

## Direct measurement of astrophysically important resonances in $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$

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**Background:** Classical novae are cataclysmic nuclear explosions occurring when a white dwarf in a binary system accretes hydrogen-rich material from its companion star. Novae are partially responsible for the galactic synthesis of a variety of nuclides up to the calcium ( $A \sim 40$ ) region of the nuclear chart. Although the structure and dynamics of novae are thought to be relatively well understood, the predicted abundances of elements near the nucleosynthesis endpoint, in particular Ar and Ca, appear to sometimes be in disagreement with astronomical observations of the spectra of nova ejecta.

**Purpose:** One possible source of the discrepancies between model predictions and astronomical observations is nuclear reaction data. Most reaction rates near the nova endpoint are estimated only from statistical model calculations, which carry large uncertainties. For certain key reactions, these rate uncertainties translate into large uncertainties in nucleosynthesis predictions. In particular, the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reaction has been identified as having a significant influence on Ar, K, and Ca production. In order to constrain the rate of this reaction, we have performed a direct measurement of the strengths of three candidate  $\ell = 0$  resonances within the Gamow window for nova burning, at  $386 \pm 10$  keV,  $515 \pm 10$  keV, and  $689 \pm 10$  keV.

**Method:** The experiment was performed in inverse kinematics using a beam of unstable  $^{38}\text{K}$  impinging on a windowless hydrogen gas target. The  $^{39}\text{Ca}$  recoils and prompt  $\gamma$  rays from  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reactions were detected in coincidence using a recoil mass separator and a bismuth-germanate scintillator array, respectively.

**Results:** For the 689 keV resonance, we observed a clear recoil- $\gamma$  coincidence signal and extracted resonance strength and energy values of  $120_{-30}^{+50}(\text{stat.})_{-60}^{+20}(\text{sys.})$  meV and  $679_{-1}^{+2}(\text{stat.}) \pm 1(\text{sys.})$  keV, respectively. We also performed a singles analysis of the recoil data alone, extracting a resonance strength of  $120 \pm 20(\text{stat.}) \pm 15(\text{sys.})$  meV, consistent with the coincidence result. For the 386 keV and 515 keV resonances, we extract 90% confidence level upper limits of 2.54 meV and 18.4 meV, respectively.

**Conclusions:** We have established a new recommended  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate based on experimental information, which reduces overall uncertainties near the peak temperatures of nova burning by a factor of  $\sim 250$ . Using the rate obtained in this work in model calculations of the hottest oxygen-neon novae reduces overall uncertainties on Ar, K, and Ca synthesis to factors of 15 or less in all cases.

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## I. INTRODUCTION

Classical novae are some of the most common explosive stellar events to occur in our galaxy, with an estimated frequency of  $35 \pm 11$  per year [1]. Novae happen when a white dwarf in a binary system accretes hydrogen-rich material from its main-sequence companion, igniting thermonuclear runaway. Observations of the spectra of ejected material indicate that two main classes of nova exist, depending on the initial composition of the underlying white dwarf: carbon-oxygen (CO) and oxygen-neon (ONE). Model calculations indicate that ONE novae, which occur on more massive white dwarves, can reach peak temperatures around 0.4 GK and synthesize

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nuclei up to the calcium region ( $A \sim 40$ ). At present, there are a number of outstanding discrepancies between astronomical observations of the spectra of nova ejecta [2–5] and nova model predictions [6,7]. In particular, the model predictions of Ref. [7] indicate Ar and Ca abundances at roughly the solar level, while in contrast the observations of Ref. [3] point towards nova ejecta with Ar and Ca abundances around an order of magnitude greater than solar. Resolution of such discrepancies requires that nova models be capable of making detailed predictions regarding the synthesis of nuclides in the Ar–Ca region. In turn, this requires improved constraints on the rates of key nuclear reactions involved in nova nucleosynthesis, in particular for reactions near the nucleosynthesis endpoint.

In 2002, Iliadis *et al.* published a seminal paper investigating the influence of nuclear reaction rate variations on nucleosynthesis in classical novae [8]. In this study, the authors varied the rates of 64 nuclear reactions within their recommended uncertainties and examined the effect of these variations on the nucleosynthesis predictions of seven different nova models. For the hottest model included in the study, reaching a peak temperature of 0.418 GK, the authors identified the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reaction as having a significant influence on the production of Ar, K, and Ca. Qualitatively, the predicted abundances of these elements were found to vary by respective factors of 24, 58, and 57 when the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate was varied within its existing uncertainties. When Ref. [8] was published, the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate was estimated entirely from statistical model predictions with no experimental nuclear physics input [9]. This rate estimate was assigned an overall uncertainty of  $10^4$ , i.e., the upper and lower limits were established at 100 and 0.01 times the central value, respectively. The importance of this reaction for nova nucleosynthesis, along with the paucity of experimental input regarding the accepted rate, prompted an attempt by the present authors to measure the strengths of the three  $\ell = 0$  resonances lying within the Gamow window for ONe novae ( $T_{\text{peak}} \simeq 0.2\text{--}0.4$  GK). The first results of this study were published in a review article [10] and a recent Letter, which recommends a new, experimentally based rate with uncertainties over two orders of magnitude smaller than before [11]. In the present article, we expand upon Ref. [11], providing significantly more detail concerning the experiment and data analysis. We also report the results of a new sensitivity study investigating the effect of our measurement on the synthesis of Ar, K, and Ca in classical novae. The results presented here supersede those published previously.

## II. EXPERIMENT

The experiment was performed in the ISAC-I [12] hall at TRIUMF, Canada’s national laboratory for particle and nuclear physics. A beam of radioactive  $^{38}\text{K}$  was produced by impinging 500 MeV protons from the TRIUMF cyclotron onto a high-power TiC production target. The  $^{38}\text{K}^{(1+)}$  ions produced by spallation reactions in the target were extracted and sent through a high-resolution mass separator. They were then charge bred to the  $7^+$  charge state in an electron cyclotron resonance (ECR) charge state booster before post-acceleration. The charge breeding is necessary because the ISAC-I radio

frequency quadrupole (RFQ) is restricted to a mass-to-charge ratio of 30 or less [13].

The  $^{38}\text{K}^{(7+)}$  beam was delivered to the Detector of Recoils and Gammas of Nuclear Reactions (DRAGON) where it impinged on a windowless extended gas target [14], filled with  $\text{H}_2$  at an average pressure and temperature of 10.6 mbar and 298 K, respectively. The  $\text{H}_2$  was cleaned by continuous recirculation through a  $\text{LN}_2$  cooled zeolite trap. The prompt  $\gamma$  rays from  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reactions were detected in array of 30 bismuth germanate (BGO) scintillators surrounding the target, while the  $^{39}\text{Ca}$  recoils were transmitted to the focal plane of DRAGON, separating them from unreacted and elastically scattered  $^{38}\text{K}$ . A timing signature for recoils was established as the time difference between signals from a pair of microchannel plates (MCPs) separated by 59 cm, which detected secondary electrons produced by the interaction of the recoil ions with a diamond-like carbon foil intersecting the beam line. The total kinetic energy and stopping power of the recoil ions was measured in a multi-anode ionization chamber (IC) [15]. Coincidences between recoils and prompt  $\gamma$  rays were identified using a timestamp-based algorithm [16]. The  $^{39}\text{Ca}$  recoils were separated from a background of scattered and charge-changed  $^{38}\text{K}$  (“leaky beam”) based primarily on the local time of flight (TOF) between the two MCPs (“MCP TOF”) and the time difference between the  $\gamma$  ray and the upstream MCP (“separator TOF”).

Laboratory beam energies of 15.58 MeV, 20.56 MeV, and 27.17 MeV were employed for measurements of the  $386 \pm 10$  keV,  $515 \pm 10$  keV, and  $689 \pm 10$  keV resonances, respectively. The beam energies were measured using the procedure given in Ref. [17]. The beam was centered on 2 mm slits downstream of DRAGON’s first magnetic dipole, and the measured field value was converted to energy by solving the relativistically correct equation

$$E/A = c_{\text{mag}}(qB/A)^2 - \frac{1}{2uc^2}(E/A)^2, \quad (1)$$

where  $E$ ,  $A$ , and  $q$  are the beam kinetic energy, mass number, and charge state, respectively, and  $u$  is the atomic mass unit. The quantity  $c_{\text{mag}}$  is related to the effective bending radius of the dipole. The recommended value from Ref. [17],  $c_{\text{mag}} = 48.15 \pm 0.07$  MeV T<sup>2</sup>, was employed for this experiment. The estimated uncertainty on this procedure is 0.17%.

The chosen beam energies cover respective center-of-mass energies in the DRAGON gas target of  $386 \pm 13$ ,  $515 \pm 13$ , and  $689 \pm 13$  keV. The resonances in question were previously identified as  $5/2^+$   $^{39}\text{Ca}$  states through  $^{40}\text{Ca}(^3\text{He},\alpha)^{39}\text{Ca}$  [18],  $^{40}\text{Ca}(d,t)^{39}\text{Ca}$  [19], and  $^{40}\text{Ca}(p,d)^{39}\text{Ca}$  [20] transfer reaction studies. Their recommended excitation energies are  $6157 \pm 10$ ,  $6286 \pm 10$ , and  $6460 \pm 10$  keV, corresponding to  $^{38}\text{K} + p$  resonances at  $386 \pm 10$ ,  $515 \pm 10$ , and  $689 \pm 10$  keV, respectively [21]. The respective  $(p,\gamma)$  cone angles for measurements at the 15.58 MeV, 20.56 MeV, and 27.17 MeV beam energies were 5.98 mrad, 5.29 mrad, and 4.73 mrad. Each of these is well within the  $\pm 21$  mrad angular acceptance of DRAGON [22].

For each beam energy, only a single charge state was transmitted to the end of DRAGON. The respective charge states were  $7^+$ ,  $9^+$ , and  $10^+$  for the 15.58 MeV, 20.56 MeV,

and 27.17 MeV beam energies. The charge state fractions and stopping powers for K and Ca ions passing through the gas target were measured separately using stable beams of  $^{39}\text{K}$  and  $^{44}\text{Ca}$ . Charge state fractions were determined by measuring the ratio  $(I_2^g/I_0^g) \times (I_0^{ng}/I_1^{ng})$ , where  $I_0$ ,  $I_1$ , and  $I_2$  represent the current on Faraday cups upstream of the gas target, downstream of the gas target, and downstream of the first magnetic dipole, respectively; and the superscripts  $g$  and  $ng$  represent currents measured with and without gas in the target, respectively. Current measurements were taken with the magnetic dipole set to accept each of the charge states that resulted in a measurable  $I_2$ . The resulting distributions were then fit with a Gaussian function (normalized to unity). The value of the Gaussian at each charge state was taken to represent the corresponding charge state fraction. Measurements were taken at three different beam energies spanning the range of beam energies employed in the experiment, and the resulting charge fractions were fit with a quadratic function. The value of this quadratic function at the various beam energies employed in the experiment was then taken as the charge state fraction to use in the recoil yield analysis. All fits were performed using MIMUIT and errors on the fit parameters were calculated using MINOS [23]. The errors on the Gaussian fit were propagated along with the errors on the quadratic interpolation to arrive at the final error on the charge state fractions used in the analysis.

The number of incoming  $^{38}\text{K}^{(7+)}$  ions was determined by counting delayed ( $t_{1/2} = 7.6$  minutes) 2.2 MeV  $\gamma$  rays emitted by the daughters of beam ions implanted into the mass slits just downstream of DRAGON's first electric dipole. These  $\gamma$  rays were detected in a NaI scintillator with an efficiency of  $(8.46 \pm 0.95) \times 10^{-6}$ . This efficiency was determined from a GEANT4 [24] simulation, which included the entire geometry of the mass slit box and NaI detector. The 11% relative uncertainty on the NaI efficiency was determined by comparing simulation results to known  $^{22}\text{Na}$  and  $^{137}\text{Cs}$  source measurements. This analysis includes an uncertainty on the source position of  $\pm 0.5$  cm. The average beam rate for each  $\sim 1$  hour run was determined by fitting the decay rate vs time curves with the expected response function

$$A(t) = I(1 - e^{-\lambda t}) + N_0\lambda e^{-\lambda t}, \quad (2)$$

where  $A(t)$  is the decay rate,  $I$  is the average beam intensity,  $N_0$  is the initial number of particles implanted in the slit, and  $\lambda = 1.5 \times 10^{-3} \text{ s}^{-1}$  is the  $^{38}\text{K}$  decay constant. In the fit, both  $I$  and  $N_0$  were allowed to vary as free parameters. Cases where the average beam rate fluctuated significantly over the course of a run were identified by a noticeable deviation from the expected response. These fluctuations in the beam rate (or the complete loss of beam delivery) arose from a number of sources upstream of the DRAGON target, for example loss of the 500 MeV proton beam or Faraday cup readings taken by the ISAC-I operators. In these cases, differing sections of the run were visually identified and independently fit to Eq. (2). Figure 1(a) shows sample fitted rate vs time curves for two runs, one with a constant beam rate and the other with a varying beam rate (and corresponding piecewise fit). Figure 1(b) shows the average beam rates determined for each run throughout the course of the experiment. The overall  $^{38}\text{K}$  rate was approximately  $2 \times 10^7$  particles per second.

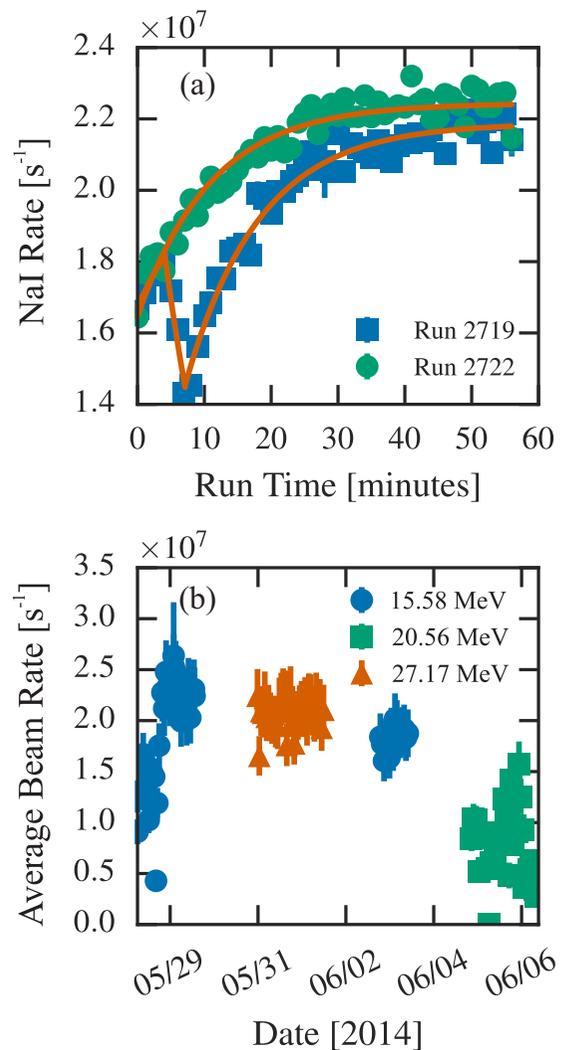


FIG. 1. (a) Sample rate vs time curves measured in the NaI detector. The green filled circles correspond to a run where the beam rate was constant, while the blue filled squares correspond to a run where there were significant changes in the rate. The solid orange lines show the fit results used to extract the average beam intensities. (b) Average beam rate determined for each  $\sim 1$  hour run taken during the experiment. The filled circles, rectangles, and triangles denote each run's beam energy, as indicated in the legend.

### A. 386(10) keV and 515(10) keV resonances

At beam energies of 15.58 MeV and 20.56 MeV (corresponding to the  $386 \pm 10$  keV and  $515 \pm 10$  keV resonances, respectively), we observed zero events in the expected recoil region. This is demonstrated in Fig. 2(a), which shows MCP vs separator TOF spectra for each of the 15.58 MeV and 20.56 MeV beam energies. The dashed and dotted ellipses included on the plots indicate the expected location of  $^{39}\text{Ca}$  recoils, based on GEANT3 simulations of the reaction and transmission through the DRAGON separator. As can be seen, in both cases no recoil events fall within this expected window. As a result, we extracted upper limits on the resonance strengths using a modification of the Rolke profile likelihood method for calculating confidence intervals in the presence

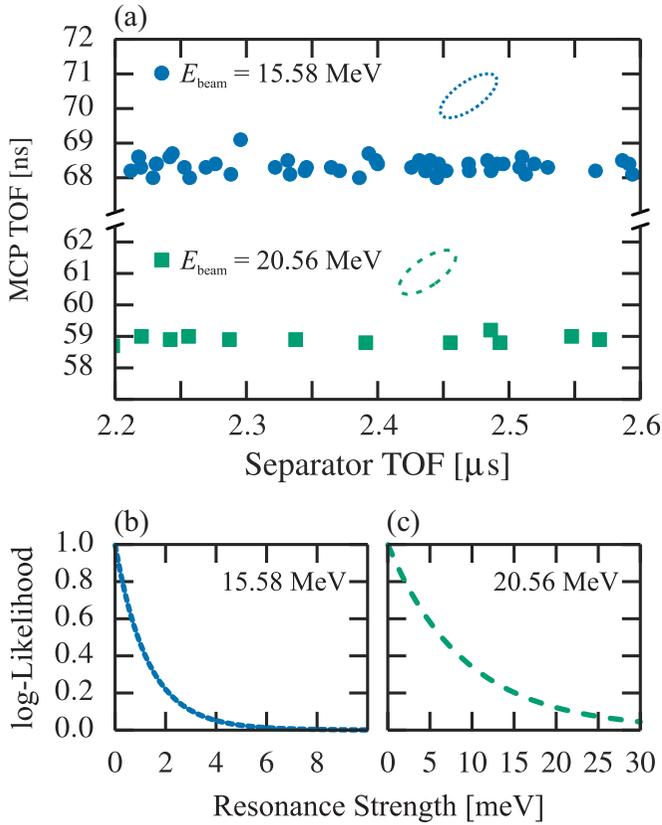


FIG. 2. (a) MCP vs separator TOF for the 15.58 MeV and 20.56 MeV beam energies. The blue dotted and green dashed ovals represent the expected location of recoils for the 15.58 MeV and 20.56 MeV beam energies, respectively. (b) Profile likelihood curve for the 15.58 MeV beam energy. (c) Profile likelihood curve for the 20.56 MeV beam energy.

of uncertain background rates and detection efficiencies [25]. In the standard Rolke treatment, the likelihood is the product of the individual likelihoods describing the signal rate  $\mu$ , background rate  $b$  (both treated as Poisson), and the detection efficiency  $\eta$  (treated as Gaussian with uncertainty  $\sigma_\eta$ ). Mathematically, this is expressed as

$$\mathcal{L}(\mu, b, \eta | x, y, z) = \left[ \frac{(\eta\mu + b)^x}{x!} e^{-\eta\mu - b} \right] \left[ \frac{(b\tau)^y}{y!} e^{-b\tau} \right] \left[ \frac{e^{-(z-\eta)^2/(2\sigma_\eta^2)}}{\sqrt{2\pi}\sigma_\eta} \right], \quad (3)$$

where  $x$  is the number of events observed in the signal region,  $y$  is the number of events observed in a background region that is  $\tau$  times as large as the signal region, and  $z$  is the observed signal rate. Equation (3) is then maximized with respect to  $b$  and  $\eta$  to construct a one-dimensional likelihood curve that is a function of only the signal strength  $\mu$  and can be analyzed to extract upper limits.

In the present analysis, we extend the Rolke method to also account for uncertainties in the resonance energy  $E_r$ , the number of incoming beam particles  $N$ , and the  $^{38}\text{K} + \text{H}_2$  stopping power  $\epsilon$ . Each of these quantities factors into the calculation of the resonance strength, and hence their uncertainties should

be included for a complete treatment of the problem. For each of these quantities, we treat the uncertainty as Gaussian (with widths  $\sigma_E$ ,  $\sigma_N$ , and  $\sigma_\epsilon$ , respectively). The complete likelihood function is then given by

$$\begin{aligned} \mathcal{L}(\omega\gamma, b, \eta, E_r, N, \epsilon | x, y, z, E_{r0}, N_0, \epsilon_0) &= \left[ \frac{(\eta\mu + b)^x}{x!} e^{-\eta\mu - b} \right] \left[ \frac{(b\tau)^y}{y!} e^{-b\tau} \right] \left[ \frac{e^{-(z-\eta)^2/(2\sigma_\eta^2)}}{\sqrt{2\pi}\sigma_\eta} \right] \\ &\times \left[ \frac{e^{-(E_{r0} - E_r)^2/(2\sigma_E^2)}}{\sqrt{2\pi}\sigma_E} \right] \left[ \frac{e^{-(N_0 - N)^2/(2\sigma_N^2)}}{\sqrt{2\pi}\sigma_N} \right] \left[ \frac{e^{-(\epsilon_0 - \epsilon)^2/(2\sigma_\epsilon^2)}}{\sqrt{2\pi}\sigma_\epsilon} \right], \end{aligned} \quad (4)$$

where  $E_{r0}$ ,  $N_0$ , and  $\epsilon_0$  are the observed central values of the resonance energy, beam ions on target, and stopping power, respectively. In Eq. (4), the signal rate  $\mu$  is no longer a constant parameter but rather a function of the resonance strength  $\omega\gamma$ , resonance energy  $E_r$ , number of incoming beam particles  $N$ , center-of-mass stopping power  $\epsilon$ , beam mass  $M$ , and target mass  $m$ :

$$\mu(\omega\gamma, E_r, N, \epsilon) = \frac{N(\omega\gamma)(hc)^2}{2\epsilon[E_r^2 + 2E_r m M / (m + M)]}. \quad (5)$$

Following the Rolke prescription, we maximize Eq. (4) with respect to the “nuisance” parameters  $\{b, \eta, E_r, N, \epsilon\}$  to arrive at a profile likelihood that is a function of only the resonance strength  $\omega\gamma$ . In practice, we first take the negative logarithm of Eq. (4) and then calculate the minimum numerically using the MINUIT package [23]. The resulting profile likelihoods for the 15.58 MeV and 20.56 MeV beam energies are shown in Figs. 2(b) and 2(c), respectively (plotted as negative log-likelihoods). To extract single-sided 68%, 90%, and 95% upper limits from the profile likelihood curves, we follow exactly the prescriptions of Ref. [25]. The resulting upper limits, along with all of the measured parameters going into the upper limit calculation are summarized in Tables I and II. It should be noted that when we refer to “68%” or “95%” confidence intervals, we mean the area under a normalized Gaussian distribution between the  $\pm 1\sigma$  or  $\pm 2\sigma$  limits. These are more precisely equal to 68.27% and 95.45%, respectively.

TABLE I. Summary of observed quantities going into the resonance strength upper limit calculations, for the 15.58 MeV beam energy.

Quantity	Value
Background rate	$4.44 \times 10^{-2}$
Beam ions on target	$(2.88 \pm 0.36) \times 10^{12}$
Stopping power (eV cm <sup>2</sup> )	$(3.78 \pm 0.14) \times 10^{-15}$
Detection efficiency	$0.093 \pm 0.016$
68% upper limit (meV)	1.16
90% upper limit (meV)	2.54
95% upper limit (meV)	3.53

TABLE II. Summary of observed quantities going into the resonance strength upper limit calculations, for the 20.56 MeV beam energy.

Quantity	Value
Background rate	$4.44 \times 10^{-3}$
Beam ions on target	$(8.8 \pm 1.2) \times 10^{11}$
Stopping power (eV cm <sup>2</sup> )	$(4.04 \pm 0.14) \times 10^{-15}$
Detection efficiency	$0.062 \pm 0.011$
68% upper limit (meV)	8.59
90% upper limit (meV)	18.4
95% upper limit (meV)	25.5

### B. 689(10) keV resonance

In contrast to the two lower-energy resonances, we observed a clear recoil signal when running with a beam energy of 27.17 MeV. This is demonstrated in the separator TOF vs MCP TOF distribution shown by the filled circles in Fig. 3(a). This spectrum exhibits a clear clustering of 27 recoil events in the region indicated by the open ellipse. The BGO  $z$ -position distribution of the identified recoil events is clustered downstream of the target center, indicating a resonance energy less than the central value of 689 keV [17]. Hence to extract a resonance strength,  $\omega\gamma$ , and a resonance energy,  $E_r$ , we use a technique similar to that employed in Ref. [26]. For a fixed beam energy of 27.17 MeV, we generate a simulated BGO  $z$ -position spectrum over the range of resonance energies contained within the gas target. For the simulations, we use the standard DRAGON GEANT3 package [27] and convolute the resulting BGO energies with a realistic hardware threshold. The hardware threshold was determined experimentally by taking long background runs with the threshold set to the value employed in the experiment, and to a reduced value of 50 mV. The resulting spectra were normalized, divided into each other, and fit with a Fermi function to arrive at the functional form used in the analysis. Following the threshold convolution, we scale the simulated spectra by the factor

$$\eta Y_{\omega\gamma} N_b / N_{\text{sim}}, \quad (6)$$

where  $\eta = 0.121 \pm 0.003$  is the heavy-ion detection efficiency,  $Y_{\omega\gamma}$  is the reaction yield at a given resonance strength  $\omega\gamma$ ,  $N_b = (2.53 \pm 0.30) \times 10^{12}$  is the number of incoming beam ions, and  $N_{\text{sim}} = 50,000$  is the number of simulated events. Scaled in this manner, the simulated spectrum represents both the magnitude and the shape of the BGO  $z$ -position distribution for a given  $\omega\gamma$  and  $E_r$ .

The  $\gamma$ -ray efficiency is implicitly included in the generation of the simulated spectra since the number of counts appearing in the spectra prior to scaling is determined by the detection efficiency, as modeled in the GEANT3 simulation. This modeling is sensitive to the branching ratios for  $\gamma$ -ray decay from the 6460 keV state in <sup>39</sup>Ca. These branching ratios have not been measured, and hence we have assumed dominant decays either directly to the ground state or through the first excited  $\frac{5}{2}^-$  state, as observed for the decay of known  $\frac{5}{2}^+$  excited states in the well-studied mirror nucleus <sup>39</sup>K

[21]. The location of the  $\frac{5}{2}^-$  state in <sup>39</sup>Ca has not been conclusively assigned, but there are a number of candidates in the  $\sim 3$ –4 MeV excitation energy region [21]. Hence, for the present analysis, we have assumed decay through a state at 3.5 MeV to represent the feeding through the  $\frac{5}{2}^-$ . To quantitatively account for the uncertainty related to the  $\gamma$ -ray decay scheme, we have utilized a profile likelihood technique to marginalize over the unknown branching ratios. Specifically, we performed separate simulations for a range of different fractional feedings directly to the ground state or through a state at 3.5 MeV. In the simulations, the ground state/excited state ratios ranged 0%–100% in steps of 10%. For each set of simulations, we took the branching with the highest likelihood value and incorporated it into the eventual likelihood surface used to extract confidence intervals on the resonance strength and energy (the calculation of likelihoods and construction of the likelihood surface is detailed later in this section). This technique of using profile likelihoods to deal with relevant, but uninteresting “nuisance” parameters is well established in the statistical literature; see, for example Refs. [25,28]. It should be noted that the uncertainty on the  $\gamma$ -ray detection efficiency is dominated by geometrical and Monte Carlo uncertainties and not the unknown branching ratios.

The yield parameter in Eq. (6),  $Y_{\omega\gamma}$ , is given by the convolution of the standard Breit-Wigner narrow-resonance cross section [29] with the gas target density profile. The density profile was measured in a previous experiment by recording the  $\gamma$ -ray yield from the <sup>3</sup>He(<sup>12</sup>C,  $p$ )<sup>14</sup>N $\gamma$  reaction in a shielded BGO detector moved along the length of the target [30]. These data (scaled to the 27.17 MeV beam energy employed in the present experiment) are shown in Fig. 3(b). The density profile was determined by fitting the data with the following function:

$$f(E) = 1/[1 + e^{(E-E_0)-\Delta E/2)/a}], \quad (7)$$

where  $E_0$  is the beam energy at the center of the gas target,  $\Delta E$  is the energy loss across the full length of the gas target, and  $a$  is a free parameter. The resulting best-fit is shown as the orange solid line in Fig. 3(b). The fitting procedure implicitly includes the stopping power,  $\epsilon = (3.95 \pm 0.14) \times 10^{-15}$  eV cm<sup>2</sup> (in the center-of-mass frame).

To extract a resonance strength and energy, we calculate the negative log-likelihood (NLL) by comparing our model (the scaled BGO  $z$ -position simulations) with experimental data, over a grid of resonance strengths and energies. We assume the counts per bin in the BGO  $z$ -position spectra are Poisson distributed, meaning the NLL is given by

$$-\ln \mathcal{L} = \sum_i \{\ln(n_i!) - n_i \ln(f_i)\} + S. \quad (8)$$

Here  $n_i$  is the number of measured counts in bin  $i$ ,  $f_i$  is the number of simulation counts in bin  $i$ , and  $S$  is the integral of the simulated distribution. The result of this likelihood analysis is shown in Fig. 3(c). This figure shows a contour plot of the NLL as a function of the resonance energy and resonance strength, which contains two local minima. The first (global) minimum is in the constant-pressure region of the target with  $E_r = 679$  keV,  $\omega\gamma = 120$  meV, and  $-\ln \mathcal{L}_0 = 16.2$ . The second

(local) minimum is far upstream in the target, where the density has not yet reached equilibrium, at  $E_r = 677$  keV,  $\omega\gamma = 650$  meV, and  $-\ln \mathcal{L}_1 = 16.9$ . Based on the NLL values, we exclude the  $E_r = 677$  keV solution at a 76% significance level. This significance level was calculated using the likelihood ratio test, wherein  $2 \ln[\mathcal{L}_0/\mathcal{L}_1]$  (here equal to 1.4) is taken to be  $\chi_1^2$  distributed [28]. The significance level is thus the value of  $X_1^2(1.4)$ , where  $X_1^2(x)$  is the  $\chi^2$  cumulative distribution function with one degree of freedom. The resulting best fits to both the BGO  $z$ -position and the  $\gamma$ -ray energy spectra are shown in Figs. 3(d) and 3(e), respectively.

Analyzing the region of the contour plot surrounding the global minimum, we extract 68% confidence intervals for the resonance energy and resonance strength of  $E_r = 679_{-1}^{+2}$  keV and  $\omega\gamma = 120_{-30}^{+50}$  meV. These quantities represent statistical uncertainties only. A number of sources of systematic uncertainty are also present, and are summarized in Table III. Note that the 0.17% systematic uncertainty on the beam energy (cf. Sec. II) is implicitly included since it is already folded into the quoted uncertainty on the stopping power. The resonance strength measurement is subject to systematic uncertainties related to each of the quantities in Table III, while the resonance

TABLE III. Sources of systematic uncertainty for the measurements at 27.17 MeV beam energy.

Quantity	Measured value	Relative uncertainty
$^{38}\text{Ar}$ background	(see text)	+0% -50%
Beam ions on target	$(2.53 \pm 0.30) \times 10^{12}$	12%
BGO efficiency	$0.541 \pm 0.054$	10%
Stopping power (eV cm <sup>2</sup> )	$(3.95 \pm 0.14) \times 10^{-15}$	3.5%
MCP transmission	$0.789 \pm 0.021$	2.7%
Charge state fraction	$0.192 \pm 0.002$	1.0%
MCP efficiency	$0.997 \pm 0.003$	0.3%
Live time	$0.79806 \pm 0.00002$	0.002%

energy measurement is affected only by the stopping power. Adding all of the relative uncertainties in quadrature, we arrive at the following resonance energy and strength values:

$$E_r = 679_{-1}^{+2}(\text{stat.}) \pm 1(\text{sys.}) \text{ keV},$$

$$\omega\gamma = 120_{-30}^{+50}(\text{stat.})_{-60}^{+20}(\text{sys.}) \text{ meV}.$$

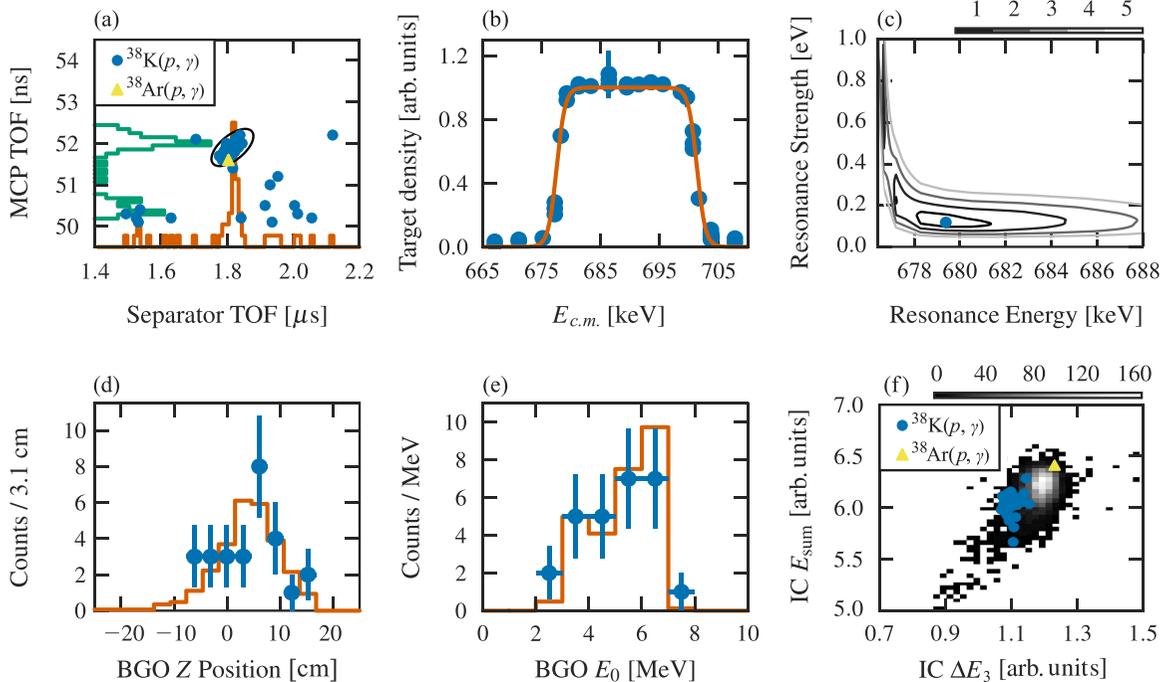


FIG. 3. Summary of the coincidence analysis for the data taken with a beam energy of 27.17 MeV. The individual descriptions of panels (a) through (f) are as follows: (a) Separator vs MCP TOF particle identification spectrum. The blue filled circles represent data collected with the radioactive  $^{38}\text{K}$  beam, while the single filled yellow triangle represents data collected with a  $^{38}\text{Ar}$  beam, for background characterization. The open ellipse outlines the expected recoil region. Projections onto the horizontal and vertical axes are also included (as the unshaded orange and green histograms). (b) Target density as a function of center-of-mass beam energy. The filled circles with error bars represent the data points and the solid line shows the fit to Eq. (7). (c) NLL contour plot, calculated by comparing simulated and measured BGO  $z$  positions as explained in the text. The solid blue point shows the location of the global minimum. (d) Measured BGO  $z$ -position distribution for recoil events (filled circles), compared with the best-fit simulation result at  $E_r = 679$  keV and  $\omega\gamma = 120$  eV (solid orange lines). (e) Same as panel (d), but showing the measured energy of the most energetic  $\gamma$ -ray hit in the BGO array. (f) Total energy deposited in the IC vs energy loss in the third (most downstream) anode. The filled (blue) circles show the location of the  $^{39}\text{Ca}$  coincidence recoils observed with the  $^{38}\text{K}$  beam. The filled yellow triangle denotes the location of the event observed with a beam of pure  $^{38}\text{Ar}$ . The greyscale map shows the location of all heavy-ion singles events observed with the 27.17 MeV  $^{38}\text{K}$  beam. This distribution is dominated by leaky beam.

The uncertainty due to potential background from reactions occurring on isobaric  $^{38}\text{Ar}$  contamination in the beam was determined through a background measurement using a stable beam of pure  $^{38}\text{Ar}$ , with a total ions on target of  $(6.9 \pm 0.6) \times 10^{11}$ . This measurement observed a single count near the edge of the expected recoil region, shown as the filled triangle in Fig. 3(a). This count is likely a random leaky beam event based on its location in the IC total energy vs energy loss spectrum. This is demonstrated in Fig. 3(f), which clearly shows that the suspected background event is well separated from the locus of  $^{38}\text{K}$  recoils and is consistent with the locus of leaky beam events. Furthermore, the known properties of  $^{38}\text{Ar} + p$  radiative capture imply that background from  $^{38}\text{Ar}$  contamination is highly unlikely. There are no known  $^{38}\text{Ar} + p$  resonances within 10 keV of the energies covered in the DRAGON gas target [21]. As a result, resonant capture is only possible through heretofore unknown proton-unbound states in the well-studied  $^{39}\text{K}$  nucleus. Concerning direct capture, we calculate an estimated cross section of 0.6 nb using the  $S$ -factor parametrization of Ref. [31]. Integrated across the length of the entire target, this results in an expected yield of only  $6 \times 10^{-5}$  recoils.

Given the small likelihood that the single event observed in the measurement with pure  $^{38}\text{Ar}$  beam is a genuine  $^{38}\text{Ar}(p,\gamma)^{39}\text{K}$  recoil, we do not alter the  $\omega\gamma = 120$  meV central value extracted from our likelihood analysis. However, for a conservative estimate of the associated uncertainties, we recommend that the lower-bound systematic uncertainty include the possibility of unforeseen contamination arising from  $^{38}\text{Ar}(p,\gamma)^{39}\text{K}$  reactions. To calculate this uncertainty, we first determine an upper limit of 2.4 events, or a yield of  $3.4 \times 10^{-12}$ , in the pure  $^{38}\text{Ar}$  beam measurement. We do this by applying the standard Rolke method [25] to the single count observed in the recoil region. In the production runs with the  $^{38}\text{K}$  radioactive beam (mixed with  $^{38}\text{Ar}$  contamination), this translates into an upper limit of 13 events. This upper limit is calculated assuming an Ar/K ratio of 1.54 in the production beam, determined by sending attenuated beam to the end of DRAGON and fitting the individual Ar and K components in the IC energy loss spectrum. Dividing by the 27 observed recoil events, we arrive at a relative uncertainty of 50%. This uncertainty applies only to the lower limit on the resonance strength since the presence of background due to beam contamination can only reduce, never increase, the measured resonance strength. We emphasize that this procedure for determining a systematic uncertainty due to potential  $^{38}\text{Ar}$  background is an *ad hoc* adjustment, not one formulated from rigorous statistical methods. Overall, it provides a conservative estimate on the total systematic uncertainty applied to the resonance strength measurement.

The beam delivered to DRAGON was also contaminated by isomeric  $^{38m}\text{K}$  ( $E_x = 130$  keV,  $t_{1/2} = 924$  ms). The ratio of  $^{38m}\text{K}$  to  $^{38g}\text{K}$  was measured to be  $7.1 \times 10^{-2}$  at the ISAC yield station. The yield measurements bypass the charge state booster, and hence some additional fraction of the isomers will decay before reaching DRAGON. The delay between production and arrival at the DRAGON target is dominated by the charge breeding time, which has been measured to be on the order of a few hundred milliseconds

[32]. Taking a nominal delay time of 400 ms, the  $^{38m}\text{K}/^{38g}\text{K}$  ratio would decrease to  $5.3 \times 10^{-2}$  by the time the beam reaches the DRAGON target. Given the small fraction of  $^{38m}\text{K}$  in the beam, no background from isomeric capture is expected.

### C. Singles analysis

In addition to the coincidence analysis of the 689 keV resonance presented in Sec. II B, we have also performed a separate extraction of the resonance strength using heavy-ion singles data alone. This analysis was *guided* by the results of the prior coincidence analysis, i.e., regions of interest in various parameter spaces were identified by the location of coincidence recoils. However, the final quantitative cuts applied to the singles data were determined from the distributions of the singles parameters alone. This singles analysis made use of the time difference between the incoming beam bunch (measured from the ISAC-I RFQ signal) and the upstream MCP to construct a separator TOF parameter without requiring prompt  $\gamma$  rays. This analysis is summarized in the plots shown in Fig. 4. Panel (a) shows the standard MCP TOF signal plotted vs the RF-MCP TOF, where the 27 events already identified as recoils in the coincidence analysis (represented by the blue filled circles) are tightly clustered in a narrow region of the plot. Continuing the analysis, we first set a gate on the entire upper-left region of the plot, which contains all of the coincidence recoils (the actual gate is included in the Fig. 4(a) as the blue dashed line). We then project these events onto the solid black diagonal axis shown in the figure.

The new projected parameter (“RF-MCP projection”) is shown in the panel (b) of Fig. 4, plotted vs two separate parameters: (1) the  $y$  position in the upstream MCP, deduced from a resistive-anode readout scheme; and (2) the energy loss in the third (most downstream) anode of the IC. In both cases, the confirmed recoil events are tightly clustered in a single region of the plot. To further separate recoil events from background, we place a one-dimensional cut on the “RF-MCP projection” parameter, including all events to the left of the black dotted line in the figure. For these events only, we then plot the IC energy loss vs the MCP  $y$  position, shown in Fig. 4(c). Here, the singles events cluster into two distinct loci, with the confirmed recoil events falling entirely within the upper-right cluster. From this, we conclude that the singles events in the upper-right locus correspond to recoils, while the events in the lower-left locus correspond to background leaky beam events.

To quantify the overlap between the recoil and leaky beam regions in Fig. 4(c), we project onto the diagonal axis indicated by the solid black line in the figure. This projection is shown in the Fig. 4(d). The measured data (shown as open circles with error bars) are well described by a double-Gaussian distribution (shown as dashed, dot-dashed, and solid lines, as indicated in the legend). The smaller Gaussian on the left of the figure corresponds to the estimated background distribution, and the larger Gaussian on the right of the figure corresponds to the recoil distribution. We take the true number of recoil events to be equal to the integral of the signal distribution,  $52.0 \pm 8.2$ . The uncertainty on this quantity comes from propagating the

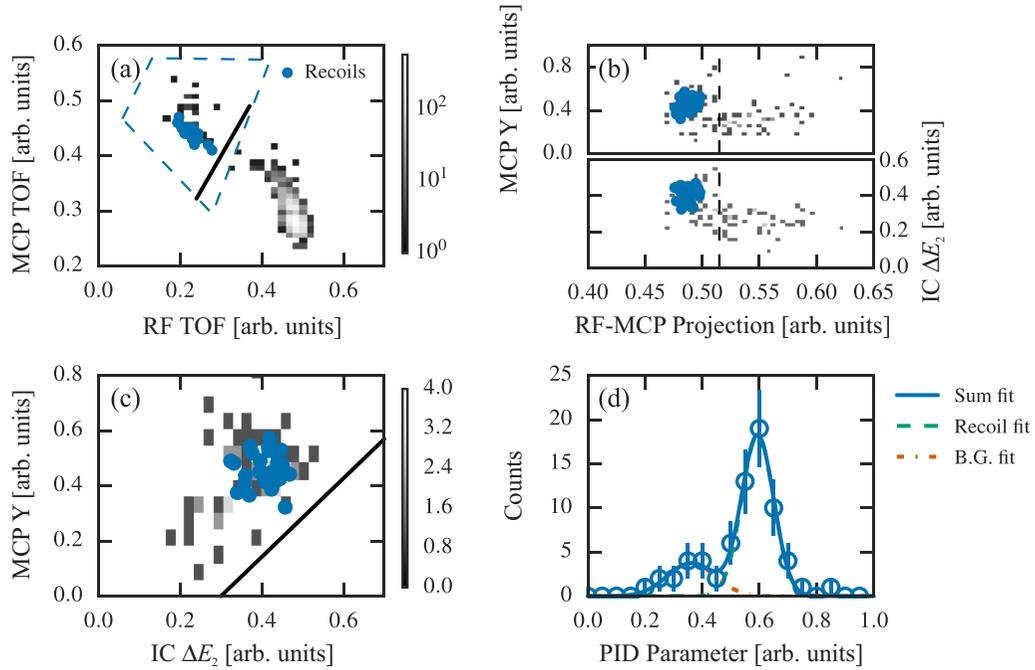


FIG. 4. Summary of the singles resonance strength analysis for the 27.17 MeV beam energy. In panels (a)–(c), the blue filled circles represent events already identified as recoils in the coincidence analysis, and the greyscale intensity maps represent all singles data. In panel (d), the open circles represent all singles data, and the various curves represent fits as indicated in the legend. The solid black lines in panels (a) and (c) represent diagonal axes onto which the two-dimensional data are projected for subsequent analysis. The dashed black lines in panel (b) represent the cut placed on the “RF-MCP Projection” parameter. The full significance of each plot is explained in the main text.

$1\sigma$  uncertainties on the individual fit parameters, which were calculated with MINUIT.

To calculate the singles resonance strength, we use the standard thick-target formula [29],

$$\omega\gamma = 2N_r\epsilon/(\eta N_b\lambda^2), \quad (9)$$

where  $N_r = 52.0 \pm 8.2$  is the number of recoil events,  $\epsilon = (3.95 \pm 0.14) \times 10^{-15}$  eV cm<sup>2</sup> is the center-of-mass stopping power,  $\eta = 0.110 \pm 0.003$  is the heavy-ion detection efficiency,  $N_b = (2.53 \pm 0.30) \times 10^{12}$  is the number of beam ions, and  $\lambda = (3.513 \pm 0.005) \times 10^{-12}$  cm is the center-of-mass de Broglie wavelength. Note that the heavy-ion detection efficiency includes the IC efficiency of  $0.913 \pm 0.003$ . This was not included in the heavy-ion efficiency used in the coincidence analysis since the IC was not used to select coincidence events. The de Broglie wavelength assumes a resonance energy of  $E_r = 679 \pm 2$  keV, as extracted from our previous maximum likelihood analysis. The influence of this assumption is minor; calculating the resonance strength using the previous resonance energy of  $689 \pm 10$  keV increases the result by less than 1 meV. The resulting resonance strength is  $\omega\gamma = 120 \pm 20$  meV (statistical uncertainty only), which is in good agreement with our coincidence result of  $120^{+50}_{-30}$  meV (the exact agreement of the central values should be considered fortuitous). The estimated singles systematic uncertainty is  $\pm 15$  meV, calculated by propagating the uncertainties for the stopping power, number of beam particles, and detection efficiency.

In practice, the singles technique frequently results in a lower systematic uncertainty than the coincidence method since there is no need to estimate the  $\gamma$ -ray detection efficiency. This efficiency typically comes with a relative uncertainty of 10% or greater, resulting from uncertainties in the GEANT3 simulation of the BGO array [27], as well as from unknown  $\gamma$ -ray decay schemes. However, for reliable application of the singles technique, it is crucial that the full width of the resonance be contained within the gas target, to ensure that the thick-target approximation of the resonance strength formula is valid. In the future, technical advances will likely improve the ability to discern resonance positions based on only a handful of recoil events. With this capability, an off-center resonance would be spotted early on during a running period, and the beam energy could be adjusted accordingly. One example presently under development is the use of a fast-timing LaBr array for  $\gamma$ -ray detection. The fast timing properties of LaBr allow the resonance position to be deduced from the time difference between the detected  $\gamma$  rays and the arrival of the corresponding beam bunch. Preliminary calculations and simulation work suggest that this method is more precise than the presently employed  $z$ -position technique and may be applied to data sets with as few as  $\sim 5$  confirmed recoils [33].

### III. DISCUSSION

As discussed in Ref. [11], the present measurements place significant constraints on the overall  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reaction rate at nova temperatures. This is demonstrated in Fig. 5,

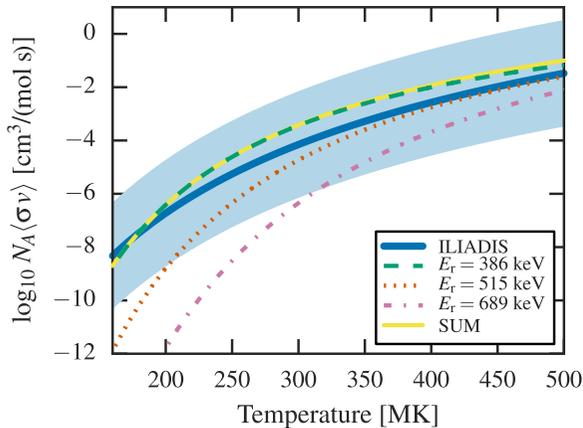


FIG. 5. Updated  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reaction rate across the temperature regime covered by classical novae, calculated assuming the rate is dominated by the three  $\ell = 0$  resonances studied in the present experiment. The contributions of the 386 and 515 keV resonances represent upper limits, while the “689 keV” resonance contribution represents our measured central value of 120 meV. Also shown is the statistical model rate of Ref. [9] (“ILIADIS”), along with its associated uncertainties (shaded region).

which shows the calculated rate vs temperature curves for the three presently reported resonances, along with their sum. Assuming the astrophysical rate is dominated by these three resonances, the lower curve (the nominal 689 keV resonance) sets a lower limit on the astrophysical rate, while the sum sets an upper limit. For comparison, the recommended rate from Iliadis *et al.*, along with the factor 100 up/down uncertainty band, is also included in the figure. At peak temperatures for ONe nova burning,  $T \simeq 0.4$  GK, the total uncertainty has been reduced from a factor of  $10^4$  to a factor of  $\sim 40$ . Applying these new, experimentally based limits to the model predictions of the Iliadis *et al.* sensitivity study [8] results in a reduction of overall uncertainties on  $^{38}\text{Ar}$ ,  $^{39}\text{K}$ , and  $^{40}\text{Ca}$  production in ONe novae from respective factors of  $\sim 25$ , 136, and 57 to factors of  $\sim 2$ , 18, and 9. Note for these calculations, the nucleosynthesis models which maximized the sensitivity to the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate were used. For  $^{38}\text{Ar}$  and  $^{40}\text{Ca}$  this corresponds to the “S1” model ( $T_{\text{peak}} = 418$  MK), while for  $^{39}\text{K}$  the “P2” model ( $T_{\text{peak}} = 356$  MK) was used.

In order to investigate the dependence of these sensitivity results to specific nova models, we have performed a separate  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  sensitivity study based on an independent calculation performed with the NUGRID package, using the single-zone “post processing network” (ppn) code [34]. The initial conditions of the calculation are a 1.3 solar mass white dwarf with a temperature of 7.0 MK. The white dwarf composition is given by the “Denisenkov” model, evolved using the Modules for Experiments in Stellar Astrophysics (MESA) code [35]. The accretion rate is  $10^{-11} M_{\odot}/\text{yr}$ , and the composition of the accreted material is assumed to be solar. The peak temperature of the model outburst is 408 MK, similar to the 418 MK peak temperature of the S1 model from Ref. [8]. A complete description of the parameters going into the NUGRID

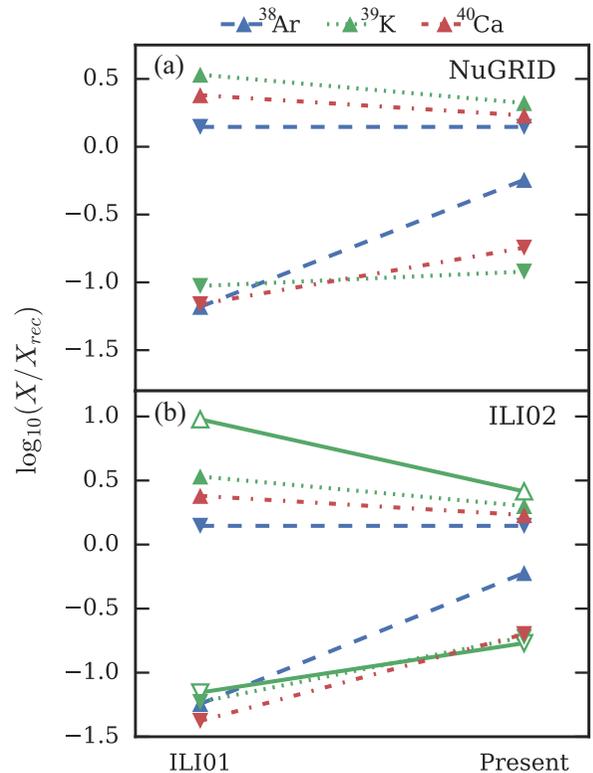


FIG. 6. Summary of the sensitivity study results presented in Table IV. The various data points represent the logarithm of the ratio  $X/X_{\text{rec}}$ , where  $X_{\text{rec}}$  is the predicted abundance of a given isotope taking the recommended  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate from Ref. [9], and  $X$  is the predicted abundance of the same isotope taking the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate to be at the upper or lower limit of various uncertainty bands. For the points on the left of the figure (labeled “ILI01”),  $X$  is evaluated at the factor of 100 up / 0.01 down uncertainty limits given in Ref. [9]. For points on the right of the figure (labeled “Present”),  $X$  is taken from the uncertainty band established in Ref. [11]. In all cases, up-turned triangles represent  $X$  calculated at the upper limit of the uncertainty band and down-turned triangles represent  $X$  at the lower limit. The various dashed, dotted, and dot-dashed lines represent abundance calculations for  $^{38}\text{Ar}$ ,  $^{39}\text{K}$ , and  $^{40}\text{Ca}$  as indicated in the legend at the top of the figure. Panel (a) shows the results of the present NUGRID sensitivity study, while panel (b) shows the results of the Iliadis *et al.* sensitivity study [8]. In panel (b), results from both the “S1” and “P2” nova models are displayed for  $^{39}\text{K}$ . The filled green triangles connected by the dotted line represent results of the “S1” model, while the open triangles connected by solid lines represent the results of the “P2” model.

calculation can be found in Ref. [34]. The results of the NUGRID sensitivity study are summarized in Table IV, with results from the Iliadis *et al.* study included for comparison; the results are also presented graphically in Fig. 6. The recommended rate utilized for this analysis is identical to the one presented in Ref. [9]. Overall, the predictions of the NUGRID and the S1 models are rather consistent, with agreement to within a factor of 2 in all cases.

Our measurements and sensitivity analyses indicate that the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate is not a likely source of significant over- or underproduction of  $^{38}\text{Ar}$ ,  $^{39}\text{K}$ , or  $^{40}\text{Ca}$  in novae (relative to

TABLE IV. Calculated abundances for Ar, K, and Ca isotopes, from the NUGRID simulations explained in the text. Results from the Iliadis *et al.* sensitivity study, model “S1” (as well as model “P2”, for  $^{39}\text{K}$ ) are included for comparison [8]. The quantity  $X_{\text{rec}}$  represents the abundance calculated using the recommended  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  rate from Ref. [9], while  $X_{100}$  and  $X_{0.01}$  represent abundances calculated with the recommended rate multiplied by factors of 100 and 0.01, respectively. The quantities  $X_{\text{up}}$  and  $X_{\text{low}}$  represent abundances calculated with the experimental upper and lower limits presented in Ref. [11].

Nuclide	$\frac{X_{100}}{X_{\text{rec}}}$	$\frac{X_{\text{up}}}{X_{\text{rec}}}$	$\frac{X_{\text{low}}}{X_{\text{rec}}}$	$\frac{X_{0.01}}{X_{\text{rec}}}$	$\frac{X_{100}}{X_{0.01}}$	$\frac{X_{\text{up}}}{X_{\text{low}}}$
NUGRID						
$^{38}\text{Ar}$	0.066	0.57	1.4	1.4	1/21	1/2.5
$^{39}\text{K}$	3.4	2.1	0.14	0.094	36	15
$^{40}\text{Ca}$	2.4	1.7	0.18	0.069	35	9.4
Iliadis <i>et al.</i> [8]						
$^{38}\text{Ar}$ (“S1”)	0.057	0.60	1.4	1.4	1/25	1/2.3
$^{39}\text{K}$ (“S1”)	3.4	2.0	0.19	0.059	58	11
$^{39}\text{K}$ (“P2”)	9.5	2.6	0.17	0.070	136	15
$^{40}\text{Ca}$ (“S1”)	2.4	1.7	0.20	0.042	57	8.5

solar abundances). Hence the over-production of Ar and Ca observed in the spectra of nova ejecta [2–5] remains unexplained. We encourage more extensive sensitivity studies and multizone model calculations to investigate the source of this anomaly.

It should be noted that the present results intrinsically depend on the veracity of previous transfer reaction studies, which have established the  $^{39}\text{Ca}$  level scheme in the  $E_x = 6\text{--}7\text{ MeV}$  region. If the spins or level energies established from these studies are incorrect or incomplete, the present experiment may have neglected to cover the most important resonance energy windows for astrophysics. For this reason, we encourage future high-resolution transfer reaction studies that are targeted specifically at measuring the properties of potential astrophysical proton capture resonances in  $^{39}\text{Ca}$ .

Although the present measurements are not directly sensitive to the spins of the measured resonances, we can still infer properties of the resonances in question based on the measured strengths. For this, we use the standard formula for the resonance strength [29],

$$\omega\gamma = \frac{2J_r + 1}{(2J_p + 1)(2J_{^{38}\text{K}} + 1)} \frac{\Gamma_\gamma \Gamma_p}{\Gamma_\gamma + \Gamma_p}, \quad (10)$$

where  $J_r = \frac{5}{2}$ ,  $J_p = \frac{1}{2}$ , and  $J_{^{38}\text{K}} = 3$  are the respective spins of the resonance, proton, and  $^{38}\text{K}$ ; and  $\Gamma_\gamma$  and  $\Gamma_p$  are the respective  $\gamma$ -ray and proton partial widths of the resonance. Assuming a “hard” upper limit on the proton spectroscopic factor for unbound states of  $C^2S \leq 0.1$  and the measured strength value of  $\omega\gamma = 120\text{ meV}$  for the  $679 \pm 2\text{ keV}$  resonance, we calculate an upper limit on the mean  $\gamma$ -decay lifetime for this state of  $\tau \leq 2.2\text{ fs}$ . For shorter lifetimes, as  $\tau \rightarrow 0$ ,  $\Gamma_\gamma/(\Gamma_\gamma + \Gamma_p) \rightarrow 1$ , and we calculate a lower limit on the spectroscopic factor of  $C^2S \geq 0.0055$ . For the  $515 \pm 10\text{ keV}$  resonance, the 90% confidence level (CL) upper limit on the strength

of  $18.4\text{ meV}$  sets an upper limit on the lifetime of  $\tau \leq 12\text{ fs}$  (again taking the “hard” upper limit on the spectroscopic factor at  $C^2S = 0.1$ ). For short lifetimes,  $\Gamma_\gamma/(\Gamma_\gamma + \Gamma_p) \simeq 1$ , we calculate an upper limit on the spectroscopic factor of  $C^2S \leq 0.022$ . For the  $386 \pm 10\text{ keV}$  resonance, the calculated upper limit on the lifetime is  $\tau \leq 38\text{ fs}$  (again taking the measured 90% upper limit on the strength of  $\omega\gamma \leq 2.54\text{ meV}$  and the “hard” spectroscopic factor limit of 0.1). For short lifetimes satisfying  $\Gamma_\gamma/(\Gamma_\gamma + \Gamma_p) \simeq 1$ , we calculate an upper limit on the spectroscopic factor of  $C^2S \leq 0.066$ . We emphasize that these limits are simply “back of the envelope” calculations and are not intended to set any rigid limits on the single-particle properties of the resonances in question.

To summarize, we have performed the first ever direct measurement of the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reaction, focusing on the three potential  $\ell = 0$  resonances within the Gamow window for classical novae, whose energies have been determined previously to be  $386 \pm 10\text{ keV}$ ,  $515 \pm 10\text{ keV}$ , and  $689 \pm 10\text{ keV}$ . For the highest-energy resonance, we observed a clear  $^{39}\text{Ca}\text{-}\gamma$  coincidence signal consisting of 27 events. We performed a two-dimensional likelihood analysis on the position distribution of the measured  $\gamma$  rays to extract a resonance strength and energy of  $\omega\gamma = 120_{-30}^{+50}(\text{stat.})_{-60}^{+20}(\text{sys.})\text{ meV}$  and  $E_r = 679_{-1}^{+2}(\text{stat.}) \pm 1(\text{sys.})\text{ keV}$ , respectively. The quoted systematic uncertainties are conservative and include the possibility of background events arising from stable  $^{38}\text{Ar}$  beam contamination. We also performed a separate analysis of  $^{39}\text{Ca}$  singles data and extracted a resonance strength of  $\omega\gamma = 120 \pm 20(\text{stat.}) \pm 15(\text{sys.})\text{ meV}$ , consistent with the coincidence result. For the lower two resonances, we observed no events consistent with recoils and used a profile likelihood technique to extract 90% CL upper limits on the resonance strengths of  $2.54\text{ meV}$  and  $18.4\text{ meV}$  for the lower and middle resonances, respectively. Based on these measurements, we have established new recommended upper and lower limits for the  $^{38}\text{K}(p,\gamma)^{39}\text{Ca}$  reaction rate which reduce uncertainties at peak nova temperatures ( $T_9 \sim 0.4$ ) from a factor of  $10^4$  to a factor  $\sim 40$ . Incorporating these new limits into two separate nova model calculations, we find that the uncertainties on the predicted abundances of  $^{38}\text{Ar}$ ,  $^{39}\text{K}$ , and  $^{40}\text{Ca}$  are reduced to a factor of 15 or below in all cases.

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