

Single quantum dot spectroscopy using a fiber taper waveguide near-field optic

Kartik Srinivasan^{a)} and Oskar Painter

Center for the Physics of Information, California Institute of Technology, Pasadena, California 91125
and Department of Applied Physics, California Institute of Technology, Pasadena, California 91125

Andreas Stintz and Sanjay Krishna

Center for High Technology Materials, University of New Mexico, Albuquerque, New Mexico 87106

(Received 13 June 2007; accepted 3 August 2007; published online 27 August 2007)

Photoluminescence spectroscopy of single InAs quantum dots at cryogenic temperatures (~ 14 K) is performed using a micron-scale optical fiber taper waveguide as a near-field optic. A lower bound on the measured collection efficiency of quantum dot spontaneous emission into the fundamental guided mode of the fiber taper is estimated at 0.1%, and spatially resolved measurements with ~ 600 nm resolution are obtained by varying the taper position with respect to the sample and using the fiber taper for both the pump and collection channels. © 2007 American Institute of Physics.
[DOI: 10.1063/1.2775811]

Semiconductor quantum dots (QDs) have drawn significant recent interest due to their potential to create optoelectronic devices such as those that can generate single photons on demand.¹ Methods to optically interrogate single QDs include physical isolation via microfabricated mesas² and diffraction-limited confocal optical microscopy.³ Despite the technical challenges involved in developing tools for use at cryogenic temperatures, solid immersion lenses⁴ have been implemented to improve the efficiency of light collection, while near-field scanning optical microscopy (NSOM) has been used to achieve sub-100 nm spatial resolution.⁵ In this letter, we examine the use of optical fiber taper waveguides as a near-field optic for performing single QD spectroscopy. These micron-scale silica waveguides have been used in many studies of optical microcavities, beginning as an efficient coupler to silica microspheres.⁶ More recently, we have shown that they can effectively probe the spatial and spectral properties of small mode volume (V_{eff}), high refractive index semiconductor cavities.⁷ Other researchers have proposed^{8,9} and realized their use as a collection tool for spontaneous emission from atomic vapors.¹⁰ Here, we show that a fiber taper may be used to channel emission from single self-assembled QDs embedded in a semiconductor slab directly into a standard single-mode fiber with high efficiency ($\sim 0.1\%$), and to provide submicron spatial resolution of QDs.

The QDs we study consist of a single layer of InAs QDs embedded in an $\text{In}_{0.15}\text{Ga}_{0.85}\text{As}$ quantum well, a so-called dot-in-a-well (DWELL) structure.¹¹ The DWELL layer is grown in the center of a GaAs waveguide (total waveguide thickness of 256 nm), which sits atop a $1.5 \mu\text{m}$ thick $\text{Al}_{0.7}\text{Ga}_{0.3}\text{As}$ buffer layer. The resulting peak of the ground state emission of the ensemble of QDs is located at $\lambda = 1.35 \mu\text{m}$ at room temperature. To limit the number of optically pumped QDs, microdisk cavities of diameter $D = 2 \mu\text{m}$ were fabricated using electron beam lithography and a series of dry and wet etching steps.¹² Although the QDs physically reside in a microcavity, they are *nonresonant* with the cavity whispering gallery modes (WGMs). In other words, our primary interest here is general single QD spectroscopy through the fiber

taper, without enhancement through interaction with the high quality factor (Q) microdisk WGMs. The samples were mounted in a continuous-flow liquid He cryostat that has been modified to allow sample probing with optical fiber tapers while being held at cryogenic temperatures ($T \sim 14$ K), as described in detail in Ref. 13. The cryostat setup provides any combination of free-space and fiber taper pumping and collection [see Fig. 1(a) for details].

The inset of Fig. 1(b) shows the emission spectrum from an ensemble of QDs in one of the microdisks. Here, the device is optically pumped through an objective lens at normal incidence (free-space pumping), with a spot size of $3 \mu\text{m}$ and wavelength $\lambda_p = 830$ nm. Clearly present are the ground and excited states (s and p shells) of the ensemble of QDs which, based on the estimated QD density of $300\text{--}500 \mu\text{m}^{-2}$, consists of ~ 1000 QDs. To study isolated emission lines from single QDs, we focus on the long-wavelength tail end of the QD distribution ($\lambda = 1290\text{--}1310$ nm). In this range, isolated emission lines from a single QD are seen for a fraction (10%) of the interrogated devices. A typical spectrum as collected through the pump lens (free-space collection) from one such device is shown in the top panel of Fig. 1(b). Under identical pumping conditions, the signal collected through a fiber taper waveguide positioned on top of, and in contact with, the microdisk is shown in the bottom panel of Fig. 1(b). The taper is a single mode optical fiber that has been heated and stretched down to a minimum radius of $a = 650$ nm, and is installed in the customized liquid He cryostat as detailed in Ref. 13. The most stark difference between the free-space and fiber taper collected spectra is the 25 times increase in fiber taper collected power. Similar improvement in collection efficiency was measured over all the QDs studied in this work.

Before further studying the fiber taper as a collection optic, we attempt to identify the different QD lines of Fig. 1(b). Of particular benefit in this assignment is the recent work of Cade *et al.*,¹⁴ who study a DWELL material very similar to that investigated here. In Fig. 2(a), we show taper-collected emission spectra as a function of pump power (free space, $\lambda_p = 830$ nm) for a fixed taper position. Emission is first seen for incident powers of a few nanowatts (estimated absorbed powers of tens of picowatts), with excitonic

^{a)}FAX: (626) 795-7258; electronic mail: kartik@caltech.edu

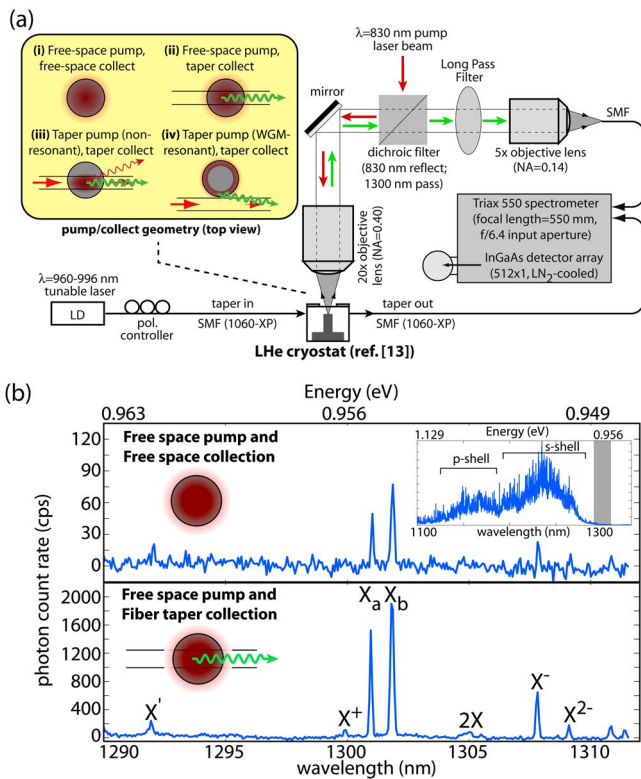


FIG. 1. (Color online) (a) Schematic of the experimental apparatus, showing the pump and collection configurations studied here. (b) Emission from a single QD using free-space collection (top) and fiber taper collection (bottom), under identical free-space pumping conditions ($\lambda_p=830$ nm, 15.7 nW incident power). The inset shows emission over a broad wavelength range from the ensemble of QDs within the microdisk. The shaded region from $\lambda=1290-1310$ nm is the spectral region where single QD emission has been observed.

lines centered at 1291.95, 1300.97, 1301.81, and 1307.75 nm. As we discuss later, spatially resolved measurements clearly indicate that the shortest wavelength emission line is unrelated to the latter three, which we identify as the polarization-split neutral exciton lines (X_a and X_b) (Refs. 14 and 15) and the negatively charged exciton (X^-). As the pump power is increased, additional emission lines appear, including the positively charged exciton (X^+) at 1299.87 nm and the biexciton ($2X$) line at 1305.11 nm. The X^-X , X^+X , and $2X-X$ splitting values of 4.6, -1.1 , and 2.8 meV match reasonably well with the 5.6, -1.1 , and 3.1 meV values measured in Ref. 14, although the fine structure splitting in the X line is significantly larger (600 versus 300 μeV) for this QD. In Fig. 2(b), we plot the emission level in each QD state against pump power. Below saturation, the emission lines all scale nearly linearly with pump power, except for the $2X$ line which scales superlinearly. The subquadratic growth rate ($n=1.54$), consistent with Ref. 14, is related to the reduced QD symmetry that generated the large X_a-X_b splitting.¹⁶ Finally, the temperature (T) dependence of the X lines is shown in Fig. 2(c), where significant broadening is seen for $T > 50$ K. Below this temperature we measure linewidths of 0.1–0.15 nm, roughly corresponding to the spectral resolution of our system (0.1 nm = 75 μeV).

A rough estimate of the absolute collection efficiency of the fiber taper is derived by considering the saturated photon count rates for the X lines in Fig. 2(b). The measured saturated photon count rate into the X_b line is 1.5×10^5 counts/s, which after considering the detector,

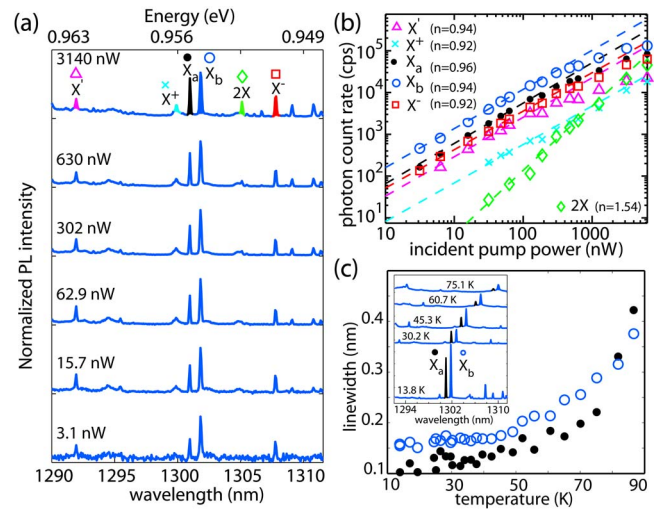


FIG. 2. (Color online) (a) Normalized emission spectra from a single QD for different free-space incident pump powers and using fiber taper collection. (b) log-log plot of the collected emission for the X_a (\bullet), X_b (\circ), X^+ (\times), X^- (\square), $2X$ (\diamond), and X' (\triangle) QD states of part (a). The dashed lines are least squares fits to the emission data below saturation, assuming $I \sim P^n$, where I is the collected photon count rate emitted into a given line and P is the pump power. (c) Linewidth of the QD exciton states X_a and X_b as a function of temperature.

taper, and spectrometer grating efficiencies (85%, 80%, and 60%), and including light in the backwards fiber channel, corresponds to a taper-collected photon count rate of 7.4×10^5 counts/s. Neglecting possible suppression or enhancement of radiation due to the presence of the microdisk, the saturated photon emission rate for InAs QDs is $5-10 \times 10^8$ counts/s [photon lifetime $\tau=1-2$ ns (Ref. 1)]. This yields an approximate fiber taper collection efficiency of $\eta_c=0.1\%$. It is important to note that this efficiency is for nonresonant collection, and does not correspond to that attainable for QDs resonant with a high- Q microdisk WGM, which one would expect to be much higher due to the Purcell-enhanced emission into a localized cavity mode¹ and the efficient taper-WGM coupling.¹² Finally, our estimate of the saturated photon emission rate neglects fluorescence intermittency due to pumping of dark states of the QD;¹⁷ 0.1% is therefore a minimum bound on η_c .

Obtaining a theoretical value for the nonresonant fiber taper collection efficiency is hindered by the complex geometry in which the QD is embedded; however, a coarse estimate can be made by comparison to Refs. 8 and 9. In these works, a spontaneous emission collection efficiency of 20%–50% is estimated for a dipole emitter on the surface of a silica fiber taper of radius $a \sim 1.4/k_0$ (k_0 is the free-space wave number). For our larger fiber taper ($a \sim 3/k_0$), and for a QD 125 nm away from its surface (corresponding to the middle of the GaAs slab), the theoretical collection efficiency is on the order of 1%. Beyond providing an upper bound on the collection efficiency, the model indicates that we are far from the optimum fiber taper diameter. Comparison with Ref. 9 indicates that an order of magnitude increase in the collection efficiency may be obtained by decreasing the fiber taper radius to $a=300$ nm, an experimentally realizable value.

An additional benefit of using the fiber taper as a near-field collection optic is the potential for spatially resolved measurements. Although the spatial resolution provided by a glass fiber taper⁷ is lower than the sub-100 nm level achieved by a scanning probe microscope,¹⁸ the resolution is still significantly better than that of a conventional microscope. AIP license of copyright, see http://apl.aip.org/apl/copyright.jsp

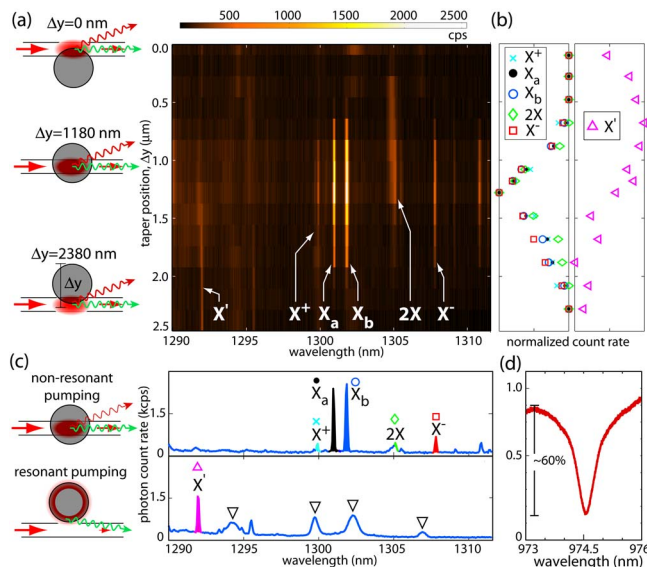


FIG. 3. (Color online) Spatially resolved measurements using fiber taper pumping and collection. (a) Collected emission spectrum as a function of taper position along the \hat{y} axis of the sample for nonresonant pumping through the fiber taper. (b) Spatial dependence of the integrated emission into each QD line. (c) Emission spectrum for nonresonant pumping with the taper at $\Delta y = 1180$ nm and an input power of 125 nW (top), and for resonant pumping of a microdisk WGM in the $\lambda_p = 980$ nm pump band at a power of 690 pW (bottom). The triangles (∇) indicate the position of WGM resonances. (d) Wavelength scan of the WGM used to pump the sample in (c).

able through NSOM,⁵ valuable information on the spatial location of QDs can be inferred from both the spatially dependent collection and excitation through the fiber taper. For the following measurements, we pump the microdisk through the fiber taper, with $\lambda_p = 978.3$ nm, where only the DWELL is significantly absorbing. Figure 3(a) shows a plot of the fiber-collected emission spectrum as a function of taper position along the \hat{y} axis of the sample (the taper position is adjusted through a piezostage on which it is mounted). Figure 3(b) plots the spatial dependence of the total photon count rate within each of the QD states identified in Fig. 2. Of note is the similar spatial dependence of the collected emission from each of the lines $\{X_a, X_b, X^+, X^-, 2X\}$, confirming that they originate from the same single QD. The full width at half maximum of the collected emission is roughly 600 nm, giving an estimate of the taper's spatial resolution transverse to its longitudinal axis. Two-dimensional mapping of a QD's position may also be obtained by rotating the sample, and repeating the measurement along the orthogonal axis as was done in Ref. 7 for the modes of photonic crystal cavities. Finally, in contrast to the other emission lines, emission from the short wavelength X' line at 1292 nm has a quite different dependence on taper position, with collected emission being strongest for the taper at the disk periphery.

Spatial selection of QDs may also be realized by resonantly pumping a microdisk WGM. This excites QDs located in a 250 nm thick annulus at the microdisk perimeter, where the pump beam resides, and efficient taper-WGM coupling allows for an accurate estimate of the *absorbed* pump power. The QDs located at the disk periphery are of course those that are of interest for cavity quantum electrodynamics (QED) studies involving high- Q , ultrasmall V_{eff} WGMs. Figure 3(d) shows a transmission scan of a pump-band WGM with a coupling depth of 60% and $Q \sim 1000$ (limited by DWELL absorption). The bottom scan of

Fig. 3(c) shows the emission spectrum when we pump on resonance with 690 pW of power at the taper input (corresponding to 330 pW of dropped/absorbed power). Emission from the centrally located QD [top scan of Fig. 3(c)] is no longer present, and has been replaced by a pronounced emission peak at $\lambda = 1291.95$ nm, corresponding to the X' line, confirming that this emission is likely due to a QD located in the disk periphery. Another difference in comparison to the nonresonant pumping spectrum is the presence of several broad emission peaks. These peaks are due to emission into relatively low- Q , higher-radial-order WGMs of the microdisk, as confirmed by fiber-taper-based transmission spectroscopy of the cavity with a tunable laser.¹² The source of such background emission into detuned cavity modes is not well understood, but has been observed to occur for even large detunings of 10–20 nm.¹⁸ In this case, it is likely that the preferential excitation of QDs that reside in the microdisk perimeter, even those that have exciton lines which are significantly detuned spectrally, results in enhanced emission into the microdisk WGMs.

In summary, we have shown that a micron-scale optical fiber taper waveguide, previously demonstrated to be an effective tool for characterization of semiconductor microcavities, can also be used to study single semiconductor quantum dots. As a near-field collection optic, the fiber taper is shown to channel quantum dot light emission directly into a single mode fiber with a high efficiency of 0.1%, and to provide a submicron spatial resolution of QDs. The ability to effectively investigate both microcavities and quantum dots suggests that these fiber tapers can serve as a very versatile tool for studying microphotonic structures, and in particular, for investigations of chip-based cavity QED.

¹P. Michler, A. Kiraz, C. Becher, W. V. Schoenfeld, P. M. Petroff, L. Zhang, E. Hu, and A. Imamoglu, *Science* **290**, 2282 (2000).

²J. Y. Marzin, J. M. Gérard, A. Izraël, D. Barrier, and G. Bastard, *Phys. Rev. Lett.* **73**, 716 (1994).

³E. Dekel, D. Gershoni, E. Ehrenfreund, D. Spektor, J. M. Garcia, and P. M. Petroff, *Phys. Rev. Lett.* **80**, 4991 (1998).

⁴Q. Wu, R. D. Grober, D. Gammon, and D. S. Katzer, *Phys. Rev. Lett.* **83**, 2652 (1999).

⁵Y. Toda and Y. Arakawa, *IEEE J. Sel. Top. Quantum Electron.* **6**, 528 (2000).

⁶J. C. Knight, G. Cheung, F. Jacques, and T. A. Birks, *Opt. Lett.* **22**, 081306 (1997).

⁷K. Srinivasan, P. E. Barclay, M. Borselli, and O. Painter, *Phys. Rev. B* **70**, 081306(R) (2004).

⁸V. V. Klimov and M. Ducloy, *Phys. Rev. A* **69**, 013812 (2004).

⁹F. L. Kien, S. D. Gupta, V. I. Balykin, and K. Hakuta, *Phys. Rev. A* **72**, 032509 (2005).

¹⁰K. Nayak, P. Melentiev, M. Morinaga, F. Kien, V. Balykin, and K. Hakuta, *Opt. Express* **15**, 5431 (2007).

¹¹G. T. Liu, A. Stintz, H. Li, T. C. Newell, A. L. Gray, P. M. Varangis, K. J. Malloy, and L. F. Lester, *IEEE J. Quantum Electron.* **36**, 1272 (2000).

¹²K. Srinivasan, M. Borselli, T. J. Johnson, P. E. Barclay, O. Painter, A. Stintz, and S. Krishna, *Appl. Phys. Lett.* **86**, 151106 (2005).

¹³K. Srinivasan and O. Painter, *Appl. Phys. Lett.* **90**, 031114 (2007).

¹⁴N. I. Cade, H. Gotoh, H. Kamada, T. Tawara, T. Sogawa, H. Nakano, and H. Okamoto, *Appl. Phys. Lett.* **87**, 172101 (2005).

¹⁵M. Bayer, G. Ortner, O. Stern, A. Kuther, A. A. Gorbunov, A. Forchel, P. Hawrylak, S. Fafard, K. Hinzer, T. L. Reinecke, S. N. Walck, J. P. Reithmaier, F. Klopff, and F. Schäfer, *Phys. Rev. B* **65**, 195315 (2002).

¹⁶G. Bacher, R. Weigand, J. Seufert, V. Kulavoskii, N. Gippius, A. Forchel, K. Leonardi, and D. Hommel, *Phys. Rev. Lett.* **83**, 4417 (1999).

¹⁷C. Santori, D. Fattal, J. Vuckovic, G. Solomon, E. Waks, and Y. Yamamoto, *Phys. Rev. B* **69**, 205324 (2004).

¹⁸S. Strauf, K. Hennessy, M. T. Rakher, Y.-S. Choi, A. Badolato, L. C. Andreani, E. L. Hu, P. M. Petroff, and D. Bouwmeester, *Phys. Rev. Lett.* **96**, 127404 (2006).