), degenerate with the wind outflow mass: higher polar angle implies higher mass, or non-axisymmetric configuration without dynamical ejecta obscuring the wind.
Figure S5 presents the light curves in the UV bands for the variations of our baseline model $\gamma A_2$.

In addition to the early data from Swift-UVOT, the Very Large Telescope (VLT) in Chile observed EM 170817 with the U filter (similar to the $u$ filter on UVOT) on the Visible wide field Imager and Multi-Object Spectrograph (VIMOS) instrument at $\Delta t = 4.47$ d. Observations were processed via the esorex software and the photometric zero-point was established using the UVOT data for field stars, transformed to the VLT U-band using the relations in (68). The (extinction corrected) magnitude at this time was $23.35 \pm 0.07$ (AB). Similar results were found by the Hubble Space Telescope (10). These points are not explained by our model, which predicts no emission at these wavelengths at late times (Figure S5D). However, since our focus is on early-time emission certain assumptions were made above (such as neglecting uncertainties on nuclear heating). Thus our which focusses only on the early-time data and reproduced this well, does not give a complete picture of the emission processes at late times.

3.3 X-ray Emission

3.3.1 Afterglow emission

To place our X-ray limits at the location of EM 170817 in context, we simulated a number of fake sGRB afterglow light curves at a distance of 40 Mpc. To do this we took the sample of sGRBs with known redshift from (13), for each of these the best-fitting power-law models from the online XRT GRB Catalog (http://www.swift.ac.uk/xrt_live_cat) (73). Occasionally we manually added extra power-law segments to more accurately reproduce the light curves. This was necessary because for the catalogue extra segments are only added when they are justified at the 4-$\sigma$ level, whereas for our purposes we simply want a model that accurately describes the data at the time of our observations. We then scaled the normalisation to correct from the observed redshift to 40 Mpc.
In Figure 4 we plot the resultant median light curve and the 25th and 75th percentiles; also shown are the *Swift* and *NuSTAR* upper limits. While the sample of sGRBs used is clearly biased, both in sample definition (it is intended to be complete in terms of BAT sensitivity) and our extra selections (objects with no XRT detection cannot be included), this nonetheless indicates that the X-ray afterglow of GRB 170817A is significantly fainter than the majority of previously-observed sGRBs.

This result is not entirely surprising, since the gamma-ray output was orders of magnitude less luminous than the current sample, and this is known to correlate with the X-ray afterglow luminosity. ([13]) presented correlations between the prompt isotropic energy output and luminosity ($E_{\gamma,iso}$, $L_{\gamma,iso}$) and the X-ray flux at 5 hours. Using the *Fermi*-GBM values of $E_{\gamma,iso} = 4.0 \times 10^{46}$ erg and $L_{\gamma,iso} = 1.6 \times 10^{47}$ erg s$^{-1}$ ([23]), and assuming a power-law decay of $L \propto t^{-1.5}$ (based on the median light curve in Figure 4), these correlations predict an X-ray luminosity at the time of our first *Swift*-XRT observation of $4.5 \times 10^{40}$ erg s$^{-1}$ (from $E_{\gamma,iso}$) or $2.1 \times 10^{41}$ erg s$^{-1}$ (from $L_{\gamma,iso}$). Our upper limit at this time was $2.2 \times 10^{40}$ erg s$^{-1}$. Thus, if the X-ray afterglow of GRB 170817A followed the ([13]) correlations exactly, we would have expected to detect it (though there is some uncertainty in the correlations).

While X-rays were detected by *Chandra* from EM 170817 at around 9 days after the trigger (no flux was reported for this proprietary observations) ([44–46, 76]), extrapolating the detected X-ray flux back to the times of the *Swift* and *NuSTAR* observations using the $t^{-1.5}$ power-law yields a flux above our upper limits ([44]). Thus our early non-detections also disfavor an on-axis afterglow as the source of the late-time X-ray emission detected by *Chandra*.

In addition to a comparison with the observed sample of sGRBs, we can translate these limits on X-ray emission from EM 170817 to physical constraints on any associated relativistic ejecta using the standard synchrotron afterglow formulation ([103]). First, we consider on-axis geometries (i.e. where the relativistic jet is pointed towards Earth).
The X-ray flux will depend on the location of the cooling frequency $\nu_c$. At the time of our first Swift-XRT observation ($\Delta t = 0.6$ d), the cooling frequency will be (104):

$$
\nu_c = 4 \times 10^{18} \left( \frac{n_0}{5 \times 10^{-3} \text{ cm}^{-3}} \right)^{-1} \left( \frac{E_{AG}}{2 \times 10^{51} \text{ erg}} \right)^{-1/2} \text{ Hz},
$$

where we have normalized the isotropic afterglow energy ($E_{AG}$) and circumburst density ($n_o$) to the median values for sGRBs derived in (21). Thus, the XRT bandpass ($0.3$–$10.0$ keV = $7.3 \times 10^{16}$–$2.4 \times 10^{18}$ Hz) will typically fall near (but slightly below) the cooling frequency at this time, while $\nu_c$ will often fall within the (broader) NuSTAR bandpass. (Here and throughout, we assume a constant-density circumburst medium, with microphysical parameters: $\epsilon_e = 0.1$ and $\epsilon_B = 0.01$, and an electron spectral index of $p = 2.5$.)

For $\nu_X < \nu_c$, our initial NuSTAR upper limit implies (104):

$$
\left( \frac{n_0}{5 \times 10^{-3} \text{ cm}^{-3}} \right)^{1/2} \left( \frac{E_{AG}}{2 \times 10^{51} \text{ erg}} \right)^{11/8} < 2 \times 10^{-4}.
$$

As a result, for any reasonable circumburst density ($n_0 > \sim 10^{-5}$ cm$^{-3}$; (21)), we place a limit on the amount of energy coupled to relativistic ejecta along our line of sight of $E_{AG} < \sim 10^{50}$ erg. Adopting $\nu_X > \nu_c$ yields similar constraints (though the result is independent of $n_0$). Utilizing the first Swift-XRT limit results in a comparable limit on the energy coupled to relativistic ejecta (for on-axis geometries).

To verify these results, we ran a series of simulations using the afterglow light curve code BOXFIT (20). Over the observed range of circumburst densities and afterglow energies for sGRBs (21), we calculated the predicted X-ray flux at the time of our initial NuSTAR observation ($\Delta t = 0.7$ d). The results are shown in the bottom panel of Figure 4. Similar to our analytic results above, we can rule out the presence of a collimated, relativistic outflow initially beamed towards Earth with $E_{AG} > \sim 10^{50}$ erg.
As discussed by (10), the prompt low γ-ray luminosity of GRB 170817A is not consistent with emission from a standard ultra-relativistic jet as seen in sGRBs to date, regardless of viewing angle. Instead it must arise from some previously unseen process, such as a heated cocoon (e.g. (24, 25)); and we are likely viewing the event off-axis (unless there was no jet at all), as indeed our modelling of the UV emission (above) suggests. In this case, if there is an ultrarelativistic jet which is oriented away from us, we can expect X-ray emission at late times. As the jet decelerates and spreads laterally, it illuminates an increasing fraction of the sky, thus eventually spreading into our line of sight. Such “orphan” afterglow emission is expected to be much fainter than the early afterglow, but due to the low distance to GW 170817 may be detectable by Swift.

We therefore repeated the afterglow light curve simulations described above for a variety of off-axis viewing geometries. For the median values of short-duration GRB afterglow energy ($E_{\text{AG}} = 2 \times 10^{51}$ erg) and circumburst density ($n_0 = 5 \times 10^{-3}$ cm$^{-3}$; (21)), the resulting light curves are plotted in Figure 5. With the XRT and NuSTAR non-detections to date, we can rule out any viewing angle with $\theta_{\text{obs}} < \sim 20^\circ$. With these afterglow parameters, the reported Chandra flux of $5 \times 10^{-15}$ erg cm$^{-2}$ s$^{-1}$ (105, 106) can be reproduced with a viewing angle of $\theta_{\text{obs}} \approx 30^\circ$. These constraints are consistent with the geometry implied by our modelling of the blue kilonova seen by the UVOT. It is also broadly consistent with the evolution of EM 170817 at radio wavelengths (9), showing no emission at early times but then later detections (107, 108).

### 3.3.2 Emission from a heated cocoon

One way of explaining the observed prompt gamma-rays may be via the ‘cocoon model’ (24, 25). In this model, if the GRB jet propagates through a baryon-contaminated site around the merger, a heated cocoon is formed from which prompt gamma-rays may be seen. Under the (25) model, the prompt luminosity of GRB 170817A can be interpreted as arising from a typical (or
rather, Swift-like) sGRB viewed $\sim 10–20^\circ$ off-axis. In their revised model, in which the cocoon is not isotropic \cite{109}, the prompt gamma-ray luminosity implies an observing angle of $\sim 40–50^\circ$ compared to the jet angle. However, the peak photon energy reported by \cite{81} (124.2 $\pm$ 52.6 keV) requires a viewing angle of $\sim 10^\circ$ according to the same model, which suggests that this model, in its current form at least, cannot explain the observed prompt emission.

Nonetheless, a heated cocoon may in principle give rise to late-time X-ray emission, as the ‘trans-relativistic’ expanding cocoon (i.e. with Lorentz factor $\Gamma \sim 2–3$) is eventually decelerated by the ambient medium which then shocks and radiates in a manner analogous to the standard on-axis GRB afterglow \cite{24}. In the main text we demonstrated that our upper limits and the Chandra detections are consistent with most of these models, which predict a peak X-ray flux several weeks to months after the merger.

### 3.3.3 Non-afterglow emission

In addition to the emission associated directly or indirectly with the jet, as above, some authors have predicted quasi-isotropic X-ray emission following a binary neutron star merger. \cite{110}, responding to an apparent X-ray excess in the light curve of GRB 130603B \cite{111} coincident with a claimed kilonova detection \cite{51, 112}, proposed that the kilonova was powered not by $r$-process nucleosynthesis but by isotropic X-ray emission, perhaps from fallback accretion via a disk. They predict typical X-ray luminosities of order $10^{41}$ erg s$^{-1}$ appearing at around 7 days after the merger. We find no evidence for such emission in our data, with upper limits of $\sim 1.5 \times 10^{40}$ erg s$^{-1}$ at 7 days after the trigger.

Chandra detected X-ray emission from EM 170817 at $\sim 9$ d after the trigger \cite{44} and in subsequent observations \cite{45, 46, 76, 105, 106, 113} with a luminosity of $\sim 9 \times 10^{38}$ erg s$^{-1}$ (0.3–8 keV), consistent with our XRT and NuSTAR upper limits. This could be indicative of the emission predicted by \cite{110}, if the accretion rate is much lower than in their fiducial cal-
culations. Alternatively, those authors note that if the ejecta initially prevent the detection of the X-rays (i.e. the X-rays are all thermalized by the optically thick ejecta) then the ejecta will still become optically thin to X-rays at tens of days, with luminosities $\sim 10^{40} - 10^{41}$ erg s$^{-1}$, still well above both the XRT/NuSTAR detection thresholds and the detected Chandra flux, thus testable with future observations.

It has also been considered that the result of the BNS merger may be a hyper-massive neutron-star which is able to support itself at least temporarily against gravitational collapse, e.g. by rotation, and which emits X-rays as it rotates. For example (19, 114–116) succesfully modelled early-time X-ray light curve plateaux as being powered by the spin-down of a magnetar. However such plateaux do not extend to the times at which our observations of EM 170817 began, and in some cases imply the collapse of the magnetar within a few ks of the merger, well before our X-ray observations of EM 170817 began.

(117) proposed three mechanisms for producing X-rays after a binary neutron star merger: the standard afterglow emission, spin-down energy extracted from a magnetar, or a ‘merger-nova’ in which the kilonova-emitting ejecta are heated also by the magnetar and emit X-rays (this latter case was also proposed by (118) in whose model this emission is extremely bright, however (117) attribute this to an error and predict dramatically less X-ray flux). In this model, the observed X-ray emission depends upon the viewing angle; with the observer either looking down the GRB jet, missing the jet but having an unobstructed view of the magnetar, or missing the jet, and viewing the magnetar through the ejecta. Based on these assumptions, (117) produced a set of example light curves for typical sGRB parameters. Almost all of their light curves predict luminosities 4–5 orders of magnitude above the X-ray detection limits presented here, thus even if the available energy in the magnetar is substatially reduced compared to typical, we would have detected this emission.

However, the X-ray flux is greatly reduced in the case that the magnetar collpases to a black
hole before our observations (consistent with the plateau durations of \((19, 114–116)\), or the collapse timescale of the magnetar is shorter than its spindown timescale. In light of these considerations, the lack of X-ray emission above \(\sim 1.5 \times 10^{40} \text{ erg s}^{-1}\) strongly suggests that if a magnetar was formed after the BNS merger, it collapsed to a black hole within the first 0.6 d of its formation.

### 3.3.4 Limits on nuclear line emission

Radioactive \(r\)-process nucleii synthesized in the merger ejecta produce gamma-ray lines through \(\beta\) decay. The emission lines have a broad spectrum from 10 keV to 1 MeV \((98, 119)\). These lines start to emerge from the ejecta when the ejecta become optically thin. We model the flux of the nuclear emission lines for the ejecta mass of 0.03 \(M_\odot\) and 0.05 \(M_\odot\) with the mean expansion velocity of \(0.2c\) \((119)\), taking photoelectric absorption and the Compton scattering into account. Here we assume a composition of nuclei that matches the solar \(r\)-process abundance pattern of stable and long-lived nuclei for \(90 \leq A \leq 140\), where \(A\) is the atomic mass number. There are three strong lines around 28, 50, and 81 keV, which are produced by \(^{129}\text{Te}\), \(^{132}\text{Te}\), and \(^{133}\text{Xe}\). However, the line nuclear emission is too weak to be detected by \(\text{NuSTAR}\), as shown in Figure S6.

### 3.4 Constraints on X-ray emission from the wide-angle search.

Our discussion so far has assumed that EM 170817, GRB 170817A and GW 170817 are all the same event. While this seems reasonable on probabilistic grounds (e.g. \((8, 120)\)) it is not certain. For this reason, even after the detection of EM 170817, \(\text{Swift}\) continued to search the GW error region for possible counterparts. In total the tiling observations continued until 4.3 d after the GW event, and XRT covered 52\% of the probability in the three-detector ‘preliminary-LALInference’ skymap \((52)\), and 92\% of the probability determined by convolving this skymap.
with GWGC, as described in (12). Therefore, if EM 170817 is not the counterpart to GW 170817 it is probable that Swift-XRT observed the true location of the event.

In Figure 4 we show a sample of sGRB X-ray light curves as they would have appeared at 40 Mpc (the best-fitting luminosity distance in the 3D GW skymap). The blue line on this plot shows the time-range spanned by our search of the GW error region and the detection limit of the observations. This shows that a sGRB afterglow typical of those observed to date would have been detectable with XRT during our follow up. As noted above, some known sGRBs have extremely fast-fading afterglows (19); such objects fall below the 25th percentile of GRB flux by the times of our observations so cannot be seen in Figure 4 and would not have been detected in our search.

To investigate in more detail, we simulated 10,000 sGRBs. Their position and distance were drawn at random from the 3-D ‘bayestar-HLV’ skymap (11) (i.e. that on which our tiles were based), and their light curves were produced by taking a random sGRB from the sample described in Section 3.3.1 and scaling it to the distance of the simulated event. We then identified for each of these simulated events whether its location was observed by XRT, and if so, whether the source would have been above our detection limit at the time of the observation. We found that in 65% of cases the source would have been detected. We then repeated the process, but instead of using the actual locations observed by XRT, which included many observations of EM 170817 which displaced planned tiling of the GW-GRB probability region, we used those which would have been observed if we had not focussed on EM 170817. In these simulations 68% of the simulated sources were recovered. Since GRB 170817A was much fainter than typical Swift-detected sGRBs, we repeated these simulations again, this time reducing the predicted XRT flux by a factor of 1,000. In this case we find we would have detected 11% of simulated events, or 10% if we had not targeted EM 170817. The slight decrease when not targeting EM 170817 – i.e. when covering more of the GW error region, arises because the focus
on EM 170817 results in longer exposures and hence the ability to detect fainter sources in that field, which in this case outweighs the benefits of searching more fields at a lower sensitivity.
Figure S1: The BAT field of view at the time of the GW/GRB trigger. The black/red/white scale indicates the BAT coding fraction; the darker colour indicates lower partial coding fraction (i.e. more toward the edge of the field of view). The GW localizations probability contours are plotted in blue. The green region shows the part of the sky occulted by the Earth and the yellow region is the Fermi-GBM 90% confidence region.
Figure S2: **Model SEDs for kilonovae compared to our data.** The SEDs were created for four models based on the $\gamma A_2$ (i.e. with the wind-2 composition), for different orientations and for two different epochs. The shaded regions display the range of model SEDs for different viewing angles over the range considered here. Thick solid lines show the flux for viewing angle $\theta \approx 30^\circ$, and thin lines represent adjacent angular bins with increasing viewing angle going up to $\approx 60^\circ$. Data points are shown as circles with error bars, or triangles (for upper limits). The models span a range of wind masses: a baseline model $\gamma A_2$ (panel A) has relatively low wind mass of $0.005 \, M_\odot$. In the following panels (B–D): $m_{\text{wind}} =$ 0.015 $M_\odot$, 0.03 $M_\odot$ and 0.1 $M_\odot$. To fit the observed data points, the wind mass needs to be at least 0.03 $M_\odot$. 

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Figure S3: **The effect of wind electron fraction on the SED.** These SEDs have the same velocities and masses of the ejecta but different composition: “wind 1” (A) with abundant iron-group and the d-shell elements vs. “wind 2” (B) with the first peak elements, largely representing the s- and p-shell elements and relatively fewer d-shell elements. Notation for the plots is the same as in the previous figure. The iron-group dominated composition not only exhibits lower brightness but also shows much more reddening in the spectrum between the two epochs. Datapoints are as in Figure S2.

Figure S4: **The effect of dynamical ejecta velocity on the SED.** These SEDs have the same velocities and masses of the wind ejecta but different velocities of the dynamical ejecta: \( v_{\text{dyn}} = 0.2 \, c \) (A) vs. \( 0.3 \, c \) (B). Notation for the plots is the same as in the previous Figure. Dynamical ejecta need to be expanding with at least \( 0.2 - 0.3 \, c \) to be responsible for the red excess around 7000-9000 Å. Datapoints are as in Figure S2.
Figure S5: **Model light curves for the variants of model configuration $\gamma A_2$.** The solid lines indicate the predictions for an on-axis viewing geometry, and the shades beneath are for off-axis geometries, with limiting angles $30^\circ$ (darker shade) and $40^\circ$ (lighter shade). Datapoints are as in Figure S2.
Figure S6: **Predicted X-ray flux from atomic line emission of high atomic mass elements.**
The predicted flux falls well below the sensitivity of our *NuSTAR* data.
Table S1: *Swift-UVOT photometry of EM 170817*. The reported magnitudes have been corrected for host contamination but not for Galactic extinction. Upper limits are quoted at the 3 $\sigma$ level. Epochs lacking sufficient SED constraints (and thus lacking estimates of the effective wavelength) are indicated with “...”.

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Table S2: X-ray pointed observations of EM 170817 with *Swift*-XRT and *NuSTAR*. For XRT we give the measured luminosity or 3-σ upper limit of NGC 4993 and 3-σ limit at the location of EM 170817. For *NuSTAR* we give only the latter.

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<td>$39.80^{+0.21}_{-0.26}$</td>
<td>&lt; 40.22</td>
</tr>
<tr>
<td>07012979009</td>
<td>Aug 28 at 10:46:17 (10.92)</td>
<td>3.37</td>
<td>&lt; 40.35</td>
<td>&lt; 40.43</td>
</tr>
<tr>
<td>07012979010</td>
<td>Aug 29 at 01:04:57 (11.52)</td>
<td>11.78</td>
<td>$39.98^{+0.15}_{-0.17}$</td>
<td>&lt; 40.06</td>
</tr>
<tr>
<td>08012979011/12</td>
<td>Aug 30 at 01:00:57 (12.51)</td>
<td>15.81</td>
<td>$39.64^{+0.18}_{-0.23}$</td>
<td>&lt; 40.02</td>
</tr>
<tr>
<td>07012979015</td>
<td>Aug 31 at 02:27:52 (13.57)</td>
<td>34.88</td>
<td>$39.90^{+0.09}_{-0.10}$</td>
<td>&lt; 39.87</td>
</tr>
<tr>
<td>08012979014/17</td>
<td>Sep 01 at 05:53:04 (14.72)</td>
<td>25.62</td>
<td>$39.63^{+0.16}_{-0.19}$</td>
<td>&lt; 39.93</td>
</tr>
<tr>
<td>08012979016/19</td>
<td>Sep 02 at 08:40:56 (15.83)</td>
<td>3.76</td>
<td>&lt; 40.68</td>
<td>&lt; 40.46</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>ObsID</th>
<th>Start time (UTC)</th>
<th>Exposure (ks)</th>
<th>$\log L_{X,NGC\ 4993}$</th>
<th>$\log L_{X,EM\ 170817}$</th>
</tr>
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<tbody>
<tr>
<td>90361003002</td>
<td>Aug 18 at 05:25 (0.7)</td>
<td>23.9</td>
<td>—</td>
<td>&lt; 39.7</td>
</tr>
<tr>
<td>90361003004</td>
<td>Aug 21 at 20:45 (4.3)</td>
<td>40.3</td>
<td>—</td>
<td>&lt; 39.6</td>
</tr>
<tr>
<td>90361003005</td>
<td>Sep 04 at 17:56 (18.2)</td>
<td>12.6</td>
<td>—</td>
<td>&lt; 40.1</td>
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<tr>
<td>90361003006</td>
<td>Sep 05 at 14:51 (19.1)</td>
<td>19.3</td>
<td>—</td>
<td>&lt; 39.9</td>
</tr>
<tr>
<td>90361003007</td>
<td>Sep 06 at 17:56 (20.1)</td>
<td>23.1</td>
<td>—</td>
<td>&lt; 39.8</td>
</tr>
<tr>
<td>90361003009</td>
<td>Sep 21 at 11:10 (34.9)</td>
<td>54.1</td>
<td>—</td>
<td>&lt; 39.5</td>
</tr>
</tbody>
</table>

* Observation year is 2017, times in parentheses are days since GW 170817. *b XRT luminosities are 0.3–10 keV, *NuSTAR* are 3–10 keV.
Table S3: *Swift* observations of the error region of GW 170817A.

<table>
<thead>
<tr>
<th>Pointing direction (J2000)</th>
<th>Start time(^a) (UTC)</th>
<th>Exposure (s)</th>
</tr>
</thead>
<tbody>
<tr>
<td>11(^h)47(^m)9.20(^s), -39 deg 50'40.2''</td>
<td>Aug 17 at 13:37:07 (0.039)</td>
<td>117</td>
</tr>
<tr>
<td>11(^h)48(^m)57.01(^s), -39 deg 50'27.9''</td>
<td>Aug 17 at 13:38:13 (0.040)</td>
<td>80</td>
</tr>
<tr>
<td>11(^h)48(^m)1.57(^s), -40 deg 08'14.1''</td>
<td>Aug 17 at 13:39:19 (0.040)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)46(^m)15.97(^s), -40 deg 08'01.1''</td>
<td>Aug 17 at 13:40:24 (0.041)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)45(^m)23.31(^s), -39 deg 50'03.5''</td>
<td>Aug 17 at 13:41:30 (0.042)</td>
<td>125</td>
</tr>
<tr>
<td>11(^h)46(^m)18.37(^s), -39 deg 32'28.2''</td>
<td>Aug 17 at 13:42:35 (0.043)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)48(^m)4.49(^s), -39 deg 32'53.9''</td>
<td>Aug 17 at 13:43:41 (0.043)</td>
<td>122</td>
</tr>
<tr>
<td>11(^h)49(^m)49.51(^s), -39 deg 32'55.3''</td>
<td>Aug 17 at 13:44:47 (0.044)</td>
<td>82</td>
</tr>
<tr>
<td>11(^h)50(^m)41.69(^s), -39 deg 50'40.1''</td>
<td>Aug 17 at 13:45:52 (0.045)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)49(^m)47.40(^s), -40 deg 08'18.8''</td>
<td>Aug 17 at 13:46:58 (0.046)</td>
<td>122</td>
</tr>
<tr>
<td>11(^h)48(^m)53.74(^s), -40 deg 25'48.7''</td>
<td>Aug 17 at 13:48:03 (0.047)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)47(^m)8.17(^s), -40 deg 25'28.4''</td>
<td>Aug 17 at 13:49:09 (0.047)</td>
<td>82</td>
</tr>
<tr>
<td>11(^h)45(^m)22.57(^s), -40 deg 25'31.7''</td>
<td>Aug 17 at 13:50:15 (0.048)</td>
<td>135</td>
</tr>
<tr>
<td>11(^h)44(^m)31.01(^s), -39 deg 07'37.1''</td>
<td>Aug 17 at 14:47:11 (0.088)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)43(^m)39.46(^s), -39 deg 49'52.3''</td>
<td>Aug 17 at 14:48:23 (0.088)</td>
<td>92</td>
</tr>
<tr>
<td>11(^h)44(^m)32.85(^s), -39 deg 32'32.2''</td>
<td>Aug 17 at 14:49:34 (0.089)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)45(^m)25.60(^s), -39 deg 14'57.5''</td>
<td>Aug 17 at 14:50:44 (0.090)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)47(^m)12.13(^s), -39 deg 15'21.6''</td>
<td>Aug 17 at 14:51:55 (0.091)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)48(^m)56.02(^s), -39 deg 15'25.5''</td>
<td>Aug 17 at 14:53:06 (0.092)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)50(^m)41.63(^s), -39 deg 15'19.1''</td>
<td>Aug 17 at 14:54:17 (0.093)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)51(^m)34.27(^s), -39 deg 33'09.8''</td>
<td>Aug 17 at 14:55:27 (0.093)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)52(^m)26.69(^s), -39 deg 50'48.2''</td>
<td>Aug 17 at 14:56:38 (0.094)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)51(^m)31.09(^s), -40 deg 08'17.0''</td>
<td>Aug 17 at 14:57:49 (0.095)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)50(^m)39.51(^s), -40 deg 25'39.4''</td>
<td>Aug 17 at 14:59:00 (0.096)</td>
<td>90</td>
</tr>
<tr>
<td>11(^h)49(^m)46.61(^s), -40 deg 43'10.6''</td>
<td>Aug 17 at 15:00:10 (0.097)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)48(^m)0.06(^s), -40 deg 42'48.2''</td>
<td>Aug 17 at 15:01:21 (0.097)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)46(^m)15.82(^s), -40 deg 42'53.4''</td>
<td>Aug 17 at 15:02:31 (0.098)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)44(^m)29.36(^s), -40 deg 42'49.7''</td>
<td>Aug 17 at 15:03:41 (0.099)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)43(^m)38.27(^s), -40 deg 25'10.7''</td>
<td>Aug 17 at 15:04:49 (0.100)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)42(^m)45.53(^s), -40 deg 07'26.9''</td>
<td>Aug 17 at 15:05:57 (0.10)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)41(^m)53.04(^s), -39 deg 49'54.0''</td>
<td>Aug 17 at 15:07:05 (0.10)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)42(^m)47.85(^s), -39 deg 32'33.8''</td>
<td>Aug 17 at 15:08:12 (0.10)</td>
<td>80</td>
</tr>
<tr>
<td>11(^h)43(^m)39.73(^s), -39 deg 15'04.2''</td>
<td>Aug 17 at 15:09:17 (0.10)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)44(^m)33.29(^s), -38 deg 57'26.2''</td>
<td>Aug 17 at 15:10:23 (0.10)</td>
<td>87</td>
</tr>
<tr>
<td>11(^h)46(^m)18.76(^s), -38 deg 57'49.4''</td>
<td>Aug 17 at 15:11:29 (0.10)</td>
<td>77</td>
</tr>
</tbody>
</table>

\(^a\) Observations were 2017, times in parentheses are days since the GW trigger.
<table>
<thead>
<tr>
<th>Pointing direction (J2000)</th>
<th>Start time(^a) (UTC)</th>
<th>Exposure (s)</th>
</tr>
</thead>
<tbody>
<tr>
<td>11(^h)48(^m)3.97(^s), -38 deg 57'51.7(^&quot;)</td>
<td>Aug 17 at 15:12:34 (0.11)</td>
<td>85</td>
</tr>
<tr>
<td>11(^h)49(^m)49.45(^s), -38 deg 57'44.9(^&quot;)</td>
<td>Aug 17 at 15:13:40 (0.11)</td>
<td>90</td>
</tr>
<tr>
<td>15(^h)02(^m)24.24(^s), -43 deg 31'11.6(^&quot;)</td>
<td>Aug 17 at 17:16:33 (0.19)</td>
<td>15</td>
</tr>
<tr>
<td>16(^h)14(^m)12.96(^s), -61 deg 03'11.2(^&quot;)</td>
<td>Aug 17 at 17:17:28 (0.19)</td>
<td>57</td>
</tr>
<tr>
<td>09(^h)16(^m)10.08(^s), -37 deg 53'03.5(^&quot;)</td>
<td>Aug 17 at 17:50:11 (0.21)</td>
<td>27</td>
</tr>
<tr>
<td>10(^h)30(^m)24.24(^s), -35 deg 03'59.0(^&quot;)</td>
<td>Aug 17 at 17:51:07 (0.22)</td>
<td>52</td>
</tr>
<tr>
<td>10(^h)40(^m)22.80(^s), -36 deg 19'07.7(^&quot;)</td>
<td>Aug 17 at 17:52:57 (0.22)</td>
<td>60</td>
</tr>
<tr>
<td>10(^h)41(^m)56.16(^s), -36 deg 19'07.7(^&quot;)</td>
<td>Aug 17 at 17:54:29 (0.22)</td>
<td>60</td>
</tr>
<tr>
<td>11(^h)08(^m)49.68(^s), -37 deg 53'03.5(^&quot;)</td>
<td>Aug 17 at 17:55:54 (0.22)</td>
<td>52</td>
</tr>
<tr>
<td>11(^h)06(^m)22.56(^s), -37 deg 15'29.2(^&quot;)</td>
<td>Aug 17 at 17:57:31 (0.22)</td>
<td>62</td>
</tr>
<tr>
<td>11(^h)37(^m)23.52(^s), -37 deg 53'03.5(^&quot;)</td>
<td>Aug 17 at 17:59:03 (0.22)</td>
<td>55</td>
</tr>
<tr>
<td>11(^h)51(^m)24.96(^s), -38 deg 30'37.4(^&quot;)</td>
<td>Aug 17 at 18:00:41 (0.22)</td>
<td>60</td>
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<tr>
<td>11(^h)52(^m)7.44(^s), -38 deg 49'24.6(^&quot;)</td>
<td>Aug 17 at 18:02:15 (0.22)</td>
<td>57</td>
</tr>
<tr>
<td>11(^h)53(^m)37.20(^s), -39 deg 26'58.9(^&quot;)</td>
<td>Aug 17 at 18:03:40 (0.22)</td>
<td>62</td>
</tr>
<tr>
<td>11(^h)55(^m)14.40(^s), -39 deg 26'58.9(^&quot;)</td>
<td>Aug 17 at 18:05:11 (0.23)</td>
<td>35</td>
</tr>
<tr>
<td>12(^h)00(^m)6.24(^s), -39 deg 26'58.9(^&quot;)</td>
<td>Aug 17 at 18:06:36 (0.23)</td>
<td>55</td>
</tr>
<tr>
<td>12(^h)50(^m)58.32(^s), -41 deg 00'54.7(^&quot;)</td>
<td>Aug 17 at 18:08:07 (0.23)</td>
<td>57</td>
</tr>
<tr>
<td>12(^h)53(^m)23.52(^s), -41 deg 38'29.0(^&quot;)</td>
<td>Aug 17 at 18:09:50 (0.23)</td>
<td>62</td>
</tr>
<tr>
<td>12(^h)52(^m)58.08(^s), -41 deg 57'16.2(^&quot;)</td>
<td>Aug 17 at 18:11:21 (0.23)</td>
<td>55</td>
</tr>
<tr>
<td>12(^h)48(^m)21.84(^s), -41 deg 38'29.0(^&quot;)</td>
<td>Aug 17 at 18:12:46 (0.23)</td>
<td>62</td>
</tr>
<tr>
<td>12(^h)47(^m)54.96(^s), -41 deg 57'16.2(^&quot;)</td>
<td>Aug 17 at 18:14:16 (0.23)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)45(^m)59.52(^s), -41 deg 00'54.7(^&quot;)</td>
<td>Aug 17 at 18:15:41 (0.23)</td>
<td>57</td>
</tr>
<tr>
<td>12(^h)44(^m)19.92(^s), -41 deg 00'54.7(^&quot;)</td>
<td>Aug 17 at 18:17:13 (0.23)</td>
<td>55</td>
</tr>
<tr>
<td>12(^h)39(^m)21.36(^s), -41 deg 00'54.7(^&quot;)</td>
<td>Aug 17 at 18:18:38 (0.23)</td>
<td>62</td>
</tr>
<tr>
<td>12(^h)37(^m)10.80(^s), -40 deg 23'20.4(^&quot;)</td>
<td>Aug 17 at 18:20:08 (0.24)</td>
<td>62</td>
</tr>
<tr>
<td>12(^h)36(^m)8.64(^s), -40 deg 04'33.2(^&quot;)</td>
<td>Aug 17 at 18:21:40 (0.24)</td>
<td>67</td>
</tr>
<tr>
<td>12(^h)25(^m)40.32(^s), -40 deg 23'20.4(^&quot;)</td>
<td>Aug 17 at 18:23:11 (0.24)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)24(^m)41.28(^s), -40 deg 04'33.2(^&quot;)</td>
<td>Aug 17 at 18:24:43 (0.24)</td>
<td>55</td>
</tr>
<tr>
<td>12(^h)25(^m)52.80(^s), -38 deg 49'24.6(^&quot;)</td>
<td>Aug 17 at 18:26:08 (0.24)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)25(^m)1.68(^s), -38 deg 30'37.4(^&quot;)</td>
<td>Aug 17 at 18:27:40 (0.24)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)47(^m)36.00(^s), -40 deg 04'33.2(^&quot;)</td>
<td>Aug 17 at 18:29:05 (0.24)</td>
<td>57</td>
</tr>
<tr>
<td>12(^h)42(^m)41.52(^s), -40 deg 04'33.2(^&quot;)</td>
<td>Aug 17 at 18:30:41 (0.24)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)42(^m)6.96(^s), -40 deg 23'20.4(^&quot;)</td>
<td>Aug 17 at 18:32:12 (0.24)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)53(^m)37.44(^s), -40 deg 23'20.4(^&quot;)</td>
<td>Aug 17 at 18:33:37 (0.24)</td>
<td>60</td>
</tr>
<tr>
<td>12(^h)54(^m)8.88(^s), -40 deg 04'33.2(^&quot;)</td>
<td>Aug 17 at 18:35:09 (0.25)</td>
<td>57</td>
</tr>
<tr>
<td>12(^h)53(^m)12.48(^s), -38 deg 49'24.6(^&quot;)</td>
<td>Aug 17 at 18:36:35 (0.25)</td>
<td>62</td>
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</table>

\(^a\) Observations were 2017, times in parentheses are days since the GW trigger.
<table>
<thead>
<tr>
<th>Pointing direction (J2000)</th>
<th>Start time&lt;sup&gt;a&lt;/sup&gt; (UTC)</th>
<th>Exposure (s)</th>
</tr>
</thead>
<tbody>
<tr>
<td>12&lt;sup&gt;h&lt;/sup&gt;53&lt;sup&gt;m&lt;/sup&gt;50.16&lt;sup&gt;a&lt;/sup&gt;, −38 deg 30′37.4″</td>
<td>Aug 17 at 18:38:07 (0.25)</td>
<td>57</td>
</tr>
<tr>
<td>12&lt;sup&gt;h&lt;/sup&gt;55&lt;sup&gt;m&lt;/sup&gt;14.88&lt;sup&gt;a&lt;/sup&gt;, −39 deg 26′58.9″</td>
<td>Aug 17 at 18:39:32 (0.25)</td>
<td>62</td>
</tr>
<tr>
<td>12&lt;sup&gt;h&lt;/sup&gt;56&lt;sup&gt;m&lt;/sup&gt;52.32&lt;sup&gt;a&lt;/sup&gt;, −39 deg 26′58.9″</td>
<td>Aug 17 at 18:41:03 (0.25)</td>
<td>60</td>
</tr>
<tr>
<td>13&lt;sup&gt;h&lt;/sup&gt;11&lt;sup&gt;m&lt;/sup&gt;52.08&lt;sup&gt;a&lt;/sup&gt;, −43 deg 31′11.6″</td>
<td>Aug 17 at 18:42:27 (0.25)</td>
<td>57</td>
</tr>
<tr>
<td>13&lt;sup&gt;h&lt;/sup&gt;20&lt;sup&gt;m&lt;/sup&gt;13.44&lt;sup&gt;a&lt;/sup&gt;, −44 deg 08′46.0″</td>
<td>Aug 17 at 18:44:05 (0.25)</td>
<td>60</td>
</tr>
<tr>
<td>14&lt;sup&gt;h&lt;/sup&gt;03&lt;sup&gt;m&lt;/sup&gt;18.00&lt;sup&gt;a&lt;/sup&gt;, −33 deg 30′03.2″</td>
<td>Aug 17 at 18:45:37 (0.25)</td>
<td>57</td>
</tr>
<tr>
<td>14&lt;sup&gt;h&lt;/sup&gt;17&lt;sup&gt;m&lt;/sup&gt;4.08&lt;sup&gt;a&lt;/sup&gt;, −48 deg 12′58.7″</td>
<td>Aug 17 at 18:47:24 (0.25)</td>
<td>57</td>
</tr>
<tr>
<td>14&lt;sup&gt;h&lt;/sup&gt;18&lt;sup&gt;m&lt;/sup&gt;56.88&lt;sup&gt;a&lt;/sup&gt;, −48 deg 12′58.7″</td>
<td>Aug 17 at 18:49:14 (0.26)</td>
<td>57</td>
</tr>
<tr>
<td>14&lt;sup&gt;h&lt;/sup&gt;57&lt;sup&gt;m&lt;/sup&gt;35.28&lt;sup&gt;a&lt;/sup&gt;, −41 deg 57′16.2″</td>
<td>Aug 17 at 18:50:39 (0.26)</td>
<td>55</td>
</tr>
<tr>
<td>14&lt;sup&gt;h&lt;/sup&gt;56&lt;sup&gt;m&lt;/sup&gt;41.28&lt;sup&gt;a&lt;/sup&gt;, −54 deg 28′41.5″</td>
<td>Aug 17 at 18:52:21 (0.26)</td>
<td>55</td>
</tr>
<tr>
<td>10&lt;sup&gt;h&lt;/sup&gt;21&lt;sup&gt;m&lt;/sup&gt;2.88&lt;sup&gt;a&lt;/sup&gt;, −34 deg 07′37.6″</td>
<td>Aug 17 at 19:26:30 (0.28)</td>
<td>40</td>
</tr>
<tr>
<td>10&lt;sup&gt;h&lt;/sup&gt;27&lt;sup&gt;m&lt;/sup&gt;26.16&lt;sup&gt;a&lt;/sup&gt;, −44 deg 08′46.0″</td>
<td>Aug 17 at 19:28:00 (0.28)</td>
<td>47</td>
</tr>
<tr>
<td>10&lt;sup&gt;h&lt;/sup&gt;29&lt;sup&gt;m&lt;/sup&gt;11.04&lt;sup&gt;a&lt;/sup&gt;, −44 deg 08′46.0″</td>
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<td>10&lt;sup&gt;h&lt;/sup&gt;27&lt;sup&gt;m&lt;/sup&gt;17.04&lt;sup&gt;a&lt;/sup&gt;, −40 deg 23′20.4″</td>
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<td>10&lt;sup&gt;h&lt;/sup&gt;26&lt;sup&gt;m&lt;/sup&gt;50.88&lt;sup&gt;a&lt;/sup&gt;, −40 deg 04′33.2″</td>
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<td>11&lt;sup&gt;h&lt;/sup&gt;31&lt;sup&gt;m&lt;/sup&gt;40.56&lt;sup&gt;a&lt;/sup&gt;, −36 deg 19′07.7″</td>
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<td>11&lt;sup&gt;h&lt;/sup&gt;54&lt;sup&gt;m&lt;/sup&gt;50.64&lt;sup&gt;a&lt;/sup&gt;, −37 deg 53′03.5″</td>
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<td>12&lt;sup&gt;h&lt;/sup&gt;14&lt;sup&gt;m&lt;/sup&gt;0.24&lt;sup&gt;a&lt;/sup&gt;, −35 deg 41′33.4″</td>
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<td>12&lt;sup&gt;h&lt;/sup&gt;20&lt;sup&gt;m&lt;/sup&gt;13.44&lt;sup&gt;a&lt;/sup&gt;, −38 deg 30′37.4″</td>
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<td>12&lt;sup&gt;h&lt;/sup&gt;06&lt;sup&gt;m&lt;/sup&gt;35.52&lt;sup&gt;a&lt;/sup&gt;, −39 deg 26′58.9″</td>
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<td>13&lt;sup&gt;h&lt;/sup&gt;33&lt;sup&gt;m&lt;/sup&gt;25.92&lt;sup&gt;a&lt;/sup&gt;, −32 deg 52′29.3″</td>
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<td>14&lt;sup&gt;h&lt;/sup&gt;21&lt;sup&gt;m&lt;/sup&gt;11.52&lt;sup&gt;a&lt;/sup&gt;, −58 deg 14′07.1″</td>
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<td>12&lt;sup&gt;h&lt;/sup&gt;42&lt;sup&gt;m&lt;/sup&gt;20.16&lt;sup&gt;a&lt;/sup&gt;, −12 deg 50′12.5″</td>
<td>Aug 17 at 21:16:58 (0.36)</td>
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<td>12&lt;sup&gt;h&lt;/sup&gt;40&lt;sup&gt;m&lt;/sup&gt;45.36&lt;sup&gt;a&lt;/sup&gt;, −12 deg 31′25.3″</td>
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<td>Aug 17 at 21:21:37 (0.36)</td>
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<sup>a</sup> Observations were 2017, times in parentheses are days since the GW trigger.
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<th>Start time&lt;sup&gt;a&lt;/sup&gt; (UTC)</th>
<th>Exposure (s)</th>
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<sup>a</sup> Observations were 2017, times in parentheses are days since the GW trigger.
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<tr>
<td>13(^h)42(^m)59.76(^s), −13 deg 27’46.8”</td>
<td>Aug 18 at 05:31:48 (0.70)</td>
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<tr>
<td>13(^h)44(^m)17.04(^s), −13 deg 27’46.8”</td>
<td>Aug 18 at 05:34:13 (0.70)</td>
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<tr>
<td>13(^h)44(^m)39.12(^s), −13 deg 46’34.0”</td>
<td>Aug 18 at 05:36:38 (0.71)</td>
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<tr>
<td>13(^h)45(^m)56.64(^s), −13 deg 46’34.0”</td>
<td>Aug 18 at 05:39:03 (0.71)</td>
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<tr>
<td>13(^h)47(^m)13.92(^s), −13 deg 46’34.0”</td>
<td>Aug 18 at 05:41:28 (0.71)</td>
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<tr>
<td>13(^h)47(^m)37.68(^s), −14 deg 05’21.1”</td>
<td>Aug 18 at 05:43:53 (0.71)</td>
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<tr>
<td>13(^h)48(^m)2.88(^s), −14 deg 24’08.3”</td>
<td>Aug 18 at 05:46:18 (0.71)</td>
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<tr>
<td>13(^h)53(^m)35.28(^s), −16 deg 54’25.2”</td>
<td>Aug 18 at 05:48:43 (0.71)</td>
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<tr>
<td>13(^h)24(^m)53.28(^s), −30 deg 22’12.0”</td>
<td>Aug 18 at 05:51:17 (0.72)</td>
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<tr>
<td>14(^h)10(^m)39.60(^s), −52 deg 17’11.4”</td>
<td>Aug 18 at 05:54:07 (0.72)</td>
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<td>14(^h)09(^m)47.28(^s), −51 deg 58’24.2”</td>
<td>Aug 18 at 05:57:07 (0.72)</td>
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<td>12(^h)40(^m)17.28(^s), −14 deg 24’08.3”</td>
<td>Aug 18 at 06:51:12 (0.76)</td>
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<td>12(^h)40(^m)43.20(^s), −14 deg 42’55.4”</td>
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<td>12(^h)42(^m)0.96(^s), −14 deg 42’55.4”</td>
<td>Aug 18 at 06:55:31 (0.76)</td>
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\(^a\) Observations were 2017, times in parentheses are days since the GW trigger.
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<th>Exposure (s)</th>
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<td>$12^h43^m45.12^s$, $-14$ deg $05'21.1''$</td>
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<td>$12^h45^m2.64^s$, $-14$ deg $05'21.1''$</td>
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<td>Aug 18 at 07:17:33 (0.78)</td>
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<td>Aug 18 at 07:20:04 (0.78)</td>
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<td>$12^h42^m4.32^s$, $-13^\circ46'34.0''$</td>
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<td>$12^h42^m27.84^s$, $-14^\circ05'21.1''$</td>
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<td>$12^h41^m10.32^s$, $-14^\circ05'21.1''$</td>
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<td>Aug 18 at 16:31:57 (1.2)</td>
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<td>$13^h07^m11.04^s$, $-23^\circ28'55.2''$</td>
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<td>$13^h08^m32.88^s$, $-23^\circ28'55.2''$</td>
<td>Aug 18 at 16:37:07 (1.2)</td>
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<td>$13^h10^m5.76^s$, $-23^\circ10'08.0''$</td>
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<td>122</td>
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<td>$13^h08^m44.16^s$, $-23^\circ10'08.0''$</td>
<td>Aug 18 at 16:42:03 (1.2)</td>
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<td>12(^h)41(^m)42.00(^s), –09 deg 42'21.2&quot;</td>
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<td>13(^h)11(^m)7.20(^s), –23 deg 47'42.4&quot;</td>
<td>Aug 18 at 18:02:06 (1.2)</td>
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<td>13(^h)09(^m)36.72(^s), –24 deg 06'29.5&quot;</td>
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<tr>
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<td>13(^h)16(^m)4.32(^s), −28 deg 48'16.6&quot;</td>
<td>Aug 19 at 05:42:53 (1.7)</td>
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<tr>
<td>13(^h)16(^m)19.92(^s), −29 deg 07'03.4&quot;</td>
<td>Aug 19 at 05:45:24 (1.7)</td>
<td>112</td>
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<tr>
<td>13(^h)29(^m)23.28(^s), −34 deg 26'24.7&quot;</td>
<td>Aug 19 at 05:47:49 (1.7)</td>
<td>115</td>
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<tr>
<td>13(^h)58(^m)32.88(^s), −43 deg 49'58.8&quot;</td>
<td>Aug 19 at 05:50:27 (1.7)</td>
<td>117</td>
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<tr>
<td>14(^h)20(^m)19.92(^s), −55 deg 06'15.8&quot;</td>
<td>Aug 19 at 05:53:11 (1.7)</td>
<td>117</td>
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<tr>
<td>12(^h)54(^m)53.76(^s), −16 deg 54'25.2&quot;</td>
<td>Aug 19 at 06:47:14 (1.8)</td>
<td>87</td>
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<tr>
<td>12(^h)55(^m)35.04(^s), −16 deg 35'38.0&quot;</td>
<td>Aug 19 at 06:49:07 (1.8)</td>
<td>117</td>
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<tr>
<td>12(^h)56(^m)53.52(^s), −16 deg 35'38.0&quot;</td>
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<td>12(^h)56(^m)28.08(^s), −15 deg 39'16.6&quot;</td>
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<td>12(^h)57(^m)46.08(^s), −15 deg 39'16.6&quot;</td>
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<td>12(^h)59(^m)4.08(^s), −15 deg 39'16.6&quot;</td>
<td>Aug 19 at 07:03:37 (1.8)</td>
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<td>12(^h)59(^m)50.40(^s), −15 deg 20'29.8&quot;</td>
<td>Aug 19 at 07:06:02 (1.8)</td>
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<td>13(^h)01(^m)8.16(^s), −15 deg 20'29.8&quot;</td>
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<td>$12^h58^m50.88^s$, $-14$ deg $42^m55.4^s$</td>
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<td>$13^h08^m7.44^s$, $-24$ deg $25^m16.7^s$</td>
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<tr>
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<td>$13^h02^m10.32^s$, $-17$ deg $50^m46.7^s$</td>
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<td>$13^h02^m46.80^s$, $-17$ deg $31^m59.5^s$</td>
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<td>$14^h11^m49.20^s$, $-51$ deg $58^m24.2^s$</td>
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<td>$12^h54^m8.64^s$, $-15$ deg $01^m42.6^s$</td>
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<tr>
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<td>12(^h)56(^m)44.40(^s), −20 deg 21'03.6&quot;</td>
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<td>13(^h)04(^m)1.20(^s), −22 deg 13'46.6&quot;</td>
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<td>13(^h)31(^m)45.60(^s), −33 deg 30'03.2&quot;</td>
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<td>14(^h)17(^m)12.72(^s), −55 deg 25'03.0&quot;</td>
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<tr>
<td>12(^h)53(^m)2.40(^s), −11 deg 53'51.0&quot; Aug 19 at 11:54:49 (2.0)</td>
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<td>12(^h)47(^m)31.84(^s), −15 deg 39'16.6&quot; Aug 19 at 13:08:56 (2.0)</td>
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<td>12(^h)47(^m)54.00(^s), −15 deg 58'03.7&quot; Aug 19 at 13:11:06 (2.0)</td>
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<td>13(^h)44(^m)5.28(^s), −41 deg 00'54.7&quot; Aug 19 at 13:48:23 (2.0)</td>
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<td>(12^h 53^m 5.76^s), (-18) deg (28'21.0&quot;)</td>
<td>Aug 19 at 18:15:02 (2.2)</td>
<td>117</td>
</tr>
<tr>
<td>(12^h 53^m 51.60^s), (-18) deg (47'08.2&quot;)</td>
<td>Aug 19 at 18:17:27 (2.2)</td>
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<tr>
<td>(12^h 54^m 39.12^s), (-19) deg (05'55.3&quot;)</td>
<td>Aug 19 at 18:19:52 (2.2)</td>
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<tr>
<td>(12^h 54^m 8.40^s), (-19) deg (24'42.5&quot;)</td>
<td>Aug 19 at 18:22:17 (2.2)</td>
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<tr>
<td>(12^h 53^m 19.44^s), (-19) deg (05'55.3&quot;)</td>
<td>Aug 19 at 18:24:42 (2.2)</td>
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<tr>
<td>(12^h 52^m 32.40^s), (-18) deg (47'08.2&quot;)</td>
<td>Aug 19 at 18:27:12 (2.2)</td>
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<tr>
<td>(13^h 07^m 59.76^s), (-20) deg (58'37.9&quot;)</td>
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<td>37</td>
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<tr>
<td>(13^h 45^m 52.80^s), (-39) deg (08'11.8&quot;)</td>
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<td>95</td>
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<td>(14^h 06^m 6.96^s), (-53) deg (32'20.0&quot;)</td>
<td>Aug 19 at 18:37:39 (2.2)</td>
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<td>(13^h 01^m 28.08^s), (-17) deg (31'59.5&quot;)</td>
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<td>(13^h 00^m 47.28^s), (-17) deg (13'12.4&quot;)</td>
<td>Aug 20 at 00:21:19 (2.5)</td>
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<td>(13^h 00^m 7.92^s), (-16) deg (54'25.2&quot;)</td>
<td>Aug 20 at 00:23:44 (2.5)</td>
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<td>(12^h 59^m 28.56^s), (-17) deg (13'12.4&quot;)</td>
<td>Aug 20 at 00:26:09 (2.5)</td>
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<td>(13^h 00^m 9.12^s), (-17) deg (31'59.5&quot;)</td>
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<td>(12^h 59^m 32.40^s), (-17) deg (50'46.7&quot;)</td>
<td>Aug 20 at 00:30:59 (2.5)</td>
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<td>(13^h 00^m 51.36^s), (-17) deg (50'46.7&quot;)</td>
<td>Aug 20 at 00:33:23 (2.5)</td>
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<tr>
<td>13(^h)18(^m)6.24(^s), −27 deg 33'07.9&quot;</td>
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<td>13(^h)18(^m)16.32(^s), −27 deg 51'55.1&quot;</td>
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<td>13(^h)18(^m)27.84(^s), −28 deg 10'42.2&quot;</td>
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<td>13(^h)43(^m)37.84(^s), −15 deg 39'16.6&quot;</td>
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<td>12(^h)43(^m)59.52(^s), −15 deg 58'03.7&quot;</td>
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<td>12(^h)44(^m)32.88(^s), −16 deg 16'50.9&quot;</td>
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<td>13(^h)45(^m)7.68(^s), −16 deg 35'38.0&quot;</td>
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<td>12(^h)46(^m)21.84(^s), −17 deg 13'12.4&quot;</td>
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<td>14(^h)04(^m)40.80(^s), −50 deg 43'16.0&quot;</td>
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<td>13(^h)16(^m)36.96(^s), −29 deg 25'50.5&quot;</td>
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<td>13&lt;sup&gt;h&lt;/sup&gt;13&lt;sup&gt;m&lt;/sup&gt;37.44&lt;sup&gt;s&lt;/sup&gt;, −26 deg 55′33.6″</td>
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<td>13&lt;sup&gt;h&lt;/sup&gt;36&lt;sup&gt;m&lt;/sup&gt;15.84&lt;sup&gt;s&lt;/sup&gt;, −33 deg 30′03.2″</td>
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<td>12(^{h})51(^{m})46.56(^{s}), −18 deg 28′21.0″</td>
<td>Aug 20 at 19:31:32 (3.3)</td>
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<td>12(^{h})51(^{m})12.96(^{s}), −18 deg 47′08.2″</td>
<td>Aug 20 at 19:33:57 (3.3)</td>
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<td>13(^{h})11(^{m})21.60(^{s}), −28 deg 10′42.2″</td>
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\(^a\) Observations were 2017, times in parentheses are days since the GW trigger.
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<td>Aug 21 at 13:09:44 (4.0)</td>
<td>120</td>
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<tr>
<td>13(^h)34(^m)55.44(^s), −32 deg 52′29.3″</td>
<td>Aug 21 at 13:12:10 (4.0)</td>
<td>115</td>
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<tr>
<td>13(^h)35(^m)34.80(^s), −33 deg 11′16.4″</td>
<td>Aug 21 at 13:14:35 (4.0)</td>
<td>112</td>
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<tr>
<td>13(^h)36(^m)24.96(^s), −32 deg 52′29.3″</td>
<td>Aug 21 at 13:17:00 (4.0)</td>
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<tr>
<td>13(^h)37(^m)54.48(^s), −32 deg 52′29.3″</td>
<td>Aug 21 at 13:19:25 (4.0)</td>
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<tr>
<td>13(^h)37(^m)4.56(^s), −33 deg 11′16.4″</td>
<td>Aug 21 at 13:21:50 (4.0)</td>
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</table>

\(^a\) Observations were 2017, times in parentheses are days since the GW trigger.
<table>
<thead>
<tr>
<th>Pointing direction (J2000)</th>
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<tbody>
<tr>
<td>13(^h)37(^m)46.08(^s), -33 deg 30'03.2&quot;</td>
<td>Aug 21 at 13:24:15 (4.0)</td>
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<tr>
<td>13(^h)38(^m)29.28(^s), -33 deg 48'50.4&quot;</td>
<td>Aug 21 at 13:26:40 (4.0)</td>
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<td>13(^h)39(^m)59.52(^s), -33 deg 48'50.4&quot;</td>
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<td>13(^h)40(^m)46.32(^s), -33 deg 30'03.2&quot;</td>
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<td>13(^h)39(^m)16.08(^s), -33 deg 30'03.2&quot;</td>
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<td>13(^h)40(^m)56.24(^s), -36 deg 37'54.8&quot;</td>
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<tr>
<td>13(^h)50(^m)45.36(^s), -48 deg 12'58.7&quot;</td>
<td>Aug 21 at 13:38:54 (4.0)</td>
<td>117</td>
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<tr>
<td>13(^h)30(^m)9.36(^s), -34 deg 45'11.9&quot;</td>
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<td>13(^h)30(^m)57.36(^s), -35 deg 03'59.0&quot;</td>
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<td>13(^h)32(^m)29.28(^s), -35 deg 03'59.0&quot;</td>
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<td>13(^h)33(^m)12.24(^s), -34 deg 45'11.9&quot;</td>
<td>Aug 21 at 14:43:22 (4.1)</td>
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<td>13(^h)32(^m)25.44(^s), -34 deg 26'24.7&quot;</td>
<td>Aug 21 at 14:45:47 (4.1)</td>
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<td>13(^h)31(^m)40.32(^s), -34 deg 07'37.6&quot;</td>
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<td>13(^h)24(^m)55.20(^s), -35 deg 48'50.4&quot;</td>
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<td>13(^h)24(^m)14.88(^s), -33 deg 30'03.2&quot;</td>
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<td>13(^h)25(^m)6.24(^s), -32 deg 11'16.4&quot;</td>
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<td>13(^h)25(^m)58.56(^s), -32 deg 52'29.3&quot;</td>
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<td>13(^h)28(^m)21.12(^s), -32 deg 33'42.1&quot;</td>
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<td>14(^h)06(^m)10.56(^s), -55 deg 25'03.0&quot;</td>
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<td>13(^h)07(^m)54.24(^s), -19 deg 05'55.3&quot;</td>
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<tr>
<td>13(^h)08(^m)24.72(^s), -18 deg 47'08.2&quot;</td>
<td>Aug 21 at 16:11:41 (4.1)</td>
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<td>13(^h)07(^m)37.20(^s), -18 deg 28'21.0&quot;</td>
<td>Aug 21 at 16:14:07 (4.1)</td>
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<td>13(^h)06(^m)18.00(^s), -18 deg 28'21.0&quot;</td>
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<td>13(^h)06(^m)51.36(^s), -18 deg 09'33.8&quot;</td>
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<td>13(^h)12(^m)24.72(^s), -21 deg 54'59.4&quot;</td>
<td>Aug 21 at 16:21:22 (4.2)</td>
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<tr>
<td>13(^h)10(^m)47.28(^s), -22 deg 13'46.6&quot;</td>
<td>Aug 21 at 16:23:57 (4.2)</td>
<td>122</td>
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<tr>
<td>13(^h)11(^m)53.52(^s), -22 deg 32'33.7&quot;</td>
<td>Aug 21 at 16:26:28 (4.2)</td>
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<td>13(^h)13(^m)1.44(^s), -22 deg 51'20.9&quot;</td>
<td>Aug 21 at 16:28:58 (4.2)</td>
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<td>13(^h)14(^m)11.04(^s), -23 deg 10'08.0&quot;</td>
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<td>13(^h)14(^m)0.48(^s), -23 deg 28'55.2&quot;</td>
<td>Aug 21 at 16:33:59 (4.2)</td>
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<tr>
<td>13(^h)15(^m)13.68(^s), -23 deg 47'42.4&quot;</td>
<td>Aug 21 at 16:36:24 (4.2)</td>
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<tr>
<td>13(^h)16(^m)28.32(^s), −24 deg 06′29.5″</td>
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<td>13(^h)16(^m)22.56(^s), −24 deg 25′16.7″</td>
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<td>13(^h)47′29.76(^s), −39 deg 08′11.8″</td>
<td>Aug 21 at 16:49:02 (4.2)</td>
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<td>14(^h)08′23.04(^s), −55 deg 25′03.0″</td>
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<tr>
<td>13(^h)03′25.44(^s), −16 deg 35′38.0″</td>
<td>Aug 21 at 17:43:57 (4.2)</td>
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<td>13(^h)02′48.96(^s), −16 deg 16′50.9″</td>
<td>Aug 21 at 17:46:42 (4.2)</td>
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<tr>
<td>13(^h)14′51.84(^s), −25 deg 02′50.6″</td>
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<td>13(^h)16′18.00(^s), −24 deg 44′03.5″</td>
<td>Aug 21 at 17:51:49 (4.2)</td>
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<td>13(^h)17′37.68(^s), −25 deg 02′50.6″</td>
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<td>13(^h)17′40.80(^s), −24 deg 44′03.5″</td>
<td>Aug 21 at 17:56:50 (4.2)</td>
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<td>13(^h)19′7.44(^s), −24 deg 25′16.7″</td>
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<td>13(^h)31′39.12(^s), −31 deg 56′07.8″</td>
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<td>13(^h)54′55.92(^s), −48 deg 31′45.8″</td>
<td>Aug 21 at 18:25:23 (4.2)</td>
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<td>12(^h)57′15.36(^s), −22 deg 13′46.6″</td>
<td>Aug 21 at 19:19:57 (4.3)</td>
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<td>12(^h)56′12.72(^s), −21 deg 54′59.4″</td>
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<td>13(^h)05′16.08(^s), −24 deg 44′03.5″</td>
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<tr>
<td>13(^h)03′53.28(^s), −24 deg 44′03.5″</td>
<td>Aug 21 at 19:35:08 (4.3)</td>
<td>142</td>
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</tbody>
</table>

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References and Notes


7. Materials and methods are available as supplementary materials.


10. M. M. Kasliwal et al., Science. 10.1126/science.aap9455


53. S. Smartt et al., *Nature* 10.1038/nature24303 (2017). [10.1038/nature24303](https://doi.org/10.1038/nature24303)


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