



Observation of $Y(1S)$ pair production in proton-proton collisions at $\sqrt{s} = 8$ TeV

The CMS Collaboration*

Abstract

Pair production of $Y(1S)$ mesons is observed at the LHC in proton-proton collisions at $\sqrt{s} = 8$ TeV by the CMS experiment in a data sample corresponding to an integrated luminosity of 20.7 fb^{-1} . Both $Y(1S)$ candidates are fully reconstructed via their decays to $\mu^+\mu^-$. The fiducial acceptance region is defined by an absolute $Y(1S)$ rapidity smaller than 2.0. The fiducial cross section for the production of $Y(1S)$ pairs, assuming that both mesons decay isotropically, is measured to be 68.8 ± 12.7 (stat) ± 7.4 (syst) ± 2.8 (\mathcal{B}) pb, where the third uncertainty comes from the uncertainty in the branching fraction of $Y(1S)$ decays to $\mu^+\mu^-$. Assuming instead that the $Y(1S)$ mesons are produced with different polarizations leads to variations in the measured cross section in the range from -38% to $+36\%$.

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1 Introduction

The measurement of quarkonium pair production in proton-proton (pp) collisions provides insight into the underlying mechanism of particle production. The production mechanisms of heavy-quarkonium pairs such as J/ψ and Y are especially valuable as they probe these interactions in both perturbative and nonperturbative regimes [1]. In this paper, the first cross section measurement of $Y(1S)$ pair production, reconstructed in the four-muon final state, in pp collisions from the CERN LHC at $\sqrt{s} = 8$ TeV is reported.

Proton collisions can be described by parton models [2] in quantum chromodynamics (QCD). In this framework, each colliding hadron is characterized as a collection of free elementary constituents. Because of the composite nature of hadrons, in a single hadron-hadron collision two partons often undergo a single interaction (single-parton scattering, SPS). However, it is also possible that multiple distinct interactions (multiple-parton interactions, MPIs) occur, the simplest case being double-parton scattering (DPS). The SPS mechanism for heavy-quarkonium pair production can be described by a nonrelativistic effective theory [3] of QCD. However, contributions from the DPS mechanism are not addressed in a simple way within the framework of perturbative QCD [4], and DPS or MPIs are widely invoked to account for the observations that cannot be explained otherwise, such as the rates for multiple heavy-flavor production [5]. There are still large uncertainties owing to possible higher-order SPS contributions and the limited knowledge of the proton transverse profile [6]. Heavy-quarkonium final states are expected to probe the distribution of gluons in a proton since their production is dominated by gluon-gluon (gg) interactions [7, 8]. Cross section measurements of quarkonium pair production are essential in understanding SPS and DPS contributions and the parton structure of the proton.

The first quarkonium pair production measurement dates back to 1982 when the NA3 experiment at CERN observed direct production of two J/ψ mesons [9] in π^- -platinum interactions with beam energy of 150 GeV and 280 GeV, where the main contribution is quark-antiquark annihilation. Early theoretical calculations of the pair production cross section assumed only color-singlet states [10–12]. However, because of the high parton flux and high center-of-mass energy at the LHC, DPS is expected to play a significant role in quarkonium pair production [13].

Quarkonium pair production in pp collisions via DPS is assumed to result from two independent SPS occurrences [7, 13]. Several DPS production processes, including final states with associated jets, are commonly described by an effective cross section (σ_{eff}) that characterizes the transverse area of the hard partonic interactions [14, 15]. It is estimated as the product of single quarkonium production cross sections divided by the corresponding DPS quarkonium pair cross section [16]. Assuming the parton distribution functions are not correlated, σ_{eff} is expected to be independent of the final state and the center-of-mass energy. Recent measurements of σ_{eff} in final states with jets are in the range 12–20 mb [17–20]. However, measurements of J/ψ pair [21] and $J/\psi + Y(1S)$ [22] production in $p\bar{p}$ collisions yield values of 4.8 and 2.2 mb, respectively. From a fit to differential cross section measurements of J/ψ pair production with CMS [23], Lansberg and Shao [24] estimate an effective cross section of 8.2 mb. In this paper, we report the first observation of $Y(1S)$ pair production and estimate σ_{eff} using this result, along with theoretical predictions.

2 The CMS detector and muon reconstruction

The central feature of the CMS apparatus is a 13 m long superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter, each composed of a barrel and two endcap sections. Extensive forward calorimetry complements the coverage provided by the barrel and endcap detectors. Muons are measured in gas-ionization detectors embedded in the steel flux-return yoke outside the solenoid. The main subdetectors used for the present analysis are the muon detection system and the silicon tracker.

Charged particle trajectories are reconstructed by the silicon tracker in the pseudorapidity range $|\eta| < 2.5$. The innermost component of the tracker consists of three cylindrical layers of pixel detectors in the barrel region and two endcap disks at each end of the barrel. The strip tracker has 10 barrel layers and 12 endcap disks at each end of the barrel. There are 66 million $100 \times 150 \mu\text{m}^2$ silicon pixels and more than 9 million silicon strips. Strip pitches vary between $80 \mu\text{m}$ and $120 \mu\text{m}$ in the inner barrel, while the inner disk sensors have a pitch between $100 \mu\text{m}$ and $141 \mu\text{m}$. For charged particles of transverse momentum $1 < p_T < 10 \text{ GeV}$ and $|\eta| < 1.4$, the relative track resolution is typically 1.5% in p_T and 25–90 (45–50) μm in the transverse (longitudinal) impact parameter [25].

The outermost part of the CMS detector is the muon system, which covers $|\eta| < 2.4$. It consists of four layers of detection planes made of drift tubes, cathode strip chambers, and resistive-plate chambers. Muon reconstruction and identification begins within the muon system using hits in the muon chambers. A matching algorithm makes the best association between a track segment in the muon system and a track segment in the inner tracker. A track fit is performed combining all hits from both subsystems to select high-purity muon candidates.

Data are collected with a two-level trigger system. The first level of the trigger system, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events that pass the trigger requirements in a fixed time interval of about $4 \mu\text{s}$, of which $1 \mu\text{s}$ is available for data processing. The second stage, the high-level trigger (HLT), is a processor farm that further reduces the event rate from around 100 kHz to less than 1 kHz before data storage. The HLT has full access to all the event information, including tracking, and therefore selections based on algorithms similar to those applied offline are used.

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in ref. [26].

3 Data, simulation, and event selection

The data sample used in this analysis was collected with the CMS detector at the LHC in proton-proton collisions at a center-of-mass energy of 8 TeV and corresponds to an integrated luminosity of 20.7 fb^{-1} . The average number of simultaneous pp collisions in a bunch crossing (pileup) during this run was 16.

Two separate simulated samples that account for the different production mechanisms of $Y(1S)$ meson pairs are used to validate the efficiency and acceptance correction methods and parameterize the kinematic distributions of the data. Strongly correlated $Y(1S)$ meson pairs produced via SPS are generated with the CASCADE 2.3.13 simulation program [27], while less-correlated $Y(1S)$ meson pairs produced via DPS are generated with the PYTHIA 8.1.58 package [28]. The CASCADE event simulation uses an off-shell matrix-element description of the

$g^*g^* \rightarrow Y(1S)Y(1S)$ production and includes the k_T -dependent parton distribution function JH-2013-set2 [29]. The parton evolution follows the CCFM prescription [30–32], which forms a bridge between the DGLAP and BFKL parton resummation models. The CMS detector response is simulated with the GEANT4 toolkit [33]. These samples are simulated with the appropriate conditions for the analysed data, including the effects of alignment, efficiency, and pileup.

For the rest of the paper, the notation YY will be used to denote any pair-wise combination of $Y(1S)$, $Y(2S)$, and $Y(3S)$ mesons. Candidate YY events are required to have four or more muons that pass the first-level muon trigger, of which at least three must be identified as muons by the HLT. There must be at least one pair of oppositely charged muons that have an invariant mass $M_{\mu\mu}$ between 8.5 and 11 GeV, and a vertex fit chi-squared (χ^2) probability greater than 0.5%, as determined from a Kalman-filter algorithm [34].

The offline selection of YY candidates begins with a search for four muons with the sum of their charges equal to 0. Selected muons are further required to pass the following quality criteria: tracks of muon candidates must have at least six hits in the silicon tracker, at least one hit in the pixel detector, and they must match at least one muon segment in the muon detector. Loose selection cuts are applied on the longitudinal and transverse impact parameters of the muons to reject muons from cosmic rays and weakly decaying hadrons. To ensure a nearly uniform single-muon acceptance and a well-defined kinematic region, the muon candidates must have $p_T^\mu > 3.5$ GeV and $|\eta^\mu| < 2.4$.

To reconstruct the $Y(1S)Y(1S) \rightarrow 4\mu$ decay, a kinematic fit [35] is performed on oppositely charged muon pairs and then on the four muons. The fit incorporates the decay and production kinematic properties to reconstruct the full decay chain of the user-defined process. To be compatible with the trigger requirements, each reconstructed Y candidate is required to have a vertex fit χ^2 probability larger than 0.5%, as determined by kinematic fitting, and an invariant mass between 8.5 and 11 GeV. To suppress contributions from pileup events, all four muons are also required to be associated with a common vertex having a fit χ^2 probability larger than 5%.

In order to measure the cross section in a well-define kinematic region, both Y candidates must have $|y^Y| < 2.0$, where y is the rapidity. Some of the events containing more than four muons result in multiple YY candidates per event since multiple dimuon pairs can pass the trigger and selection requirements. These events constitute about 4% of the selected events and are discarded from further analysis. After all selection criteria are applied, a total of 313 YY candidates are found. The major contribution of background comes from miscombined muons from Drell–Yan production and semileptonic $b \rightarrow \mu + X$ or $c \rightarrow \mu + X$ decays that pass the trigger and selection requirements.

To extract the signal yield of $Y(1S)Y(1S)$ events, a two-dimensional (2D) unbinned extended maximum-likelihood fit to the invariant masses of the two reconstructed $\mu^+\mu^-$ combinations is used. Two kinematic variables are employed to discriminate the $Y(1S)Y(1S)$ signal from the background: the invariant mass of the higher-mass Y candidate ($M_{\mu\mu}^{(1)}$) and the invariant mass of the lower-mass Y candidate ($M_{\mu\mu}^{(2)}$). This choice of variables helps to resolve the ambiguity of having a dimuon pair in which an $Y(2S)$ candidate appears. There is no visible signal for $Y(3S)$, which is not considered further in this study. Figure 1 shows the distribution of $M_{\mu\mu}^{(1)}$ versus $M_{\mu\mu}^{(2)}$ after all selection criteria are applied. Five categories of events are considered to represent the sample of events:

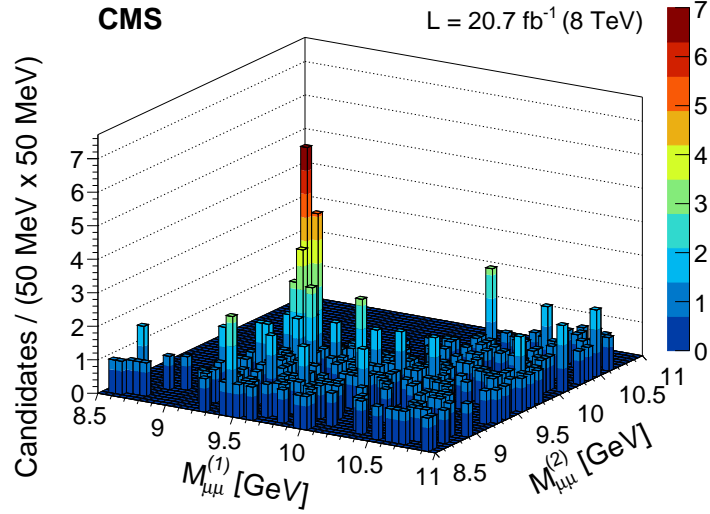


Figure 1: Two-dimensional distribution of $M_{\mu\mu}^{(1)}$ vs. $M_{\mu\mu}^{(2)}$.

1. events containing a genuine $Y(1S)Y(1S)$ pair (labeled signal),
2. events containing a genuine $Y(2S)Y(1S)$ pair (labeled $Y(2S)-Y(1S)$),
3. events containing a genuine $Y(1S)$ and two unassociated muons that contribute to the $M_{\mu\mu}^{(2)}$ distribution (labeled $Y(1S)$ -combinatorial),
4. events containing a genuine $Y(2S)$ and two unassociated muons that contribute to the $M_{\mu\mu}^{(2)}$ distribution (labeled $Y(2S)$ -combinatorial), and
5. events containing four unassociated muons (labeled combinatorial).

The contributions from other categories (combinatorial- $Y(1S)$ and combinatorial- $Y(2S)$), where two unassociated muons form the higher-mass Y candidate, are found to be consistent with zero. There is no indication of $Y(2S)Y(2S)$ events.

Each opposite-sign dimuon invariant mass distribution contains the contributions from Y candidates and a nonresonant background. The Y signal is modeled as a sum of two Crystal Ball functions [36] with a common mean, labeled DCB, to capture the variation in dimuon resolution as a function of y^Y . The DCB parameters are extracted from the signal SPS simulated sample and fixed in the fit to the data. The combinatorial background components are modeled with first-order Chebyshev polynomials (Poly). The distributions of the complete combinatorial background are found to be compatible with those of like-sign dimuon combinations. The yields of the five event categories and the two slope parameters for the combinatorial background are the free parameters in the fit to data. The likelihood function ℓ_k for an event k (ignoring the normalisation term for simplicity) is given as:

$$\begin{aligned}
 \ell_k = & N_{\text{signal}} [\text{DCB}_{Y(1S)}(M_{\mu\mu}^{(1)}) \text{DCB}_{Y(1S)}(M_{\mu\mu}^{(2)})] \\
 & + N_{Y(2S)-Y(1S)} [\text{DCB}_{Y(2S)}(M_{\mu\mu}^{(1)}) \text{DCB}_{Y(1S)}(M_{\mu\mu}^{(2)})] \\
 & + N_{Y(1S)\text{-combinatorial}} [\text{DCB}_{Y(1S)}(M_{\mu\mu}^{(1)}) \text{Poly}(M_{\mu\mu}^{(2)})] \\
 & + N_{Y(2S)\text{-combinatorial}} [\text{DCB}_{Y(2S)}(M_{\mu\mu}^{(1)}) \text{Poly}(M_{\mu\mu}^{(2)})] \\
 & + N_{\text{combinatorial}} [\text{Poly}(M_{\mu\mu}^{(1)}) \text{Poly}(M_{\mu\mu}^{(2)})].
 \end{aligned} \tag{1}$$

The yields N_i for the five categories are determined using the RooFit package [37] by mini-

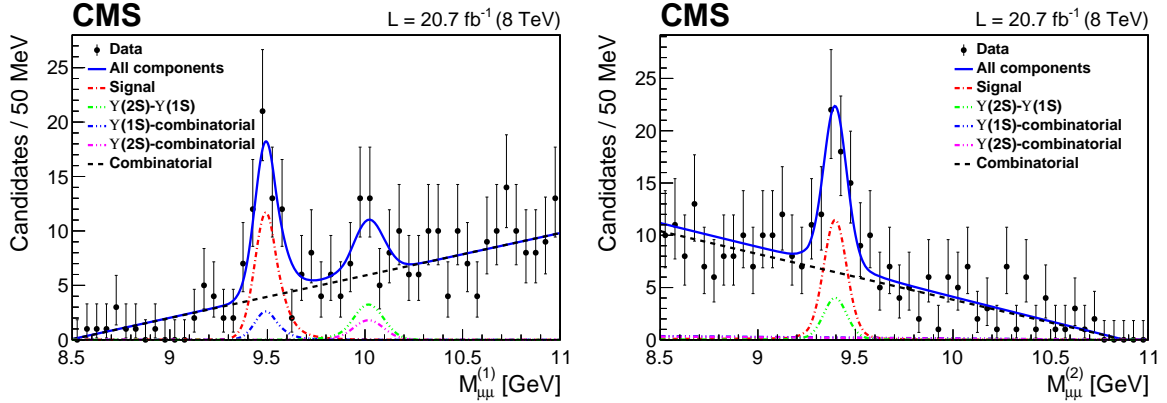


Figure 2: Invariant mass distributions of the higher-mass muon pair (left) and the lower-mass muon pair (right). The data are shown by the points. The different curves show the contributions of the various event categories from the fit.

mizing the quantity $-\ln L$, where $L = \prod_k \ell_k$ is the total extended likelihood function. The extracted yields with this procedure are $N_{\text{signal}} = 38 \pm 7$ and $N_{Y(2S)-Y(1S)} = 13^{+6}_{-5}$. To evaluate the significance of the $Y(1S)Y(1S)$ signal, the consistency of the data is checked relative to two hypotheses: a background-only hypothesis (the categories containing $Y(2S)-Y(1S)$, $Y(1S)$ -combinatorial, $Y(2S)$ -combinatorial, and combinatorial), and a signal hypothesis (the categories containing background and signal). The maximum-likelihood values returned from the fits to the signal and background-only hypotheses are denoted by L_S and L_0 , respectively. The resulting statistical significance of the signal is calculated as $\sqrt{-2(L_S/L_0)}$ [38], and is found to be well above five standard deviations. The statistical significance of the $Y(2S)-Y(1S)$ peak, evaluated using the same procedure, is 2.6 standard deviations. The invariant mass distributions of the higher- and lower-mass muon pairs in the data, with the result of the 2D fit superimposed, are shown in figure 2. The fit is validated by applying it to simulated samples generated using the probability distribution functions (pdf) that describe the data. No bias in the yields is found, and the likelihood value from the fit to the data closely agrees with the mean likelihood value obtained from the simulated samples.

4 Efficiency and acceptance

The muon acceptance is defined by the geometry of the CMS detector and the kinematic requirements of the analysis, which are established based on studies of simulated $Y(1S)Y(1S)$ events.

The kinematic variable distributions of the $Y(1S)Y(1S)$ system change with the underlying production mechanism. Therefore, to minimize the model dependence, the acceptance and efficiency corrections are calculated on an event-by-event basis using the measured Y meson and muon momenta. This method was utilized in a similar CMS analysis [23], and is intended to minimize the model dependence of the correction factors.

The probability of an event passing the muon acceptance requirements is obtained by repeatedly generating the four-momenta of the four decay muons of the measured Y pairs in the event. The acceptance correction for each selected event is calculated separately by simulating a large sample of such decays. The acceptance correction factor, a_i for an event i , is defined as the number of times all four muons pass the acceptance criteria divided by the total number

of trials. The angular distribution of the decay muons with respect to the direction of flight of the parent $Y(1S)$ meson, in the $Y(1S)$ meson rest frame, is assumed to be isotropic. Different polarization scenarios of $Y(1S)$ mesons are considered and discussed later.

The efficiency correction for the four-muon system is determined with a data-embedding method that repeatedly substitutes the measured muon four-momenta into different simulated events, which are then subjected to the complete CMS detector simulation and reconstruction chain. Each event is simulated with the same pileup scenario as in the corresponding data event. The efficiency correction factor ϵ_i for the i^{th} event in the data is the number of simulated events that pass the offline reconstruction and trigger requirements, divided by the total number of simulated events. For each data event, 1000 different simulated events are generated. The number of efficiency-corrected signal events, N_{true} , is then given as $\sum_i 1/\epsilon_i$, where the summation is over the data events that satisfy all the selection requirements.

Because of the finite $\mu^+\mu^-$ mass resolution, typically near the kinematic acceptance of the detector, the measured momenta of reconstructed Y candidates will be distorted from the original sample. To account for this, scaling factors are determined for both the SPS and DPS simulated samples. First, an average efficiency using a data-embedding method, $\bar{\epsilon}_1$, is calculated as the number of events surviving the trigger and reconstruction requirements, divided by the number of efficiency-corrected events N_{true} . Alternatively, an average efficiency using the generator information, $\bar{\epsilon}_2$, is determined as the number of events satisfying the trigger and reconstruction criteria, divided by the number of events generated within the muon and Y meson acceptance of the detector. Compared to the first method, the latter efficiency is based on the originally generated muon momenta. Then, the scaling factor is defined as the ratio $\bar{\epsilon}_2 / \bar{\epsilon}_1$. Scaling factors of 98% and 96% are determined for the SPS and DPS models, respectively. The average of these scaling factors is applied to the data-embedding efficiency.

5 Systematic uncertainties

Several sources of systematic uncertainty in the cross section measurement are considered and discussed below.

To estimate the uncertainty caused by the imperfect modeling of the data, several variations are checked: (i) the parameters for the $Y(1S)$ resonance shape that are fixed for the maximum-likelihood fit are varied by their uncertainty, as determined from the simulated samples; (ii) the $Y(1S)$ resonance shape is alternatively parameterized with a relativistic Breit-Wigner convolved with a Gaussian function; (iii) $Y(1S)$ mesons are alternatively parameterized using signal parameters obtained from the DPS simulation samples. The largest difference in the signal yield among these fits of 7.9% is taken as the corresponding systematic uncertainty.

Evaluation of the trigger and reconstruction efficiencies for Y pair production rely on detector simulation. To estimate the systematic uncertainty owing to the detector simulation, the four-muon efficiency is also calculated from the single-muon efficiencies measured in data with a "tag-and-probe" method [39] that utilizes data samples containing a single J/ψ meson decaying into a muon pair. To calculate the four-muon efficiency using the tag-and-probe method, the correlations among the muons is separately derived from the simulation. A total correlation factor, determined from simulation as a function of the p_T and rapidity of the four-muon system, is multiplied by the product of the four single-muon efficiencies, also measured versus η and p_T . The difference of the corrected signal yield with the tag-and-probe method and the data-embedding method of 4.9% is taken as the systematic uncertainty from the simulation.

The efficiency correction method is evaluated using the full detector simulation and reconstruction event samples generated according to the two production mechanisms: SPS and DPS. The systematic uncertainty from using the efficiency-correction method, a data-embedding efficiency correction method is applied to obtain the per-event efficiency, ϵ_i . The efficiency-corrected yield is compared to the number of generated events, and the fractional difference between the two is used to estimate the uncertainty in the correction method. This procedure is repeated for both simulated samples separately. The systematic uncertainty is then taken to be the larger of the differences in the two samples and is found to be 3.7%.

Signal SPS and DPS simulated samples are used to estimate the uncertainty in the event-by-event acceptance correction method. A sample of N_{gen} simulated events are subjected to the acceptance criteria. The event-based acceptance correction is then applied to the events passing the full selection to arrive at a corrected yield, N'_{gen} . The uncertainty is taken as the relative deviation of N'_{gen} from N_{gen} . The larger deviation of 2.8% between the SPS and DPS samples is taken as a systematic uncertainty.

The luminosity delivered to CMS is measured based on cluster counting from the pixel detector. The overall systematic uncertainty in the integrated luminosity is estimated to be 2.6% [40].

Table 1 provides a summary of the individual systematic uncertainties. The total systematic uncertainty in the $Y(1S)Y(1S)$ cross section measurement of 10.7% is calculated as the sum in quadrature of the individual components. In this analysis, $Y(1S)$ mesons are assumed to decay

Table 1: Summary of the relative systematic uncertainties in the $Y(1S)$ pair production cross section.

Component	Uncertainty (%)
Resonance shape	7.9
Simulation	4.9
Efficiency	3.7
Acceptance	2.8
Integrated luminosity	2.6
Total	10.7

isotropically. To check the sensitivity of the fiducial cross section measurement on the $Y(1S)$ meson decay angular distribution, studies are performed for several extreme polarization scenarios. With θ^* defined as the angle between the μ^+ direction in the $Y(1S)$ meson rest frame and the direction of the $Y(1S)$ in the lab frame, the angular distribution is parameterized as $1 + \lambda_\theta \cos^2(\theta^*)$, where λ_θ is the polar anisotropic parameter [41]. The extreme cases $\lambda_\theta = 0$ or $\lambda_\theta = +1$ (-1) corresponds to an isotropic $Y(1S)$ meson decay or 100% longitudinal (transverse) polarization. The polar anisotropic parameter for each $Y(1S)$ candidate, labeled as λ_{θ_1} and λ_{θ_2} , are varied simultaneously, and Table 2 shows a summary of this study. Compared to an isotropic $Y(1S)$ meson decay, the corrected yield is 38% lower for $\lambda_{\theta_1} = \lambda_{\theta_2} = -1$ and 36% higher for $\lambda_{\theta_1} = \lambda_{\theta_2} = +1$. This variation is not included in the total systematic uncertainty, but is quoted separately.

Table 2: Relative change in percent of the acceptance-corrected $Y(1S)Y(1S)$ yield for different polarization assumptions with respect to that for $\lambda_{\theta_1} = \lambda_{\theta_2} = 0$.

λ_{θ_1}	+1	+1	+0.5	-0.5	+0.5	-1	+1	-0.5	-0.5	-1
λ_{θ_2}	+1	+0.5	+0.5	+0.5	-0.5	+1	-1	-0.5	-1	-1
Change (%)	+36	+26	+18	-2	-3	-9	-9	-19	-29	-38

6 Results

The fiducial cross section of $Y(1S)$ pair production, σ_{fid} , is measured for events in which both $Y(1S)$ mesons have $|y^Y| < 2.0$. It is derived from the measured $Y(1S)$ pair yield, N_{signal} , using the expression:

$$\sigma_{\text{fid}} = \frac{N_{\text{signal}} \bar{\omega}}{\mathcal{L} [\mathcal{B}(Y(1S) \rightarrow \mu^+ \mu^-)]^2}, \quad (2)$$

where $\bar{\omega}$ is the average correction factor that accounts for the inefficiencies stemming from reconstruction, identification, and trigger requirements, as well as the fraction of events lost because of the muon acceptance criteria, \mathcal{L} is the integrated luminosity, and $\mathcal{B}(Y(1S) \rightarrow \mu^+ \mu^-) = (2.48 \pm 0.05)\%$ is the world-average branching fraction of the $Y(1S)$ to $\mu^+ \mu^-$ [42].

The factor $\bar{\omega}$ is the inverse product of per-event efficiencies, ϵ_i , and acceptances, a_i , averaged over the events that pass selection cuts. For the calculation of the total cross section, values of $\mathcal{L} = 20.7 \text{ fb}^{-1}$, $N_{\text{signal}} = 38 \pm 7$, and $\bar{\omega} = 23.06$ are used. Under the assumption that both $Y(1S)$ mesons decay isotropically, the total cross section for $Y(1S)$ pair production is measured to be $\sigma_{\text{fid}} = 68.8 \pm 12.7 \text{ (stat)} \pm 7.4 \text{ (syst)} \pm 2.8 \text{ (}\mathcal{B}\text{)} \text{ pb}$, where the uncertainties are statistical, systematic, and related to the $Y(1S) \rightarrow \mu^+ \mu^-$ branching fraction, respectively.

In quarkonium pair production, the measurement of the effective cross section σ_{eff} depends on the fraction of DPS, which is usually estimated either as a residual to the SPS prediction or as the result of a fit to the rapidity or azimuthal angle difference between quarkonia pairs. To estimate σ_{eff} using the fiducial $Y(1S)$ pair production cross section σ_{fid} , the following formula is used [7]:

$$\sigma_{\text{eff}} = \frac{[\sigma(Y)]^2}{2 f_{\text{DPS}} \sigma_{\text{fid}} [\mathcal{B}(Y(1S) \rightarrow \mu^+ \mu^-)]^2}, \quad (3)$$

where $\sigma(Y)$ is the cross section for single- $Y(1S)$ production [43] extrapolated to the fiducial region ($|y^Y| < 2.0$) and f_{DPS} is the fraction of the DPS contribution. To estimate the effective cross section, we use $\sigma(Y) = 7.5 \pm 0.6 \text{ nb}$ and a value of $f_{\text{DPS}} \approx 10\%$ [8], which gives $\sigma_{\text{eff}} \approx 6.6 \text{ mb}$. The cross section for $Y(1S)$ pair production, assuming SPS with feed-down from higher Y states, is 48 pb [8]. This result, combined with our measurement, gives $f_{\text{DPS}} \simeq 30\%$, and with this estimation, σ_{eff} is calculated to be $\approx 2.2 \text{ mb}$. These two estimates of σ_{eff} are consistent with the range of values from the heavy-quarkonium measurements (2–8 mb) [21, 22], but are smaller than that from multijet studies (12–20 mb) [17–20]. The quarkonium final states are dominantly produced from gluon-gluon interactions, while the jet-related channels that correspond to higher values of σ_{eff} are produced by quark-antiquark and quark-gluon parton interactions. The trend in the measured effective cross section might indicate that the average transverse distance between gluons in the proton is smaller than between quarks, or between quarks and gluons.

7 Summary

The first observation of $Y(1S)$ pair production has been performed with the CMS detector in proton-proton collisions at $\sqrt{s} = 8 \text{ TeV}$ from a sample corresponding to an integrated luminosity of 20.7 fb^{-1} . A signal yield of 38 ± 7 $Y(1S)Y(1S)$ events is measured with a significance exceeding five standard deviations. Using the measured muon kinematic properties in the data, an event-based acceptance and efficiency-correction method is applied, minimizing the model dependence of the cross section measurement. The fiducial cross section of $Y(1S)$ pair production measured within the single- Y acceptance $|y^Y| < 2.0$ is found to be $\sigma_{\text{fid}} = 68.8 \pm 12.7 \text{ (stat)} \pm 7.4 \text{ (syst)} \pm 2.8 \text{ (}\mathcal{B}\text{)} \text{ pb}$, where the $Y(1S)$ decay into muons is assumed

to be isotropic. Assuming instead that the $Y(1S)$ mesons are produced with different polarizations leads to variations in the measured cross section in the range from -38% to $+36\%$. An effective cross section is also estimated using our result that is in agreement with the values from heavy-quarkonium measurements, but is smaller than that from multijet studies.

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Facilities Council, UK; the US Department of Energy, and the US National Science Foundation.

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A The CMS Collaboration

Yerevan Physics Institute, Yerevan, Armenia

V. Khachatryan, A.M. Sirunyan, A. Tumasyan

Institut für Hochenergiephysik, Wien, Austria

W. Adam, E. Asilar, T. Bergauer, J. Brandstetter, E. Brondolin, M. Dragicevic, J. Erö, M. Flechl, M. Friedl, R. Frühwirth¹, V.M. Ghete, C. Hartl, N. Hörmann, J. Hrubec, M. Jeitler¹, A. König, I. Krätschmer, D. Liko, T. Matsushita, I. Mikulec, D. Rabadý, N. Rad, B. Rahbaran, H. Rohringer, J. Schieck¹, J. Strauss, W. Waltenberger, C.-E. Wulz¹

Institute for Nuclear Problems, Minsk, Belarus

O. Dvornikov, V. Makarenko, V. Zykunov

National Centre for Particle and High Energy Physics, Minsk, Belarus

V. Mossolov, N. Shumeiko, J. Suarez Gonzalez

Universiteit Antwerpen, Antwerpen, Belgium

S. Alderweireldt, E.A. De Wolf, X. Janssen, J. Lauwers, M. Van De Klundert, H. Van Haevermaet, P. Van Mechelen, N. Van Remortel, A. Van Spilbeek

Vrije Universiteit Brussel, Brussel, Belgium

S. Abu Zeid, F. Blekman, J. D'Hondt, N. Daci, I. De Bruyn, K. Deroover, S. Lowette, S. Moortgat, L. Moreels, A. Olbrechts, Q. Python, S. Tavernier, W. Van Doninck, P. Van Mulders, I. Van Parijs

Université Libre de Bruxelles, Bruxelles, Belgium

H. Brun, B. Clerbaux, G. De Lentdecker, H. Delannoy, G. Fasanella, L. Favart, R. Goldouzian, A. Grebenyuk, G. Karapostoli, T. Lenzi, A. Léonard, J. Luetic, T. Maerschalk, A. Marinov, A. Randle-conde, T. Seva, C. Vander Velde, P. Vanlaer, R. Yonamine, F. Zenoni, F. Zhang²

Ghent University, Ghent, Belgium

A. Cimmino, T. Cornelis, D. Dobur, A. Fagot, G. Garcia, M. Gul, I. Khvastunov, D. Poyraz, S. Salva, R. Schöfbeck, A. Sharma, M. Tytgat, W. Van Driessche, E. Yazgan, N. Zaganidis

Université Catholique de Louvain, Louvain-la-Neuve, Belgium

H. Bakhshiansohi, C. Beluffi³, O. Bondu, S. Brochet, G. Bruno, A. Caudron, S. De Visscher, C. Delaere, M. Delcourt, B. Francois, A. Giammanco, A. Jafari, P. Jez, M. Komm, V. Lemaitre, A. Magitteri, A. Mertens, M. Musich, C. Nuttens, K. Piotrkowski, L. Quertenmont, M. Selvaggi, M. Vidal Marono, S. Wertz

Université de Mons, Mons, Belgium

N. Bely

Centro Brasileiro de Pesquisas Físicas, Rio de Janeiro, Brazil

W.L. Aldá Júnior, F.L. Alves, G.A. Alves, L. Brito, C. Hensel, A. Moraes, M.E. Pol, P. Rebello Teles

Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil

E. Belchior Batista Das Chagas, W. Carvalho, J. Chinellato⁴, A. Custódio, E.M. Da Costa, G.G. Da Silveira⁵, D. De Jesus Damiao, C. De Oliveira Martins, S. Fonseca De Souza, L.M. Huertas Guativa, H. Malbouisson, D. Matos Figueiredo, C. Mora Herrera, L. Mundim, H. Nogima, W.L. Prado Da Silva, A. Santoro, A. Sznajder, E.J. Tonelli Manganote⁴, A. Vilela Pereira

Universidade Estadual Paulista ^a, Universidade Federal do ABC ^b, São Paulo, Brazil

S. Ahuja^a, C.A. Bernardes^b, S. Dogra^a, T.R. Fernandez Perez Tomei^a, E.M. Gregores^b,

P.G. Mercadante^b, C.S. Moon^a, S.F. Novaes^a, Sandra S. Padula^a, D. Romero Abad^b, J.C. Ruiz Vargas

Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria

A. Aleksandrov, R. Hadjiiska, P. Iaydjiev, M. Rodozov, S. Stoykova, G. Sultanov, M. Vutova

University of Sofia, Sofia, Bulgaria

A. Dimitrov, I. Glushkov, L. Litov, B. Pavlov, P. Petkov

Beihang University, Beijing, China

W. Fang⁶

Institute of High Energy Physics, Beijing, China

M. Ahmad, J.G. Bian, G.M. Chen, H.S. Chen, M. Chen, Y. Chen⁷, T. Cheng, C.H. Jiang, D. Leggat, Z. Liu, F. Romeo, S.M. Shaheen, A. Spiezia, J. Tao, C. Wang, Z. Wang, H. Zhang, J. Zhao

State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China

Y. Ban, G. Chen, Q. Li, S. Liu, Y. Mao, S.J. Qian, D. Wang, Z. Xu

Universidad de Los Andes, Bogota, Colombia

C. Avila, A. Cabrera, L.F. Chaparro Sierra, C. Florez, J.P. Gomez, C.F. González Hernández, J.D. Ruiz Alvarez, J.C. Sanabria

University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia

N. Godinovic, D. Lelas, I. Puljak, P.M. Ribeiro Cipriano, T. Sculac

University of Split, Faculty of Science, Split, Croatia

Z. Antunovic, M. Kovac

Institute Rudjer Boskovic, Zagreb, Croatia

V. Brigljevic, D. Ferencek, K. Kadija, S. Micanovic, L. Sudic, T. Susa

University of Cyprus, Nicosia, Cyprus

A. Attikis, G. Mavromanolakis, J. Mousa, C. Nicolaou, F. Ptochos, P.A. Razis, H. Rykaczewski, D. Tsiakkouri

Charles University, Prague, Czech Republic

M. Finger⁸, M. Finger Jr.⁸

Universidad San Francisco de Quito, Quito, Ecuador

E. Carrera Jarrin

Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt

A. Ellithi Kamel⁹, M.A. Mahmoud^{10,11}, A. Radi^{11,12}

National Institute of Chemical Physics and Biophysics, Tallinn, Estonia

B. Calpas, M. Kadastik, M. Murumaa, L. Perrini, M. Raidal, A. Tiko, C. Veelken

Department of Physics, University of Helsinki, Helsinki, Finland

P. Eerola, J. Pekkanen, M. Voutilainen

Helsinki Institute of Physics, Helsinki, Finland

J. Härkönen, T. Järvinen, V. Karimäki, R. Kinnunen, T. Lampén, K. Lassila-Perini, S. Lehti, T. Lindén, P. Luukka, J. Tuominiemi, E. Tuovinen, L. Wendland

Lappeenranta University of Technology, Lappeenranta, Finland

J. Talvitie, T. Tuuva

IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France

M. Besancon, F. Couderc, M. Dejardin, D. Denegri, B. Fabbro, J.L. Faure, C. Favaro, F. Ferri, S. Ganjour, S. Ghosh, A. Givernaud, P. Gras, G. Hamel de Monchenault, P. Jarry, I. Kucher, E. Locci, M. Machet, J. Malcles, J. Rander, A. Rosowsky, M. Titov, A. Zghiche

Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France

A. Abdulsalam, I. Antropov, S. Baffioni, F. Beaudette, P. Busson, L. Cadamuro, E. Chapon, C. Charlot, O. Davignon, R. Granier de Cassagnac, M. Jo, S. Lisniak, P. Miné, M. Nguyen, C. Ochando, G. Ortona, P. Paganini, P. Pigard, S. Regnard, R. Salerno, Y. Sirois, T. Strebler, Y. Yilmaz, A. Zabi

Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France

J.-L. Agram¹³, J. Andrea, A. Aubin, D. Bloch, J.-M. Brom, M. Buttignol, E.C. Chabert, N. Chanon, C. Collard, E. Conte¹³, X. Coubez, J.-C. Fontaine¹³, D. Gelé, U. Goerlach, A.-C. Le Bihan, K. Skovpen, P. Van Hove

Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France

S. Gadrat

Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France

S. Beauceron, C. Bernet, G. Boudoul, E. Bouvier, C.A. Carrillo Montoya, R. Chierici, D. Contardo, B. Courbon, P. Depasse, H. El Mamouni, J. Fan, J. Fay, S. Gascon, M. Gouzevitch, G. Grenier, B. Ille, F. Lagarde, I.B. Laktineh, M. Lethuillier, L. Mirabito, A.L. Pequegnot, S. Perries, A. Popov¹⁴, D. Sabes, V. Sordini, M. Vander Donckt, P. Verdier, S. Viret

Georgian Technical University, Tbilisi, Georgia

A. Khvedelidze⁸

Tbilisi State University, Tbilisi, Georgia

D. Lomidze

RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany

C. Autermann, S. Beranek, L. Feld, A. Heister, M.K. Kiesel, K. Klein, M. Lipinski, A. Ostapchuk, M. Preuten, F. Raupach, S. Schael, C. Schomakers, J. Schulz, T. Verlage, H. Weber, V. Zhukov¹⁴

RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

A. Albert, M. Brodski, E. Dietz-Laursonn, D. Duchardt, M. Endres, M. Erdmann, S. Erdweg, T. Esch, R. Fischer, A. Güth, M. Hamer, T. Hebbeker, C. Heidemann, K. Hoepfner, S. Knutzen, M. Merschmeyer, A. Meyer, P. Millet, S. Mukherjee, M. Olschewski, K. Padeken, T. Pook, M. Radziej, H. Reithler, M. Rieger, F. Scheuch, L. Sonnenschein, D. Teyssier, S. Thüer

RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany

V. Cherepanov, G. Flügge, F. Hoehle, B. Kargoll, T. Kress, A. Künsken, J. Lingemann, T. Müller, A. Nehr Korn, A. Nowack, I.M. Nugent, C. Pistone, O. Pooth, A. Stahl¹⁵

Deutsches Elektronen-Synchrotron, Hamburg, Germany

M. Aldaya Martin, T. Arndt, C. Asawatangtrakuldee, K. Beernaert, O. Behnke, U. Behrens, A.A. Bin Anuar, K. Borras¹⁶, A. Campbell, P. Connor, C. Contreras-Campana, F. Costanza, C. Diez Pardos, G. Dolinska, G. Eckerlin, D. Eckstein, T. Eichhorn, E. Eren, E. Gallo¹⁷,

J. Garay Garcia, A. Geiser, A. Gizhko, J.M. Grados Luyando, P. Gunnellini, A. Harb, J. Hauk, M. Hempel¹⁸, H. Jung, A. Kalogeropoulos, O. Karacheban¹⁸, M. Kasemann, J. Keaveney, C. Kleinwort, I. Korol, D. Krücker, W. Lange, A. Lelek, J. Leonard, K. Lipka, A. Lobanov, W. Lohmann¹⁸, R. Mankel, I.-A. Melzer-Pellmann, A.B. Meyer, G. Mittag, J. Mnich, A. Mussgiller, E. Ntomari, D. Pitzl, R. Placakyte, A. Raspereza, B. Roland, M.Ö. Sahin, P. Saxena, T. Schoerner-Sadenius, C. Seitz, S. Spannagel, N. Stefaniuk, G.P. Van Onsem, R. Walsh, C. Wissing

University of Hamburg, Hamburg, Germany

V. Blobel, M. Centis Vignali, A.R. Draeger, T. Dreyer, E. Garutti, D. Gonzalez, J. Haller, M. Hoffmann, A. Junkes, R. Klanner, R. Kogler, N. Kovalchuk, T. Lapsien, T. Lenz, I. Marchesini, D. Marconi, M. Meyer, M. Niedziela, D. Nowatschin, F. Pantaleo¹⁵, T. Peiffer, A. Perieanu, J. Poehlsen, C. Sander, C. Scharf, P. Schleper, A. Schmidt, S. Schumann, J. Schwandt, H. Stadie, G. Steinbrück, F.M. Stober, M. Stöver, H. Tholen, D. Troendle, E. Usai, L. Vanelderen, A. Vanhoefer, B. Vormwald

Institut für Experimentelle Kernphysik, Karlsruhe, Germany

M. Akbiyik, C. Barth, S. Baur, C. Baus, J. Berger, E. Butz, R. Caspart, T. Chwalek, F. Colombo, W. De Boer, A. Dierlamm, S. Fink, B. Freund, R. Friese, M. Giffels, A. Gilbert, P. Goldenzweig, D. Haitz, F. Hartmann¹⁵, S.M. Heindl, U. Husemann, I. Katkov¹⁴, S. Kudella, P. Lobelle Pardo, H. Mildner, M.U. Mozer, Th. Müller, M. Plagge, G. Quast, K. Rabbertz, S. Röcker, F. Roscher, M. Schröder, I. Shvetsov, G. Sieber, H.J. Simonis, R. Ulrich, J. Wagner-Kuhr, S. Wayand, M. Weber, T. Weiler, S. Williamson, C. Wöhrmann, R. Wolf

Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece

G. Anagnostou, G. Daskalakis, T. Gerasis, V.A. Giakoumopoulou, A. Kyriakis, D. Loukas, I. Topsis-Giotis

National and Kapodistrian University of Athens, Athens, Greece

S. Kesisoglou, A. Panagiotou, N. Saoulidou, E. Tziaferi

University of Ioánnina, Ioánnina, Greece

I. Evangelou, G. Flouris, C. Foudas, P. Kokkas, N. Loukas, N. Manthos, I. Papadopoulos, E. Paradis

MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary

N. Filipovic

Wigner Research Centre for Physics, Budapest, Hungary

G. Bencze, C. Hajdu, P. Hidas, D. Horvath¹⁹, F. Sikler, V. Veszpremi, G. Vesztergombi²⁰, A.J. Zsigmond

Institute of Nuclear Research ATOMKI, Debrecen, Hungary

N. Beni, S. Czellar, J. Karancsi²¹, A. Makovec, J. Molnar, Z. Szillasi

University of Debrecen, Debrecen, Hungary

M. Bartók²⁰, P. Raics, Z.L. Trocsanyi, B. Ujvari

National Institute of Science Education and Research, Bhubaneswar, India

S. Bahinipati, S. Choudhury²², P. Mal, K. Mandal, A. Nayak²³, D.K. Sahoo, N. Sahoo, S.K. Swain

Panjab University, Chandigarh, India

S. Bansal, S.B. Beri, V. Bhatnagar, R. Chawla, U. Bhawandeep, A.K. Kalsi, A. Kaur, M. Kaur, R. Kumar, P. Kumari, A. Mehta, M. Mittal, J.B. Singh, G. Walia

University of Delhi, Delhi, India

Ashok Kumar, A. Bhardwaj, B.C. Choudhary, R.B. Garg, S. Keshri, S. Malhotra, M. Naimuddin, N. Nishu, K. Ranjan, R. Sharma, V. Sharma

Saha Institute of Nuclear Physics, Kolkata, India

R. Bhattacharya, S. Bhattacharya, K. Chatterjee, S. Dey, S. Dutt, S. Dutta, S. Ghosh, N. Majumdar, A. Modak, K. Mondal, S. Mukhopadhyay, S. Nandan, A. Purohit, A. Roy, D. Roy, S. Roy Chowdhury, S. Sarkar, M. Sharan, S. Thakur

Indian Institute of Technology Madras, Madras, India

P.K. Behera

Bhabha Atomic Research Centre, Mumbai, India

R. Chudasama, D. Dutta, V. Jha, V. Kumar, A.K. Mohanty¹⁵, P.K. Netrakanti, L.M. Pant, P. Shukla, A. Topkar

Tata Institute of Fundamental Research-A, Mumbai, India

T. Aziz, S. Dugad, G. Kole, B. Mahakud, S. Mitra, G.B. Mohanty, B. Parida, N. Sur, B. Sutar

Tata Institute of Fundamental Research-B, Mumbai, India

S. Banerjee, S. Bhowmik²⁴, R.K. Dewanjee, S. Ganguly, M. Guchait, Sa. Jain, S. Kumar, M. Maity²⁴, G. Majumder, K. Mazumdar, T. Sarkar²⁴, N. Wickramage²⁵

Indian Institute of Science Education and Research (IISER), Pune, India

S. Chauhan, S. Dube, V. Hegde, A. Kapoor, K. Kotheekar, S. Pandey, A. Rane, S. Sharma

Institute for Research in Fundamental Sciences (IPM), Tehran, Iran

H. Behnamian, S. Chenarani²⁶, E. Eskandari Tadavani, S.M. Etesami²⁶, A. Fahim²⁷, M. Khakzad, M. Mohammadi Najafabadi, M. Naseri, S. Paktinat Mehdiabadi²⁸, F. Rezaei Hosseinabadi, B. Safarzadeh²⁹, M. Zeinali

University College Dublin, Dublin, Ireland

M. Felcini, M. Grunewald

INFN Sezione di Bari ^a, Università di Bari ^b, Politecnico di Bari ^c, Bari, Italy

M. Abbrescia^{a,b}, C. Calabria^{a,b}, C. Caputo^{a,b}, A. Colaleo^a, D. Creanza^{a,c}, L. Cristella^{a,b}, N. De Filippis^{a,c}, M. De Palma^{a,b}, L. Fiore^a, G. Iaselli^{a,c}, G. Maggi^{a,c}, M. Maggi^a, G. Miniello^{a,b}, S. My^{a,b}, S. Nuzzo^{a,b}, A. Pompili^{a,b}, G. Pugliese^{a,c}, R. Radogna^{a,b}, A. Ranieri^a, G. Selvaggi^{a,b}, L. Silvestris^{a,15}, R. Venditti^{a,b}, P. Verwilligen^a

INFN Sezione di Bologna ^a, Università di Bologna ^b, Bologna, Italy

G. Abbiendi^a, C. Battilana, D. Bonacorsi^{a,b}, S. Braibant-Giacomelli^{a,b}, L. Brigliadori^{a,b}, R. Campanini^{a,b}, P. Capiluppi^{a,b}, A. Castro^{a,b}, F.R. Cavallo^a, S.S. Chhibra^{a,b}, G. Codispoti^{a,b}, M. Cuffiani^{a,b}, G.M. Dallavalle^a, F. Fabbri^a, A. Fanfani^{a,b}, D. Fasanella^{a,b}, P. Giacomelli^a, C. Grandi^a, L. Guiducci^{a,b}, S. Marcellini^a, G. Masetti^a, A. Montanari^a, F.L. Navarria^{a,b}, A. Perrotta^a, A.M. Rossi^{a,b}, T. Rovelli^{a,b}, G.P. Siroli^{a,b}, N. Tosi^{a,b,15}

INFN Sezione di Catania ^a, Università di Catania ^b, Catania, Italy

S. Albergo^{a,b}, M. Chiorboli^{a,b}, S. Costa^{a,b}, A. Di Mattia^a, F. Giordano^{a,b}, R. Potenza^{a,b}, A. Tricomi^{a,b}, C. Tuve^{a,b}

INFN Sezione di Firenze ^a, Università di Firenze ^b, Firenze, Italy

G. Barbagli^a, V. Ciulli^{a,b}, C. Civinini^a, R. D'Alessandro^{a,b}, E. Focardi^{a,b}, V. Gori^{a,b}, P. Lenzi^{a,b}, M. Meschini^a, S. Paoletti^a, G. Sguazzoni^a, L. Viliani^{a,b,15}

INFN Laboratori Nazionali di Frascati, Frascati, Italy

L. Benussi, S. Bianco, F. Fabbri, D. Piccolo, F. Primavera¹⁵

INFN Sezione di Genova ^a, Università di Genova ^b, Genova, Italy

V. Calvelli^{a,b}, F. Ferro^a, M. Lo Vetere^{a,b}, M.R. Monge^{a,b}, E. Robutti^a, S. Tosi^{a,b}

INFN Sezione di Milano-Bicocca ^a, Università di Milano-Bicocca ^b, Milano, Italy

L. Brianza^{a,b,15}, M.E. Dinardo^{a,b}, S. Fiorendi^{a,b,15}, S. Gennai^a, A. Ghezzi^{a,b}, P. Govoni^{a,b}, M. Malberti^{a,b}, S. Malvezzi^a, R.A. Manzoni^{a,b}, D. Menasce^a, L. Moroni^a, M. Paganoni^{a,b}, D. Pedrini^a, S. Pigazzini^{a,b}, S. Ragazzi^{a,b}, T. Tabarelli de Fatis^{a,b}

INFN Sezione di Napoli ^a, Università di Napoli 'Federico II' ^b, Napoli, Italy, Università della Basilicata ^c, Potenza, Italy, Università G. Marconi ^d, Roma, Italy

S. Buontempo^a, N. Cavallo^{a,c}, G. De Nardo, S. Di Guida^{a,d,15}, M. Esposito^{a,b}, F. Fabozzi^{a,c}, F. Fienga^{a,b}, A.O.M. Iorio^{a,b}, G. Lanza^a, L. Lista^a, S. Meola^{a,d,15}, P. Paolucci^{a,15}, C. Sciacca^{a,b}, F. Thyssen

INFN Sezione di Padova ^a, Università di Padova ^b, Padova, Italy, Università di Trento ^c, Trento, Italy

P. Azzi^{a,15}, N. Bacchetta^a, L. Benato^{a,b}, D. Bisello^{a,b}, A. Boletti^{a,b}, R. Carlin^{a,b}, P. Checchia^a, M. Dall'Osso^{a,b}, P. De Castro Manzano^a, T. Dorigo^a, U. Dosselli^a, F. Gasparini^{a,b}, U. Gasparini^{a,b}, A. Gozzelino^a, S. Lacaprara^a, M. Margoni^{a,b}, G. Maron^{a,30}, A.T. Meneguzzo^{a,b}, M. Michelotto^a, J. Pazzini^{a,b}, M. Pegoraro^a, N. Pozzobon^{a,b}, P. Ronchese^{a,b}, F. Simonetto^{a,b}, M. Zanetti, G. Zumerle^{a,b}

INFN Sezione di Pavia ^a, Università di Pavia ^b, Pavia, Italy

A. Braghieri^a, A. Magnani^{a,b}, P. Montagna^{a,b}, S.P. Ratti^{a,b}, V. Re^a, C. Riccardi^{a,b}, P. Salvini^a, I. Vai^{a,b}, P. Vitulo^{a,b}

INFN Sezione di Perugia ^a, Università di Perugia ^b, Perugia, Italy

L. Alunni Solestizi^{a,b}, G.M. Bilei^a, D. Ciangottini^{a,b}, L. Fanò^{a,b}, P. Lariccia^{a,b}, R. Leonardi^{a,b}, G. Mantovani^{a,b}, M. Menichelli^a, A. Saha^a, A. Santocchia^{a,b}

INFN Sezione di Pisa ^a, Università di Pisa ^b, Scuola Normale Superiore di Pisa ^c, Pisa, Italy

K. Androsov^{a,31}, P. Azzurri^{a,15}, G. Bagliesi^a, J. Bernardini^a, T. Boccali^a, R. Castaldi^a, M.A. Ciocci^{a,31}, R. Dell'Orso^a, S. Donato^{a,c}, G. Fedi, A. Giassi^a, M.T. Grippo^{a,31}, F. Ligabue^{a,c}, T. Lomtadze^a, L. Martini^{a,b}, A. Messineo^{a,b}, F. Palla^a, A. Rizzi^{a,b}, A. Savoy-Navarro^{a,32}, P. Spagnolo^a, R. Tenchini^a, G. Tonelli^{a,b}, A. Venturi^a, P.G. Verdini^a

INFN Sezione di Roma ^a, Università di Roma ^b, Roma, Italy

L. Barone^{a,b}, F. Cavallari^a, M. Cipriani^{a,b}, D. Del Re^{a,b,15}, M. Diemoz^a, S. Gelli^{a,b}, E. Longo^{a,b}, F. Margaroli^{a,b}, B. Marzocchi^{a,b}, P. Meridiani^a, G. Organtini^{a,b}, R. Paramatti^a, F. Preiato^{a,b}, S. Rahatlou^{a,b}, C. Rovelli^a, F. Santanastasio^{a,b}

INFN Sezione di Torino ^a, Università di Torino ^b, Torino, Italy, Università del Piemonte Orientale ^c, Novara, Italy

N. Amapane^{a,b}, R. Arcidiacono^{a,c,15}, S. Argiro^{a,b}, M. Arneodo^{a,c}, N. Bartosik^a, R. Bellan^{a,b}, C. Biino^a, N. Cartiglia^a, M. Costa^{a,b}, G. Cotto^{a,b}, R. Covarelli^{a,b}, P. De Remigis^a, A. Degano^{a,b}, N. Demaria^a, L. Finco^{a,b}, B. Kiani^{a,b}, C. Mariotti^a, S. Maselli^a, E. Migliore^{a,b}, V. Monaco^{a,b}, E. Monteil^{a,b}, M.M. Obertino^{a,b}, L. Pacher^{a,b}, N. Pastrone^a, M. Pelliccioni^a, G.L. Pinna

Angioni^{a,b}, F. Ravera^{a,b}, A. Romero^{a,b}, M. Ruspa^{a,c}, R. Sacchi^{a,b}, V. Sola^a, A. Solano^{a,b},
A. Staiano^a, P. Traczyk^{a,b}

INFN Sezione di Trieste ^a, Università di Trieste ^b, Trieste, Italy

S. Belforte^a, M. Casarsa^a, F. Cossutti^a, G. Della Ricca^{a,b}, A. Zanetti^a

Kyungpook National University, Daegu, Korea

D.H. Kim, G.N. Kim, M.S. Kim, S. Lee, S.W. Lee, Y.D. Oh, S. Sekmen, D.C. Son, Y.C. Yang

Chonbuk National University, Jeonju, Korea

A. Lee

Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea

H. Kim

Hanyang University, Seoul, Korea

J.A. Brochero Cifuentes, T.J. Kim

Korea University, Seoul, Korea

S. Cho, S. Choi, Y. Go, D. Gyun, S. Ha, B. Hong, Y. Jo, Y. Kim, B. Lee, K. Lee, K.S. Lee, S. Lee,
J. Lim, S.K. Park, Y. Roh

Seoul National University, Seoul, Korea

J. Almond, J. Kim, H. Lee, S.B. Oh, B.C. Radburn-Smith, S.h. Seo, U.K. Yang, H.D. Yoo, G.B. Yu

University of Seoul, Seoul, Korea

M. Choi, H. Kim, J.H. Kim, J.S.H. Lee, I.C. Park, G. Ryu, M.S. Ryu

Sungkyunkwan University, Suwon, Korea

Y. Choi, J. Goh, C. Hwang, J. Lee, I. Yu

Vilnius University, Vilnius, Lithuania

V. Dudenas, A. Juodagalvis, J. Vaitkus

National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia

I. Ahmed, Z.A. Ibrahim, J.R. Komaragiri, M.A.B. Md Ali³³, F. Mohamad Idris³⁴, W.A.T. Wan
Abdullah, M.N. Yusli, Z. Zolkapli

Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

H. Castilla-Valdez, E. De La Cruz-Burelo, I. Heredia-De La Cruz³⁵, A. Hernandez-Almada,
R. Lopez-Fernandez, R. Magaña Villalba, J. Mejia Guisao, A. Sanchez-Hernandez

Universidad Iberoamericana, Mexico City, Mexico

S. Carrillo Moreno, C. Oropeza Barrera, F. Vazquez Valencia

Benemerita Universidad Autonoma de Puebla, Puebla, Mexico

S. Carpinteyro, I. Pedraza, H.A. Salazar Ibarguen, C. Uribe Estrada

Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico

A. Morelos Pineda

University of Auckland, Auckland, New Zealand

D. Krofcheck

University of Canterbury, Christchurch, New Zealand

P.H. Butler

National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

A. Ahmad, M. Ahmad, Q. Hassan, H.R. Hoorani, W.A. Khan, A. Saddique, M.A. Shah, M. Shoaib, M. Waqas

National Centre for Nuclear Research, Swierk, Poland

H. Bialkowska, M. Bluj, B. Boimska, T. Frueboes, M. Górski, M. Kazana, K. Nawrocki, K. Romanowska-Rybinska, M. Szleper, P. Zalewski

Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland

K. Bunkowski, A. Byszuk³⁶, K. Doroba, A. Kalinowski, M. Konecki, J. Krolikowski, M. Misiura, M. Olszewski, M. Walczak

Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal

P. Bargassa, C. Beirão Da Cruz E Silva, A. Di Francesco, P. Faccioli, P.G. Ferreira Parracho, M. Gallinaro, J. Hollar, N. Leonardo, L. Lloret Iglesias, M.V. Nemallapudi, J. Rodrigues Antunes, J. Seixas, O. Toldaiev, D. Vadrucio, J. Varela, P. Vischia

Joint Institute for Nuclear Research, Dubna, Russia

S. Afanasiev, P. Bunin, M. Gavrilenko, I. Golutvin, I. Gorbunov, A. Kamenev, V. Karjavin, A. Lanev, A. Malakhov, V. Matveev^{37,38}, V. Palichik, V. Perelygin, S. Shmatov, S. Shulha, N. Skatchkov, V. Smirnov, N. Voytishin, A. Zarubin

Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia

L. Chtchypounov, V. Golovtsov, Y. Ivanov, V. Kim³⁹, E. Kuznetsova⁴⁰, V. Murzin, V. Oreshkin, V. Sulimov, A. Vorobyev

Institute for Nuclear Research, Moscow, Russia

Yu. Andreev, A. Dermenev, S. Gninenko, N. Golubev, A. Karneyev, M. Kirsanov, N. Krasnikov, A. Pashenkov, D. Tlisov, A. Toropin

Institute for Theoretical and Experimental Physics, Moscow, Russia

V. Epshteyn, V. Gavrilov, N. Lychkovskaya, V. Popov, I. Pozdnyakov, G. Safronov, A. Spiridonov, M. Toms, E. Vlasov, A. Zhokin

Moscow Institute of Physics and Technology

A. Bylinkin³⁸

National Research Nuclear University 'Moscow Engineering Physics Institute' (MEPhI), Moscow, Russia

R. Chistov⁴¹, M. Danilov⁴¹, V. Rusinov

P.N. Lebedev Physical Institute, Moscow, Russia

V. Andreev, M. Azarkin³⁸, I. Dremin³⁸, M. Kirakosyan, A. Leonidov³⁸, S.V. Rusakov, A. Terkulov

Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia

A. Baskakov, A. Belyaev, E. Boos, M. Dubinin⁴², L. Dudko, A. Ershov, A. Gribushin, V. Klyukhin, O. Kodolova, I. Lokhtin, I. Miagkov, S. Obraztsov, S. Petrushanko, V. Savrin, A. Snigirev

Novosibirsk State University (NSU), Novosibirsk, Russia

V. Blinov⁴³, Y.Skovpen⁴³, D. Shtol⁴³

State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia

I. Azhgirey, I. Bayshev, S. Bitioukov, D. Elumakhov, V. Kachanov, A. Kalinin, D. Konstantinov, V. Krychkin, V. Petrov, R. Ryutin, A. Sobol, S. Troshin, N. Tyurin, A. Uzunian, A. Volkov

University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia

P. Adzic⁴⁴, P. Cirkovic, D. Devetak, M. Dordevic, J. Milosevic, V. Rekovic

Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain

J. Alcaraz Maestre, M. Barrio Luna, E. Calvo, M. Cerrada, M. Chamizo Llatas, N. Colino, B. De La Cruz, A. Delgado Peris, A. Escalante Del Valle, C. Fernandez Bedoya, J.P. Fernández Ramos, J. Flix, M.C. Fouz, P. Garcia-Abia, O. Gonzalez Lopez, S. Goy Lopez, J.M. Hernandez, M.I. Josa, E. Navarro De Martino, A. Pérez-Calero Yzquierdo, J. Puerta Pelayo, A. Quintario Olmeda, I. Redondo, L. Romero, M.S. Soares

Universidad Autónoma de Madrid, Madrid, Spain

J.F. de Trocóniz, M. Missiroli, D. Moran

Universidad de Oviedo, Oviedo, Spain

J. Cuevas, J. Fernandez Menendez, I. Gonzalez Caballero, J.R. González Fernández, E. Palencia Cortezon, S. Sanchez Cruz, I. Suárez Andrés, J.M. Vizán Garcia

Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain

I.J. Cabrillo, A. Calderon, J.R. Castiñeiras De Saa, E. Curras, M. Fernandez, J. Garcia-Ferrero, G. Gomez, A. Lopez Virto, J. Marco, C. Martinez Rivero, F. Matorras, J. Piedra Gomez, T. Rodrigo, A. Ruiz-Jimeno, L. Scodellaro, N. Trevisani, I. Vila, R. Vilar Cortabitarte

CERN, European Organization for Nuclear Research, Geneva, Switzerland

D. Abbaneo, E. Auffray, G. Auzinger, M. Bachtis, P. Baillon, A.H. Ball, D. Barney, P. Bloch, A. Bocci, A. Bonato, C. Botta, T. Camporesi, R. Castello, M. Cepeda, G. Cerminara, M. D'Alfonso, D. d'Enterria, A. Dabrowski, V. Daponte, A. David, M. De Gruttola, A. De Roeck, E. Di Marco⁴⁵, M. Dobson, B. Dorney, T. du Pree, D. Duggan, M. Dünser, N. Dupont, A. Elliott-Peisert, S. Fartoukh, G. Franzoni, J. Fulcher, W. Funk, D. Gigi, K. Gill, M. Girone, F. Glege, D. Gulhan, S. Gundacker, M. Guthoff, J. Hammer, P. Harris, J. Hegeman, V. Innocente, P. Janot, J. Kieseler, H. Kirschenmann, V. Knünz, A. Kornmayer¹⁵, M.J. Kortelainen, K. Kousouris, M. Krammer¹, C. Lange, P. Lecoq, C. Lourenço, M.T. Lucchini, L. Malgeri, M. Mannelli, A. Martelli, F. Meijers, J.A. Merlin, S. Mersi, E. Meschi, P. Milenovic⁴⁶, F. Moortgat, S. Morovic, M. Mulders, H. Neugebauer, S. Orfanelli, L. Orsini, L. Pape, E. Perez, M. Peruzzi, A. Petrilli, G. Petrucciani, A. Pfeiffer, M. Pierini, A. Racz, T. Reis, G. Rolandi⁴⁷, M. Rovere, M. Ruan, H. Sakulin, J.B. Sauvan, C. Schäfer, C. Schwick, M. Seidel, A. Sharma, P. Silva, P. Sphicas⁴⁸, J. Steggemann, M. Stoye, Y. Takahashi, M. Tosi, D. Treille, A. Triossi, A. Tsirou, V. Veckalns⁴⁹, G.I. Veres²⁰, N. Wardle, H.K. Wöhri, A. Zagozdinska³⁶, W.D. Zeuner

Paul Scherrer Institut, Villigen, Switzerland

W. Bertl, K. Deiters, W. Erdmann, R. Horisberger, Q. Ingram, H.C. Kaestli, D. Kotlinski, U. Langenegger, T. Rohe

Institute for Particle Physics, ETH Zurich, Zurich, Switzerland

F. Bachmair, L. Bäni, L. Bianchini, B. Casal, G. Dissertori, M. Dittmar, M. Donegà, C. Grab, C. Heidegger, D. Hits, J. Hoss, G. Kasieczka, P. Lecomte[†], W. Lustermann, B. Mangano, M. Marionneau, P. Martinez Ruiz del Arbol, M. Masciovecchio, M.T. Meinhard, D. Meister,

F. Micheli, P. Musella, F. Nessi-Tedaldi, F. Pandolfi, J. Pata, F. Pauss, G. Perrin, L. Perrozzi, M. Quitnat, M. Rossini, M. Schönenberger, A. Starodumov⁵⁰, V.R. Tavolaro, K. Theofilatos, R. Wallny

Universität Zürich, Zurich, Switzerland

T.K. Aarrestad, C. AMSler⁵¹, L. Caminada, M.F. Canelli, A. De Cosa, C. Galloni, A. Hinzmann, T. Hreus, B. Kilminster, J. Ngadiuba, D. Pinna, G. Rauco, P. Robmann, D. Salerno, Y. Yang, A. Zucchetta

National Central University, Chung-Li, Taiwan

V. Candelise, T.H. Doan, Sh. Jain, R. Khurana, M. Konyushikhin, C.M. Kuo, W. Lin, Y.J. Lu, A. Pozdnyakov, S.S. Yu

National Taiwan University (NTU), Taipei, Taiwan

Arun Kumar, P. Chang, Y.H. Chang, Y.W. Chang, Y. Chao, K.F. Chen, P.H. Chen, C. Dietz, F. Fiori, W.-S. Hou, Y. Hsiung, Y.F. Liu, R.-S. Lu, M. Miñano Moya, E. Paganis, A. Psallidas, J.f. Tsai, Y.M. Tzeng

Chulalongkorn University, Faculty of Science, Department of Physics, Bangkok, Thailand

B. Asavapibhop, G. Singh, N. Srimanobhas, N. Suwonjandee

Cukurova University, Adana, Turkey

A. Adiguzel, S. Damarseckin, Z.S. Demiroglu, C. Dozen, E. Eskut, S. Girgis, G. Gokbulut, Y. Guler, I. Hos, E.E. Kangal⁵², O. Kara, A. Kayis Topaksu, U. Kiminsu, M. Oglakci, G. Onengut⁵³, K. Ozdemir⁵⁴, S. Ozturk⁵⁵, A. Polatoz, B. Tali⁵⁶, S. Turkcapar, I.S. Zorbakir, C. Zorbilmez

Middle East Technical University, Physics Department, Ankara, Turkey

B. Bilin, S. Bilmis, B. Isildak⁵⁷, G. Karapinar⁵⁸, M. Yalvac, M. Zeyrek

Bogazici University, Istanbul, Turkey

E. Gülmez, M. Kaya⁵⁹, O. Kaya⁶⁰, E.A. Yetkin⁶¹, T. Yetkin⁶²

Istanbul Technical University, Istanbul, Turkey

A. Cakir, K. Cankocak, S. Sen⁶³

Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkov, Ukraine

B. Grynyov

National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine

L. Levchuk, P. Sorokin

University of Bristol, Bristol, United Kingdom

R. Aggleton, F. Ball, L. Beck, J.J. Brooke, D. Burns, E. Clement, D. Cussans, H. Flacher, J. Goldstein, M. Grimes, G.P. Heath, H.F. Heath, J. Jacob, L. Kreczko, C. Lucas, D.M. Newbold⁶⁴, S. Paramesvaran, A. Poll, T. Sakuma, S. Seif El Nasr-storey, D. Smith, V.J. Smith

Rutherford Appleton Laboratory, Didcot, United Kingdom

K.W. Bell, A. Belyaev⁶⁵, C. Brew, R.M. Brown, L. Calligaris, D. Cieri, D.J.A. Cockerill, J.A. Coughlan, K. Harder, S. Harper, E. Olaiya, D. Petyt, C.H. Shepherd-Themistocleous, A. Thea, I.R. Tomalin, T. Williams

Imperial College, London, United Kingdom

M. Baber, R. Bainbridge, O. Buchmuller, A. Bundock, D. Burton, S. Casasso, M. Citron, D. Colling, L. Corpe, P. Dauncey, G. Davies, A. De Wit, M. Della Negra, R. Di Maria, P. Dunne,

A. Elwood, D. Futyan, Y. Haddad, G. Hall, G. Iles, T. James, R. Lane, C. Laner, R. Lucas⁶⁴, L. Lyons, A.-M. Magnan, S. Malik, L. Mastrolorenzo, J. Nash, A. Nikitenko⁵⁰, J. Pela, B. Penning, M. Pesaresi, D.M. Raymond, A. Richards, A. Rose, C. Seez, S. Summers, A. Tapper, K. Uchida, M. Vazquez Acosta⁶⁶, T. Virdee¹⁵, J. Wright, S.C. Zenz

Brunel University, Uxbridge, United Kingdom

J.E. Cole, P.R. Hobson, A. Khan, P. Kyberd, D. Leslie, I.D. Reid, P. Symonds, L. Teodorescu, M. Turner

Baylor University, Waco, USA

A. Borzou, K. Call, J. Dittmann, K. Hatakeyama, H. Liu, N. Pastika

The University of Alabama, Tuscaloosa, USA

O. Charaf, S.I. Cooper, C. Henderson, P. Rumerio, C. West

Boston University, Boston, USA

D. Arcaro, A. Avetisyan, T. Bose, D. Gastler, D. Rankin, C. Richardson, J. Rohlf, L. Sulak, D. Zou

Brown University, Providence, USA

G. Benelli, E. Berry, D. Cutts, A. Garabedian, J. Hakala, U. Heintz, J.M. Hogan, O. Jesus, K.H.M. Kwok, E. Laird, G. Landsberg, Z. Mao, M. Narain, S. Piperov, S. Sagir, E. Spencer, R. Syarif

University of California, Davis, Davis, USA

R. Breedon, G. Breto, D. Burns, M. Calderon De La Barca Sanchez, S. Chauhan, M. Chertok, J. Conway, R. Conway, P.T. Cox, R. Erbacher, C. Flores, G. Funk, M. Gardner, W. Ko, R. Lander, C. Mclean, M. Mulhearn, D. Pellett, J. Pilot, S. Shalhout, J. Smith, M. Squires, D. Stolp, M. Tripathi, S. Wilbur, R. Yohay

University of California, Los Angeles, USA

C. Bravo, R. Cousins, A. Dasgupta, P. Everaerts, A. Florent, J. Hauser, M. Ignatenko, N. Mccoll, D. Saltzberg, C. Schnaible, E. Takasugi, V. Valuev, M. Weber

University of California, Riverside, Riverside, USA

K. Burt, R. Clare, J. Ellison, J.W. Gary, S.M.A. Ghiasi Shirazi, G. Hanson, J. Heilman, P. Jandir, E. Kennedy, F. Lacroix, O.R. Long, M. Olmedo Negrete, M.I. Paneva, A. Shrinivas, W. Si, H. Wei, S. Wimpenny, B. R. Yates

University of California, San Diego, La Jolla, USA

J.G. Branson, G.B. Cerati, S. Cittolin, M. Derdzinski, R. Gerosa, A. Holzner, D. Klein, V. Krutelyov, J. Letts, I. Macneill, D. Olivito, S. Padhi, M. Pieri, M. Sani, V. Sharma, S. Simon, M. Tadel, A. Vartak, S. Wasserbaech⁶⁷, C. Welke, J. Wood, F. Würthwein, A. Yagil, G. Zevi Della Porta

University of California, Santa Barbara - Department of Physics, Santa Barbara, USA

N. Amin, R. Bhandari, J. Bradmiller-Feld, C. Campagnari, A. Dishaw, V. Dutta, K. Flowers, M. Franco Sevilla, P. Geffert, C. George, F. Golf, L. Gouskos, J. Gran, R. Heller, J. Incandela, S.D. Mullin, A. Ovcharova, J. Richman, D. Stuart, I. Suarez, J. Yoo

California Institute of Technology, Pasadena, USA

D. Anderson, A. Apresyan, J. Bendavid, A. Bornheim, J. Bunn, Y. Chen, J. Duarte, J.M. Lawhorn, A. Mott, H.B. Newman, C. Pena, M. Spiropulu, J.R. Vlimant, S. Xie, R.Y. Zhu

Carnegie Mellon University, Pittsburgh, USA

M.B. Andrews, V. Azzolini, T. Ferguson, M. Paulini, J. Russ, M. Sun, H. Vogel, I. Vorobiev, M. Weinberg

University of Colorado Boulder, Boulder, USA

J.P. Cumalat, W.T. Ford, F. Jensen, A. Johnson, M. Krohn, T. Mulholland, K. Stenson, S.R. Wagner

Cornell University, Ithaca, USA

J. Alexander, J. Chaves, J. Chu, S. Dittmer, K. Mcdermott, N. Mirman, G. Nicolas Kaufman, J.R. Patterson, A. Rinkevicius, A. Ryd, L. Skinnari, L. Soffi, S.M. Tan, Z. Tao, J. Thom, J. Tucker, P. Wittich, M. Zientek

Fairfield University, Fairfield, USA

D. Winn

Fermi National Accelerator Laboratory, Batavia, USA

S. Abdullin, M. Albrow, G. Apollinari, S. Banerjee, L.A.T. Bauerdick, A. Beretvas, J. Berryhill, P.C. Bhat, G. Bolla, K. Burkett, J.N. Butler, H.W.K. Cheung, F. Chlebana, S. Cihangir[†], M. Cremonesi, V.D. Elvira, I. Fisk, J. Freeman, E. Gottschalk, L. Gray, D. Green, S. Grünendahl, O. Gutsche, D. Hare, R.M. Harris, S. Hasegawa, J. Hirschauer, Z. Hu, B. Jayatilaka, S. Jindariani, M. Johnson, U. Joshi, B. Klima, B. Kreis, S. Lammel, J. Linacre, D. Lincoln, R. Lipton, T. Liu, R. Lopes De Sá, J. Lykken, K. Maeshima, N. Magini, J.M. Marraffino, S. Maruyama, D. Mason, P. McBride, P. Merkel, S. Mrenna, S. Nahn, C. Newman-Holmes[†], V. O'Dell, K. Pedro, O. Prokofyev, G. Rakness, L. Ristori, E. Sexton-Kennedy, A. Soha, W.J. Spalding, L. Spiegel, S. Stoynev, N. Strobbe, L. Taylor, S. Tkaczyk, N.V. Tran, L. Uplegger, E.W. Vaandering, C. Vernieri, M. Verzocchi, R. Vidal, M. Wang, H.A. Weber, A. Whitbeck

University of Florida, Gainesville, USA

D. Acosta, P. Avery, P. Bortignon, D. Bourilkov, A. Brinkerhoff, A. Carnes, M. Carver, D. Curry, S. Das, R.D. Field, I.K. Furic, J. Konigsberg, A. Korytov, J.F. Low, P. Ma, K. Matchev, H. Mei, G. Mitselmakher, D. Rank, L. Shchutska, D. Sperka, L. Thomas, J. Wang, S. Wang, J. Yelton

Florida International University, Miami, USA

S. Linn, P. Markowitz, G. Martinez, J.L. Rodriguez

Florida State University, Tallahassee, USA

A. Ackert, J.R. Adams, T. Adams, A. Askew, S. Bein, B. Diamond, S. Hagopian, V. Hagopian, K.F. Johnson, A. Khatiwada, H. Prosper, A. Santra

Florida Institute of Technology, Melbourne, USA

M.M. Baarmand, V. Bhopatkar, S. Colafranceschi⁶⁸, M. Hohlmann, D. Noonan, T. Roy, F. Yumiceva

University of Illinois at Chicago (UIC), Chicago, USA

M.R. Adams, L. Apanasevich, D. Berry, R.R. Betts, I. Bucinskaite, R. Cavanaugh, O. Evdokimov, L. Gauthier, C.E. Gerber, D.J. Hofman, K. Jung, P. Kurt, C. O'Brien, I.D. Sandoval Gonzalez, P. Turner, N. Varelas, H. Wang, Z. Wu, M. Zakaria, J. Zhang

The University of Iowa, Iowa City, USA

B. Bilki⁶⁹, W. Clarida, K. Dilsiz, S. Durgut, R.P. Gandrajula, M. Haytmyradov, V. Khristenko, J.-P. Merlo, H. Mermerkaya⁷⁰, A. Mestvirishvili, A. Moeller, J. Nachtman, H. Ogul, Y. Onel, F. Ozok⁷¹, A. Penzo, C. Snyder, E. Tiras, J. Wetzel, K. Yi

Johns Hopkins University, Baltimore, USA

I. Anderson, B. Blumenfeld, A. Cocoros, N. Eminizer, D. Fehling, L. Feng, A.V. Gritsan, P. Maksimovic, C. Martin, M. Osherson, J. Roskes, U. Sarica, M. Swartz, M. Xiao, Y. Xin, C. You

The University of Kansas, Lawrence, USA

A. Al-bataineh, P. Baringer, A. Bean, S. Boren, J. Bowen, C. Bruner, J. Castle, L. Forthomme, R.P. Kenny III, A. Kropivnitskaya, D. Majumder, W. Mcbrayer, M. Murray, S. Sanders, R. Stringer, J.D. Tapia Takaki, Q. Wang

Kansas State University, Manhattan, USA

A. Ivanov, K. Kaadze, S. Khalil, Y. Maravin, A. Mohammadi, L.K. Saini, N. Skhirtladze, S. Toda

Lawrence Livermore National Laboratory, Livermore, USA

F. Rebassoo, D. Wright

University of Maryland, College Park, USA

C. Anelli, A. Baden, O. Baron, A. Belloni, B. Calvert, S.C. Eno, C. Ferraioli, J.A. Gomez, N.J. Hadley, S. Jabeen, R.G. Kellogg, T. Kolberg, J. Kunkle, Y. Lu, A.C. Mignerey, F. Ricci-Tam, Y.H. Shin, A. Skuja, M.B. Tonjes, S.C. Tonwar

Massachusetts Institute of Technology, Cambridge, USA

D. Abercrombie, B. Allen, A. Apyan, R. Barbieri, A. Baty, R. Bi, K. Bierwagen, S. Brandt, W. Busza, I.A. Cali, Z. Demiragli, L. Di Matteo, G. Gomez Ceballos, M. Goncharov, D. Hsu, Y. Iiyama, G.M. Innocenti, M. Klute, D. Kovalskyi, K. Krajczar, Y.S. Lai, Y.-J. Lee, A. Levin, P.D. Luckey, B. Maier, A.C. Marini, C. McGinn, C. Mironov, S. Narayanan, X. Niu, C. Paus, C. Roland, G. Roland, J. Salfeld-Nebgen, G.S.F. Stephans, K. Sumorok, K. Tatar, M. Varma, D. Velicanu, J. Veverka, J. Wang, T.W. Wang, B. Wyslouch, M. Yang, V. Zhukova

University of Minnesota, Minneapolis, USA

A.C. Benvenuti, R.M. Chatterjee, A. Evans, A. Finkel, A. Gude, P. Hansen, S. Kalafut, S.C. Kao, Y. Kubota, Z. Lesko, J. Mans, S. Nourbakhsh, N. Ruckstuhl, R. Rusack, N. Tambe, J. Turkewitz

University of Mississippi, Oxford, USA

J.G. Acosta, S. Oliveros

University of Nebraska-Lincoln, Lincoln, USA

E. Avdeeva, R. Bartek, K. Bloom, D.R. Claes, A. Dominguez⁷², C. Fangmeier, R. Gonzalez Suarez, R. Kamalieddin, I. Kravchenko, A. Malta Rodrigues, F. Meier, J. Monroy, J.E. Siado, G.R. Snow, B. Stieger

State University of New York at Buffalo, Buffalo, USA

M. Alyari, J. Dolen, J. George, A. Godshalk, C. Harrington, I. Iashvili, J. Kaisen, A. Kharchilava, A. Kumar, A. Parker, S. Rappoccio, B. Roozbahani

Northeastern University, Boston, USA

G. Alverson, E. Barberis, A. Hortiangtham, A. Massironi, D.M. Morse, D. Nash, T. Orimoto, R. Teixeira De Lima, D. Trocino, R.-J. Wang, D. Wood

Northwestern University, Evanston, USA

S. Bhattacharya, K.A. Hahn, A. Kubik, A. Kumar, N. Mucia, N. Odell, B. Pollack, M.H. Schmitt, K. Sung, M. Trovato, M. Velasco

University of Notre Dame, Notre Dame, USA

N. Dev, M. Hildreth, K. Hurtado Anampa, C. Jessop, D.J. Karmgard, N. Kellams, K. Lannon,

N. Marinelli, F. Meng, C. Mueller, Y. Musienko³⁷, M. Planer, A. Reinsvold, R. Ruchti, G. Smith, S. Taroni, M. Wayne, M. Wolf, A. Woodard

The Ohio State University, Columbus, USA

J. Alimena, L. Antonelli, J. Brinson, B. Bylsma, L.S. Durkin, S. Flowers, B. Francis, A. Hart, C. Hill, R. Hughes, W. Ji, B. Liu, W. Luo, D. Puigh, B.L. Winer, H.W. Wulsin

Princeton University, Princeton, USA

S. Cooperstein, O. Driga, P. Elmer, J. Hardenbrook, P. Hebda, D. Lange, J. Luo, D. Marlow, J. McDonald, T. Medvedeva, K. Mei, M. Mooney, J. Olsen, C. Palmer, P. Piroué, D. Stickland, C. Tully, A. Zuranski

University of Puerto Rico, Mayaguez, USA

S. Malik

Purdue University, West Lafayette, USA

A. Barker, V.E. Barnes, S. Folgueras, L. Gutay, M.K. Jha, M. Jones, A.W. Jung, D.H. Miller, N. Neumeister, J.F. Schulte, X. Shi, J. Sun, A. Svyatkovskiy, F. Wang, W. Xie, L. Xu

Purdue University Calumet, Hammond, USA

N. Parashar, J. Stupak

Rice University, Houston, USA

A. Adair, B. Akgun, Z. Chen, K.M. Ecklund, F.J.M. Geurts, M. Guilbaud, W. Li, B. Michlin, M. Northup, B.P. Padley, R. Redjimi, J. Roberts, J. Rorie, Z. Tu, J. Zabel

University of Rochester, Rochester, USA

B. Betchart, A. Bodek, P. de Barbaro, R. Demina, Y.t. Duh, T. Ferbel, M. Galanti, A. Garcia-Bellido, J. Han, O. Hindrichs, A. Khukhunaishvili, K.H. Lo, P. Tan, M. Verzetti

Rutgers, The State University of New Jersey, Piscataway, USA

A. Agapitos, J.P. Chou, E. Contreras-Campana, Y. Gershtein, T.A. Gómez Espinosa, E. Halkiadakis, M. Heindl, D. Hidas, E. Hughes, S. Kaplan, R. Kunnawalkam Elayavalli, S. Kyriacou, A. Lath, K. Nash, H. Saka, S. Salur, S. Schnetzer, D. Sheffield, S. Somalwar, R. Stone, S. Thomas, P. Thomassen, M. Walker

University of Tennessee, Knoxville, USA

A.G. Delannoy, M. Foerster, J. Heideman, G. Riley, K. Rose, S. Spanier, K. Thapa

Texas A&M University, College Station, USA

O. Bouhali⁷³, A. Celik, M. Dalchenko, M. De Mattia, A. Delgado, S. Dildick, R. Eusebi, J. Gilmore, T. Huang, E. Juska, T. Kamon⁷⁴, R. Mueller, Y. Pakhotin, R. Patel, A. Perloff, L. Perniè, D. Rathjens, A. Rose, A. Safonov, A. Tatarinov, K.A. Ulmer

Texas Tech University, Lubbock, USA

N. Akchurin, C. Cowden, J. Damgov, F. De Guio, C. Dragoiu, P.R. Duderod, J. Faulkner, E. Gурpinar, S. Kunori, K. Lamichhane, S.W. Lee, T. Libeiro, T. Peltola, S. Undleeb, I. Volobouev, Z. Wang

Vanderbilt University, Nashville, USA

S. Greene, A. Gurrola, R. Janjam, W. Johns, C. Maguire, A. Melo, H. Ni, P. Sheldon, S. Tuo, J. Velkovska, Q. Xu

University of Virginia, Charlottesville, USA

M.W. Arenton, P. Barria, B. Cox, J. Goodell, R. Hirosky, A. Ledovskoy, H. Li, C. Neu, T. Sinthuprasith, X. Sun, Y. Wang, E. Wolfe, F. Xia

Wayne State University, Detroit, USA

C. Clarke, R. Harr, P.E. Karchin, J. Sturdy

University of Wisconsin - Madison, Madison, WI, USA

D.A. Belknap, C. Caillol, S. Dasu, L. Dodd, S. Duric, B. Gomber, M. Grothe, M. Herndon, A. Hervé, P. Klabbers, A. Lanaro, A. Levine, K. Long, R. Loveless, I. Ojalvo, T. Perry, G.A. Pierro, G. Polese, T. Ruggles, A. Savin, N. Smith, W.H. Smith, D. Taylor, N. Woods

†: Deceased

1: Also at Vienna University of Technology, Vienna, Austria

2: Also at State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China

3: Also at Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France

4: Also at Universidade Estadual de Campinas, Campinas, Brazil

5: Also at Universidade Federal de Pelotas, Pelotas, Brazil

6: Also at Université Libre de Bruxelles, Bruxelles, Belgium

7: Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany

8: Also at Joint Institute for Nuclear Research, Dubna, Russia

9: Also at Cairo University, Cairo, Egypt

10: Also at Fayoum University, El-Fayoum, Egypt

11: Now at British University in Egypt, Cairo, Egypt

12: Now at Ain Shams University, Cairo, Egypt

13: Also at Université de Haute Alsace, Mulhouse, France

14: Also at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia

15: Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland

16: Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

17: Also at University of Hamburg, Hamburg, Germany

18: Also at Brandenburg University of Technology, Cottbus, Germany

19: Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary

20: Also at MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary

21: Also at University of Debrecen, Debrecen, Hungary

22: Also at Indian Institute of Science Education and Research, Bhopal, India

23: Also at Institute of Physics, Bhubaneswar, India

24: Also at University of Visva-Bharati, Santiniketan, India

25: Also at University of Ruhuna, Matara, Sri Lanka

26: Also at Isfahan University of Technology, Isfahan, Iran

27: Also at University of Tehran, Department of Engineering Science, Tehran, Iran

28: Also at Yazd University, Yazd, Iran

29: Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran

30: Also at Laboratori Nazionali di Legnaro dell'INFN, Legnaro, Italy

31: Also at Università degli Studi di Siena, Siena, Italy

32: Also at Purdue University, West Lafayette, USA

33: Also at International Islamic University of Malaysia, Kuala Lumpur, Malaysia

34: Also at Malaysian Nuclear Agency, MOSTI, Kajang, Malaysia

35: Also at Consejo Nacional de Ciencia y Tecnología, Mexico city, Mexico

36: Also at Warsaw University of Technology, Institute of Electronic Systems, Warsaw, Poland

- 37: Also at Institute for Nuclear Research, Moscow, Russia
- 38: Now at National Research Nuclear University 'Moscow Engineering Physics Institute' (MEPhI), Moscow, Russia
- 39: Also at St. Petersburg State Polytechnical University, St. Petersburg, Russia
- 40: Also at University of Florida, Gainesville, USA
- 41: Also at P.N. Lebedev Physical Institute, Moscow, Russia
- 42: Also at California Institute of Technology, Pasadena, USA
- 43: Also at Budker Institute of Nuclear Physics, Novosibirsk, Russia
- 44: Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia
- 45: Also at INFN Sezione di Roma; Università di Roma, Roma, Italy
- 46: Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia
- 47: Also at Scuola Normale e Sezione dell'INFN, Pisa, Italy
- 48: Also at National and Kapodistrian University of Athens, Athens, Greece
- 49: Also at Riga Technical University, Riga, Latvia
- 50: Also at Institute for Theoretical and Experimental Physics, Moscow, Russia
- 51: Also at Albert Einstein Center for Fundamental Physics, Bern, Switzerland
- 52: Also at Mersin University, Mersin, Turkey
- 53: Also at Cag University, Mersin, Turkey
- 54: Also at Piri Reis University, Istanbul, Turkey
- 55: Also at Gaziosmanpasa University, Tokat, Turkey
- 56: Also at Adiyaman University, Adiyaman, Turkey
- 57: Also at Ozyegin University, Istanbul, Turkey
- 58: Also at Izmir Institute of Technology, Izmir, Turkey
- 59: Also at Marmara University, Istanbul, Turkey
- 60: Also at Kafkas University, Kars, Turkey
- 61: Also at Istanbul Bilgi University, Istanbul, Turkey
- 62: Also at Yildiz Technical University, Istanbul, Turkey
- 63: Also at Hacettepe University, Ankara, Turkey
- 64: Also at Rutherford Appleton Laboratory, Didcot, United Kingdom
- 65: Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom
- 66: Also at Instituto de Astrofísica de Canarias, La Laguna, Spain
- 67: Also at Utah Valley University, Orem, USA
- 68: Also at Facoltà Ingegneria, Università di Roma, Roma, Italy
- 69: Also at Argonne National Laboratory, Argonne, USA
- 70: Also at Erzincan University, Erzincan, Turkey
- 71: Also at Mimar Sinan University, Istanbul, Istanbul, Turkey
- 72: Now at The Catholic University of America, Washington, USA
- 73: Also at Texas A&M University at Qatar, Doha, Qatar
- 74: Also at Kyungpook National University, Daegu, Korea